



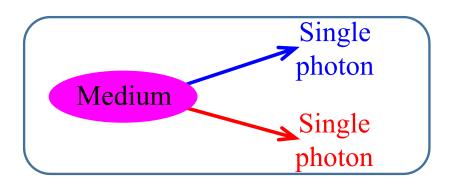
Ite A. Yu

Department of Physics
National Tsing Hua University

Center for Quantum Technology Hsinchu, Taiwan

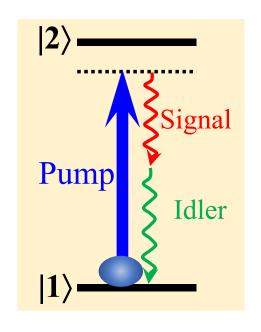
### Introduction

## What are biphotons or heralded single photons? Why are they useful in quantum information processing?



- The **biphoton** is a pair of time-correlated single photons.
- Single photons appear randomly in time. It is difficult to use them in the random timing.
- Biphotons also appear randomly in time. Nevertheless, we can use the first photon of the biphoton to trigger a quantum operation, and employ the second one in the quantum operation. So, the second one is called the heralded single photons.

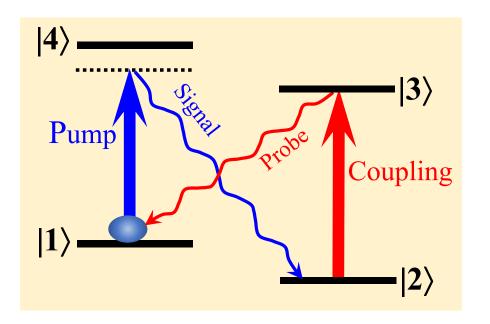
### Mechanisms for Generation of Biphotons — SPDC



### **Spontaneous Parametric Down Conversion**

- SPDC: A pump photon is converted to a signal and an idler photons induced by the vacuum fluctuation.
- Typical media are nonlinear crystals.
- Since 1970.

### Mechanisms for Generation of Biphotons — SFWM

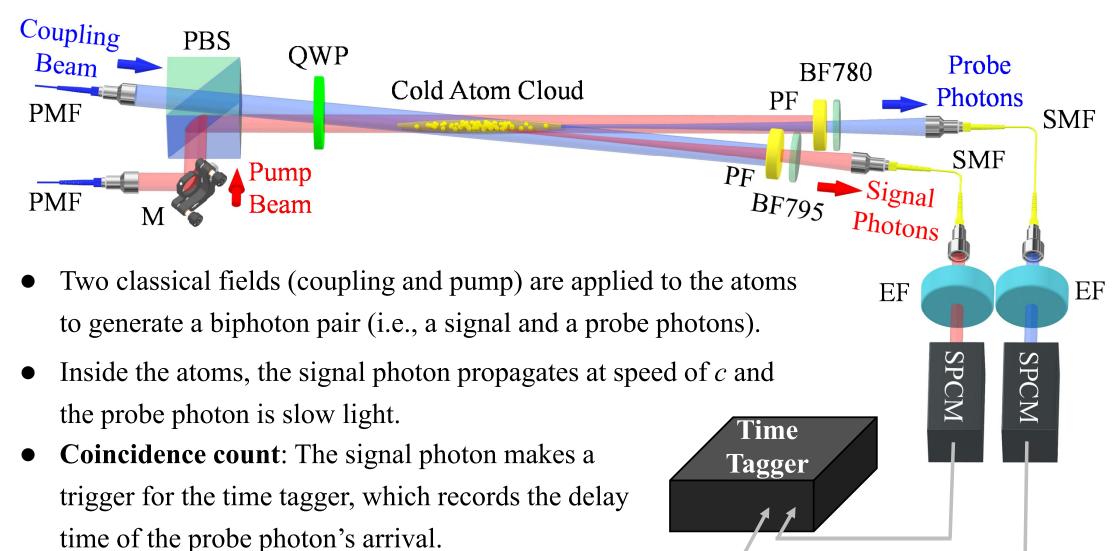


**Spontaneous Four-Wave Mixing** 

- SFWM: The vacuum fluctuation induces a Raman transition to generate the signal photon and also the coherence between states 1 and 2. The coupling field utilizes the coherence to generate the probe photon similar to the electromagnetically induced transparency (EIT) effect.
- The EIT effect makes the probe photon become slow light.
- Typical media: cold atoms since 2005 and room-temperature or hot atomic vapors since 2016.

