



First lecture:

Precision measurement in atomic physics
---selected topics

Second lecture:

Laser spectroscopy of simple atoms
---work performed at NTHU

王立邦 Li-Bang Wang

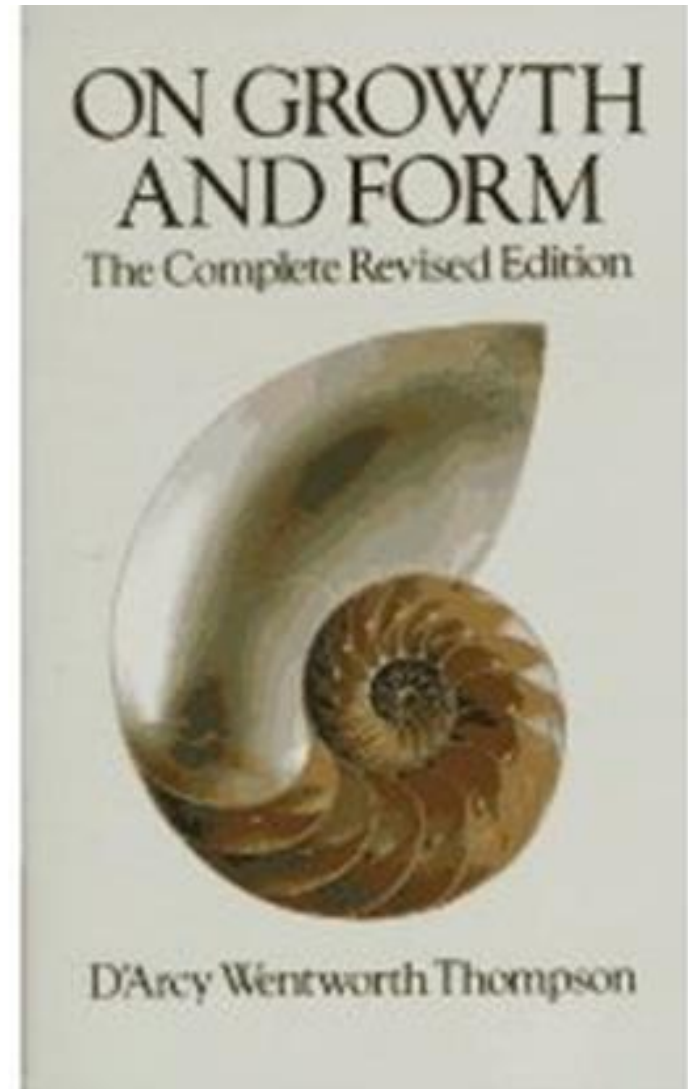
清華大學物理系 NTHU Physics



Thompson, D'Arcy Wentworth

On Growth and Form, 1917

...numerical precision is the
very soul of science...



How precise?



- Magnetic moment of electron,

$$g_e (\text{exp}) = 2.00231930436146(56)$$

- Rydberg constant = $\frac{2\pi^2 m_e e^4}{h^3 c}$
= 109,737.31568160(21)

- Electric dipole moment of electron

$$|d_e| < 1.1 \times 10^{-29} \text{ e} \cdot \text{cm}$$

- Time variation of fine structure constant

$$\delta\alpha/\alpha = (-1.6 \pm 2.3) \times 10^{-17} / \text{year} \quad \alpha = e^2 / \hbar c \approx 1/137$$

[1] D. Hanneke *et al.*, *Phys. Rev. Lett.* **100**, 120801 (2008)

[2] 2018 CODATA value (2019)

[3] The ACME Collaboration, *Nature* **562**, 355 (2018)

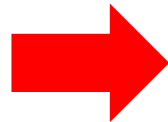
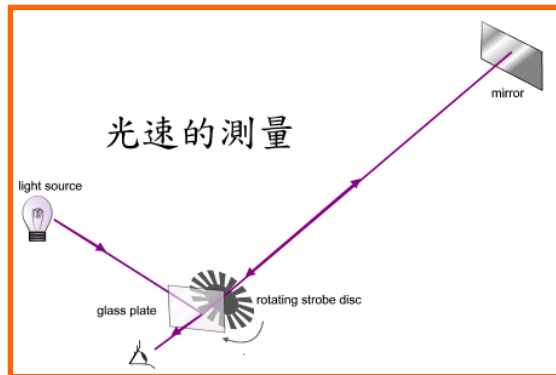
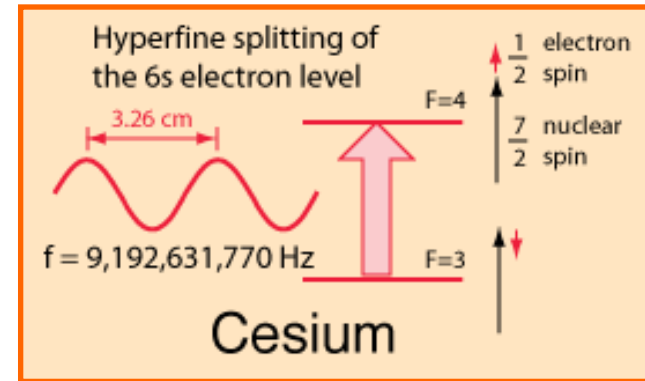
[4] T. Rosenband *et al.*, *Science* **319**, 1808 (2008)

Classical measurements

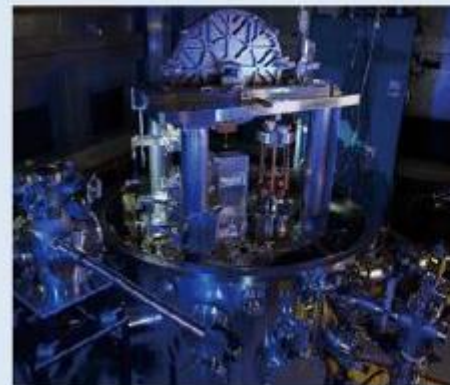


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- Time, Length, Mass, Charge
- Definition of Unit
- Good Tools



Kibble balance



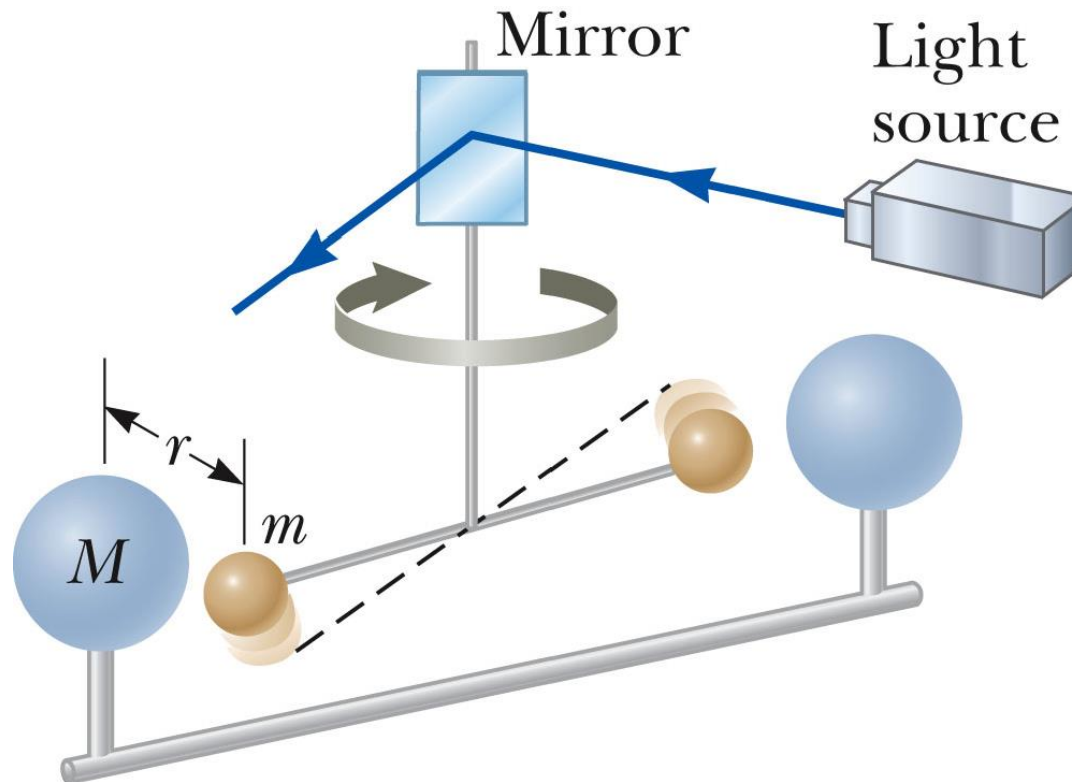
Silicon sphere



Example



- Cavendish 1798, $G = 6.74 \times 10^{-11}$
- $G = (6.674 \pm 0.001) \times 10^{-11}$, $\delta G/G \sim 100$ ppm

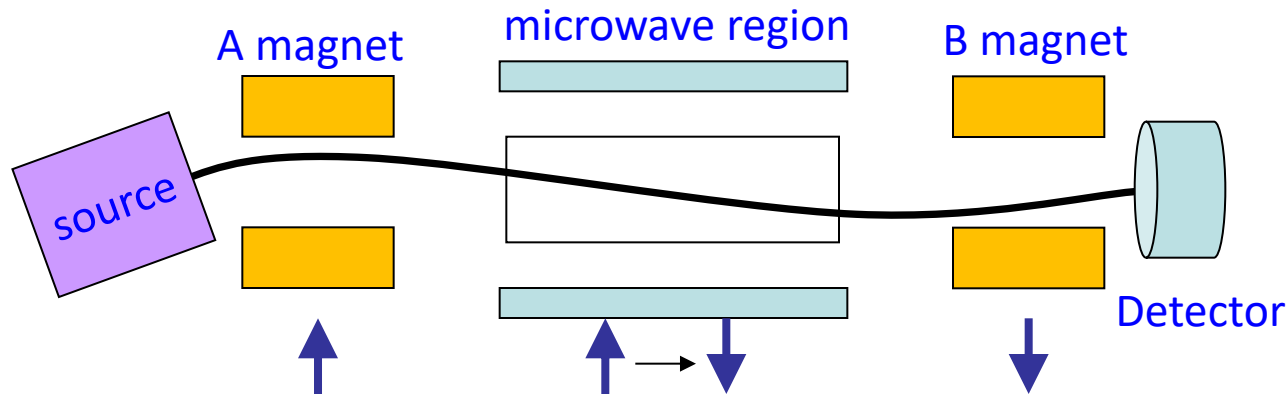


- Historical Review:
 - Precision measurement lead to new physics
- Modern approach: resonance and frequency measurement
- Selected topics

Resonance Method



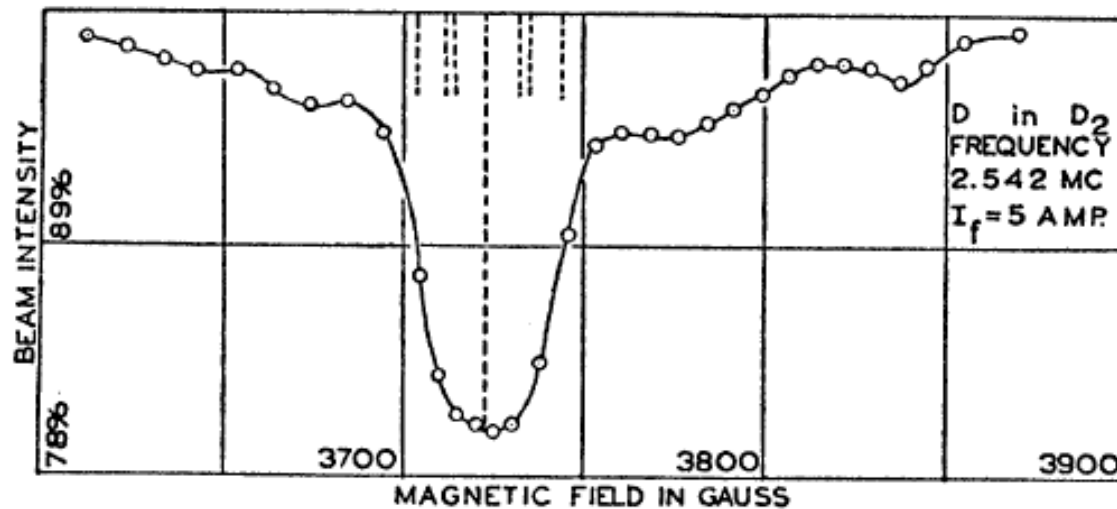
- First Motivated by Stern-Gerlach experiment, 1922
 - First atomic spectroscopy by **resonance** method
- NMR, Rabi and Ramsey, 1934-1939



Isidor Isaac Rabi



Norman F. Ramsey

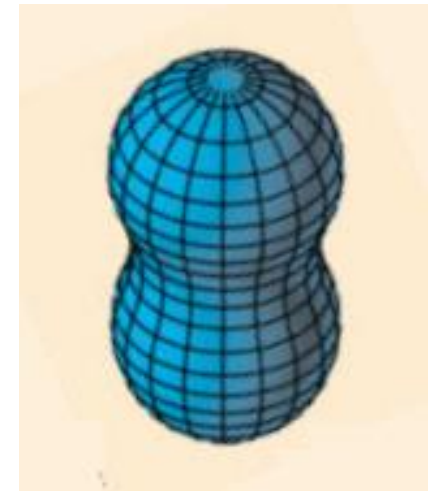
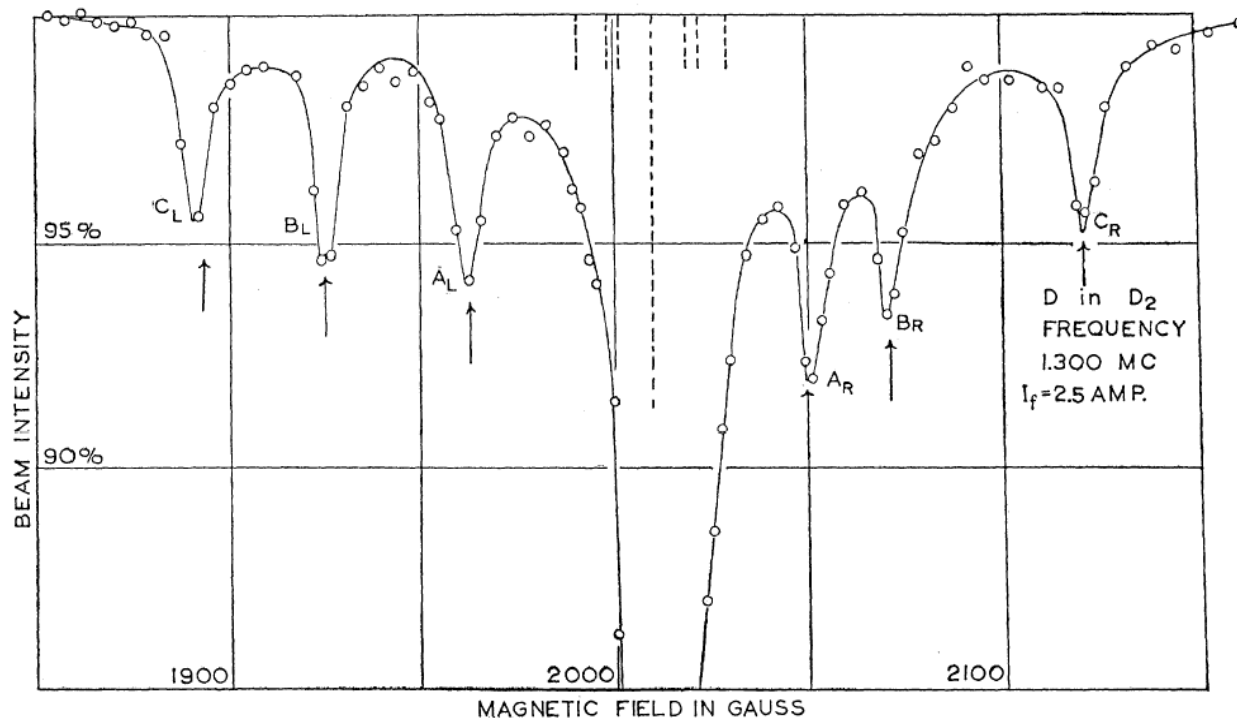


NMR for Deuteron

Electric Quadrupole Moment

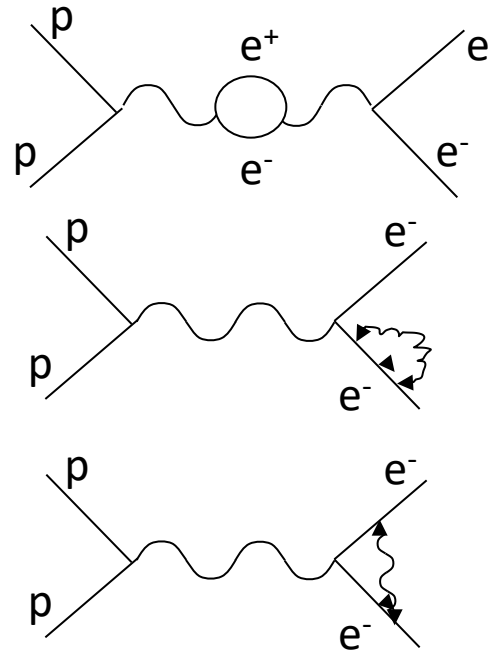
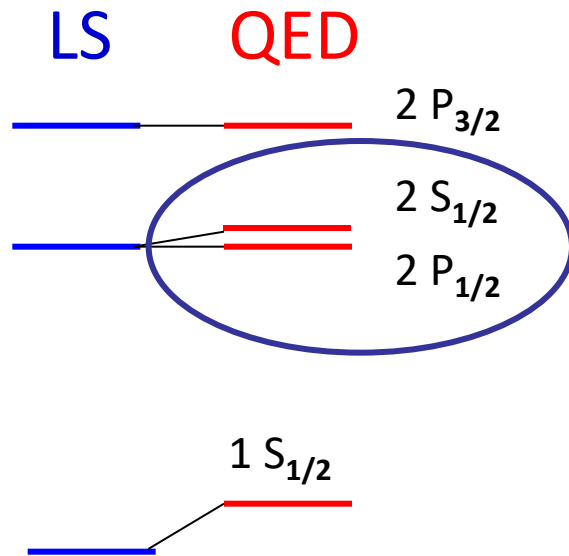


- Electric Quadrupole moment of Deuteron discovered
- Nuclear force is **Non-central!** Tensor force



J. M. Kellogg, I. I. Rabi, N. F. Ramsey, and J. R. Zacharias Phys. Rev. 57, 677 (1940)

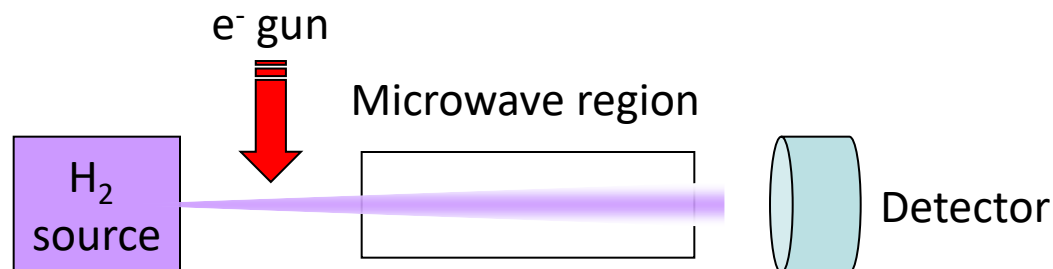
The Lamb Shift



Willis Eugene Lamb
Nobel Prize in 1955

"for his discoveries concerning
the fine structure of the
hydrogen spectrum"

1947 by Lamb ~ 1060 MHz
now: $1057.846(4)$ MHz



- Quantum Electrodynamics
- The most precisely tested theory



**Sin-ichiro Tomonaga (朝永振一郎), Julian Schwinger, Richard P. Feynman
Nobel Prize in Physics 1965**

"for their fundamental work in quantum electrodynamics, with deep-ploughing consequences for the physics of elementary particles"

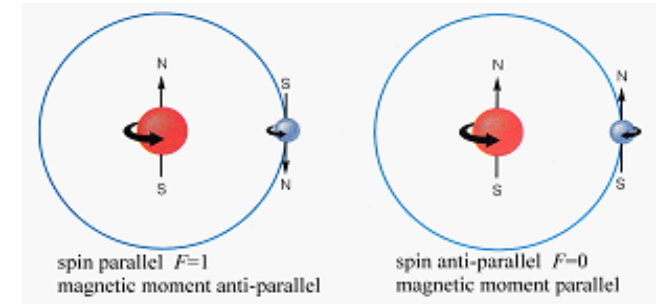


- Hyperfine structure: probing nuclear moment
- Test QED: the $g-2$ of electron and muon
- Strong interaction: size of proton
- Time and frequency standard: making a better clock
- Constancy of fundamental constant: can speed of light change?
- New source of CP violation: electric dipole moment of electron and nucleon

Hyperfine structure

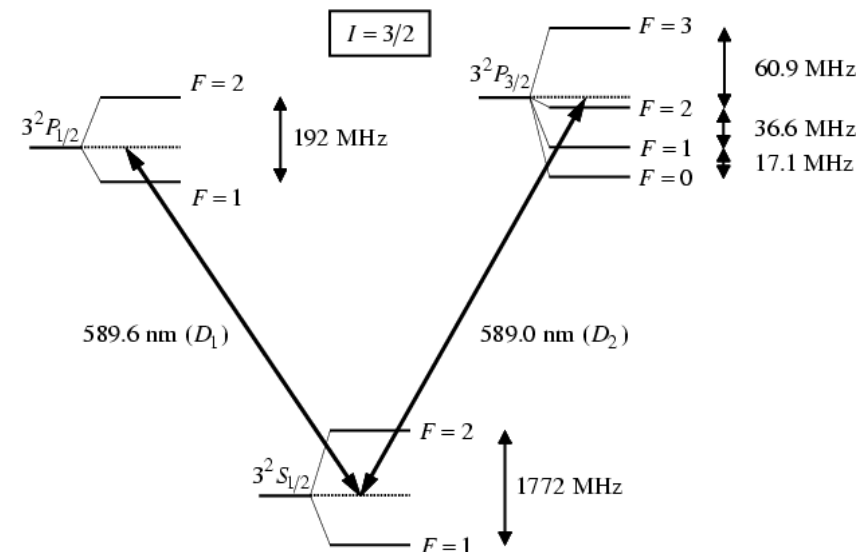


- Nuclear magnetic moment
- $F = I + J$
- a and b constant



$$\Delta E_{HFS} = \frac{1}{2} aK + \frac{1}{4} b \frac{\frac{3}{2} K(K+1) - 2I(I+1)J(J+1)}{I(2I-1)J(2J-1)}$$

$$K = F(F+1) - J(J+1) - I(I+1)$$

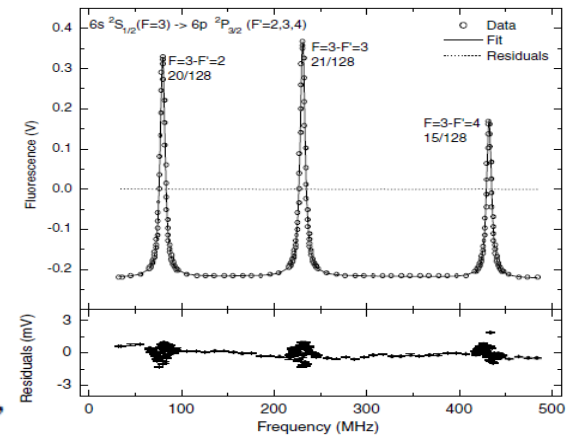
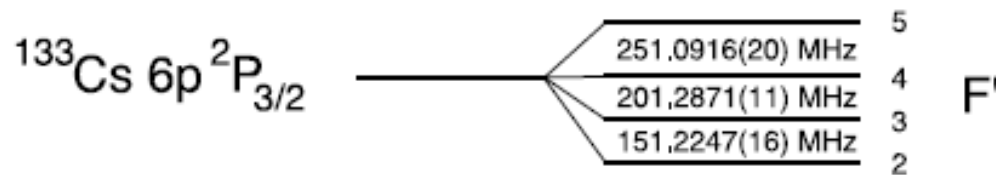


Magnetic octupole moment



● Observe magnetic octupole moment of ^{133}Cs

V. Gerginov, A. Derevianko, and C. E. Tanner, PRL 91, 072501 (2003)



$$H_{\text{dipole}} = aI \cdot J, \quad H_{\text{quadrupole}} = b \frac{3(I \cdot J)^2 + \frac{3}{2}(I \cdot J) - I(I+1)J(J+1)}{2I(2I-1)J(2J-1)},$$

$$H_{\text{octupole}} = c \frac{\{10(I \cdot J)^3 + 20(I \cdot J)^2 + 2(I \cdot J)[-3I(I+1)J(J+1) + I(I+1) + J(J+1) + 3] - 5I(I+1)J(J+1)\}}{I(I-1)(2I-1)J(J-1)(2J-1)}$$

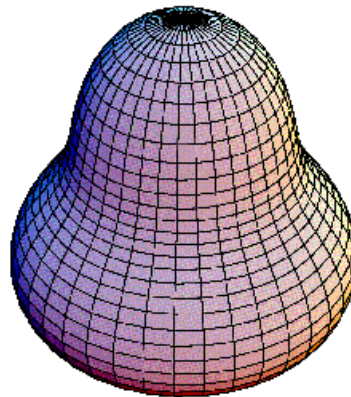
$$\Delta\nu_{54} = 5a + \frac{5}{7}b + \frac{40}{7}c + (-0.000520 \text{ MHz}),$$

$$\Delta\nu_{43} = 4a - \frac{2}{7}b - \frac{88}{7}c + (+0.000119 \text{ MHz}),$$

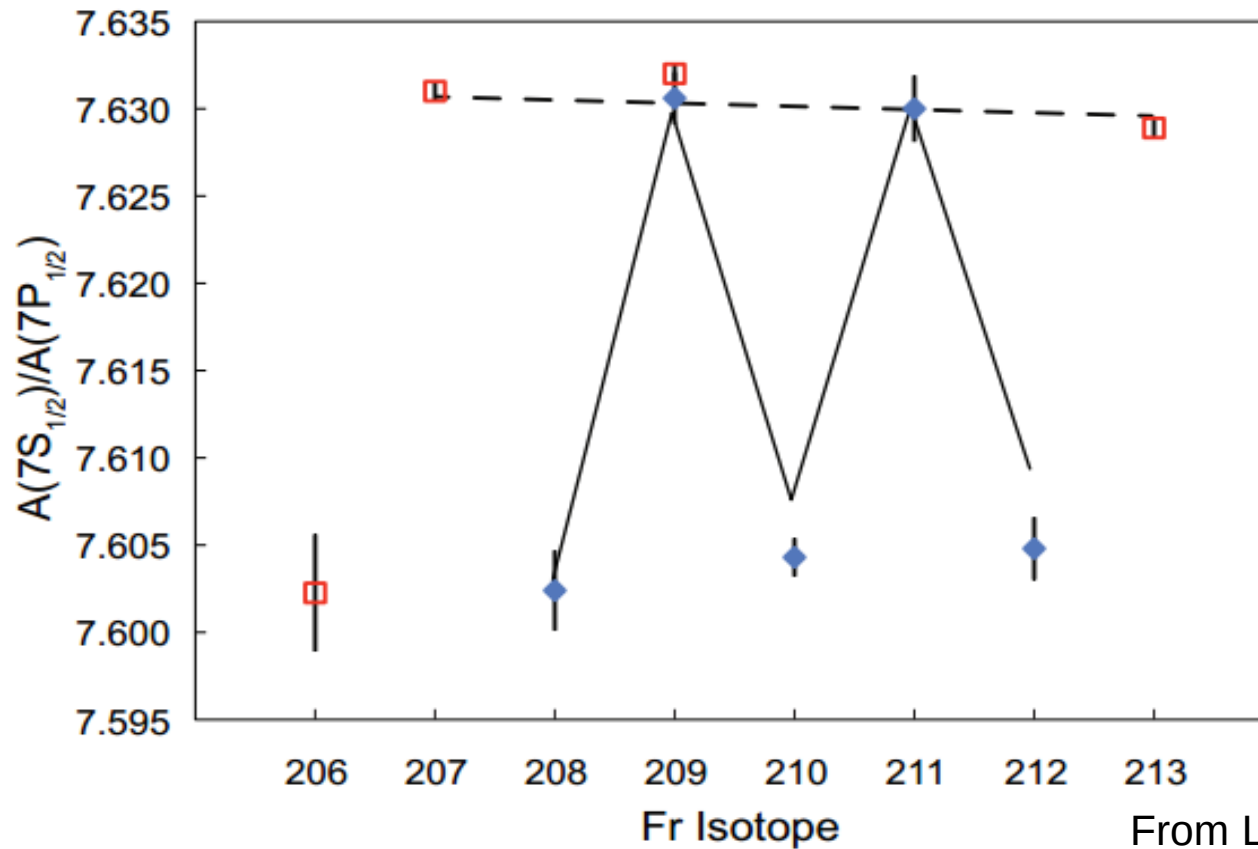
$$\Delta\nu_{32} = 3a - \frac{5}{7}b + \frac{88}{7}c + (+0.000401 \text{ MHz}),$$

a (MHz)	b (MHz)	c (kHz)
50.288 27(23)	-0.4934(17)	0.56(7)

- Observe magnetic octupole moment of ^{87}Rb from spectroscopy of hyperfine intervals
V Gerginov, C E Tanner, W R Johnson, Can. J. Phys., 87(1): 101-104 (2009)
- Large deformation for some heavy nuclei
- Comparison to nuclear shell model



Hyperfine anomaly



From L. Orozco, UMD

Blue points: Grossman et al., PRL **83**, 935–938 (1999)

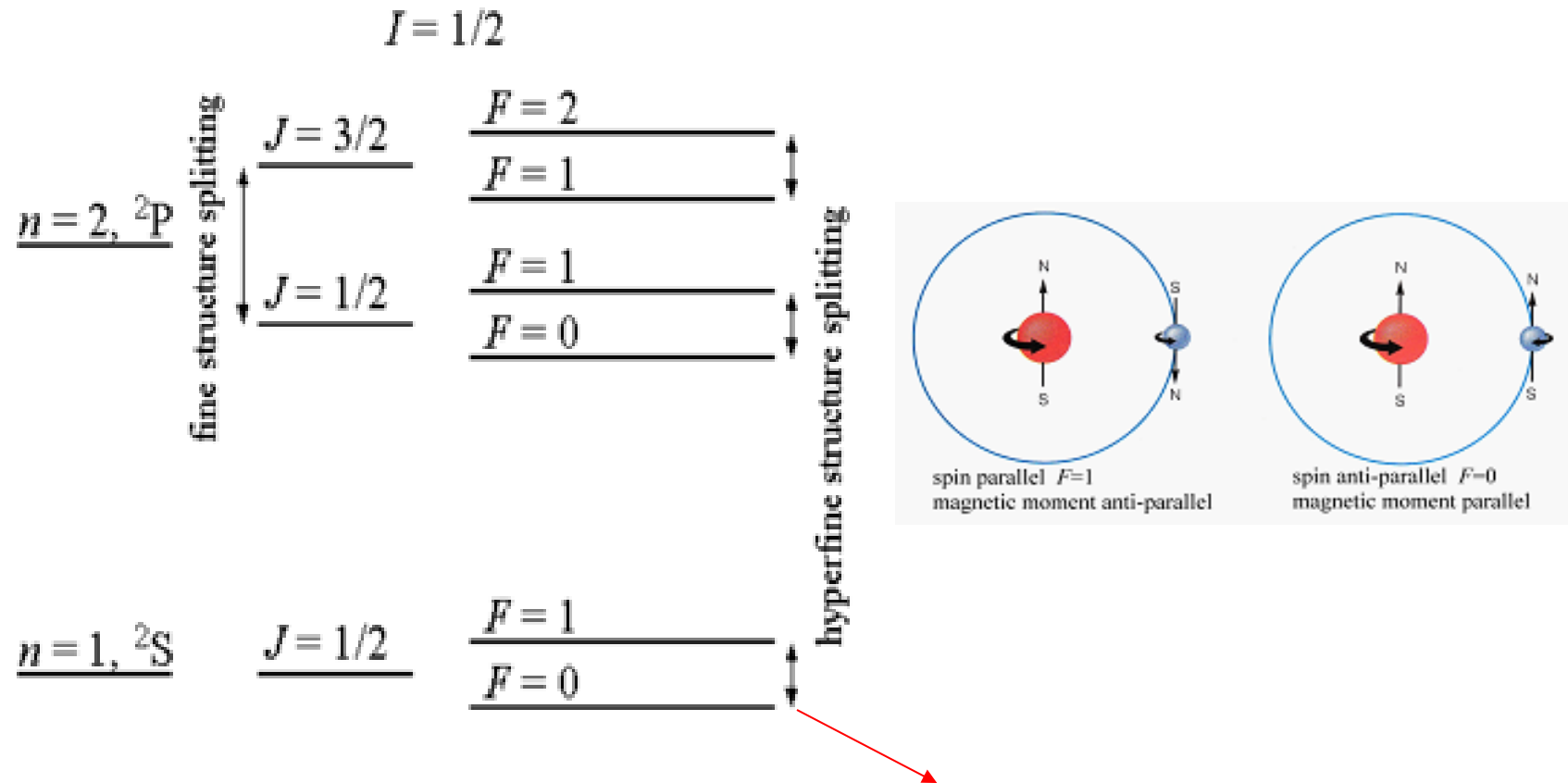
Red points: new measurements at TRIUMF (error as of May 30, 2013)

Dashed line: magnetization radius = charge radius, normalized to ^{211}Fr

^{213}Fr ground state: Duong, et al. Europhys. Lett., **3**(2), 175-182 (1987)

^{206}Fr ground state: Voss, et al. To be published.

- Measure hyperfine constant of isotopes
- Magnetic radius $\neq$ electric radius



$E_{\text{HFS}} = 1420.4057517667(9) \text{ MHz} = 21 \text{ cm}$

$$E_F^N = \frac{8}{3\pi} \alpha^3 \mu_B \mu_N \frac{m_e^3 m_N^3}{(m_N + m_e)^3}$$

$$E_{HFS} = (1 + \Delta_{QED} + \Delta_R + \Delta_S) E_F$$

calculable (pointing to Δ_{QED} and Δ_R)
 unknown (pointing to Δ_S)

where $E_F = \frac{8}{3\pi} \alpha^3 \mu_B \mu_p \frac{m_e^3 m_p^3}{(m_e + m_p)^3} = \text{known}$

$$\Delta_S = \Delta_Z + \Delta_{pol}$$

$$= -2\alpha m_e \langle r \rangle_Z (1 + \delta_Z^{rad}) + \Delta_{pol}$$

calculable (pointing to δ_Z^{rad} and Δ_{pol})

$$\langle r \rangle_Z = -\frac{4}{\pi} \int_0^\infty \frac{dQ}{Q^2} \left[\frac{2G_E G_M}{g} - 1 \right]$$

Zemach moment,
to be precisely determined

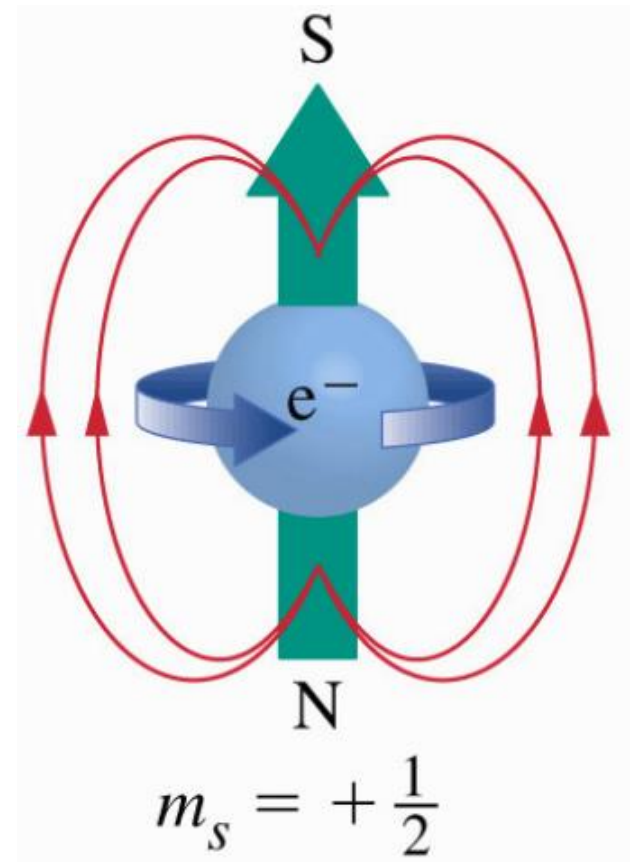
What is g-factor?

The magnetic moment of the electron is proportional to the electron spin

$$\frac{|\bar{\mu}|}{\mu_B} = g \cdot \frac{|\bar{s}|}{\hbar}$$

- Classical non-relativistic $g=1$
- In QM, Dirac Theory predict $g=2$

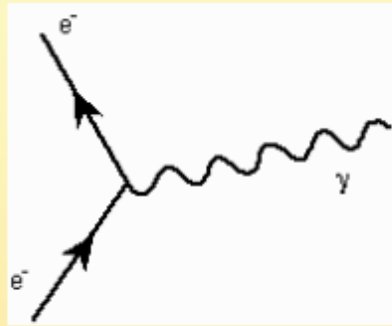
In reality, $g= 2.002319304$



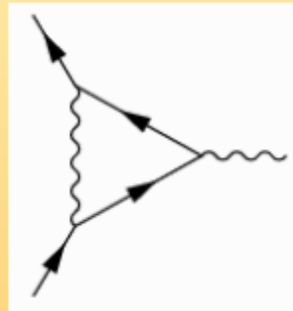
● Vacuum is not real vacuum

and a lot more....

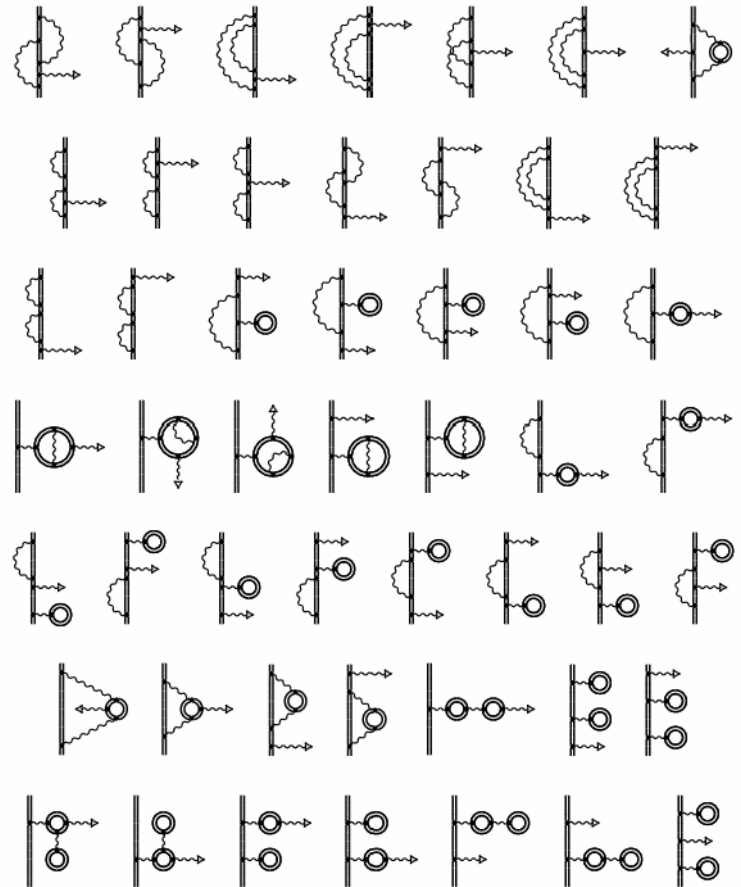
0th order:



1st order:



+ ...



- Test QED
- Search for structure of electrons
- g can be expanded as a function of α

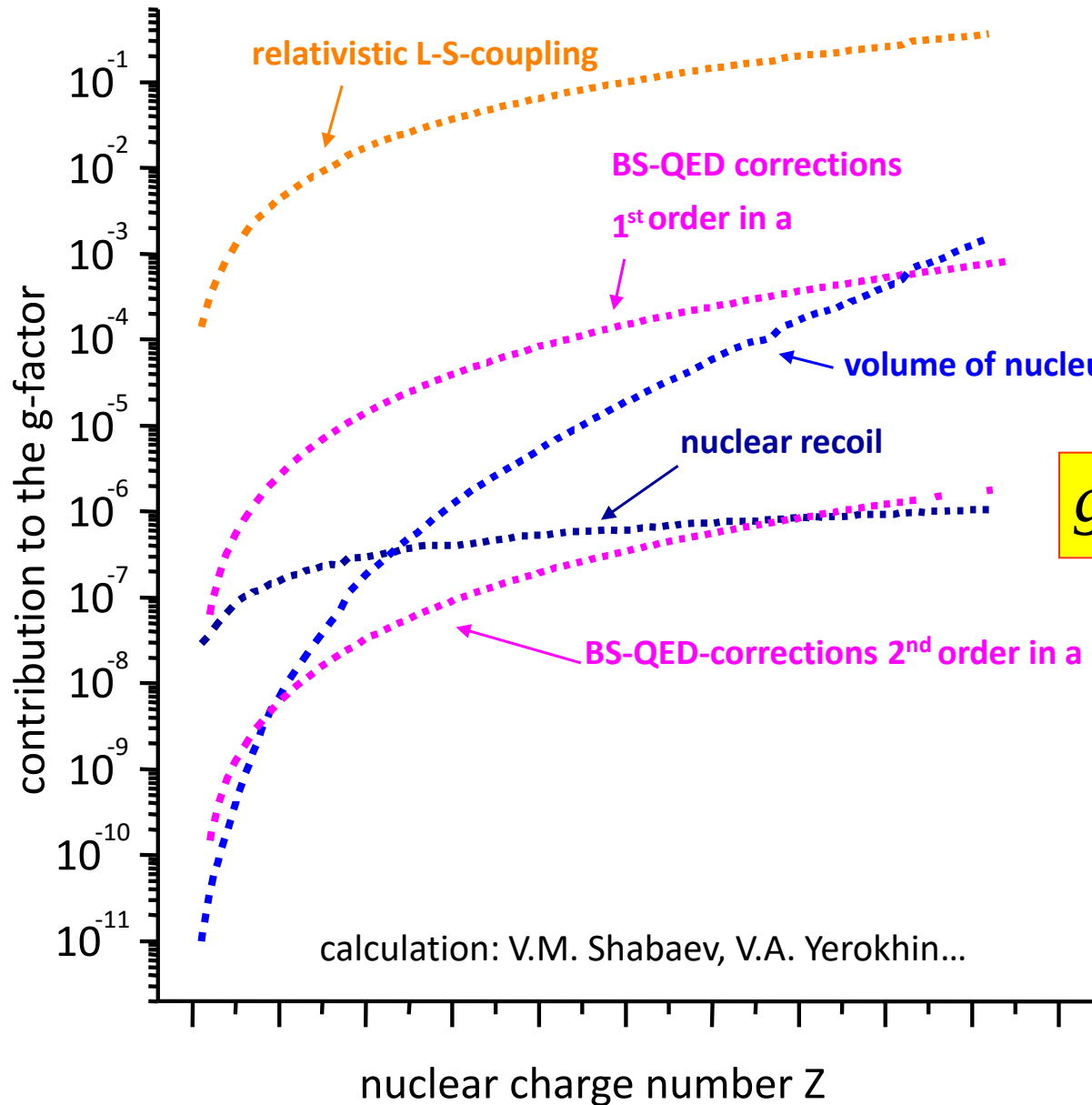
$$\frac{g}{2} = 1 + C_2 \left(\frac{\alpha}{\pi} \right) + C_4 \left(\frac{\alpha}{\pi} \right)^2 + C_6 \left(\frac{\alpha}{\pi} \right)^3 + \dots$$
$$+ a_{\mu\tau} + a_{hadronic} + a_{weak}$$

$$\alpha = e^2 / \hbar c \approx 1/137$$

$$\alpha^{-1} = 137.035\,999\,074\,(44)$$

D. Hanneke, S. Fogwell, and G. Gabrielse, *PRL* **100**, 120801 (2008)

Measure electron in bound system



QED of free electron
(exp. Dehmelt et al., Gabrielse
et al. theo. Kinoshita)

Dirac value

Nuclear
correction

$$g_J = 2 + a_e + a_b + b + c$$

Relativistic
correction

QED of bound
electron

Spin precession (Larmor) frequency

$$\hbar\omega_L = m_s g \cdot \mu_B \cdot B$$

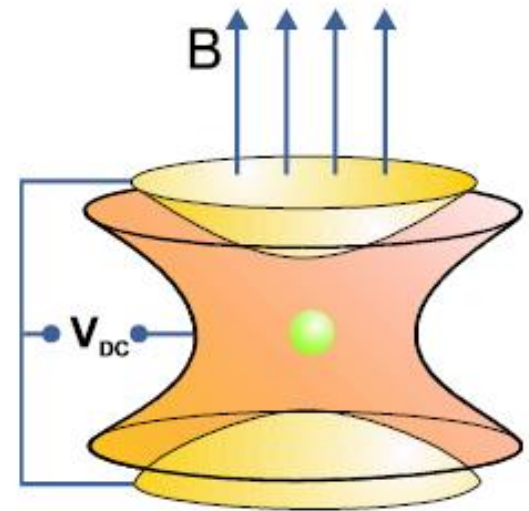
$$\mu_B = \frac{e\hbar}{2m}$$

Calibration of magnetic field by cyclotron frequency:

$$\hbar\omega_C = \frac{q}{M} B$$

$$g = 2 \frac{\omega_L}{\omega_C} \frac{q}{e} \frac{m}{M} = 2 \frac{\omega_L}{\omega_C}$$

Measurement performed on a single electron in Penning trap



$g-2$ experiment of electron and muon:

$$g_e(\text{exp}) = 2.00231930436146(56)$$

$g_e(\text{th})$ to determine fundamental constant

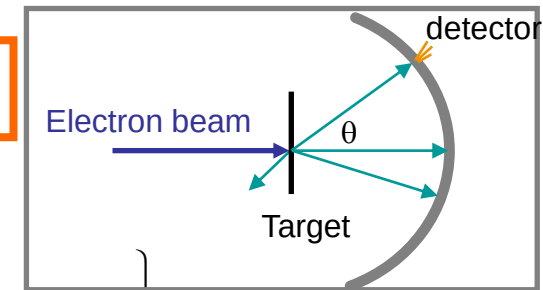
$$\left. \begin{array}{l} g_\mu(\text{exp}) = 2.0023318416(12)^* \\ g_\mu(\text{th}) = 2.0023318367(13)^* \end{array} \right\} 5 \sigma \text{ deviation}$$

G.W. Bennett et al., *Phys Rev Lett* 92, 161802 (2004)

Muon $g-2$ at Fermilab data taken

- Fundamental quantity bothering for decades
- Very important for nuclear calculation
- Lattice QCD **not** able to calculate precisely

Charge radius of proton by electron scattering

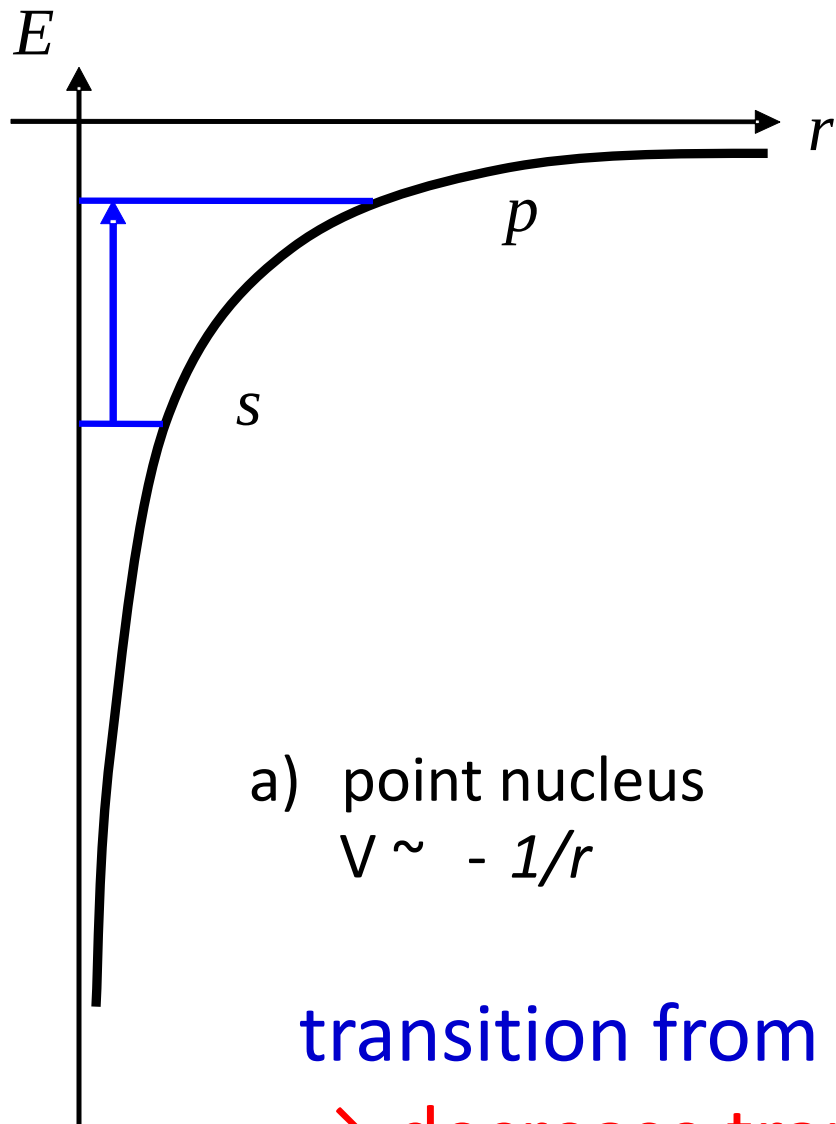


$$\left(\frac{d\sigma}{d\Omega}\right)_{\text{Rosenbluth}} = \left(\frac{d\sigma}{d\Omega}\right)_{\text{Mott}} \left\{ \frac{G_E^2(Q^2) + \frac{Q^2}{4M^2} G_M^2(Q^2)}{1 + \frac{Q^2}{4M^2}} + \frac{Q^2}{2M^2} \cdot G_M^2(Q^2) \cdot \tan^2(\theta/2) \right\}$$

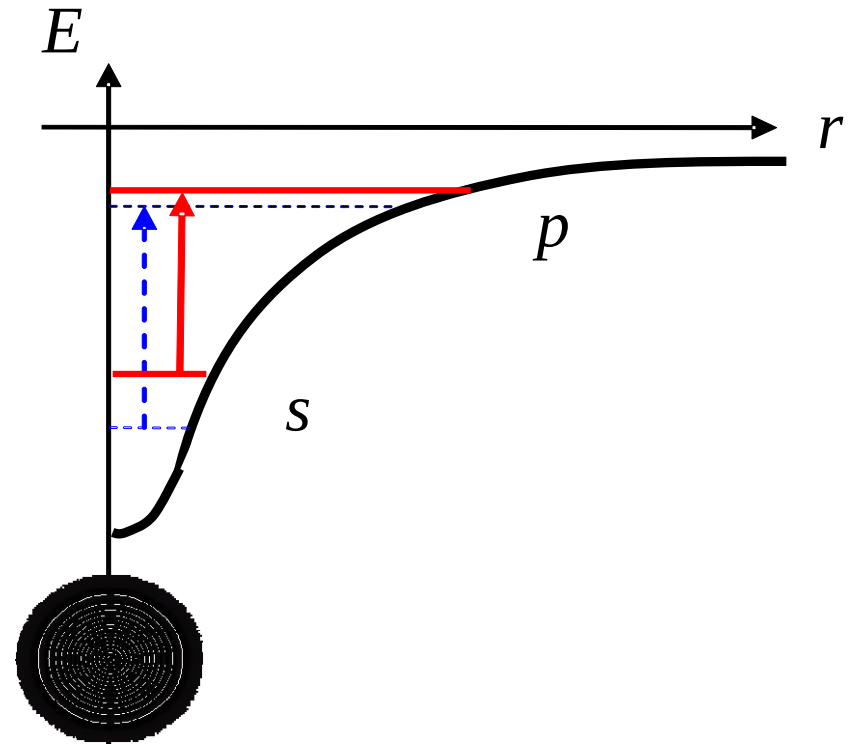
➔ $\langle r_c^2 \rangle = -6 \frac{dG_E(Q^2)}{dQ^2} \Big|_{Q^2=0}$

- Simon et al, (1980) $R_{\text{rms}} = 0.862(12) \text{ fm}$
- I. Sick, (2003). $R_{\text{rms}} = 0.895(18) \text{ fm}$

Nuclear size effect



a) point nucleus
 $V \sim -1/r$



b) finite size nucleus

transition from s to p state

→ decrease transition frequency

- Measure H transition frequency $1s \rightarrow 2s$ at 121 nm, uncertainty: (th) 16 kHz, (exp) 22 kHz

→ charge radius of proton $R_{\text{rms}} = 0.883(14) \text{ fm}$

Lack of precision (only 1~2%) very annoying!