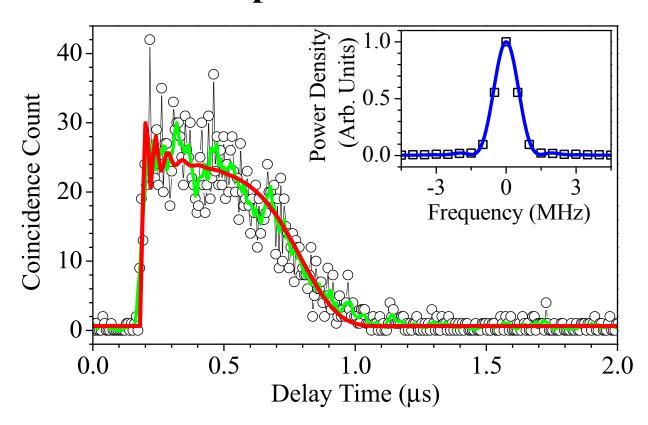
# Illustration of Biphoton Generation and Wave Packet

### **Biphoton Generation from Cold Atoms**



### **Temporal Profile of the Biphoton Wave Packet in Cold Atoms**

Data represent biphoton wave packet  $|\psi|^2$ , where  $\psi$  is the wave function of two-photon correlation.

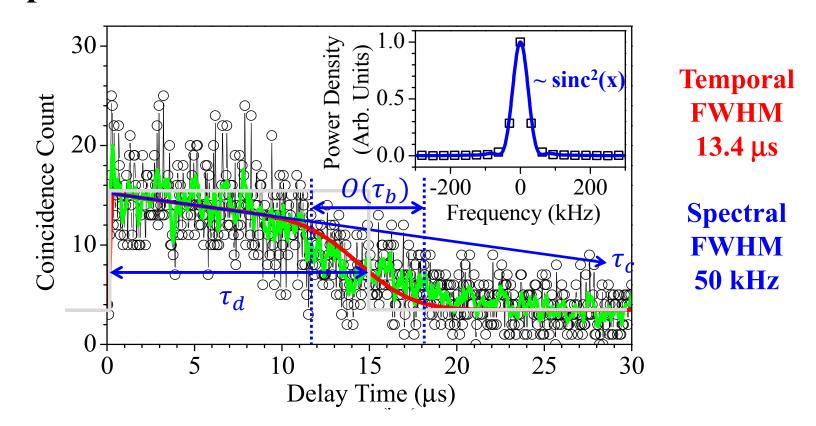


- Temporal FWHM 0.57 µs
- Spectral FWHM 1.2 MHz

- Circles are the two-photon coincidence counts, green line is the result of 4-point moving average of the circles, and red line is the theoretical prediction.
- In the inset, squares and blue line are the Fourier transforms of the data and the prediction.

#### The Longest Biphoton with the Narrowest Linewidth in the World

Y.-S. Wang, K.-B. Li, C.-F. Chang, T.-W. Lin, J.-Q. Li, S.-S. Hsiao, J.-M. Chen, Y.-H. Lai, Y.-C. Chen, Y.-F. Chen, C.-S. Chuu, and I. A. Yu, arXiv:2205.13778.