Kirill Melnikov and Timo van Ritbergen, Phys. Rev. Lett. **84**, 1673(2000)

$$E(nS) \approx \frac{-R_\infty}{n^2} + \frac{L_{1S}}{n^3} \quad L_{1S} = (8172 + 1.56 \cdot r_p^2) \text{MHz}$$

- Two unknowns: R_∞ and r_p

0.883(14) fm, hydrogen spectroscopy, ENS Paris

C. Schwob et al., Phys. Rev. Lett. **82**, 4960 (1999).

K. Melnikov and T. van Ritbergen, Phys. Rev. Lett. **84**, 1673(2000).

0.890(14) fm, hydrogen spectroscopy, MPI Garching

T. Udem et al. Phys. Rev. Lett. **79**, 2646, (1997).

0.862(12) fm, original electron scattering result

G.G. Simon et.al. Nucl. Phys. A **333**, 381 (1980)

0.895(18) fm, re-analysis of world data

I. Sick, Phys. Lett. **B 576**, 62-67 (2003)

0.879(8) fm, new experiment by GSI

J. C. Bernauer et al., Phys. Rev. Lett. **105**, 242001 (2010)

Precision Spectroscopy of Atomic Hydrogen

23

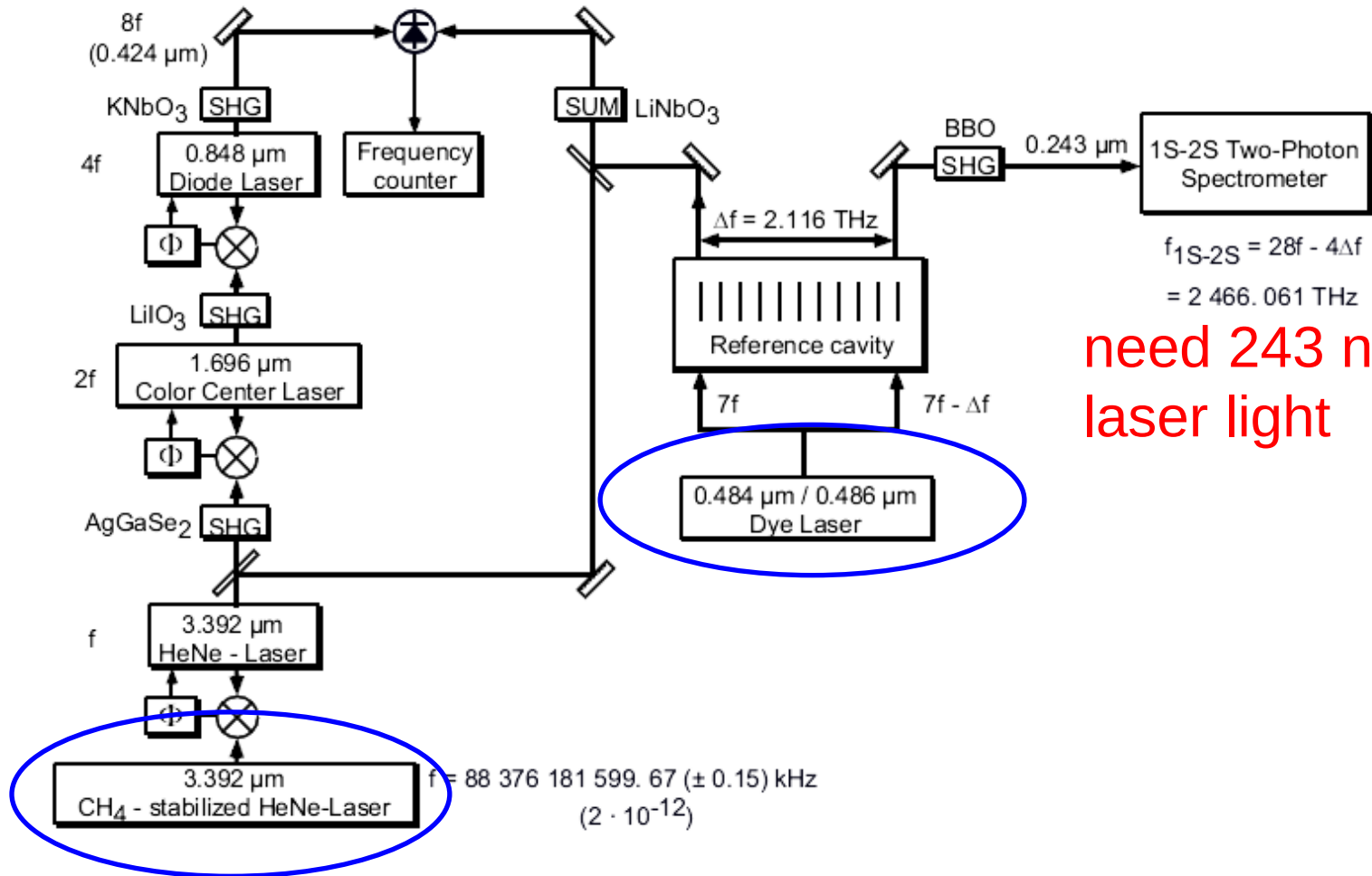


Fig. 4. The first 1992 Garching frequency chain for the measurement of the $1S - 2S$ transition in atomic hydrogen (Φ : phase-locked loop, SHG: second harmonic generation)

Frequency chain

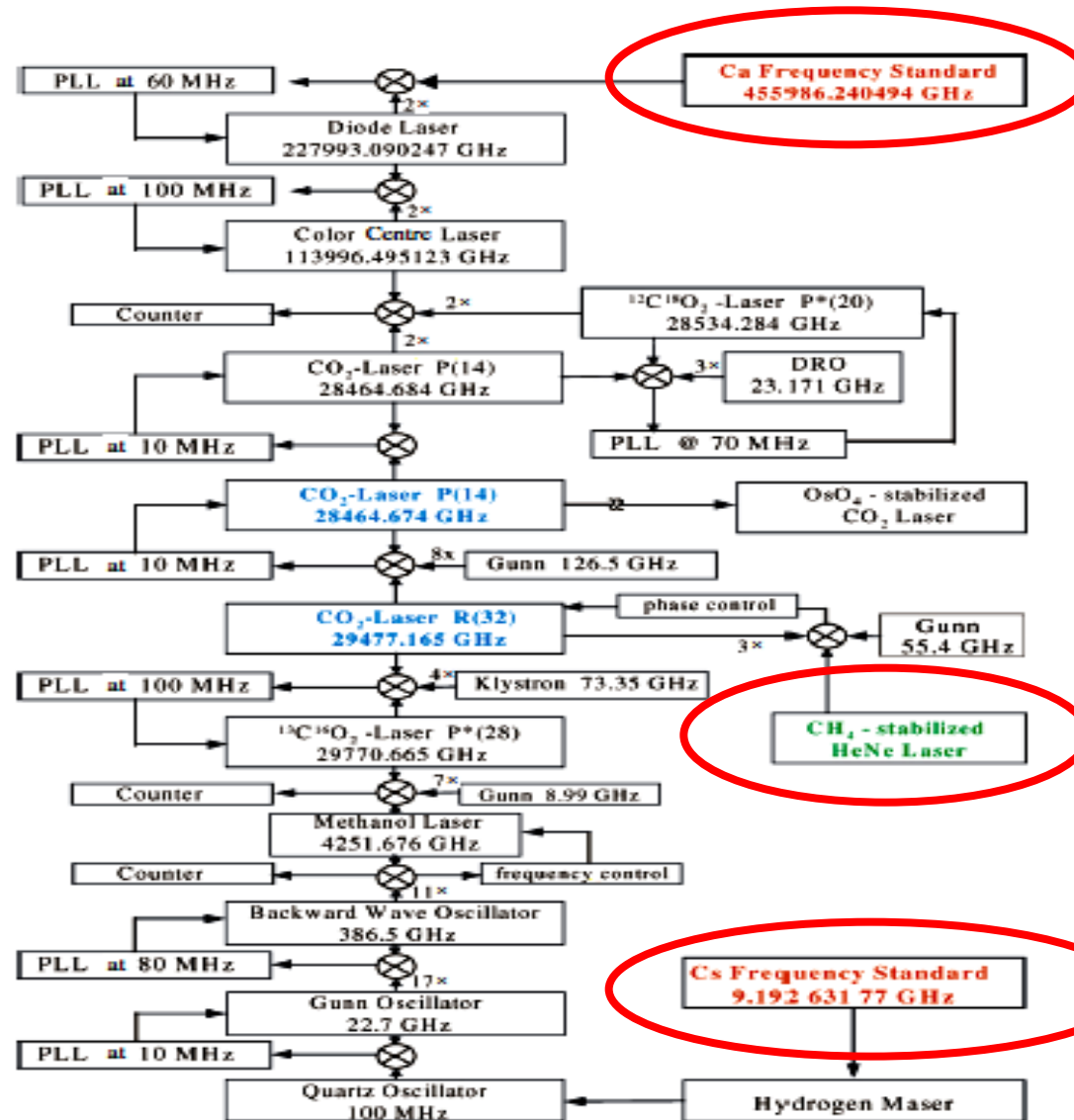
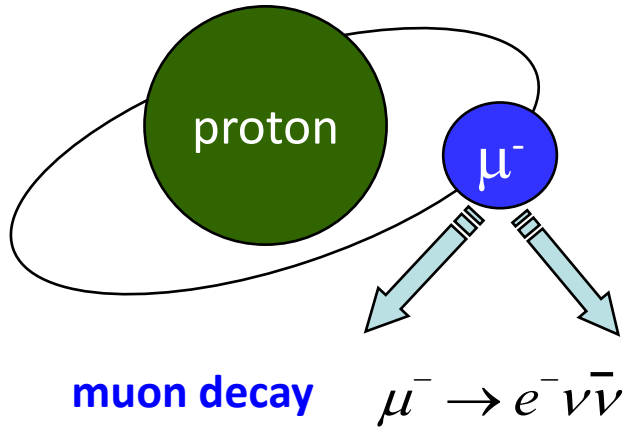


Figure 5. The PTB harmonic frequency chain was the first to achieve a phase-coherent connection between the caesium primary standard and an optical-frequency reference in the visible range. In this case, the target was the calcium optical frequency standard at 457 THz (657 nm). (Adapted with permission from [96].)

Muonic atom result



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- $m_{\mu}/m_e \sim 200$

- Bohr radius

- Energy level

- Wave function

- Energy shift due to nuclear size $\sim |\Psi(0)|^2 \langle r^2 \rangle$

- Sensitivity $\sim (m_{\mu}/m_e)^2$

$$a_0 = \frac{4\pi\epsilon_0\hbar^2}{m_e e^2}$$

$$E_n = -\frac{e^4}{2(4\pi\epsilon_0\hbar)^2} \frac{m_e}{n^2}$$

$$\Psi(r) \sim a_0^{-3/2} e^{-r/a_0}$$

$$r_p = 0.84184(67) \text{ fm}$$

Pohl et al., Nature **466**, 213, 2010

$$r_p = 0.84087(39) \text{ fm}$$

Antognini et al., Science **339**, 417, 2013

Problem in old hydrogen experiments?

Thomas et. al, arXiv:1903.04252 (2019)

$$r_p = 0.877(13) \text{ fm by } 1S-3S \text{ spectroscopy}$$



Frequency standard

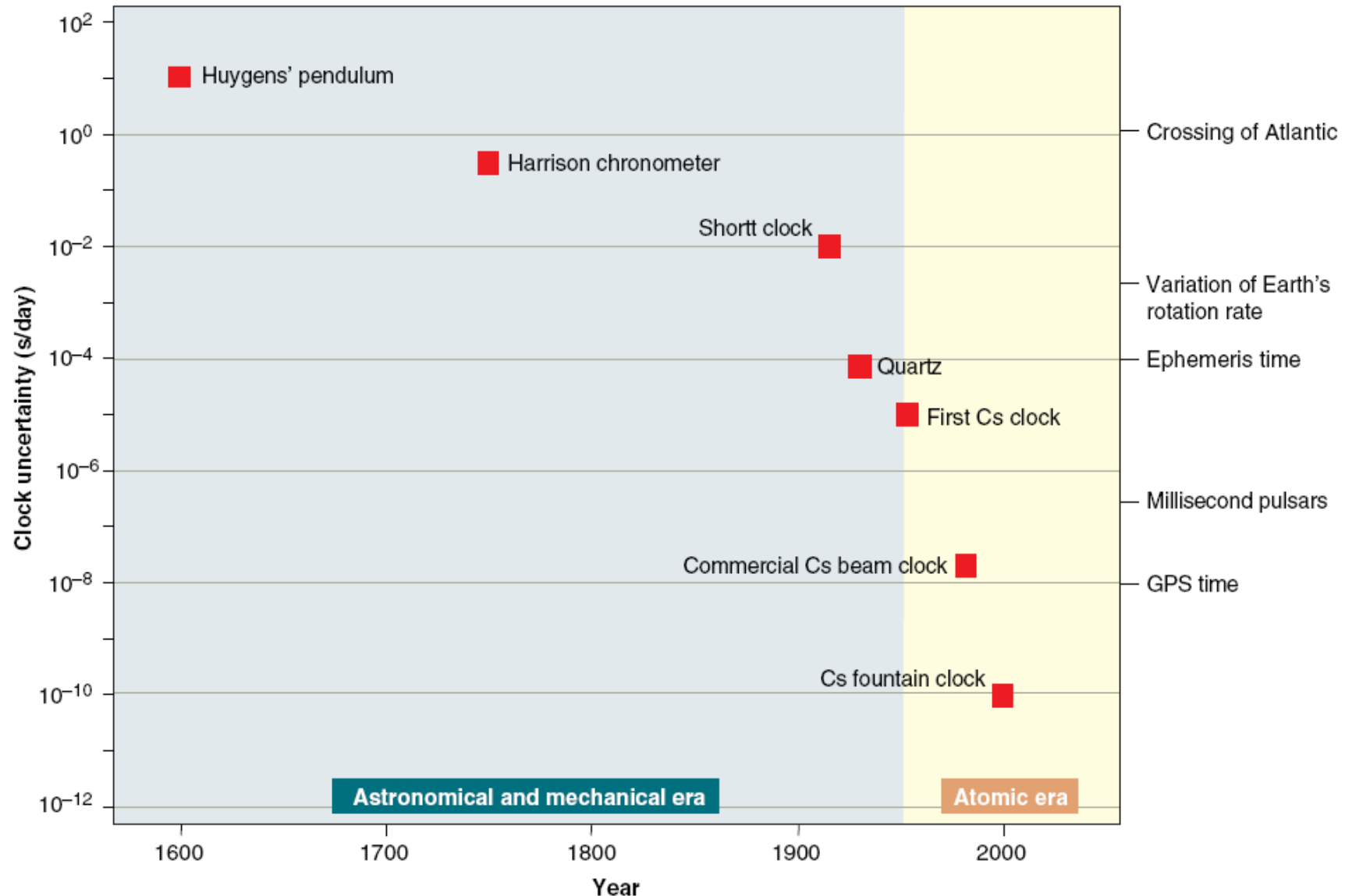
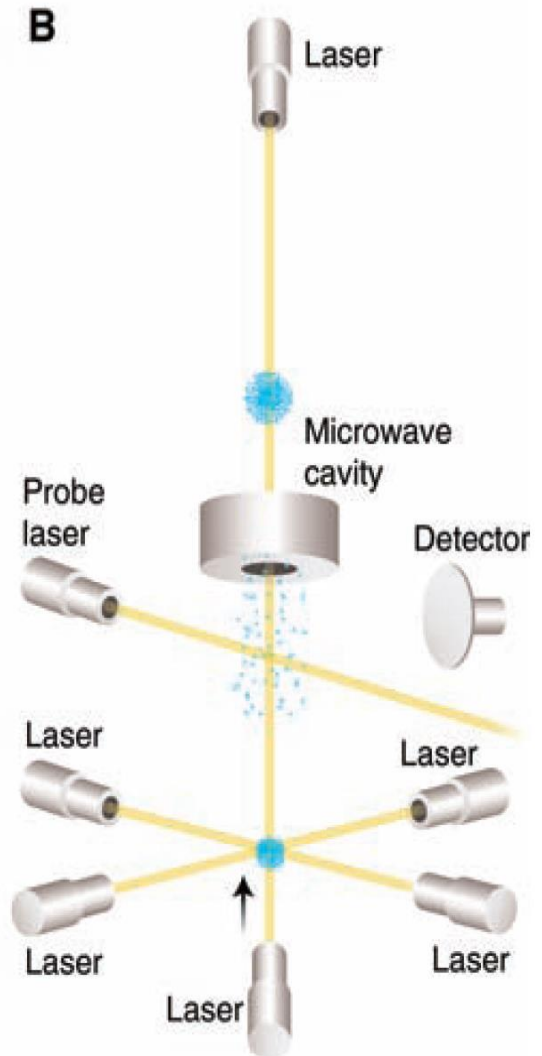
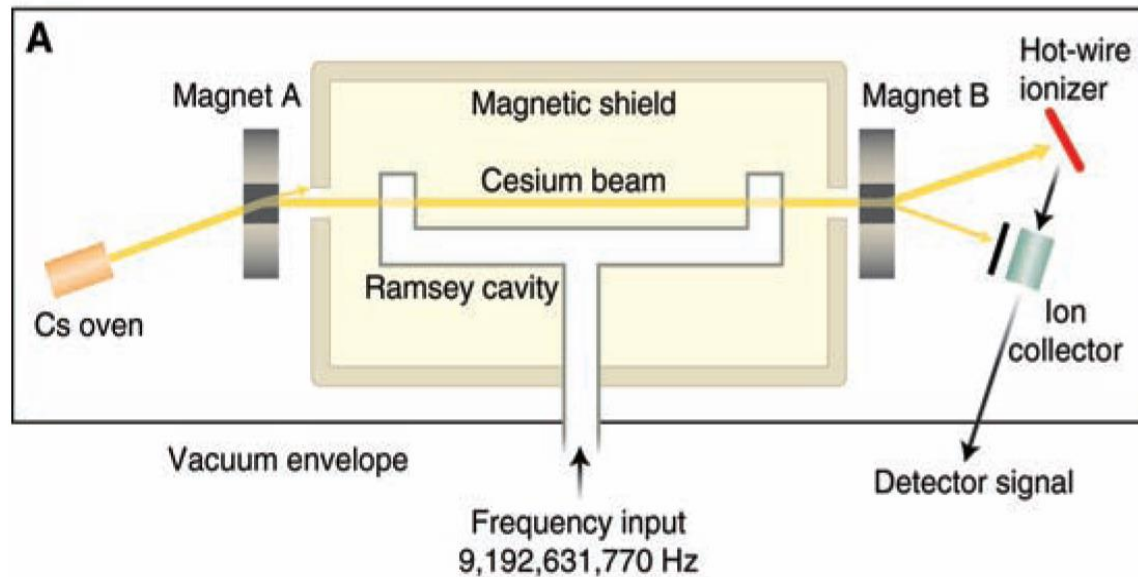


Figure: S.A. Diddams et al., Science 306, 1318, 2004

- When you measure something **very precisely**, **new physics** will come out by itself!
- Applications:
 - GPS
 - Test special relativity, gravitational red-shift
 - Gravitational wave radiation in binary pulsar
 - Test linearity of quantum mechanics
 - Variation of fundamental constant

Microwave Cs clock



NIST Cs fountain

A: Microwave resonance of Cs beam

$$\delta\nu/\nu_0 = 3 \times 10^{-15}$$

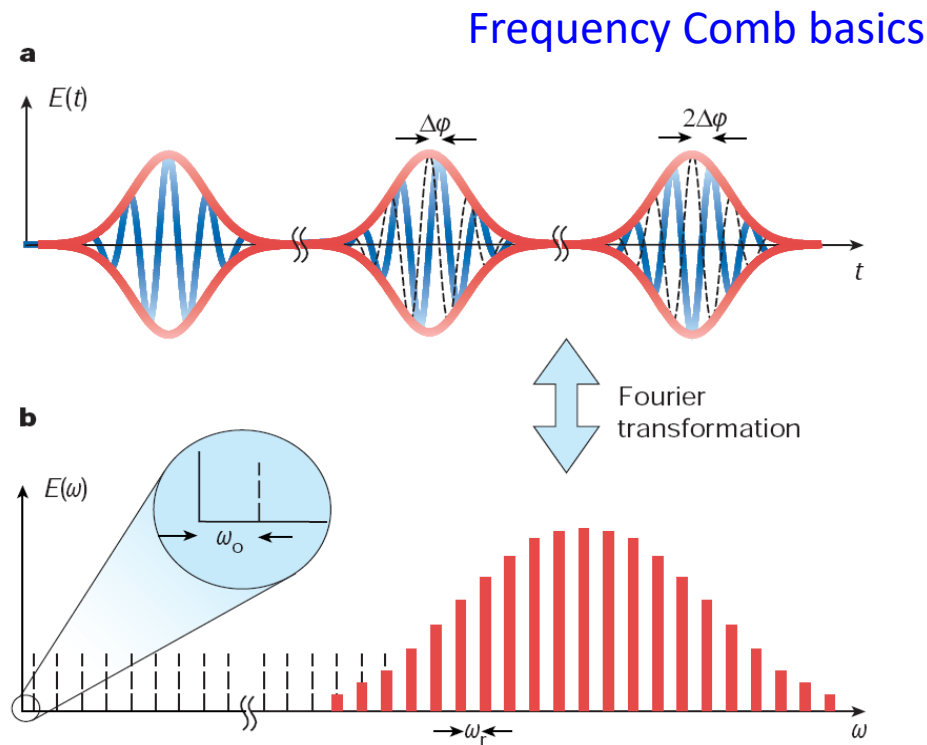
state selection, interaction time

B: Cs fountain clock

$$\delta\nu/\nu_0 = 4 \times 10^{-16}$$

> To improve $\delta\nu/\nu_0$?

- Frequency comb: pioneering work by T.W. Hänsch and J. Hall made frequency measurement possible



$$\omega_n = n \omega_r + \omega_{\text{offset}}$$

$\omega_r : 50\text{MHz} \sim 1\text{GHz}$

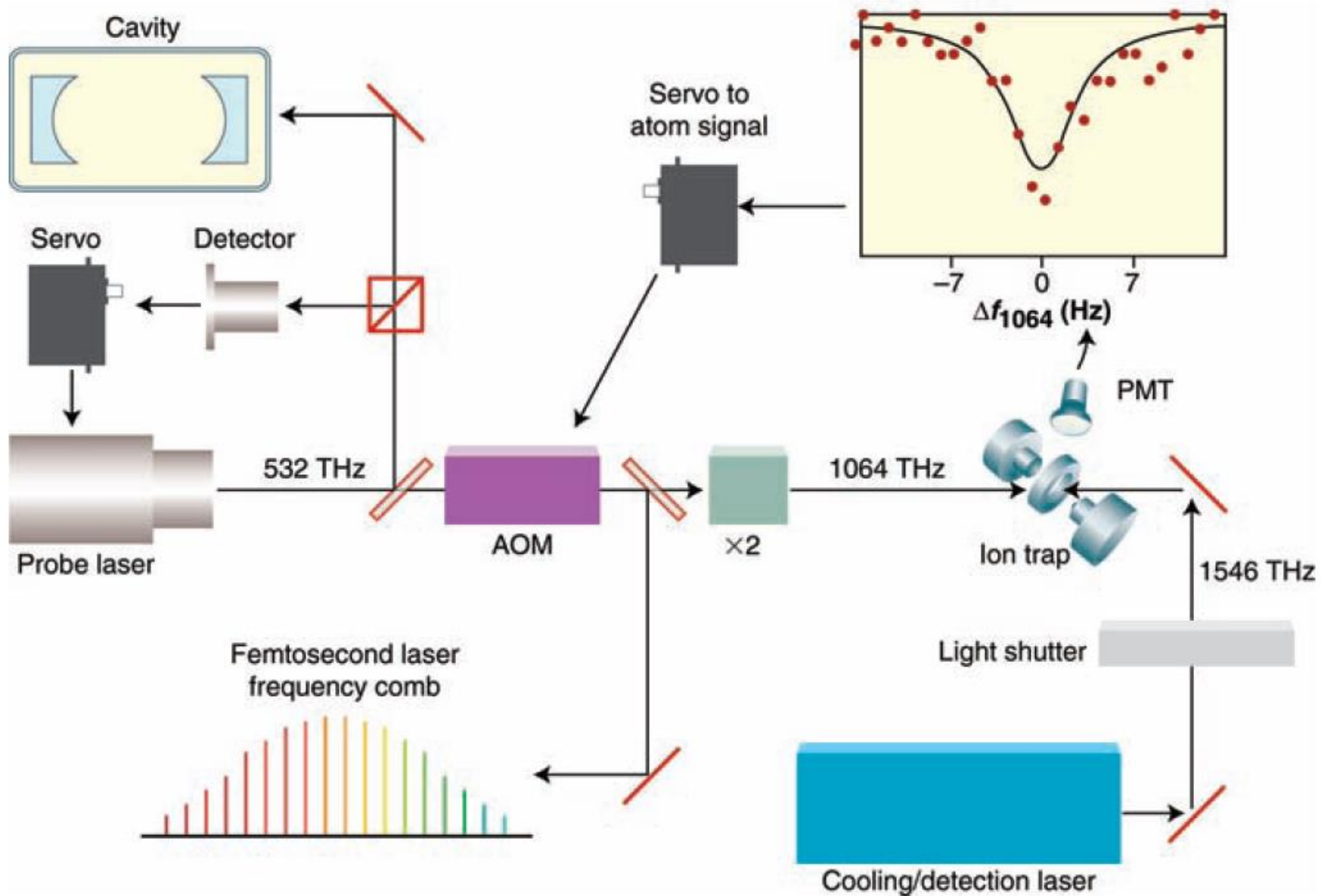
ω_r :
Determined by cavity
length

ω_{offset} :
measured by f-2f
technique

Picture from Th. Udem, R. Holzwarth,
T.W. Hänsch, Nature 416, 233, 2002

- Trapped ions: Hg^+ , Ca^+ , Sr^+ , Yb^+ , Ba^+ , In^+ , Al^+
Long interaction time, can be laser cooled
Single ion detection, poor S/N
- Trapped neutral atoms: Ca, Sr, Yb, Mg, H.....
Large number of atoms, very high S/N
Complicated systematics
- Precision $\sim 10^{-18} \text{ (Hz}^{1/2}\text{)}$