- A system with a low decoherence rate of  $2\pi \times 1.8$  kHz and a high optical depth of 110 makes an ultralong biphoton wave packet.
- Coherence time  $\tau_c$ : 43 µs; propagation delay time  $\tau_d$ : 16 µs; (EIT bandwidth)<sup>-1</sup>  $\tau_b$ : 2 µs.

### Biphoton Sources with Spectral Profiles of FWHM below 1 MHz

Process	Medium	Type <sup>§</sup>	Temporal FWHM (µs)	Spectral FWHM (kHz)	Reference
SPDC	nonlinear crystal	MM	$0.33^{\dagger}$	670 <sup>‡</sup>	5
SPDC	nonlinear crystal	MM	$0.83^\dagger$	$265^{\ddagger}$	10
SFWM	cold atom cloud	SM	1.7	430	11
SFWM	cold atom cloud	SM	2.1	380	12
SFWM	cold atom cloud	SM	2.9	250	13
SFWM	cold atom cloud	SM	13.4	50	this work
SFWM	hot atomic vapor	SM	0.66	320	22, 24

<sup>§</sup>MM denotes multimode, and SM denotes single mode.

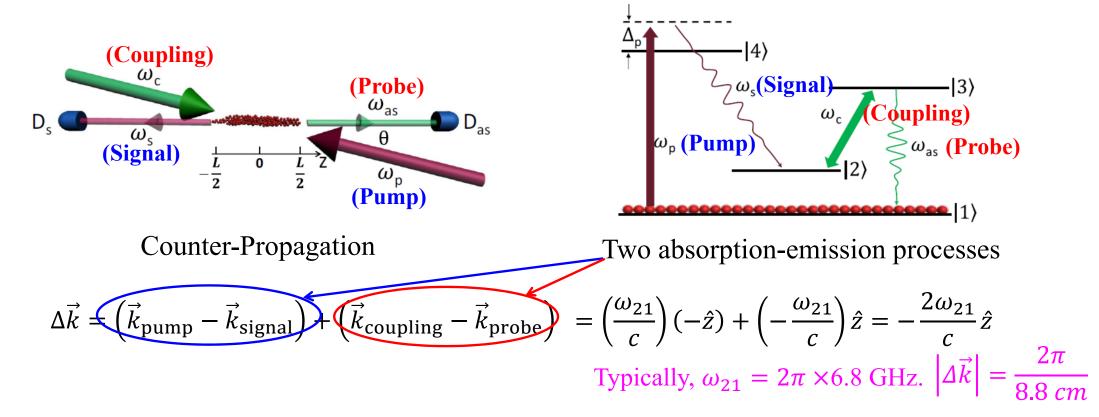
- However, the duty cycle of the biphoton generation with cold atoms is low, e.g., 0.8% in our case and the average generation rate is merely about 30 pairs/s.
- The data in each figure were accumulated by about one hour.

<sup>&</sup>lt;sup>†</sup>The FWHM of the envelope formed by all peaks in a temporally comb-like structure.

<sup>&</sup>lt;sup>‡</sup>The spectral FWHM of the envelope.

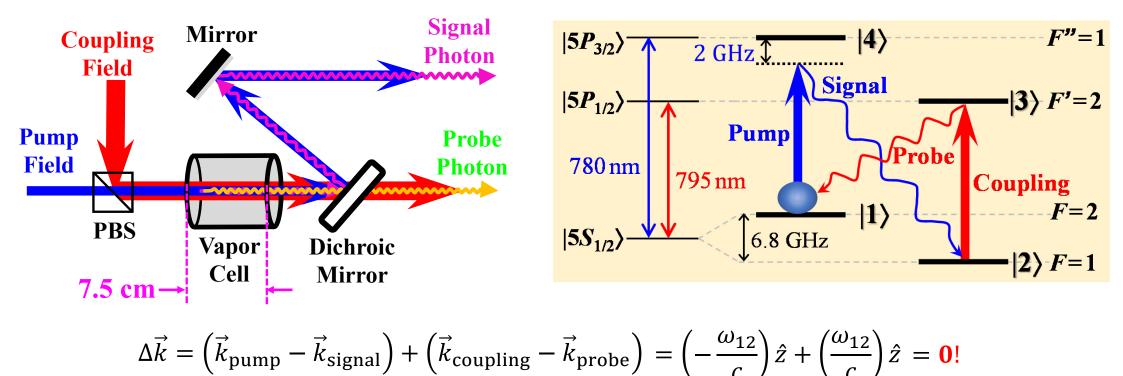
# Features of Our Hot-Atom Biphoton Source