How optical clock work





$$\alpha = e^2 / \hbar c \approx 1/137$$

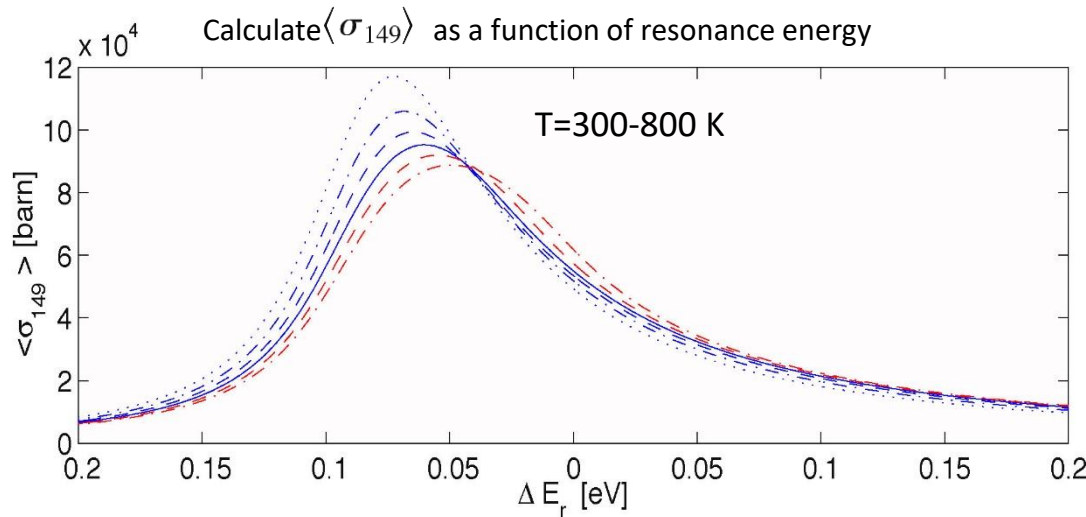
$$\alpha^{-1} = 137.035\,999\,074\,(44)$$

- How constant is the constant?
- Two reports show the constant may change in time
- Oklo phenomenon
- Quasar spectra

The Oklo Natural Nuclear Reactor



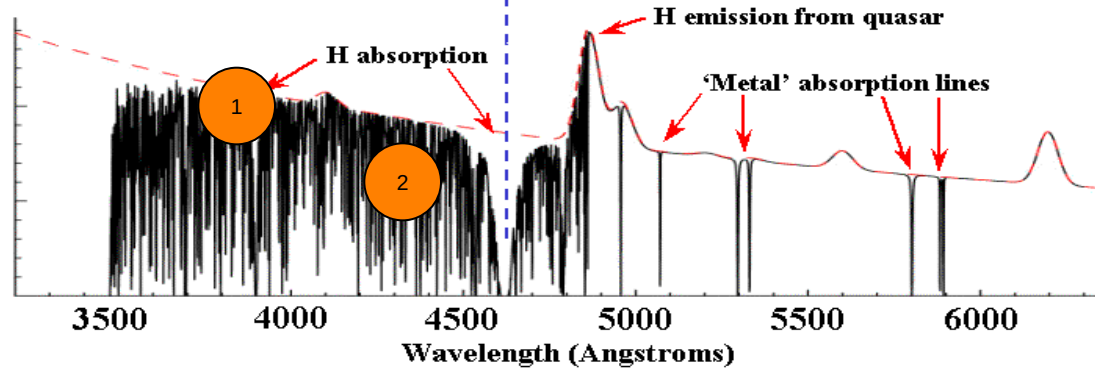
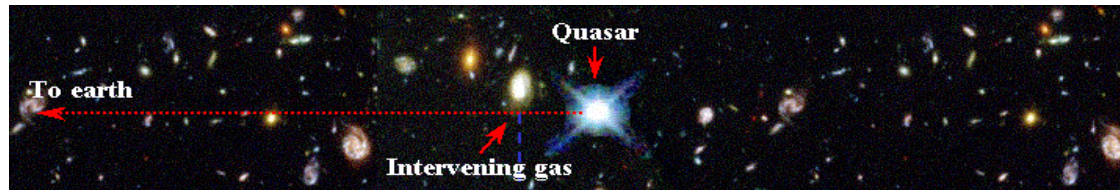
- Natural nuclear reactor 1.9 billion years ago
- ^{235}U abundance very low
- Fission of ^{235}U , neutron slowed by water
- $n + ^{149}\text{Sm} \rightarrow ^{150}\text{Sm} + \gamma, E_r = 97.3 \text{ meV}$



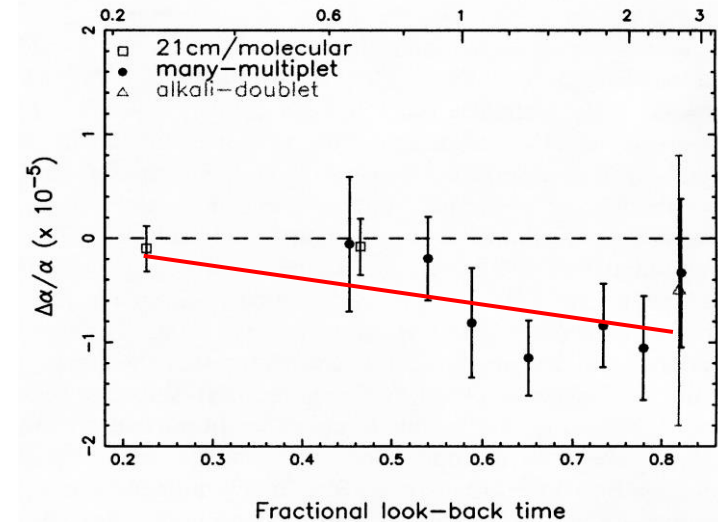
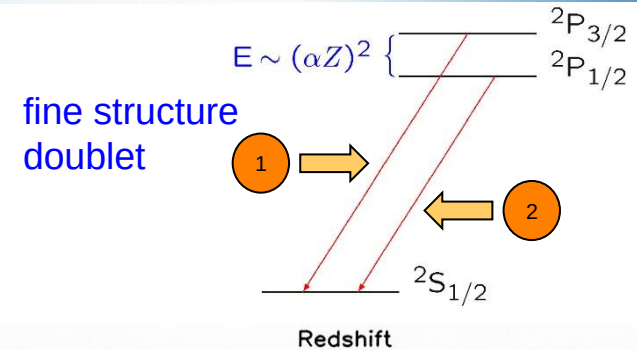
- ^{149}Sm isotope abundance, $\rightarrow$ neutron absorption cross section σ_{149}
- From $\sigma_{149} \rightarrow \Delta E_r$, assume reasonable T

$$\frac{\Delta E_r}{1.1 \text{ MeV}} = \frac{\Delta \alpha}{\alpha} \quad \Delta E_r < 0.1 \text{ eV}, \Delta \alpha / \alpha < 10^{-7}$$

Quasar Absorption Spectra



$\Delta\alpha/\alpha$	z	t
$(7.2 \pm 1.8) \cdot 10^{-6}$	0.5–2.6	0.38–0.82
$\rightarrow (6.4 \pm 1.4) \cdot 10^{-16}/\text{yr}$		



- (1) Atoms located in different regions with different velocities
→ different Doppler red shift
- (2) Isotopic abundance in the past may change
→ shift the position of absorption line
- (3) Magnetic field shift the energy levels by Zeeman effect

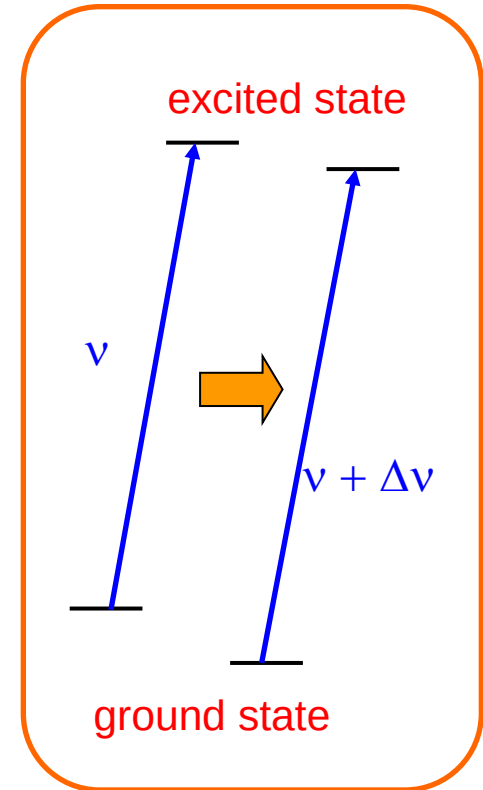
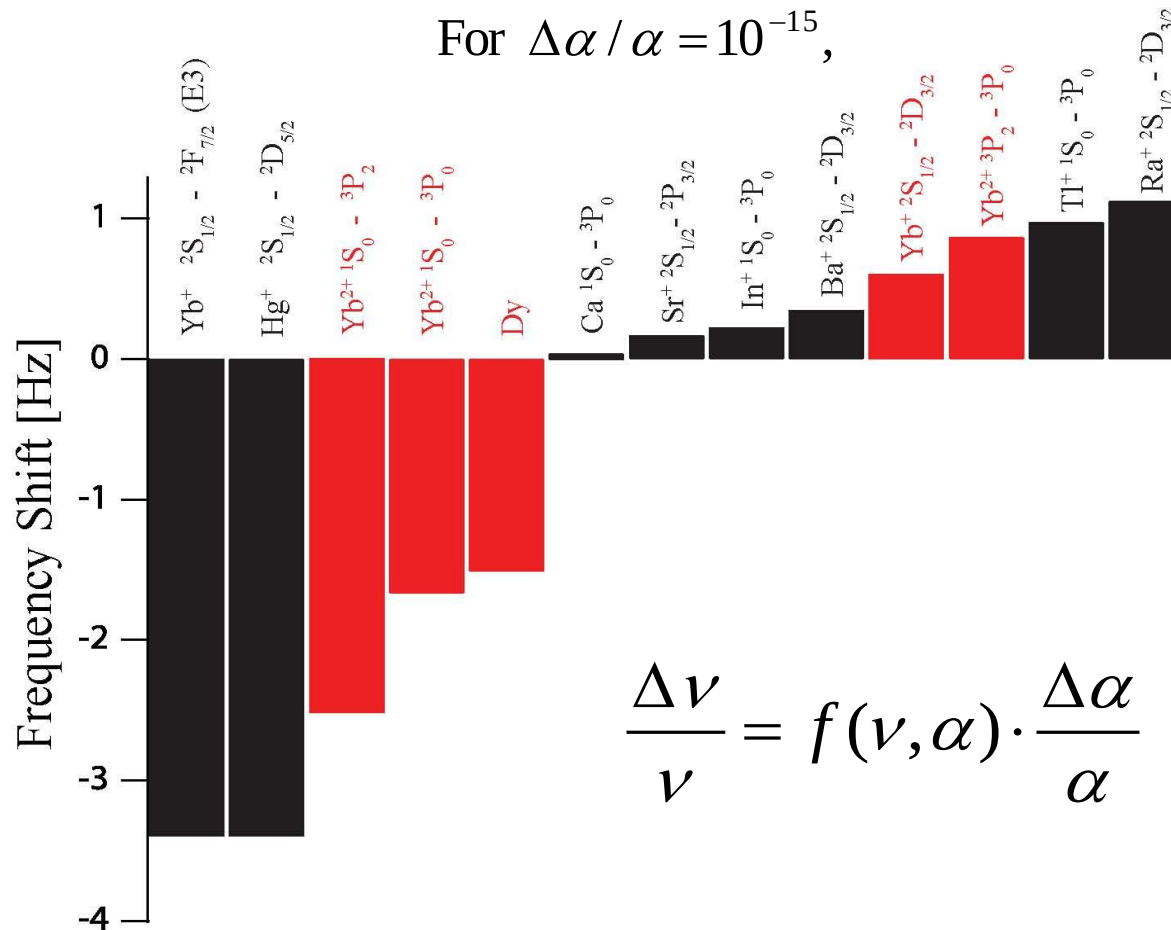
Results From Searches for $d\alpha/dt$



$\Delta\alpha/\alpha$ per year	method	quantity	ref
$-(2.3 +0.7/-0.3) 10^{-17}$	Oklo nuclear reactor	$\alpha(m_\pi/\Lambda_{\text{qcd}})^{0.5}$	1
$< (0.8) 10^{-17}$	Oklo nuclear reactor	$\alpha(m_\pi/\Lambda_{\text{qcd}})^{0.5}$	2
$(7.2 \pm 1.8) 10^{-16}$	Quasar spectra	α	3
$(6.4 \pm 1.4) 10^{-16}$	Quasar spectra	α	4
$(0.6 \pm 0.6) 10^{-16}$	Quasar spectra	α	5
$(4.4 \pm 1.6) 10^{-16}$	Quasar spectra	α	6
$(-2.7 \pm 2.6) 10^{-15}$	Dy atomic beam	α	7
$(-0.9 \pm 2.9) 10^{-15}$	Opt-H $\leftrightarrow$ Cs	$g_{\text{Cs}}m_e/m_p\alpha^{2.8}$	8
$(0.3 \pm 2.0) 10^{-15}$	Opt-Yb ⁺ $\leftrightarrow$ Cs	$g_{\text{Cs}}m_e/m_p\alpha^{1.9}$	9
$(0.02 \pm 0.7) 10^{-15}$	Rb $\leftrightarrow$ Cs	$g_{\text{Cs}}/g_{\text{Rb}}\alpha^{0.44}$	10
$(-6.2 \pm 6.5) 10^{-17}$	Opt-Hg ⁺ $\leftrightarrow$ Cs	$g_{\text{Cs}}m_e/m_p\alpha^{6.0}$	11
$(-1.6 \pm 2.3)10^{-17}$	Opt-Hg ⁺ $\leftrightarrow$ Opt-Al ⁺	α	12

- [1] S. K. Lamoreaux and J. R. Torgerson, Phys. Rev. D **69**, 121701R (2004).
- [2] Y. Fujii *et al.*, Nuclear Physics B **573**, 377 (2000).
- [3] J. K. Webb *et al.*, Phys. Rev. Lett. **87**, 091301 (2001).
- [4] M. T. Murphy, J. K. Webb, and V. V. Flambaum, Mon. Not. R. Astron. Soc. **345**, 609 (2003).
- [5] R. Srianand *et al.*, Phys. Rev. Lett. **92**, 121302 (2004).
- [6] M. T. Murphy, J. K. Webb, and V. V. Flambaum, astro-ph/0612407 (2006).
- [7] A. Cingöz *et al.*, Phys. Rev. Lett. **98**, 040801 (2007).
- [8] M. Fischer *et al.*, Phys. Rev. Lett. **92**, 230802 (2004).
- [9] E. Peik *et al.*, Phys. Rev. Lett. **93**, 170801 (2004).
- [10] H. Marion *et al.*, Phys. Rev. Lett. **90**, 150801 (2003).
- [11] T. M. Fortier *et al.*, Phys. Rev. Lett. **98**, 070801 (2007).
- [12] T. Rosenband *et al.*, Science **319**, 1808 (2008)

- Laboratory and controllable test strongly desired
- Ultra-cold trapped ions and atoms best testing ground



Results From Searches for $d\alpha/dt$



$\Delta\alpha/\alpha$ per year	method	quantity	ref
$-(2.3 \pm 0.7/-0.3) 10^{-17}$	Oklo nuclear reactor	$\alpha(m_\pi/\Lambda_{\text{qcd}})^{0.5}$	1
$< 0.8 10^{-17}$	Oklo nuclear reactor	$\alpha(m_\pi/\Lambda_{\text{qcd}})^{0.5}$	2
$(7.2 \pm 1.8) 10^{-16}$	Quasar spectra	α	3
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$(0.6 \pm 0.6) 10^{-16}$	Quasar spectra	α	5
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$(-2.7 \pm 2.6) 10^{-15}$	Dy atomic beam	α	7
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[1] S. K. Lamoreaux and J. R. Torgerson, Phys. Rev. D **69**, 121701R (2004).

[2] Y. Fujii *et al.*, Nuclear Physics B **573**, 377 (2000).

[3] J. K. Webb *et al.*, Phys. Rev. Lett. **87**, 091301 (2001).

[4] M. T. Murphy, J. K. Webb, and V. V. Flambaum, Mon. Not. R. Astron. Soc. **345**, 609 (2003).

[5] R. Srianand *et al.*, Phys. Rev. Lett. **92**, 121302 (2004).

[6] M. T. Murphy, J. K. Webb, and V. V. Flambaum, astro-ph/0612407 (2006).

[7] A. Cingöz *et al.*, Phys. Rev. Lett. **98**, 040801 (2007).

[8] M. Fischer *et al.*, Phys. Rev. Lett. **92**, 230802 (2004).

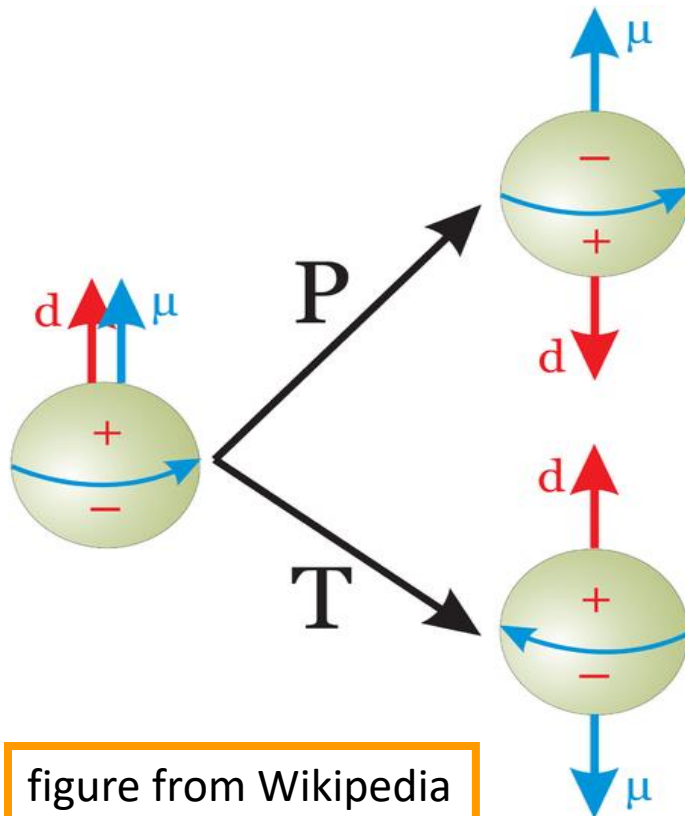
[9] E. Peik *et al.*, Phys. Rev. Lett. **93**, 170801 (2004).

[10] H. Marion *et al.*, Phys. Rev. Lett. **90**, 150801 (2003).

[11] T. M. Fortier *et al.*, Phys. Rev. Lett. **98**, 070801 (2007).

[12] T. Rosenband *et al.*, Science **319**, 1808 (2008)

- 1957 Landau first pointed out the electric dipole moment (EDM) of a fundamental particle would suggest P and T violation



$\mathbf{d}$: electric dipole moment
= vector

$\mathbf{\mu}$: spin or magnetic moment
= pseudo-vector, (like $\mathbf{r} \times \mathbf{p}$)

figure from Wikipedia

Neutron EDM: history

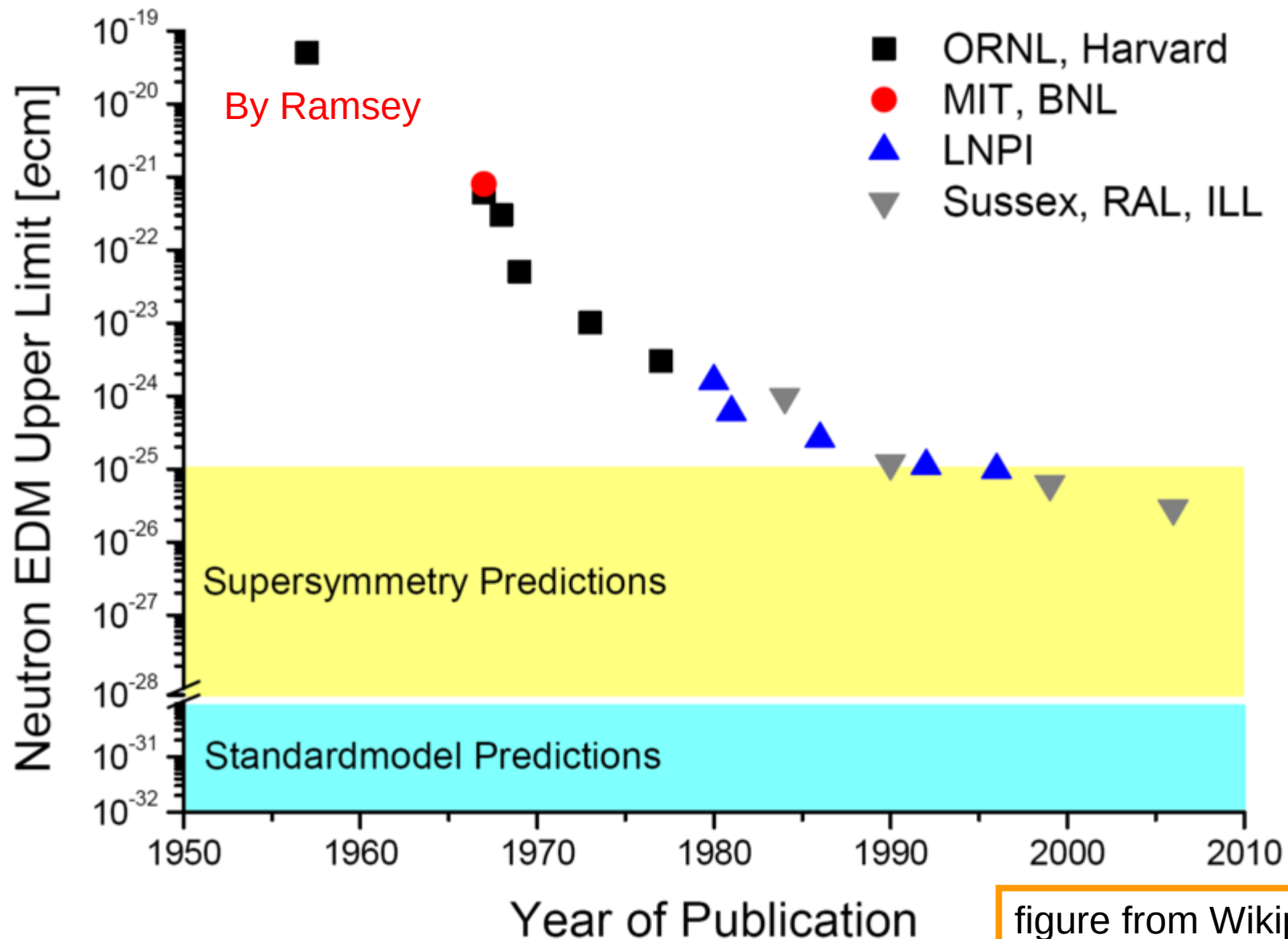


figure from Wikipedia

- Matter–antimatter asymmetry suggest another source of CP violation
 - Electron: **intrinsic ?**
 - Quark: **intrinsic ?**
 - Neutron/proton:
from quark EDM ? New interaction?
 - Atoms, molecules:
large enhancement factors

Atoms: Tl, Cs, Hg, Xe, Rn, Ra, Fr

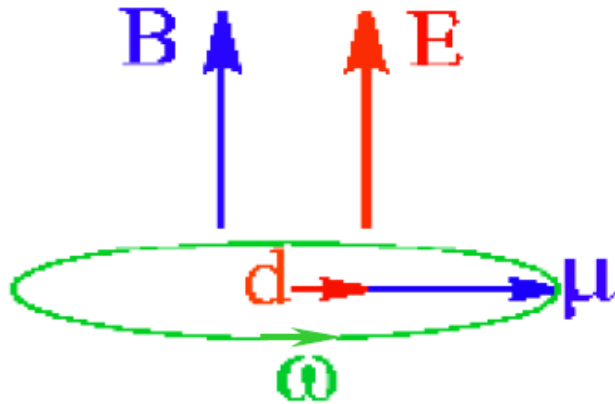
Molecule: PbO, **YbF**, **ThO**

Ions, solid state systems, etc...