### Phase Mismatch in Others' Biphoton Experiments



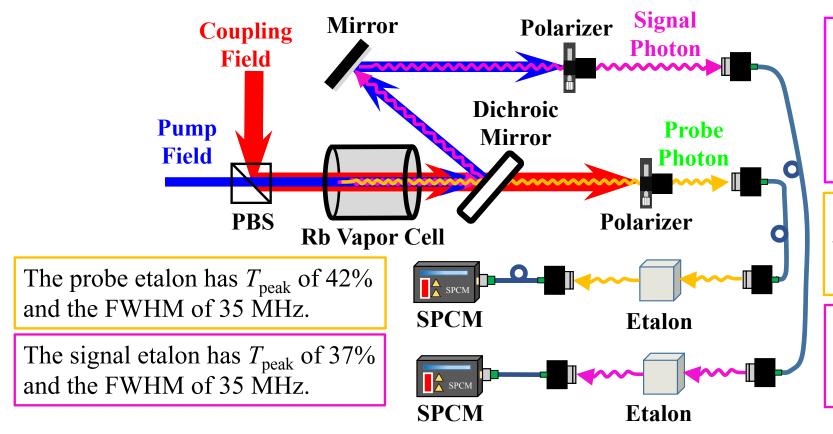
- Previously, SFWM biphoton sources utilized the counter-propagation scheme.
- The degree of phase mismatch is given by  $L \left| \Delta \vec{k} \right|$  (L: the medium length). At L = 7.5 cm, the phase mismatch reduces the generation rate by 1000 folds!

### Phase-Mismatch-Free in Our Biphoton Experiment



- Our biphoton source utilized the all-copropagation scheme.
- The all-copropagation scheme ensures the phase match, and also maintains a low decoherence rate, which enables a narrow linewidth.

### **High Extinction for Laser Light**



The polarization filter provides an extinction ratio (ER) of 60 (48) dB to block the pump (coupling) field.

The probe etalon blocks the coupling field with an ER of 88 dB.

The signal etalon blocks the pump field with an ER of 74 dB.

- Laser light of a power up to 40 mW, and the single-photon pulse of a power as low as 0.4 pW.
- Fortunately, an overall ER of  $\sim 135$  dB to block the pump and coupling fields.

# A Narrow-Linewidth, High-Generation Rate, High-Spectral Brightness Biphoton Source

# Biphoton Wave Packet and EIT Spectrum of Hot Atoms

S.-S. Hsiao, W.-K. Huang, Y.-M. Lin, J.-M. Chen, C.-Y. Hsu, and I. A. Yu, Phys. Rev. A 106, 023709 (2022).

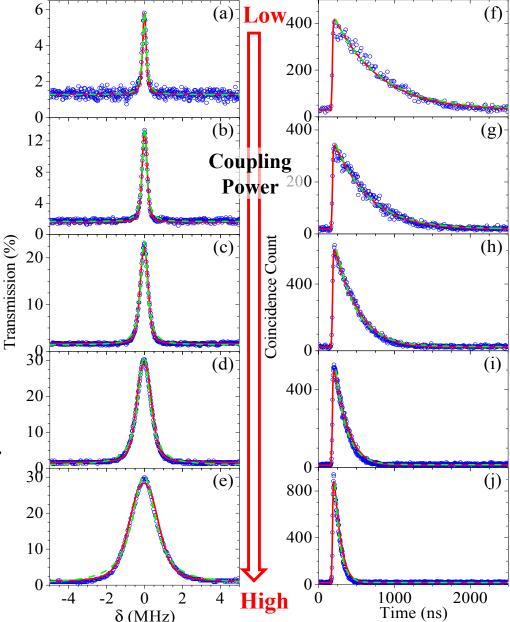
In hot atoms, the biphoton wave packet is mainly

determined by the EIT spectrum; in cold atoms,

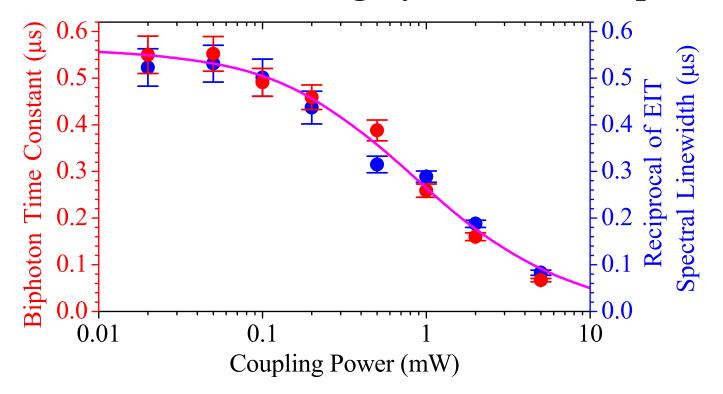
it is strongly influenced by the propagation delay time as shown by the earlier slide.

A higher coupling power makes a broader EIT

- linewidth & a narrower biphoton temporal width.
- The EIT spectral profile is a Lorentz function.
- The biphoton temporal profile is an exponential-decay function.



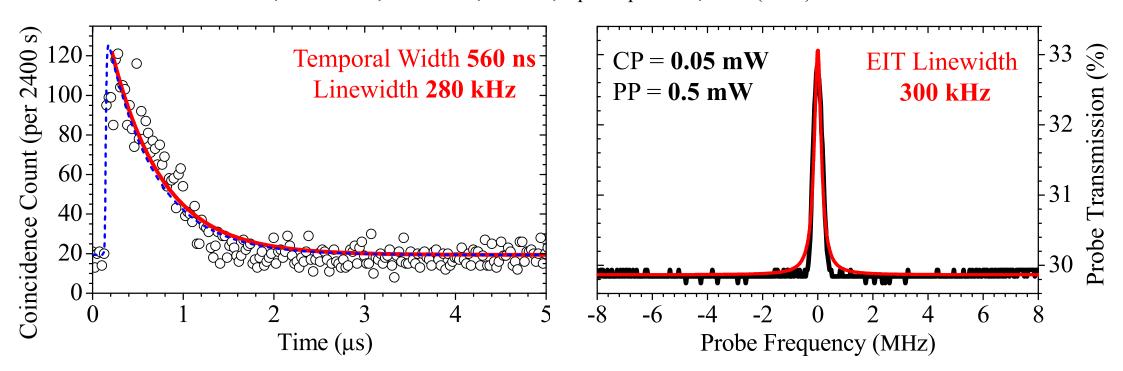
### A Biphoton Source with a Highly Tunable Temporal Width



- Data of biphoton temporal widths (red) are consistent with results calculated from those of EIT linewidths (blue), and also in agreement with the theoretical predictions (line).
- The temporal width or spectral linewidth of biphotons can be tuned by about one order of magnitude (from 60 to 560 ns). A higher laser power can make the tuning range larger.