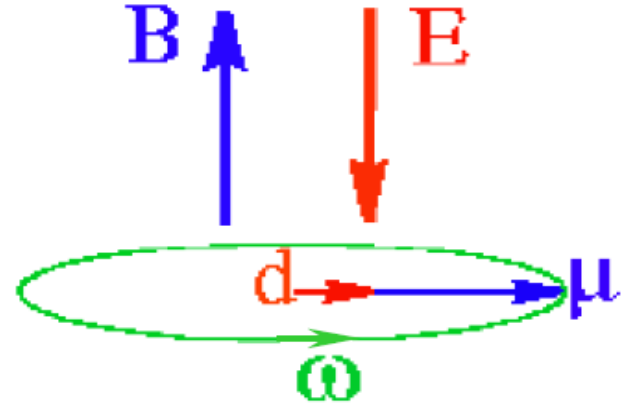
How to measure EDM



$$H = -\vec{\mu} \cdot \vec{B} - \vec{d} \cdot \vec{E}$$



$$\omega_1 = \frac{2\mu B + 2dE}{\hbar}$$



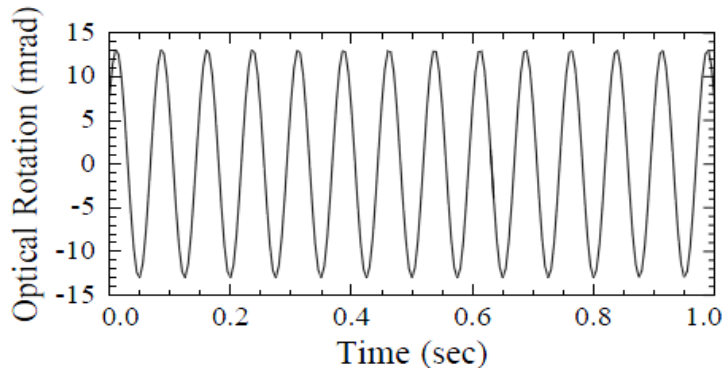
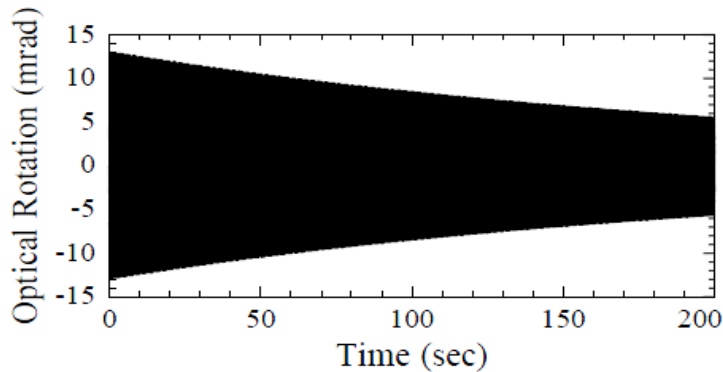
$$\omega_2 = \frac{2\mu B - 2dE}{\hbar}$$

$$\omega_1 - \omega_2 = \frac{4dE}{\hbar}$$

figure from M. Romalis, Princeton

Problem:
reverse E field cause leakage current,
huge noise from B field

^{199}Hg Precession Data



Typical value

$E \sim 10 \text{ kV/cm}$

$\omega \sim 10 \text{ Hz}, \delta\omega \sim \text{nHz}$

Sensitivity approaching
some Standard Model extension

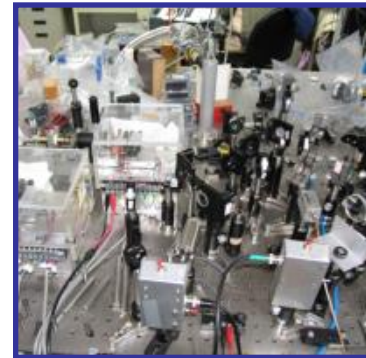
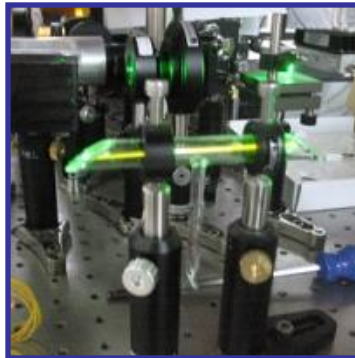
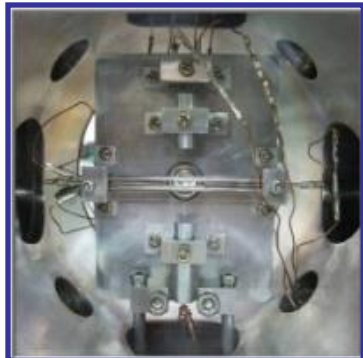
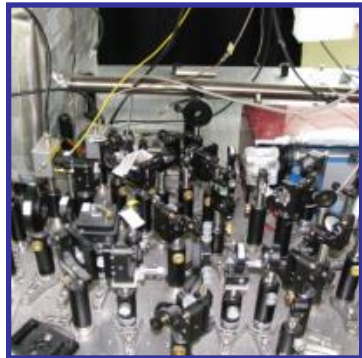
What if someone observes EDM?
Theories ready with many parameters

figure from M. Romalis, Princeton

Second lecture:

Laser spectroscopy of simple atoms

work performed at NTHU Physics



Precision spectroscopy of simple atoms

- **Hydrogen**: exact solution exist
- **Helium**: numerical method
- **Lithium**: numerical method

Remain as input for atomic calculation

- **Magnetic moment** of nucleus
- **Charge radius** of nucleus

- Calculable **atomic structure**: precision measurement test fundamental physics
- Advances in atomic theories, e.g. in H, He, Li, He^+ , Li^+ , Be^+ etc...
- New experimental techniques, e.g.: cold atoms by atom and ion trap, frequency comb, exotic atoms, etc...
- What **fundamental physics** to study?

Atomic structure (Hydrogen)



- non-relativistic
- relativistic correction
- spin-orbit interaction (L·S, fine structure)

● nuclear moment (HFS)

● QED effect (Lamb shift)

● nuclear size effect

$$En = -\alpha^2 mc^2 \left(\frac{1}{2n^2} \right) = \frac{-13.6 \text{ eV}}{n^2}$$
$$- \alpha^4 mc^2 \frac{1}{4n^2} \left[\frac{2n}{(l+1/2)} - \frac{3}{2} \right]$$
$$- \alpha^4 mc^2 \frac{1}{4n^2} \left[\frac{2n}{(j+1/2)} - \frac{3}{2} \right]$$

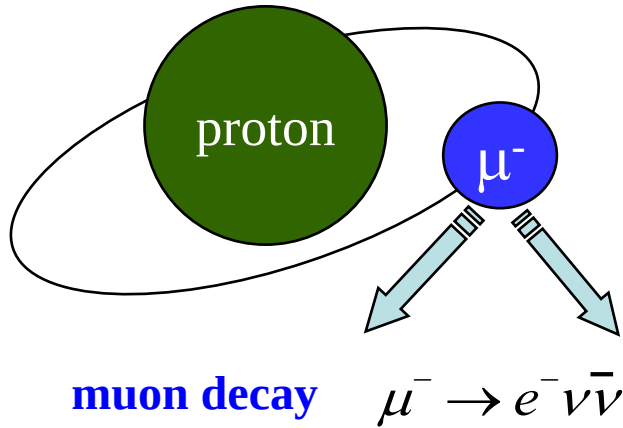
$$\left(\frac{m}{m_p} \right) \alpha^4 mc^2 \frac{4\gamma_p}{3n^2} [f(f+1) - 3/2]$$
$$\alpha^5 mc^2 \frac{1}{4n^3} \left\{ k(n, l) \pm \frac{1}{\pi(j+1/2)(l+1/2)} \right\}$$

$$\frac{2\pi}{3} Ze^2 |\psi(0)|^2 \langle r^2 \rangle_{\text{proton}}$$

Muonic atom result



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- $m_\mu/m_e \sim 200$

- Bohr radius

- Energy level

- Wave function

- Energy shift due to nuclear size $\sim$

- Sensitivity $\sim (m_\mu/m_e)^2$

$$a_0 = \frac{4\pi\epsilon_0\hbar^2}{m_e e^2}$$

$$E_n = -\frac{e^4}{2(4\pi\epsilon_0\hbar)^2} \frac{m_e}{n^2}$$

$$\Psi(r) \sim a_0^{-3/2} e^{-r/a_0}$$

$$|\Psi(0)|^2 \langle r^2 \rangle$$

$$r_p = 0.84184(67) \text{ fm}$$

Pohl et al., Nature **466**, 213, 2010

$$r_p = 0.84087(39) \text{ fm}$$

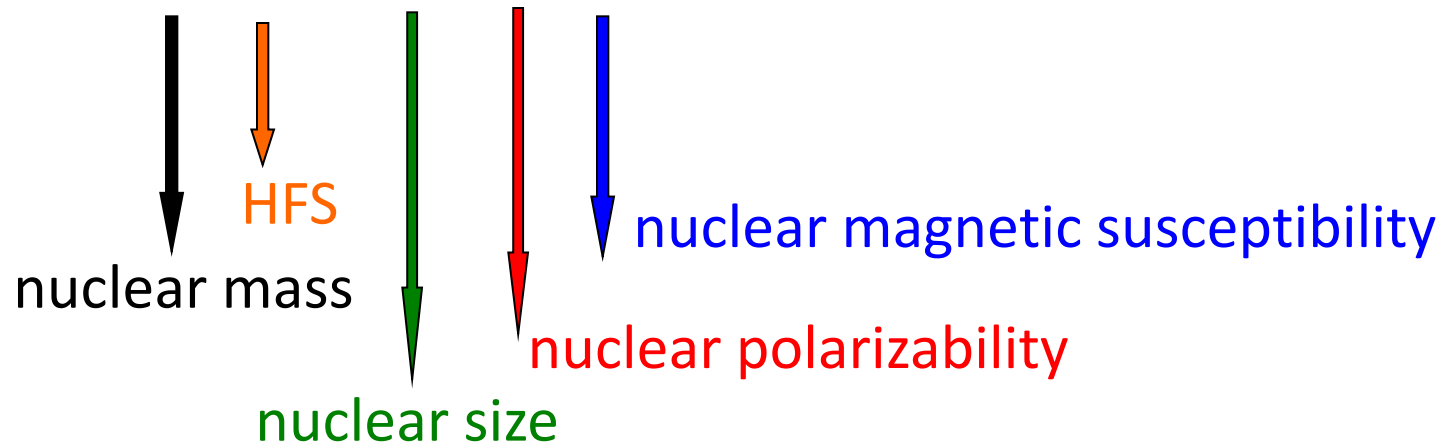
Antognini et al., Science **339**, 417, 2013

Need more experimental verification of
atomic QED calculations



- Measure isotope shift (H-D), 1s - 2s at 121 nm

$$\Delta\nu = 670\,994.33464(15) \text{ MHz}$$



Reference:

- [1] Kirill Melnikov and Timo van Ritbergen, Phys. Rev. Lett. **84**, 1673(2000)
- [2] A. Huber, Th. Udem, B. Gross, J. Reichert, M. Kourogi, K. Pachucki, M. Weitz, and T.W. Hänsch. Phys. Rev. Lett. **80**, 468 (1998)

Total transition frequency (optical):

- QED effects
- H: ~ 10 kHz, He: ~ 1 MHz, Li: > 10 MHz

Isotope shift:

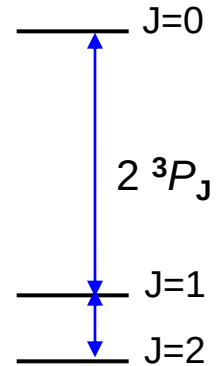
- Nuclear charge radius
- H: ~ 1 kHz, He: ~ 1 kHz, Li: ~ 10 kHz

Fine and hyperfine structure :

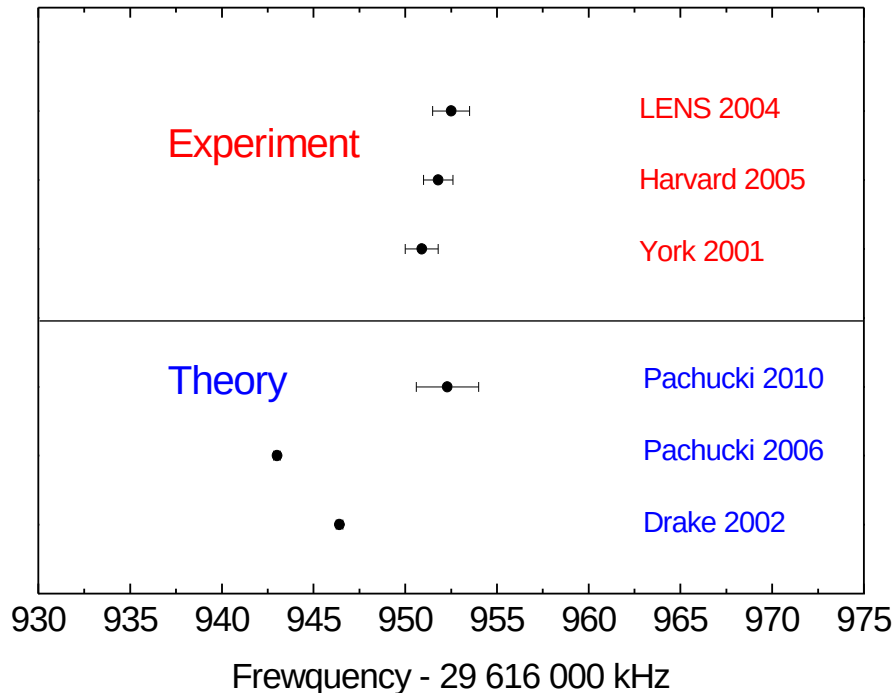
- State mixing, find hyperfine constant
- H: < 1 kHz, He: < 1 kHz, Li: several kHz

Previous attempt for helium FS

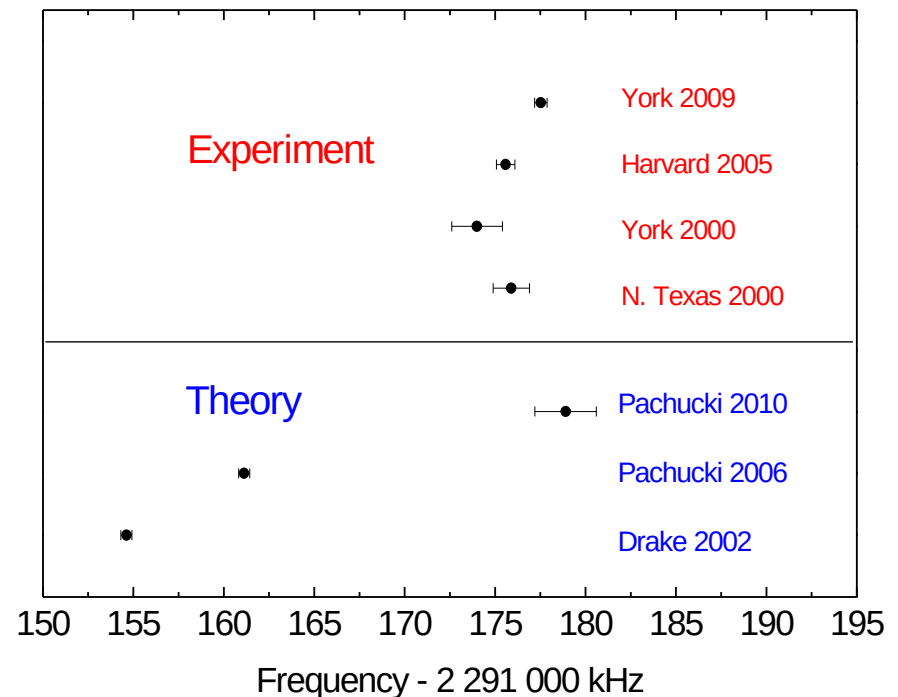
- uncertainty of theory and exp ~ 1 kHz
- Discrepancy = 10 kHz and 20 kHz for 10 years
- Include high-order terms ($\alpha^5 \text{Ry}$)



2^3P_{0-1}

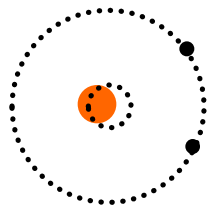


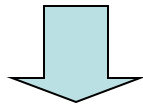
2^3P_{1-2}



Isotope Shift $\delta\nu = \delta\nu_{\text{MS}} + \delta\nu_{\text{FS}}$

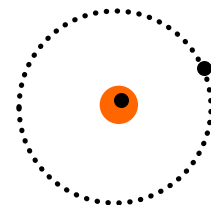
Mass shift: due to nucleus recoil

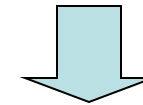

$$\delta\nu_{\text{MS}} \propto \frac{A - A'}{AA'}$$



Nuclear mass precisely known

Field shift: due to nuclear size

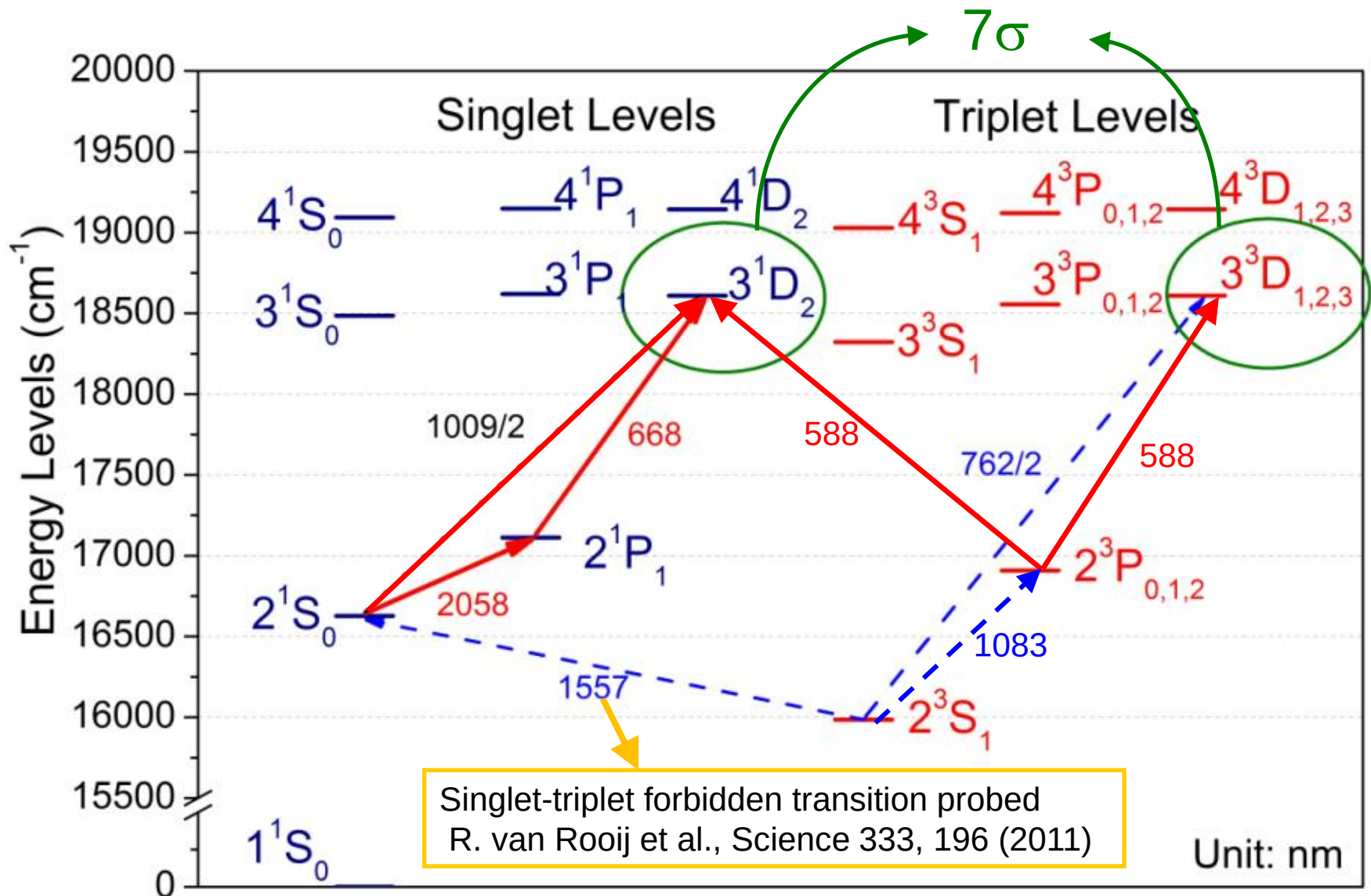

$$\delta\nu_{\text{FS}} \propto [\Psi(0)]^2 \times \delta\langle r_c^2 \rangle$$



Atomic structure $\Psi(0)$ calculable,
H, He, Li.

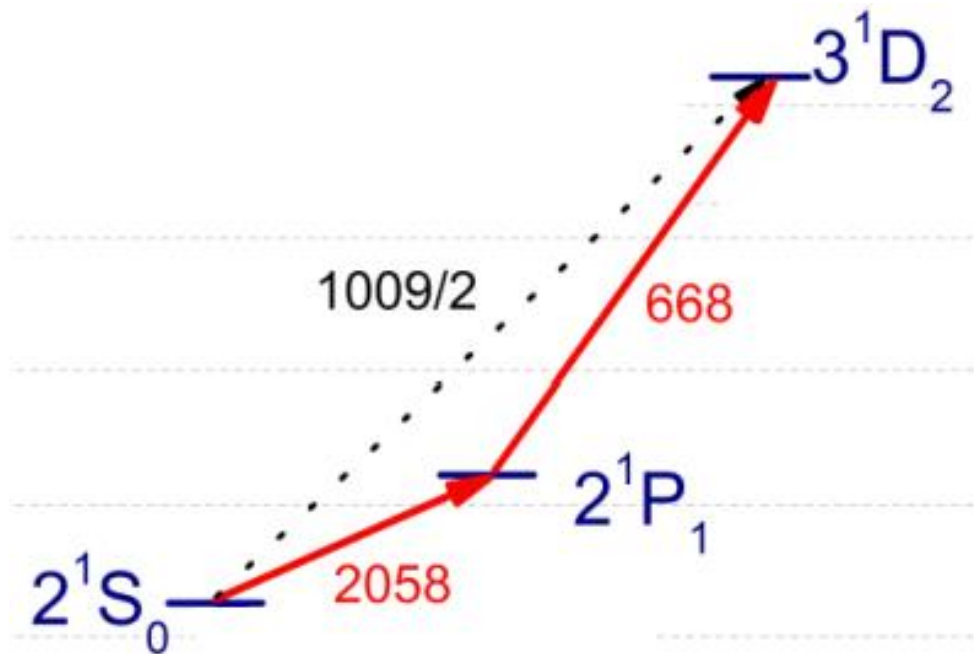
This method widely used in H-D, ^4He - ^3He , and ^6Li - ^7Li and
determine the difference in the nuclear size

Helium energy levels



- Better theory
- Experiments mostly performed in triplet state
- Singlet-triplet forbidden transition probed

R. van Rooij et al., Science **333**, 196 (2011)



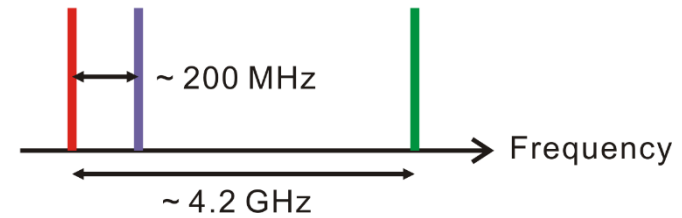
2058 nm transition



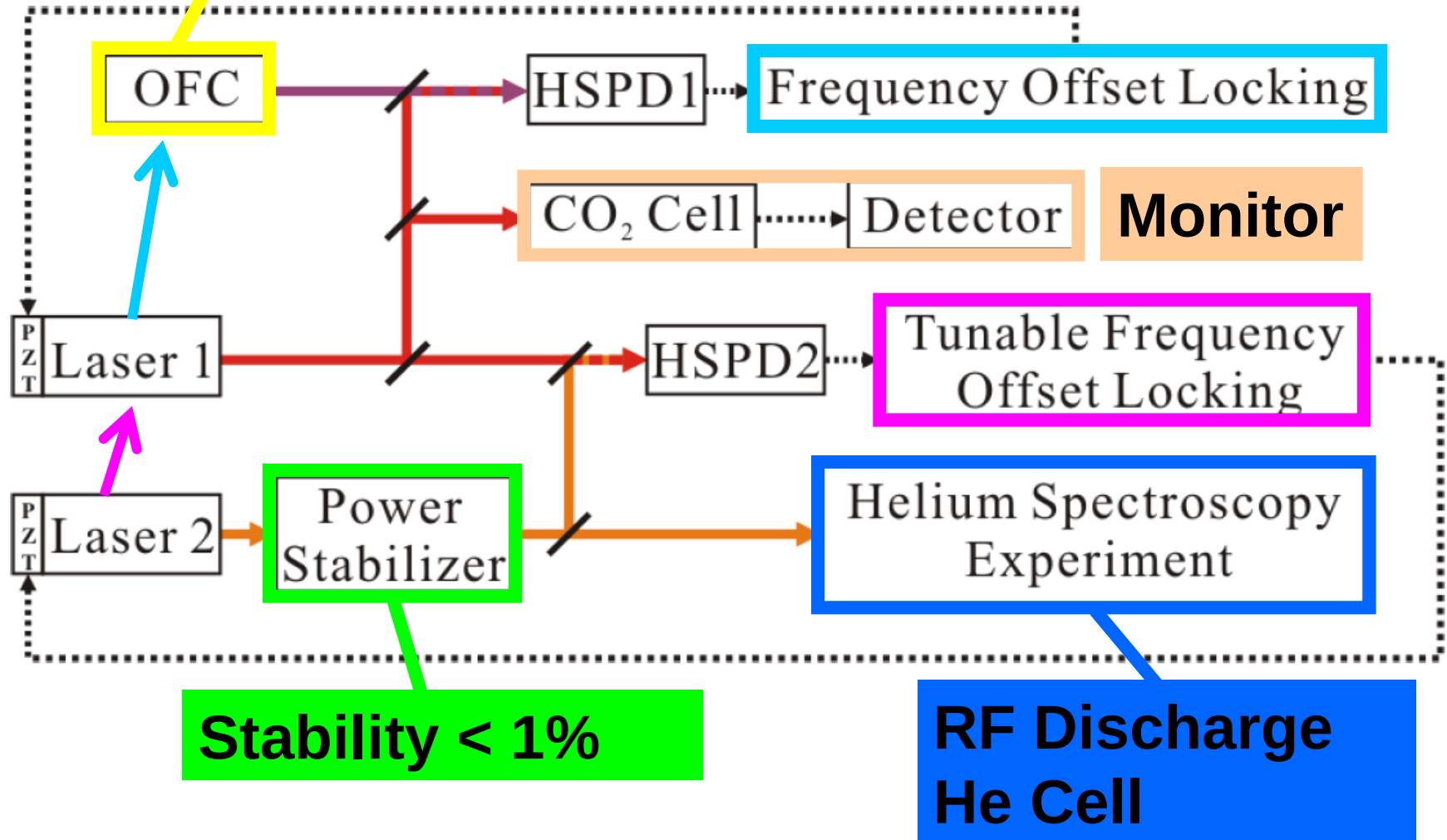
CO₂ 00001-20013 R(4)

⁴He 2¹S₀-2¹P₁

³He 2¹S₀-2¹P₁



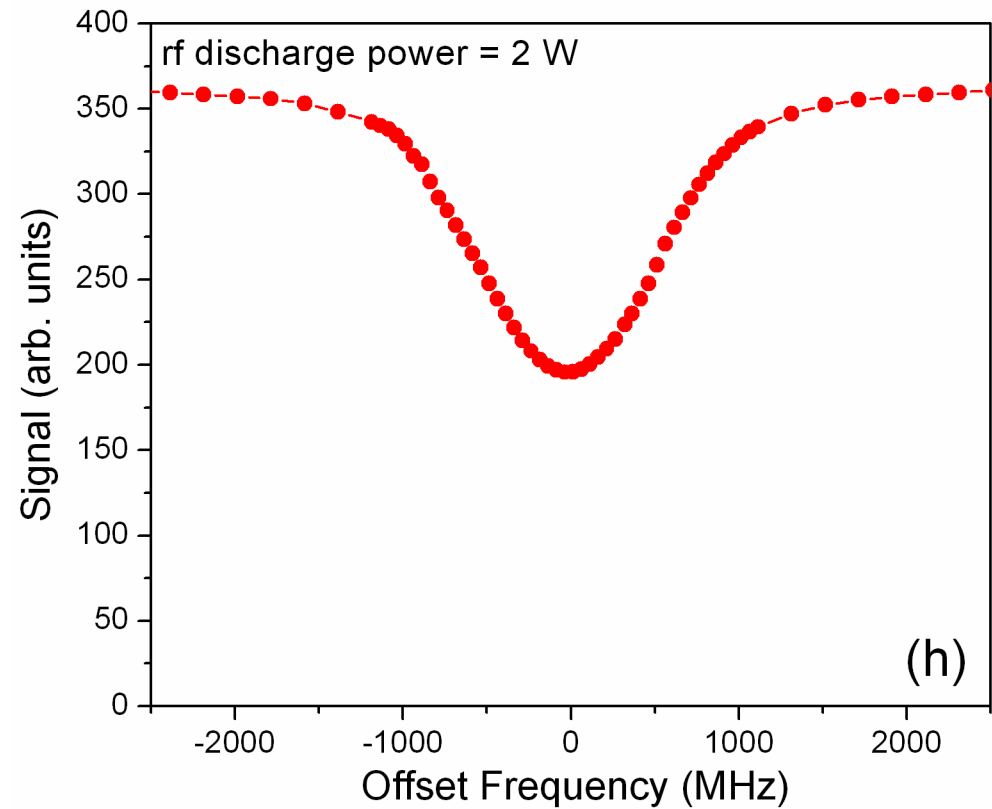
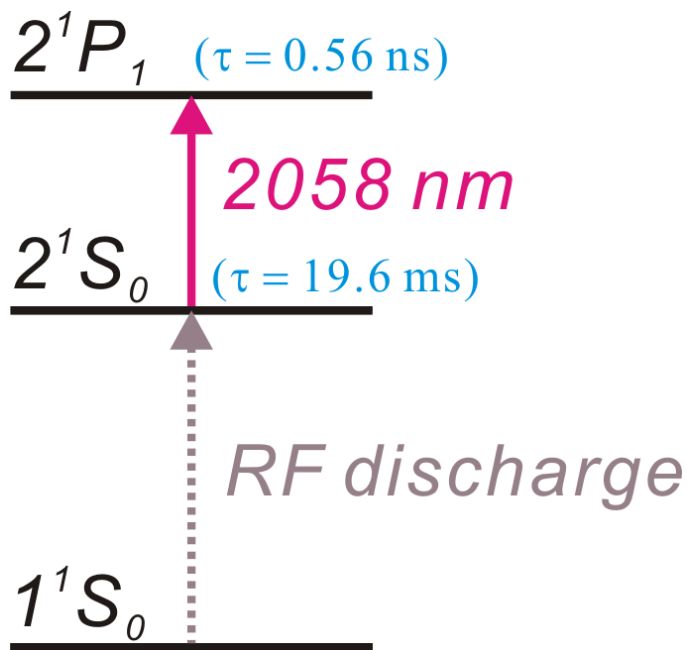
Accuracy = 10^{-12} @ 1000-sec



Direct Absorption Spectrum



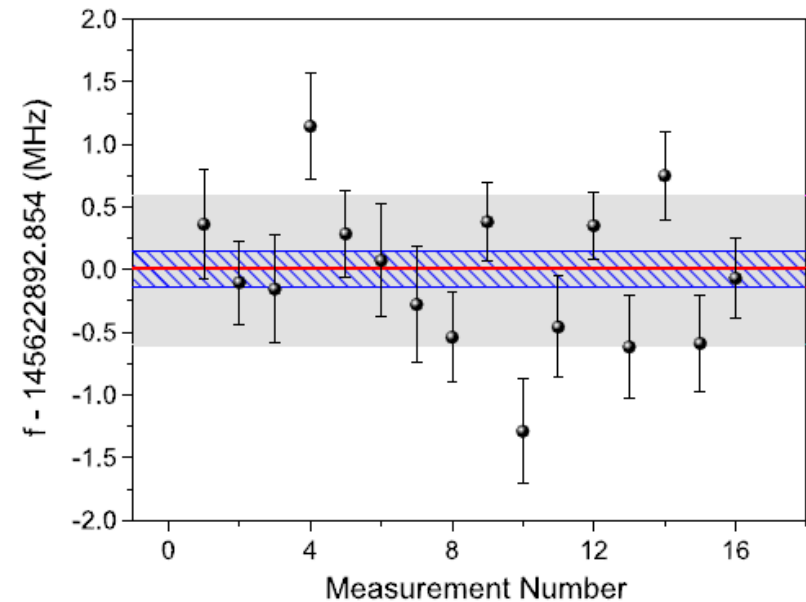
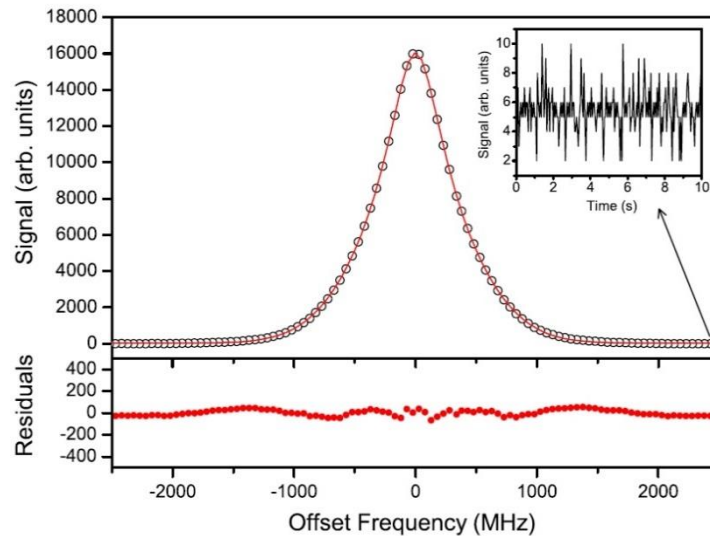
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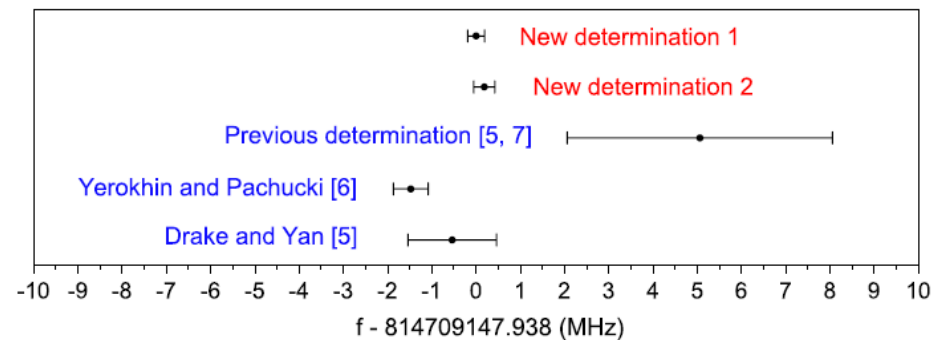
Spectroscopic results



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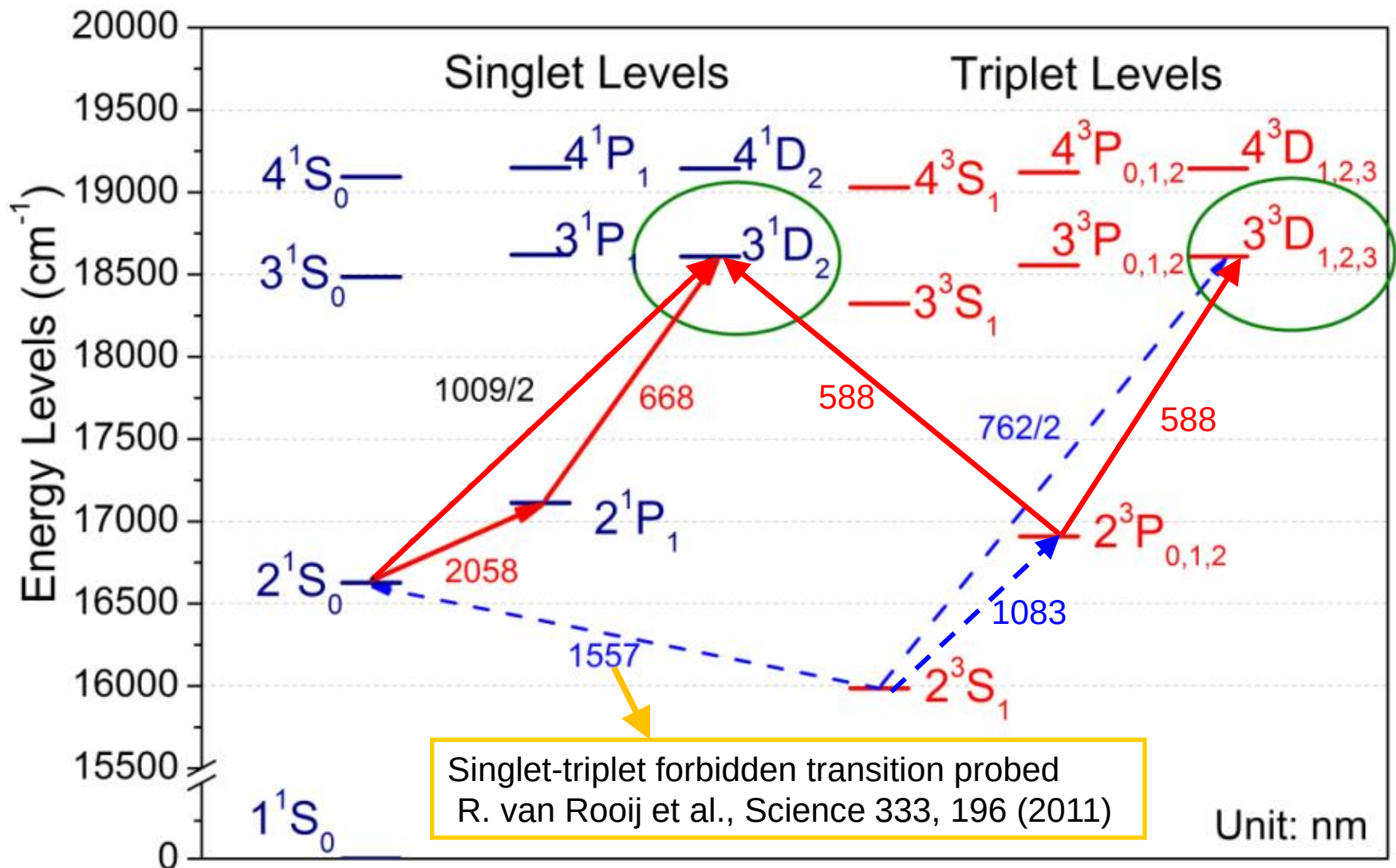
Source	kHz
Statistical error	141
OFC accuracy	<5
Frequency locking stability of laser 1 and OFC	<5
Frequency offset locking stability of two lasers	<2
Pressure shift	9.3
Light intensity	66
rf discharge power	82
Zeeman effect	<7
Overall uncertainty	183



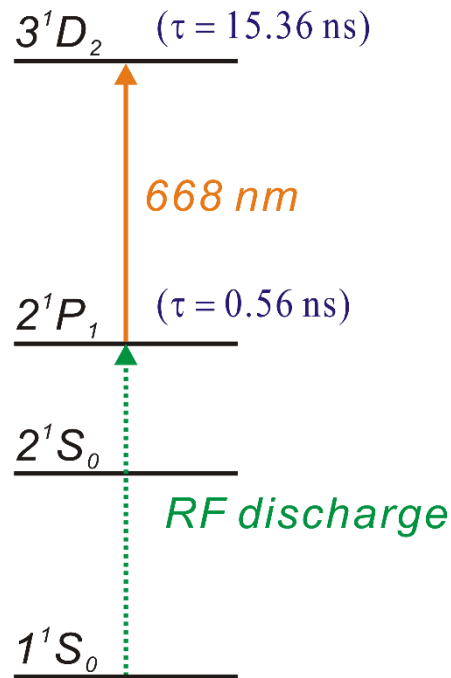
P.-L. Luo et al., *Phys. Rev. Lett.* **111**, 013002 (2013)

- First absolute frequency measurement on this transition
- 10 time more precise in isotope shift

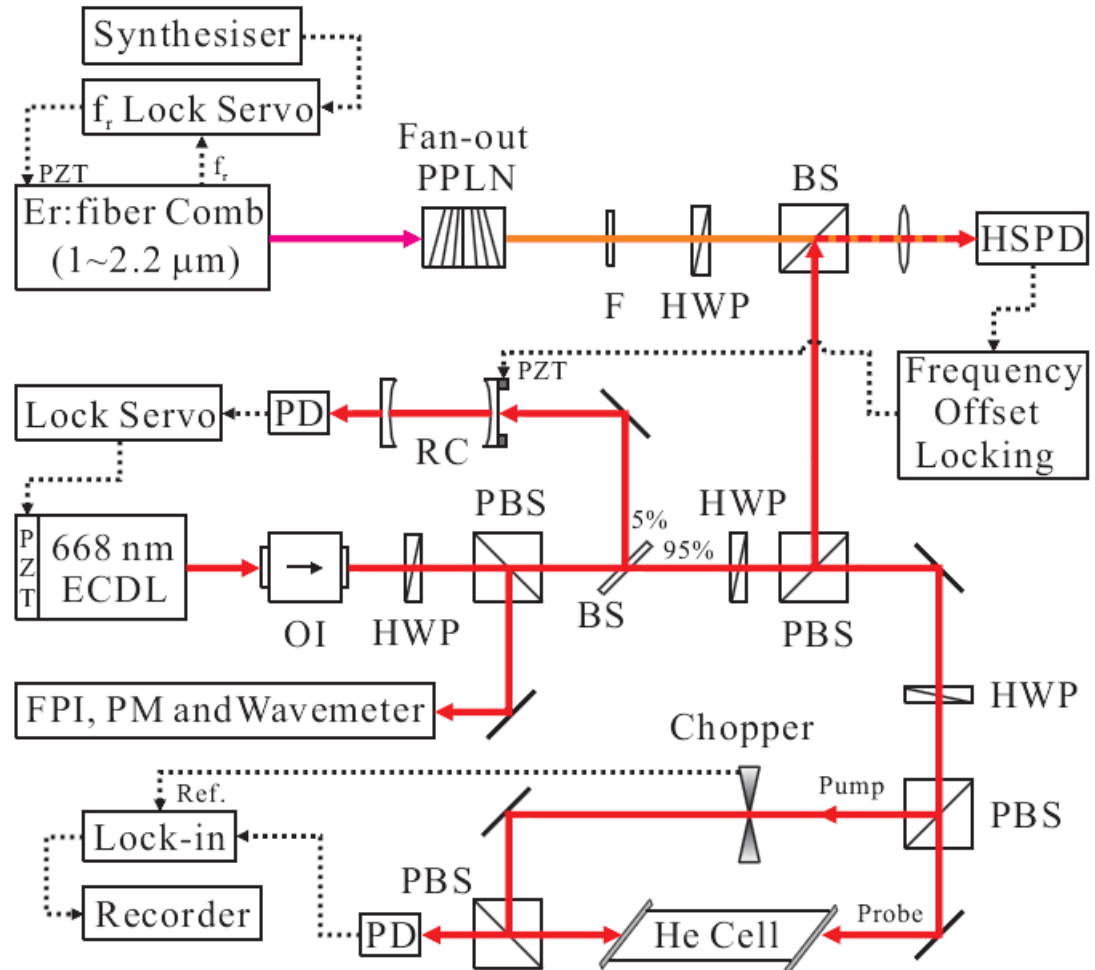
Helium energy levels

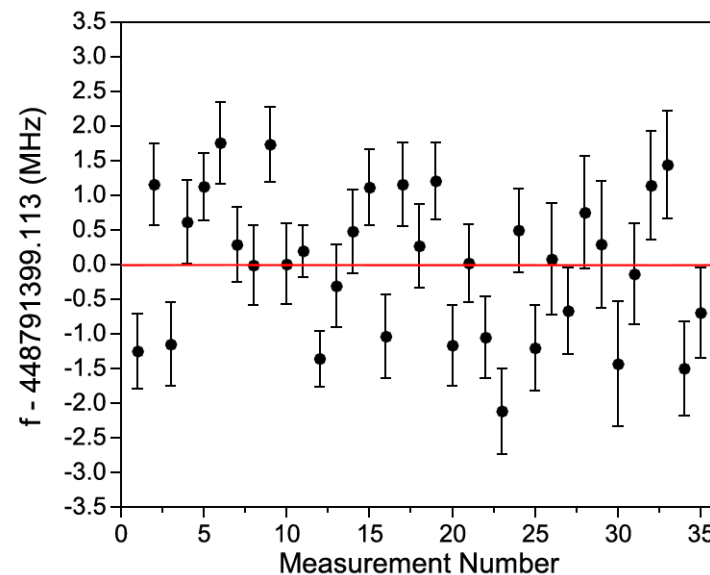
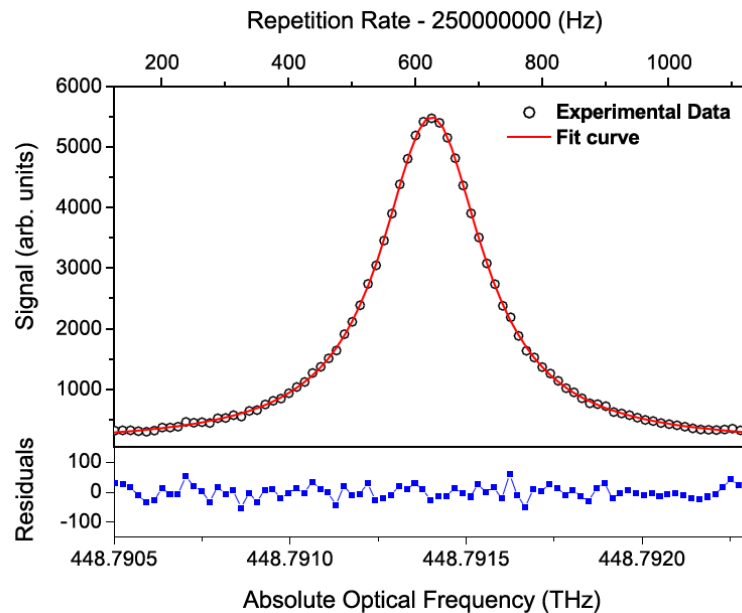


- ECDL at 668 nm locked to frequency comb
- Comb $\rightarrow$ PPLN, change rep rate



$$f_L = n \times f_r + 2 \times f_o \pm f_b$$



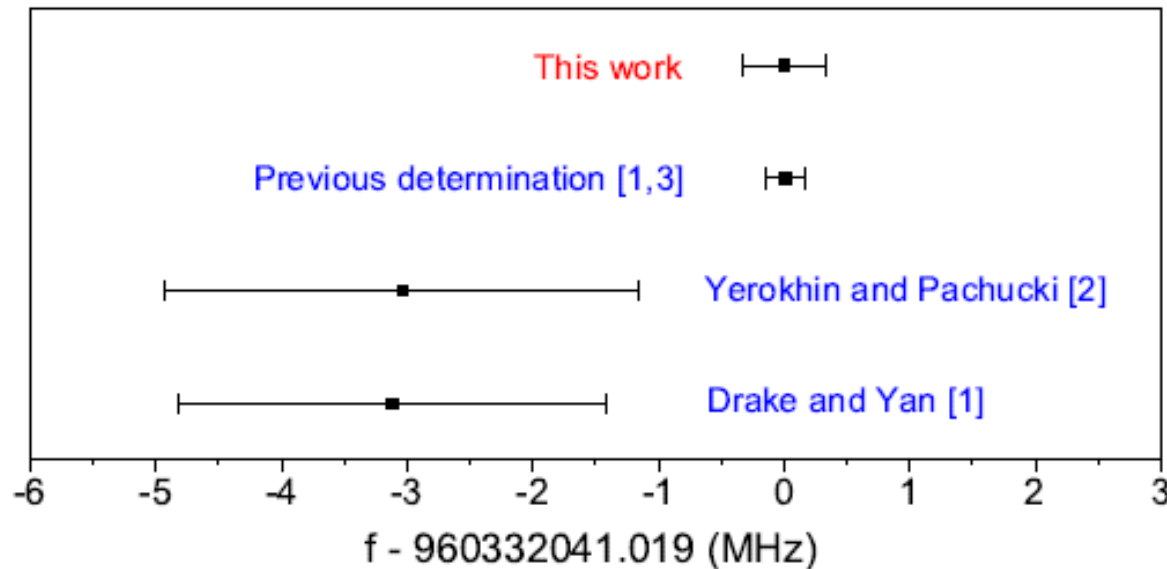
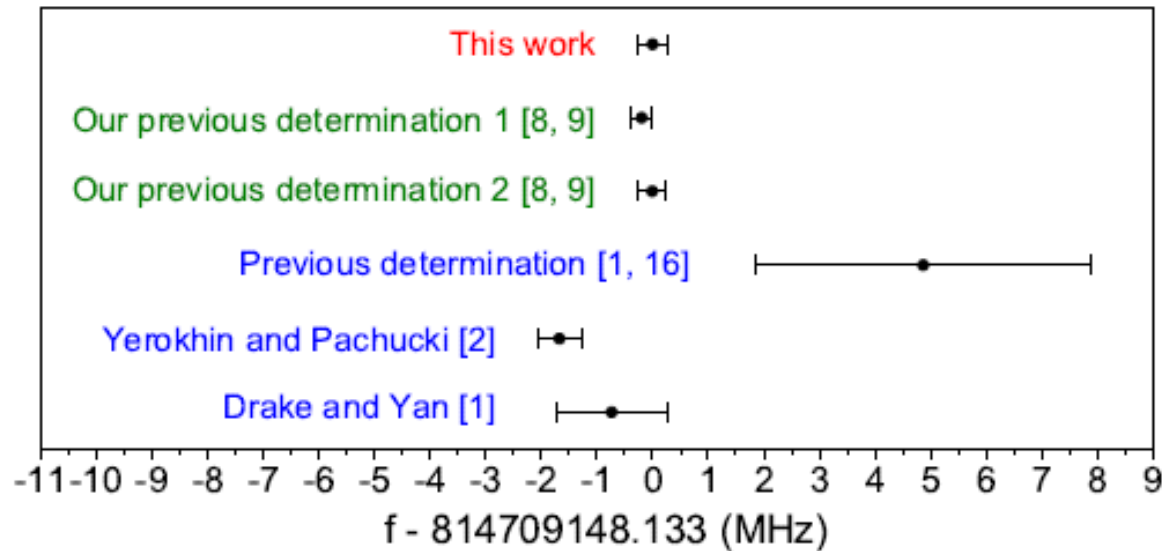


Source	shift	uncertainty
Statistical error	0	177
OFC accuracy	0	<5
Frequency locking of ECDL and OFC	0	<5
Pressure shift	0	192
Zeeman effect	0	<7
Overall	0	268

P.-L. Luo et al., *Phys. Rev. A* **88**, 054501 (2013)

10 time more precise than previous measurement

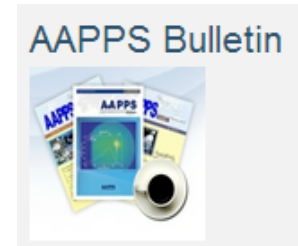
Ionization energy





OCTOBER 2013 VOL.23 NO.5

Reported in Physics Focus of AAPPS Bulletin



PHYSICS FOCUS

Precision Laser Spectroscopy of Helium Testing QED Atomic Calculations

Helium is the simplest multi-electron atom and has been the best testing ground for many-body QED atomic calculations. Unlike the hydrogen atom, for which analytical solutions exist, similar studies of helium require extensive numerical calculations in order to determine its electronic structures. Precision laser spectroscopy of He can improve the theoretical value of Lamb shift and be used to determine the nuclear charge radii of helium.

Recently, researchers from a multi-

had to control precisely many experimental parameters, e.g., laser power stability, magnetic field shielding, discharge condition, and helium gas pressure. Otherwise, any systematic error will limit the final precision of the measurement. A typical spectrum is shown in Fig. 2. This represents the first Doppler-free measurement of the $2^1S_0 \rightarrow 2^1P_1$ transition.

For the ionization energy of the 2^1P_1 state, a discrepancy of 3.5σ with the most precise theoretical value is found. This is shown in Fig. 3. This

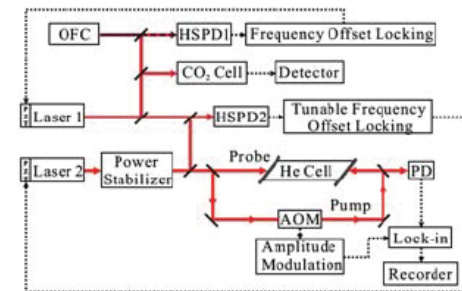
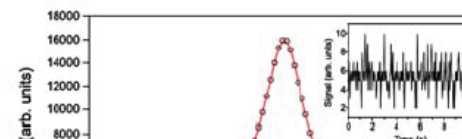
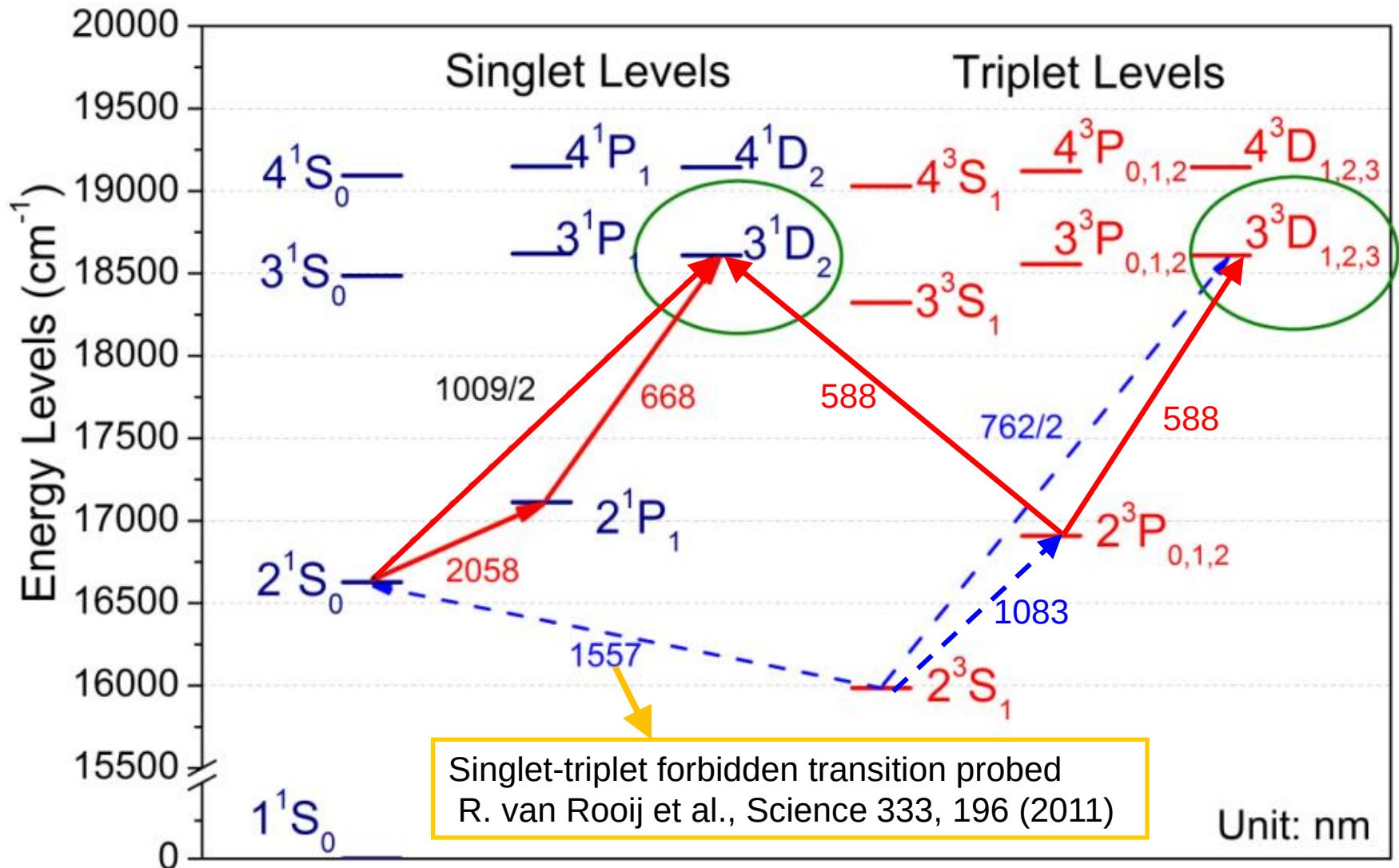


Fig. 1: Schematic of the experimental setup. OFC: optical frequency comb; PD: photodetector; HSPD: high-speed photodetector; AOM: acousto-optical modulator.