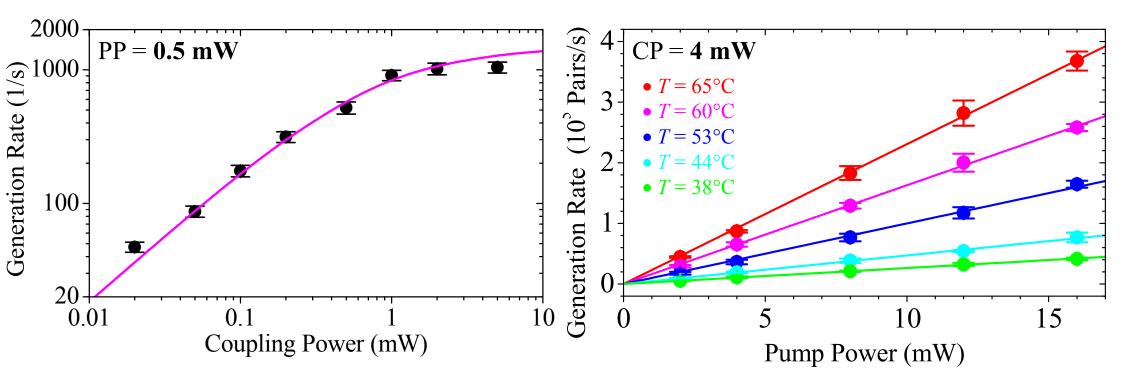
### The Narrowest-Linewidth Biphotons Generated from Hot Atoms [18] in the World

C.-Y. Hsu, Y.-S. Wang, J.-M. Chen, F.-C. Huang, Y.-T. Ke, E. K. Huang, W. Hung, K.-L. Chao, S.-S. Hsiao, Y.-H. Chen, C.-S. Chuu, Y.-C. Chen, Y.-F. Chen, I. A. Yu, Opt. Express 29, 4632 (2021). Editors' Pick.



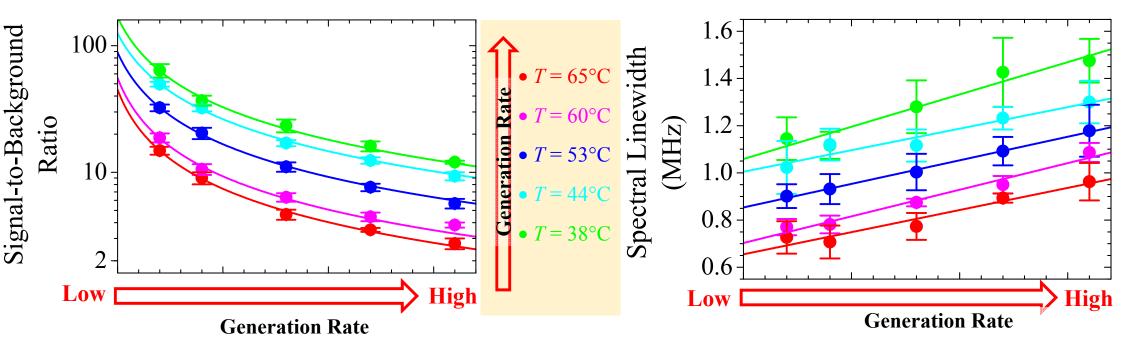
- Biphoton wave packet (left) and EIT spectrum (right) were measured at the same condition.
- The decoherence rate in the experimental system limits the narrowest linewidth.

### **Generation Rate of the Biphoton Source (1/2)**



- At large coupling powers, the generation rate saturates.
- The generate rate is linearly proportional to the pump power, and increases with the atom temperature, i.e., the optical depth.
- Experimental data are consistent with theoretical predictions.

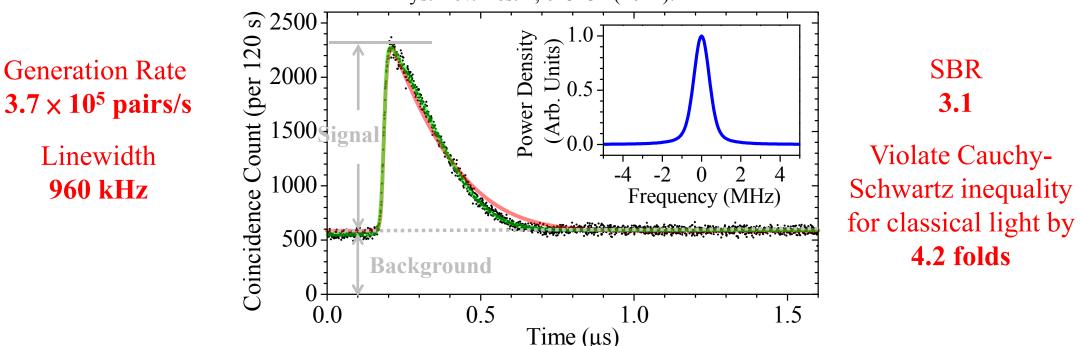
### **Generation Rate of the Biphoton Source (2/2)**



- A larger the signal-to-background ratio (SBR) makes a higher purity of the single photon. However, as the generate increases, the SBR decreases.
- A narrower biphoton's linewidth makes a higher success probability of a quantum process, e.g. quantum memory. However, as the generate increases, the linewidth becomes broadened.

### The Highest Spectral Brightness in the World

J.-M. Chen, C.-Y. Hsu, W.-K. Huang, S.-S. Hsiao, F.-C. Huang, Y.-H. Chen, C.-S. Chuu, Y.-C. Chen, Y.-F. Chen, and I. A. Yu, Phys. Rev. Res. 4, 023132 (2022).



- The spectral brightness, i.e., generation rate per linewidth, is the measure of success rate of a quantum information process.
- The high generation rate, together with the narrow linewidth, results in a spectral brightness of  $3.8 \times 10^5$  pairs/s/MHz, better than all known results with all kinds of media.