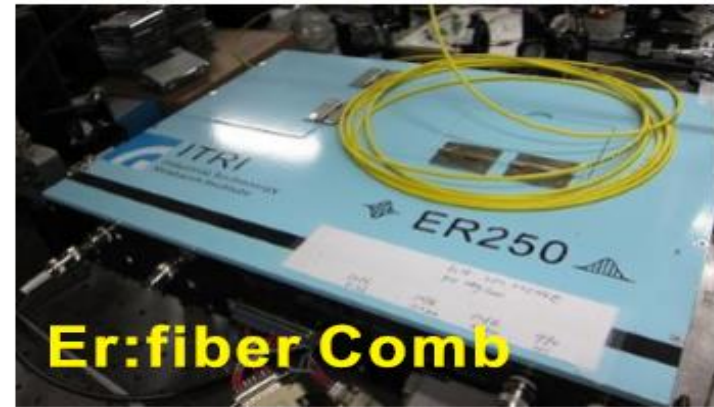
Helium energy levels



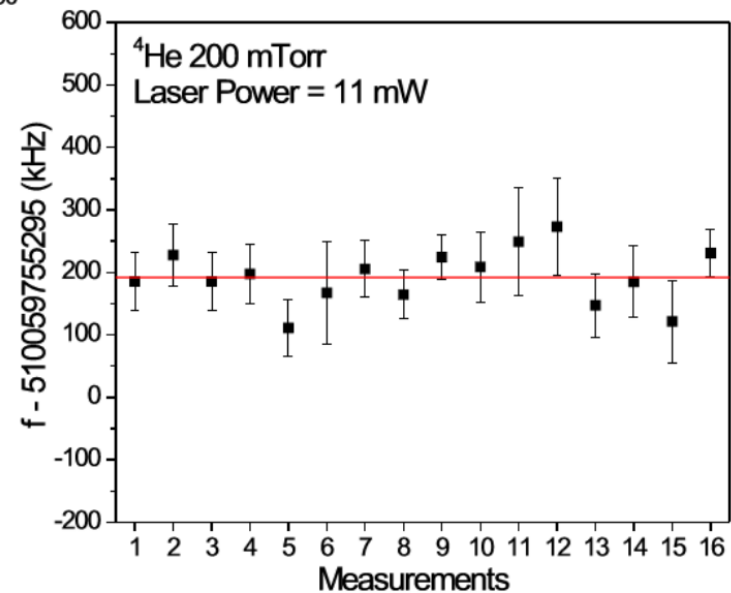
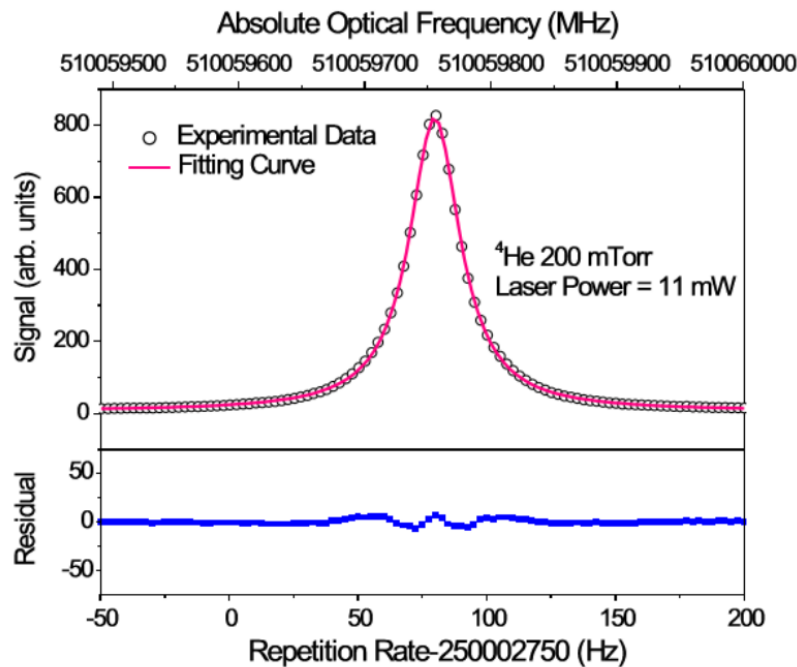
588 nm light source



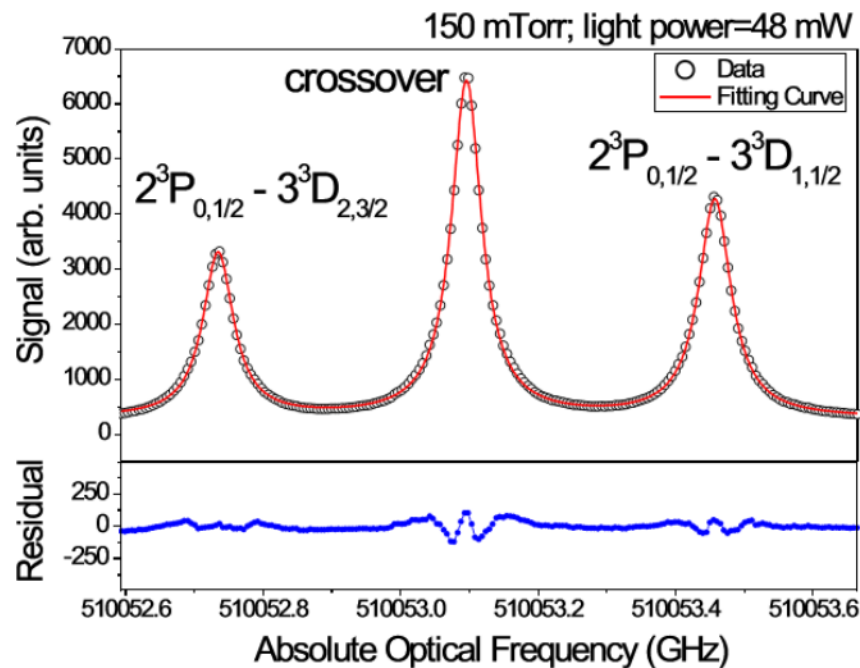
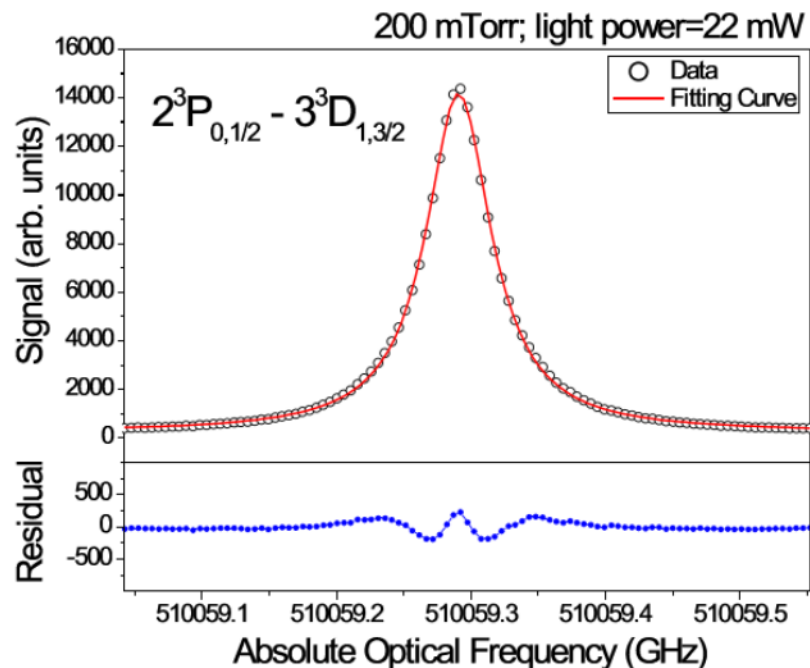
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P.-L. Luo et al., Applied Phys. B **120**, 279 (2015)

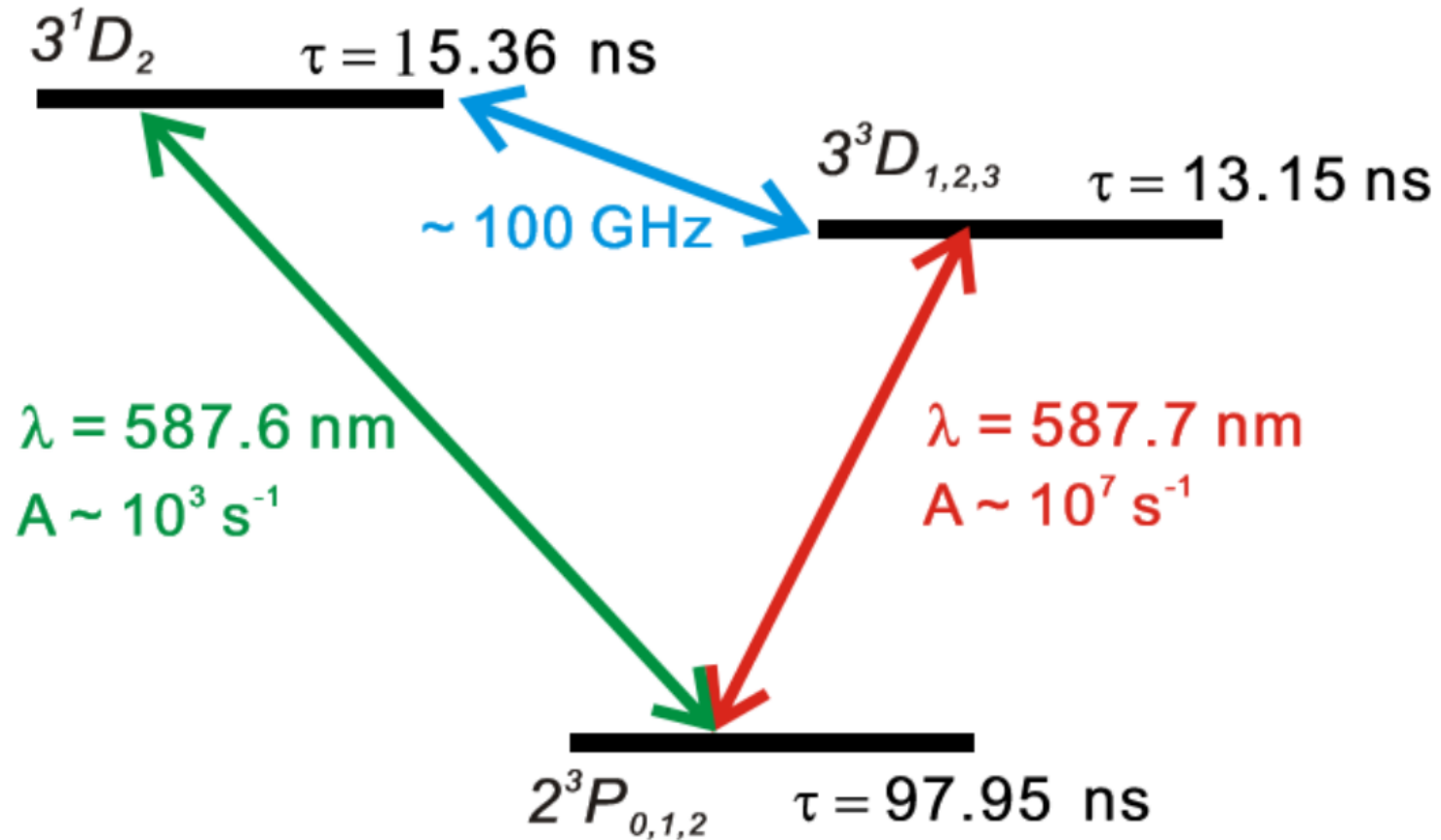


Helium-3



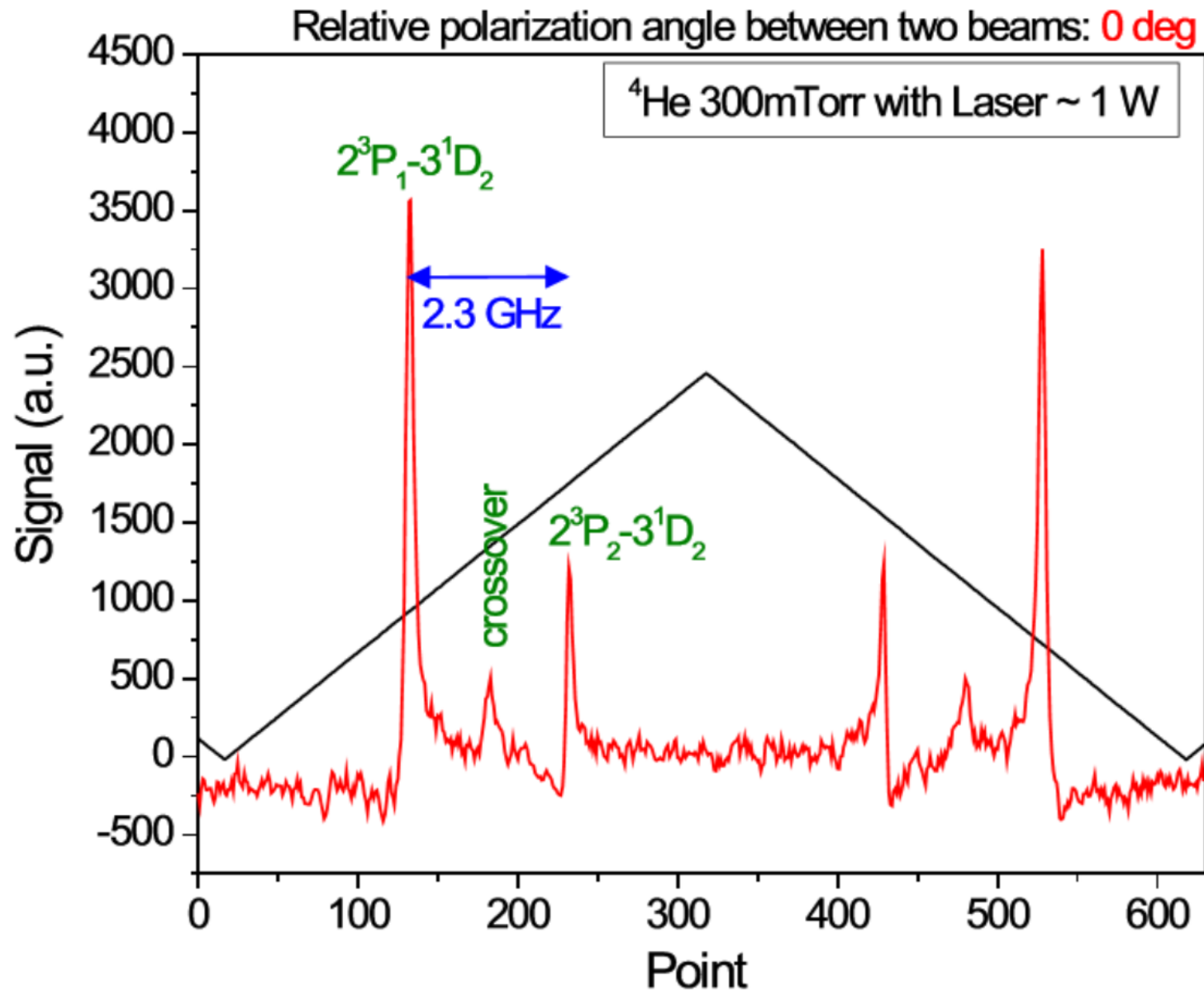
Separation	Our work	Theory [5]
$^3\text{He } 3^3D_{1,3/2} - 3^3D_{1,1/2}$	5834.073 (0.134)	5834.03 (0.04)
$^3\text{He } 3^3D_{1,3/2} - 3^3D_{2,3/2}$	6556.458 (0.149)	6556.42 (0.03)
$^3\text{He } 3^3D_{1,1/2} - 3^3D_{2,3/2}$	722.385 (0.183)	722.39 (0.04)

Spin-forbidden transition



- Only Doppler-broadened spectrum observed before

Spin-forbidden transition

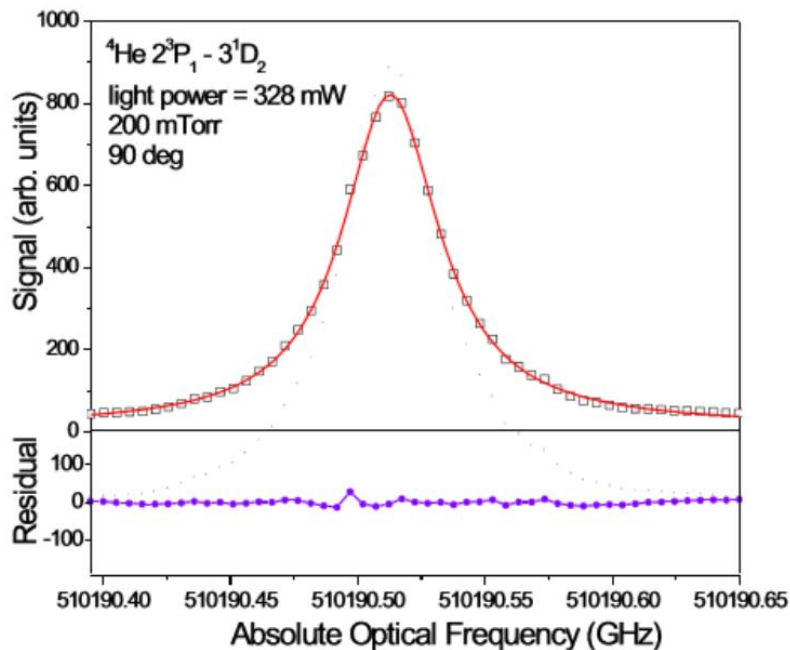


Spin-forbidden transition

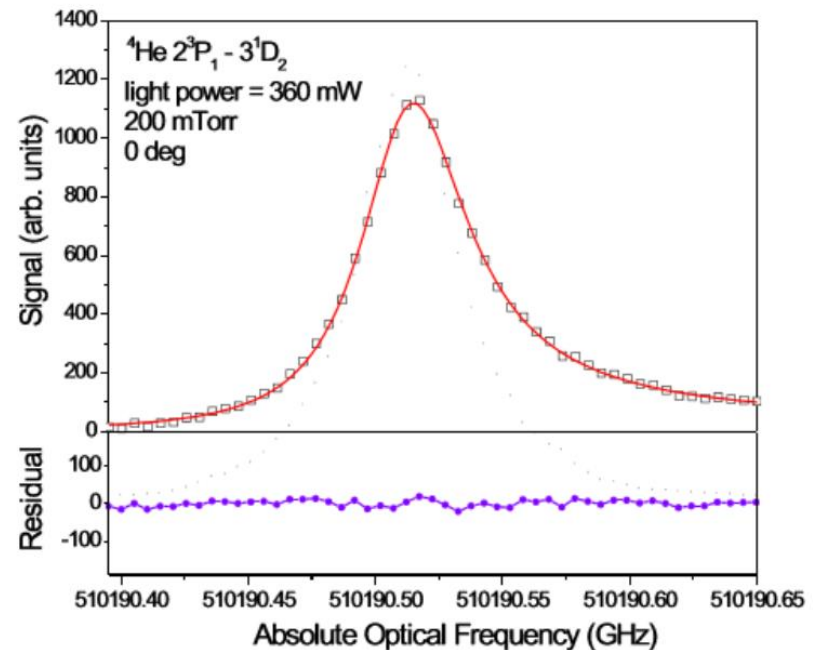


$^4\text{He } 2^3\text{P}_1 - 3^1\text{D}_2 @ 200 \text{ mTorr}$

$\theta = 90 \text{ deg}$



$\theta = 0 \text{ deg}$



Fitting Formula

$$S = A[L(\nu - \nu_c) + B \times D(\nu - \nu_c)]$$

Lorentzian

Dispersion

Analysis helium-4



$^4\text{He } 2^3\text{P}_1\text{-}3^1\text{D}_2$ Transition	MHz
Statistical value (zero light power at 300 mTorr)	510 190 516.631 (0.796)
Pressure shift	0 (1.0)
OFC accuracy	0 (0.002)
Frequency locking of ECDL and OFC	0 (0.002)
Second-order Doppler shift	0.004 (0.001)
Zeeman effect (residual magnetic fields < 1 mG)	0 (1)
Final result (this work)	510 190 516.635 (1.796)
Previous measurement (Doppler-limited) [5]	510 190 516.5 (12.0)
Experimental determination [1-4]	510 190 516.412 (0.325)
Theoretical calculation [6]	510 190 515.98 (2.0)

Analysis helium-3

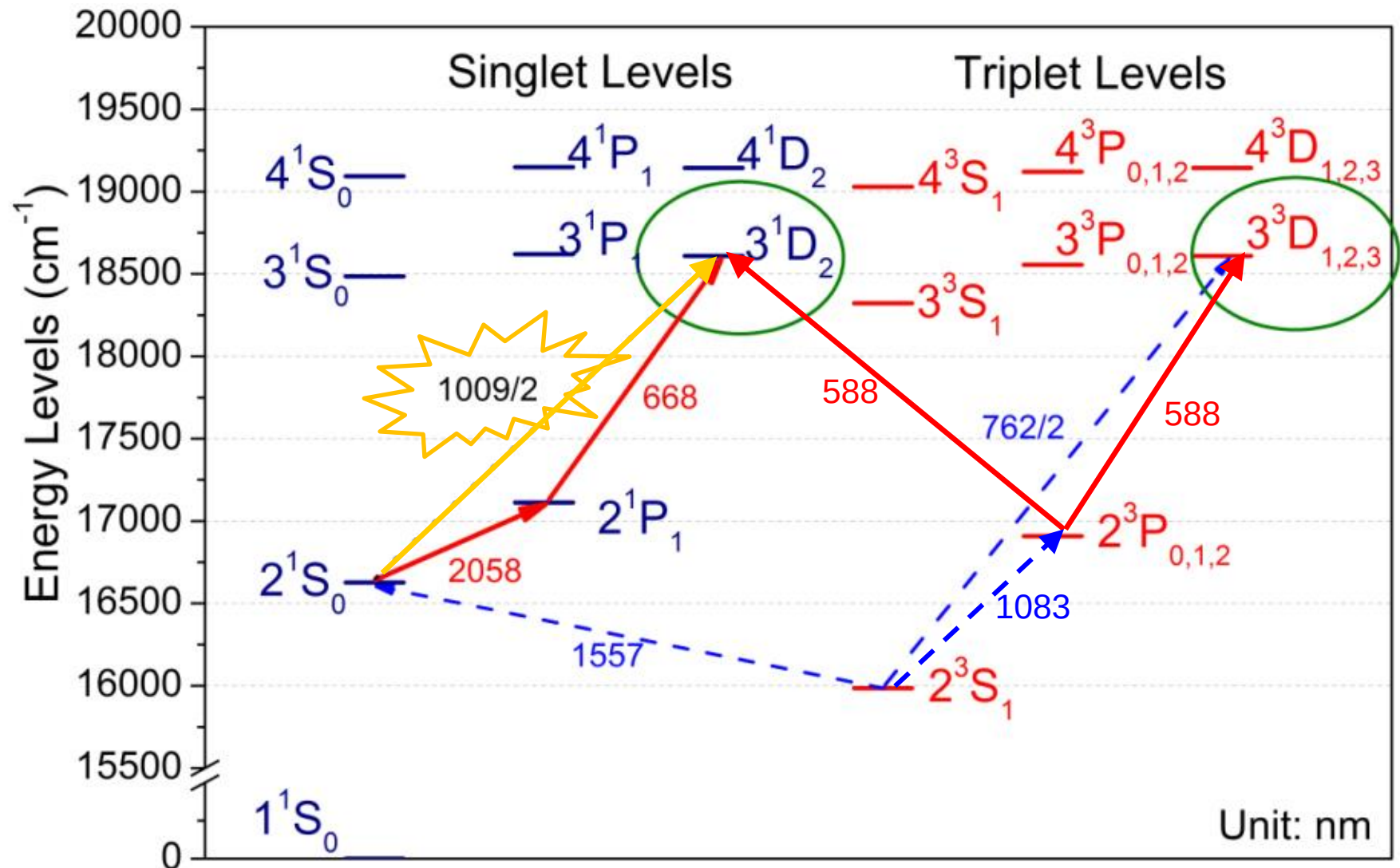


$^3\text{He } 2^3\text{P}_{0,1/2} - 3^1\text{D}_{2,3/2}$ Transition	MHz
Statistical value (zero light power at 300 mTorr)	510 156 738.642 (1.074)
Pressure shift	0 (1.0)
OFC accuracy	0 (0.002)
Frequency locking of ECDL and OFC	0 (0.002)
Second-order Doppler shift	0.004 (0.001)
Zeeman effect (residual magnetic fields < 1 mG)	0 (1)
Final result (this work)	510 156 738.646 (2.074)
Theoretical calculation [6]	510 156 732.24 (2.0)

Separation	Our work	Theory [6]
$^4\text{He } 3^1\text{D}_2 - 3^3\text{D}_1$	-101144.331 (1.796)	-101143.95 (0.03)
$^3\text{He } 3^1\text{D}_{2,3/2} - 3^3\text{D}_{1,3/2}$	97448.153 (2.075)	97448.15 (0.03)
$^3\text{He } 3^1\text{D}_{2,3/2} - 3^3\text{D}_{1,1/2}$	103282.226 (2.078)	103282.18 (0.04)
$^3\text{He } 3^1\text{D}_{2,3/2} - 3^3\text{D}_{2,3/2}$	104004.611 (2.079)	104004.57 (0.03)

- The 37 MHz (7 σ) discrepancy not seen here.