### Comparison between Different Kinds of Biphoton Sources

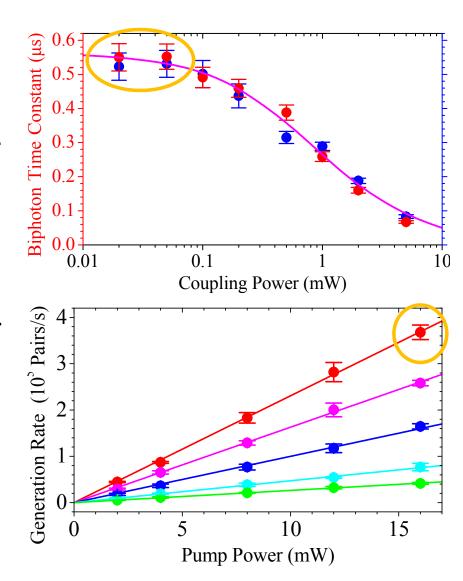
	and the second s		1	Notes
DC 3 MH	7 1	N A	a few GHz	
OC 265 kF	4,300 hz <sup>[2]</sup> pairs/s/MH	$\mathbf{z}^{[6]}$ N.A.	a few GHz	The values refer to one of the frequency modes.
M 50 kH	<b>/</b>   11   ``		f N.A.	Duty cycle ≤ 10%.
es 92 MH	$Iz^{[3]}$ 1.4×10 <sup>5</sup> pairs/s/MH	$\mathbf{z}^{[3]}$ N.A.	160 MHz <sup>[3]</sup>	Micro-ring resonator <sup>[3]</sup> with $Q$ of $\sim 10^6$ .
2 MH	7 4	гот	f 600 MHz	The frequency tunability is determined by width of the Doppler broadening.
			(Rb atoms)	
	Linewicks  Linewicks  3 MH  265 kH  M 50 kH  92 MH  2 MH	Linewidth         Brightness           DC         3 MHz <sup>[1]</sup> 3.5×10 <sup>5</sup> pairs/s/MH           DC         265 kHz <sup>[2]</sup> 4,300 pairs/s/MH           VM         50 kHz <sup>[11]</sup> 4,700(×10 <sup>9</sup> pairs/s/MH           ees         92 MHz <sup>[3]</sup> 1.4×10 <sup>5</sup> pairs/s/MH           ees         2 MHz <sup>[4]</sup> 1.4×10 <sup>4</sup> pairs/s/MH           ese         290 kHz <sup>[9]</sup> 3.8×10 <sup>5</sup>	Linewidth Brightness Tunability $3.5 \times 10^5$ $pairs/s/MHz^{[5]}$ N.A. $2.6 \times 10^5$ $pairs/s/MHz^{[5]}$ N.A. $2.6 \times 10^5$ $pairs/s/MHz^{[6]}$ N.A. $2.6 \times 10^5$ $pairs/s/MHz^{[6]}$ N.A. $2.6 \times 10^5$ $pairs/s/MHz^{[7]}$ N.A. $2.6 \times 10^5$ $pairs/s/MHz^{[3]}$ N.A. $2.6 \times 10^5$ $pairs/s/MHz^{[8]}$ N.A. $2.6 \times 10^5$ $pairs/s/MHz^{[8]}$ N.A. $2.6 \times 10^5$ $pairs/s/MHz^{[8]}$ One order of magnitude or	Linewidth Brightness Tunability Tunability  DC $3 \text{ MHz}^{[1]}$ $3.5 \times 10^5 \text{ pairs/s/MHz}^{[5]}$ N.A. a few GHz  DC $265 \text{ kHz}^{[2]}$ $4,300 \text{ pairs/s/MHz}^{[6]}$ N.A. a few GHz  MM $50 \text{ kHz}^{[11]}$ $4,700(\times 10\%) \text{ pairs/s/MHz}^{[7]}$ one order of magnitude  PS $92 \text{ MHz}^{[3]}$ $1.4 \times 10^5 \text{ pairs/s/MHz}^{[3]}$ N.A. $160 \text{ MHz}^{[3]}$ The rightness ier $2 \text{ MHz}^{[4]}$ $1.4 \times 10^4 \text{ pairs/s/MHz}^{[8]}$ one order of magnitude  Tunability Tunability  N.A. a few GHz  N.A. one order of magnitude  Tunability  N.A. a few GHz  N.A. one order of magnitude  Tunability  N.A. a few GHz  N.A. one order of magnitude  Tunability  N.A. a few GHz  N.A. one order of magnitude  Tunability

- [1] New J. Phys. 18, 123013 (2018).
- [2] APL Photon. 5, 066105 (2020).
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- [4] Nat. Commun. 7, 12783 (2016).
- [5] Phys. Rev. A 92, 063827 (2015). [6] APL Photon. 4, 090804 (2019).
  - [7] Optica 1, 84 (2014).
  - [8] Appl. Phys. Lett. 110, 161101 (2017).
- [9] Opt. Express 29, 4632 (2021). [10] Phys. Rev. Res. 4, 023132 (2022).
- [11] arXiv:2205.13778.

### Conclusion

#### **Conclusion**

- The linewidth of our hot-atom biphotons can be as narrow as 290 kHz, which is the narrowest among all kinds of single-mode biphotons generated from room-temperature or hot media.
- The spectral brightness of our hot-atom biphoton source can be as high as 3.7×10<sup>5</sup> pairs/s/MHz, which the highest spectral brightness of all biphoton sources.
- Our biphoton source not only surpasses the sources produced with the hot atoms in the previous works, but also competes with the sources produced with the nonlinear crystals and integrated photonic devices.



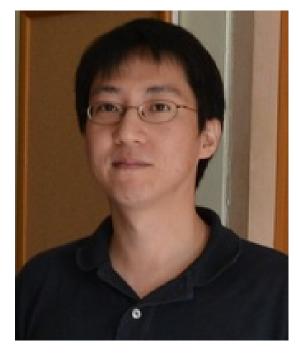
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## Thank you for your attention