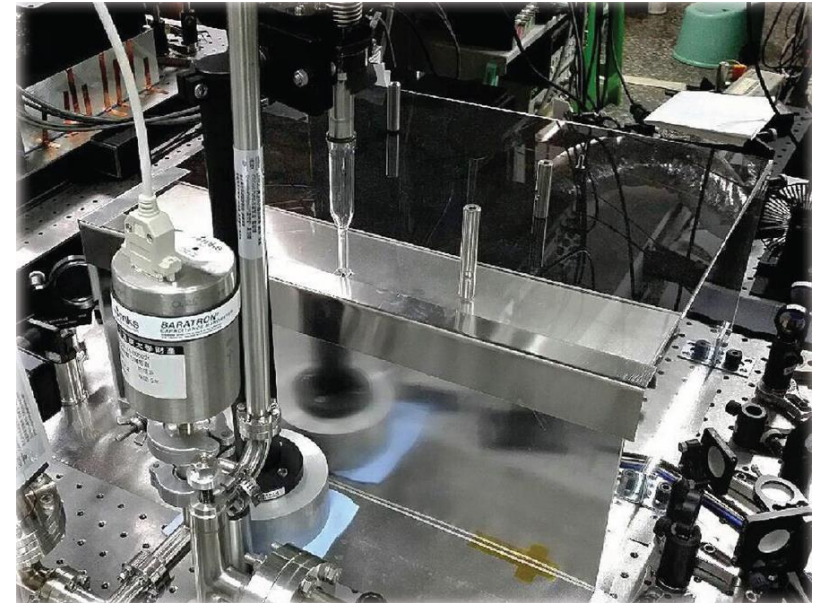
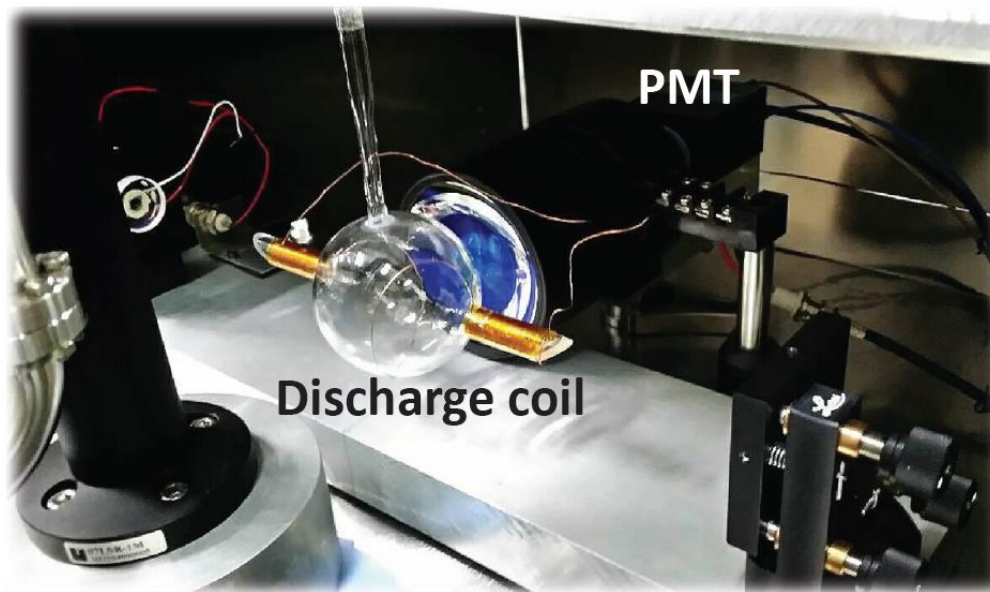
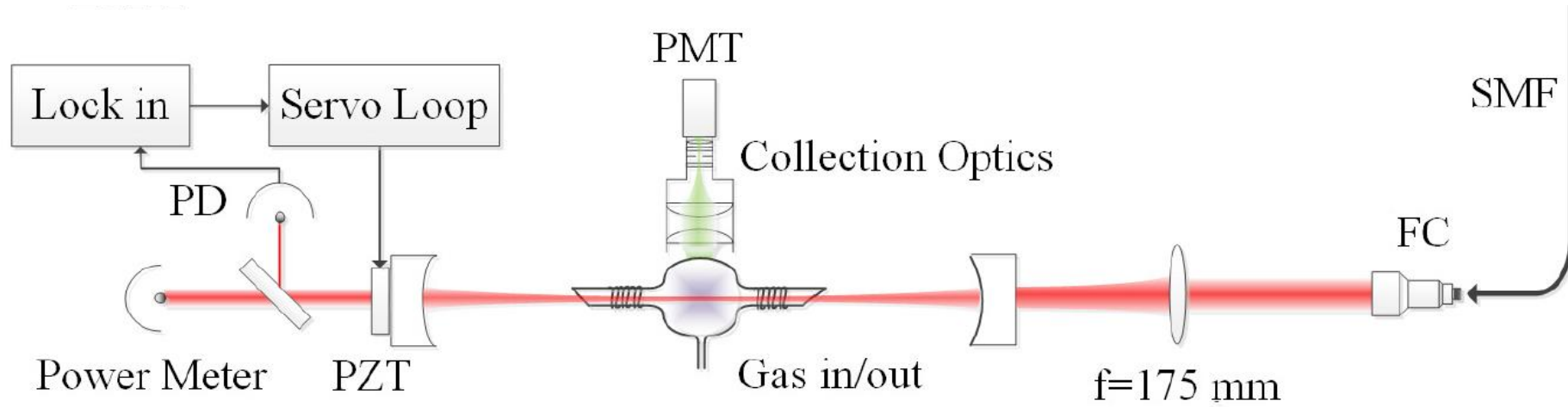
Helium energy levels

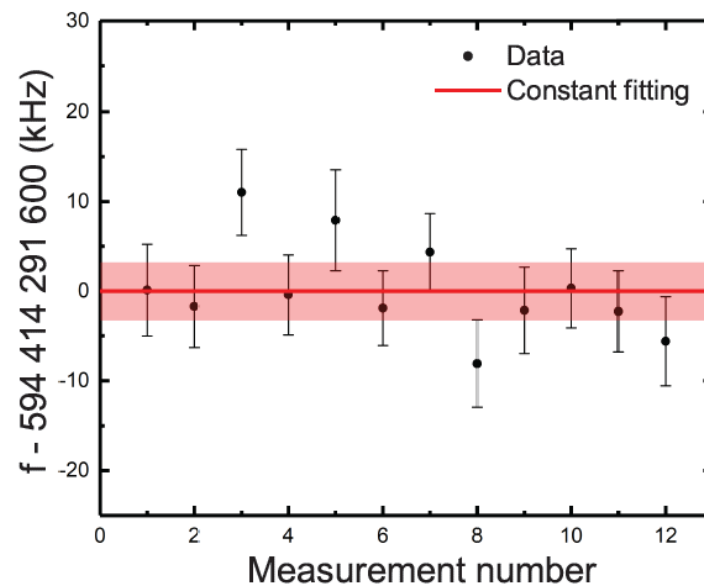
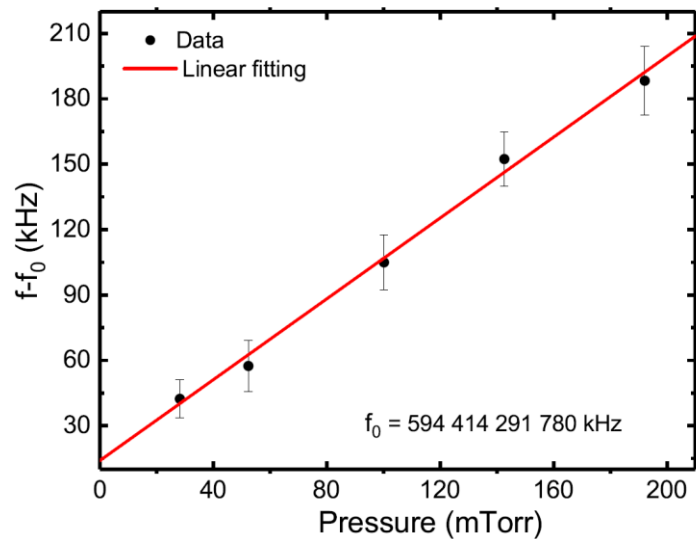
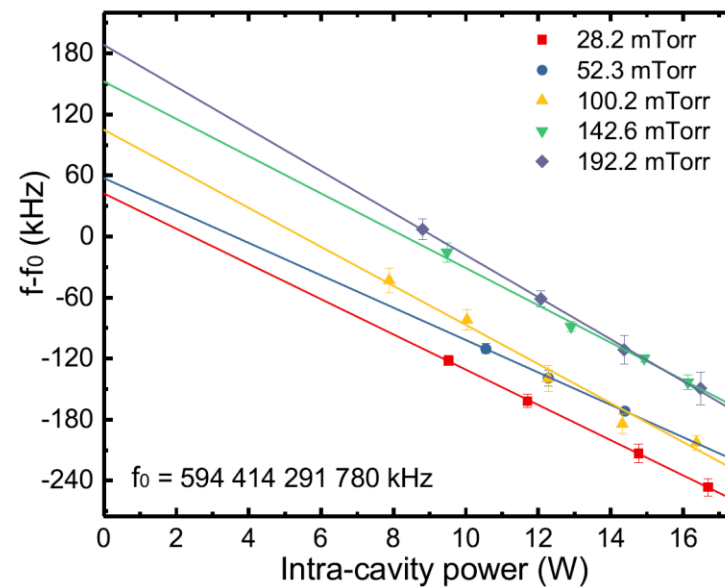
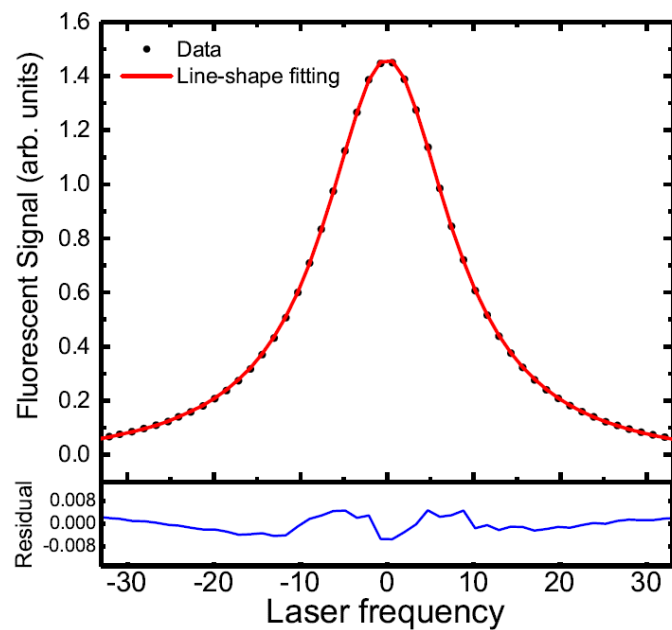


1009 nm two-photon transition



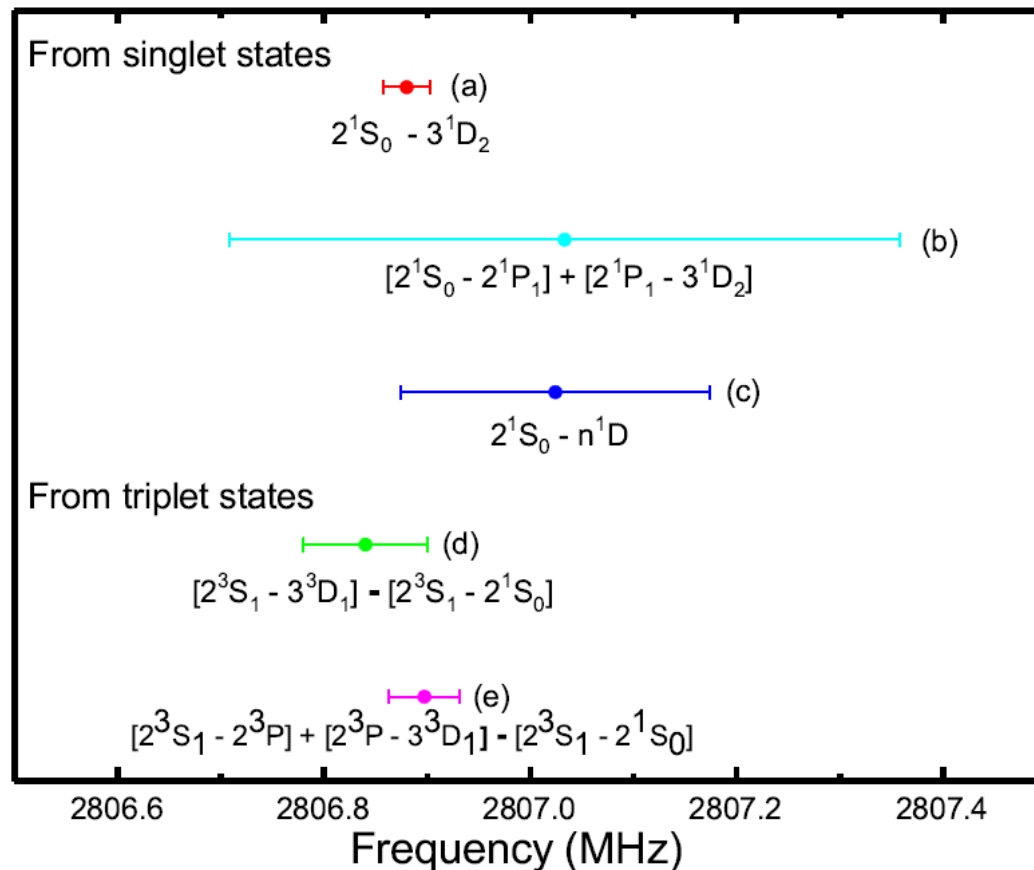
DEPARTMENT OF PHYSICS
National Tsing Hua University





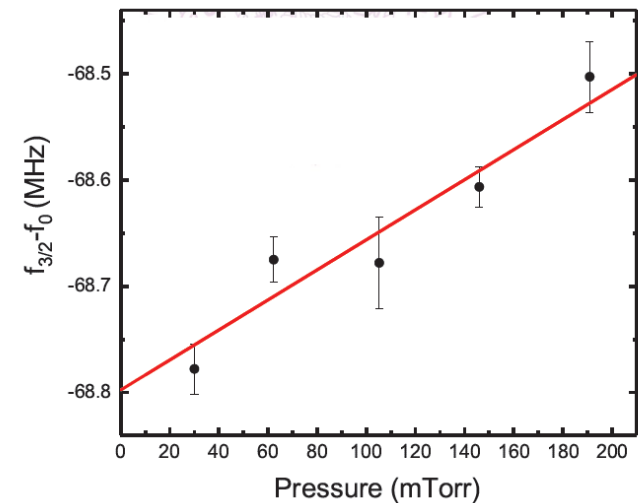
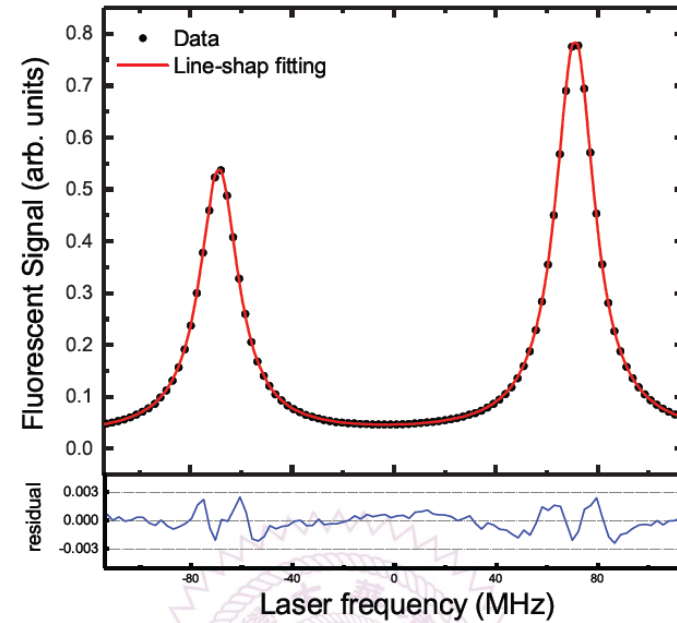
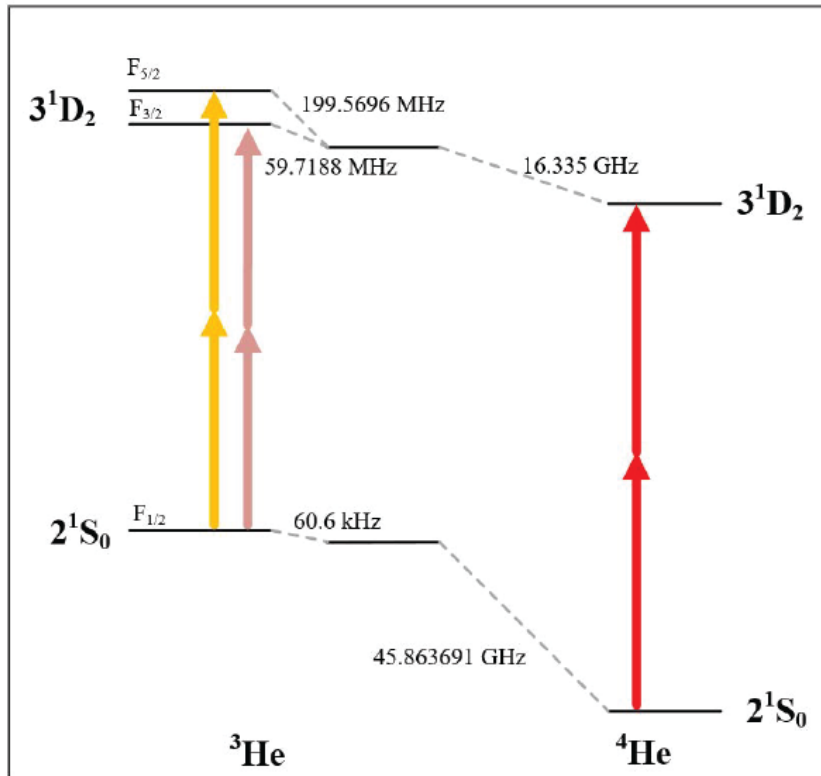


Lamb shift of the singlet 2S state



Y.-J. Huang, Y.-C. Guan, Y.-C. Huang, T.-H. Suen, J.-L. Peng, L.-B. Wang, and J.-T. Shy,, *Phys. Rev. A* **97**, 032516 (2018)

Isotope shift and helium-3



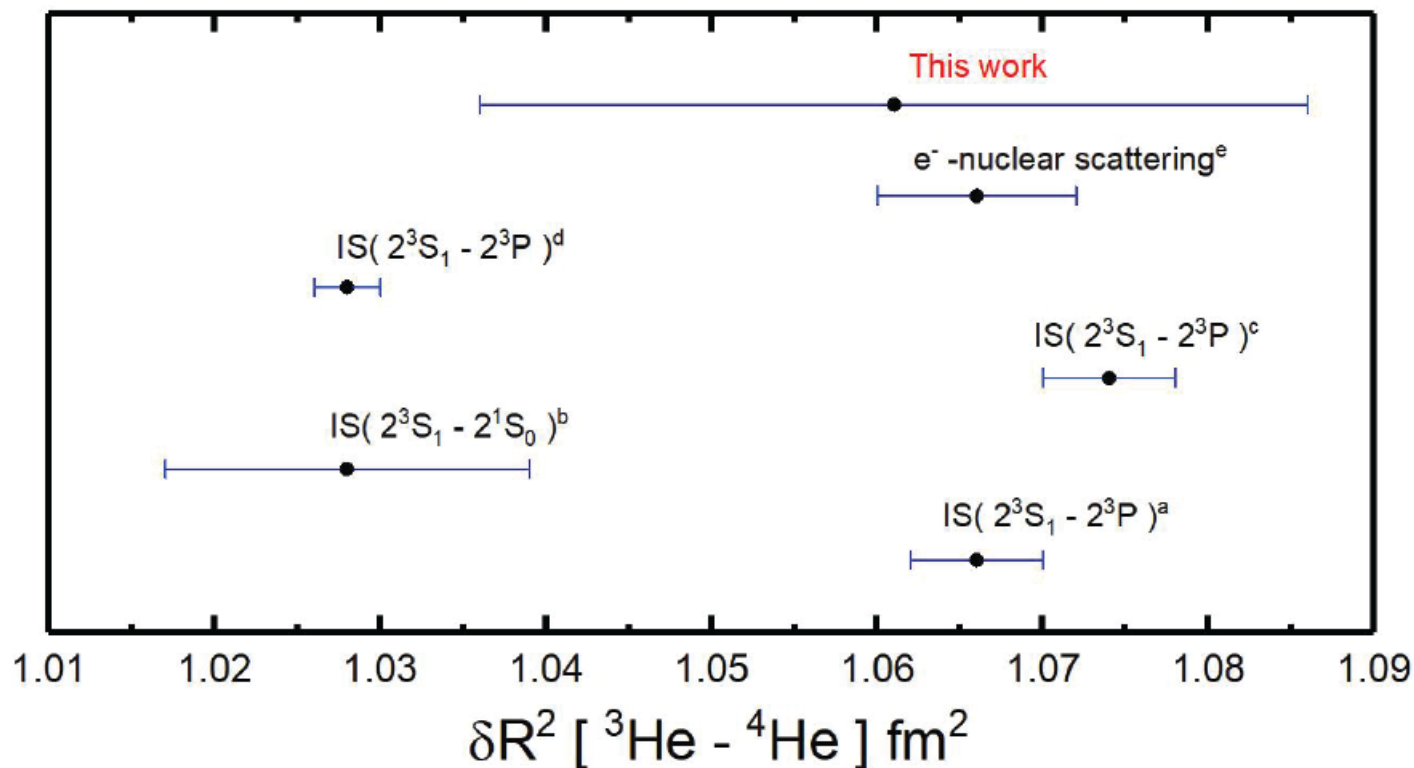


Figure 5.13: The comparison of the determination of the δR^2 from our work and other measurements. (a) Shiner et al. [61]. (b) Cancio Pastor et al. [18, 19, 56]. (c) Van Rooij et al. [42]. (d) X. Zeng et al. [46]. (e) Ingo Sick [62].

Work completed

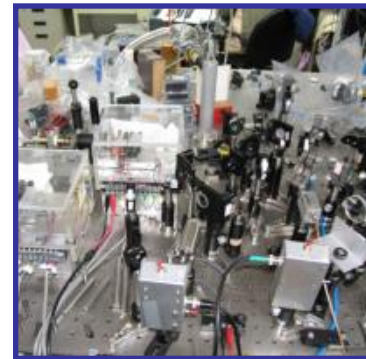
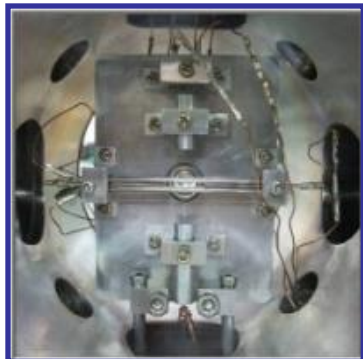
- Helium singlet 2S-2P at 2058nm, 2P-3D at 668 nm
- Helium triplet 2P to 3D (singlet and triplet) at 588 nm
- Helium singlet 2S-3D two photon transition
- Either resolve discrepancy or test theories

Work in progress and future work

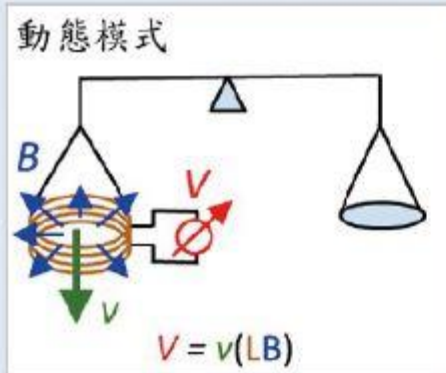
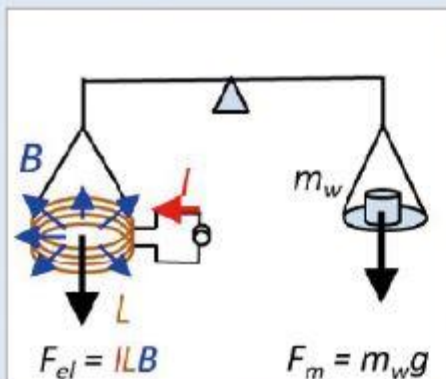
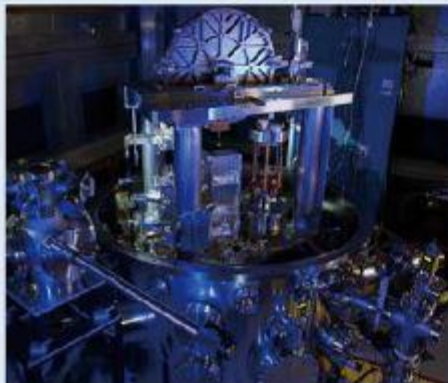
- Helium Rydberg spectroscopy
- Construction of metastable helium beam
- Upgrade of optical frequency comb

清華大學物理系: 施宙聰老師、孫德輝、官鈺禪、黃一展
原分所: 羅佩凌博士
工研院量測中心: 彭錦龍博士

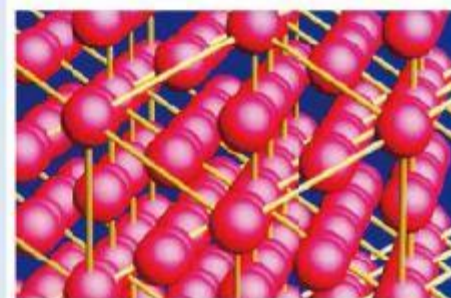
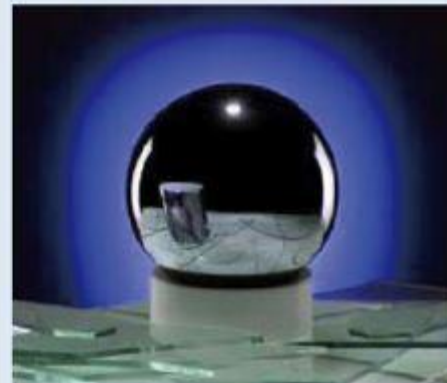
Supported by National Science Council, and Minister
of Education



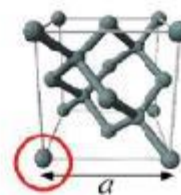
瓦特天平法



矽晶球法



矽原子數量 $N = \frac{V_{core}}{a^3} \times 8$



矽原子質量

$$m_a(^{28}\text{Si}) = \left(\frac{A_r(^{28}\text{Si})}{A_r(e)} \frac{2R_e}{c\alpha^2} \right) h$$