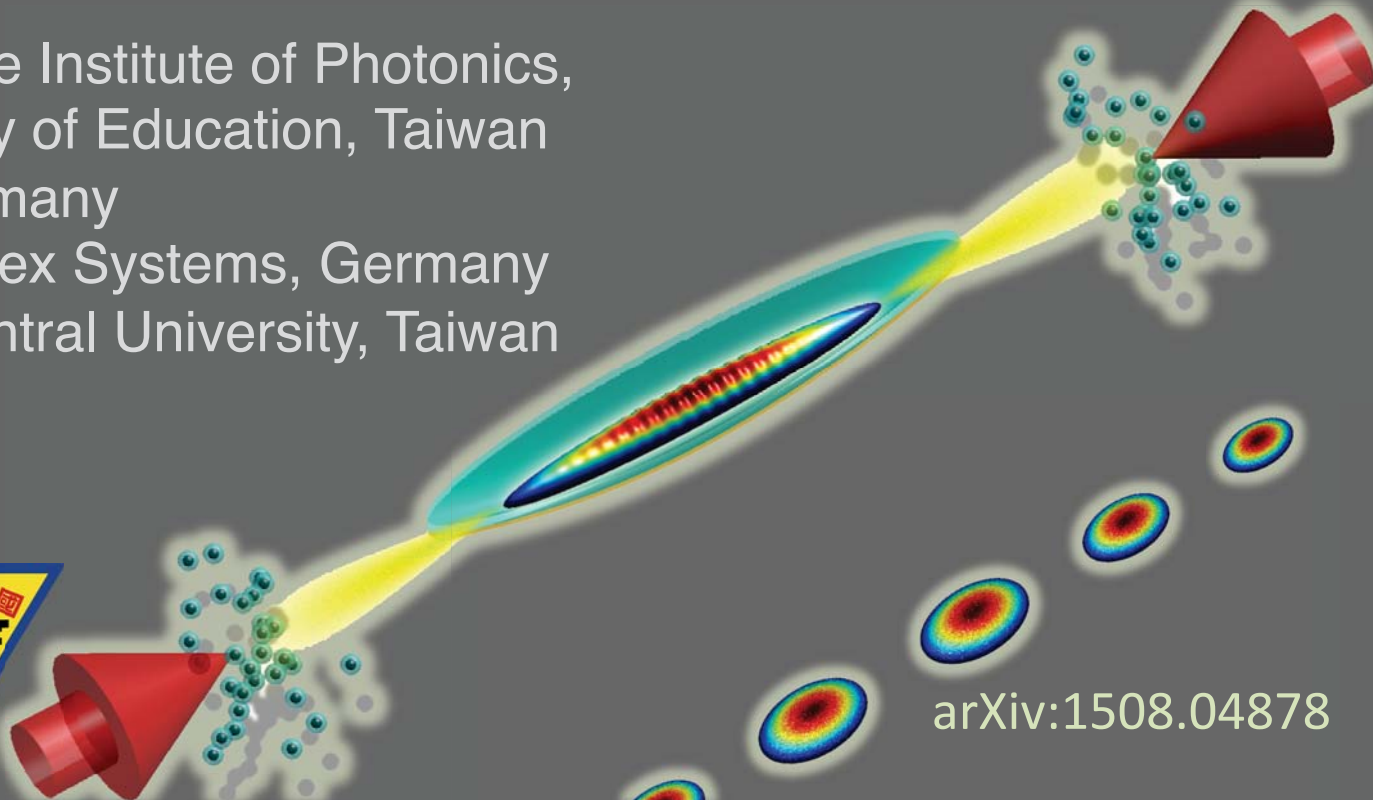


All-optical cavity using atomic mirrors

Shih-Wei Su(蘇士煒)¹, Zhen-Kai Lu(呂振凱)²,
Shih-Chuan Gou(郭西川)¹ and Wen-Te Liao(廖文德)^{3,4}

1. Dep. of Physics and Graduate Institute of Photonics, National Changhua University of Education, Taiwan
2. MPI of Quantum Optics, Germany
3. MPI for the Physics of Complex Systems, Germany
4. Dep. of Physics, National Central University, Taiwan

2015 AMO Summer School
2015/8/27 NTHU Taiwan

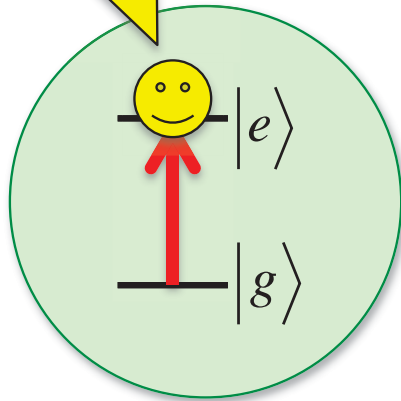


arXiv:1508.04878

Why Cavity?

Is eigenstate stationary?

Be here forever!



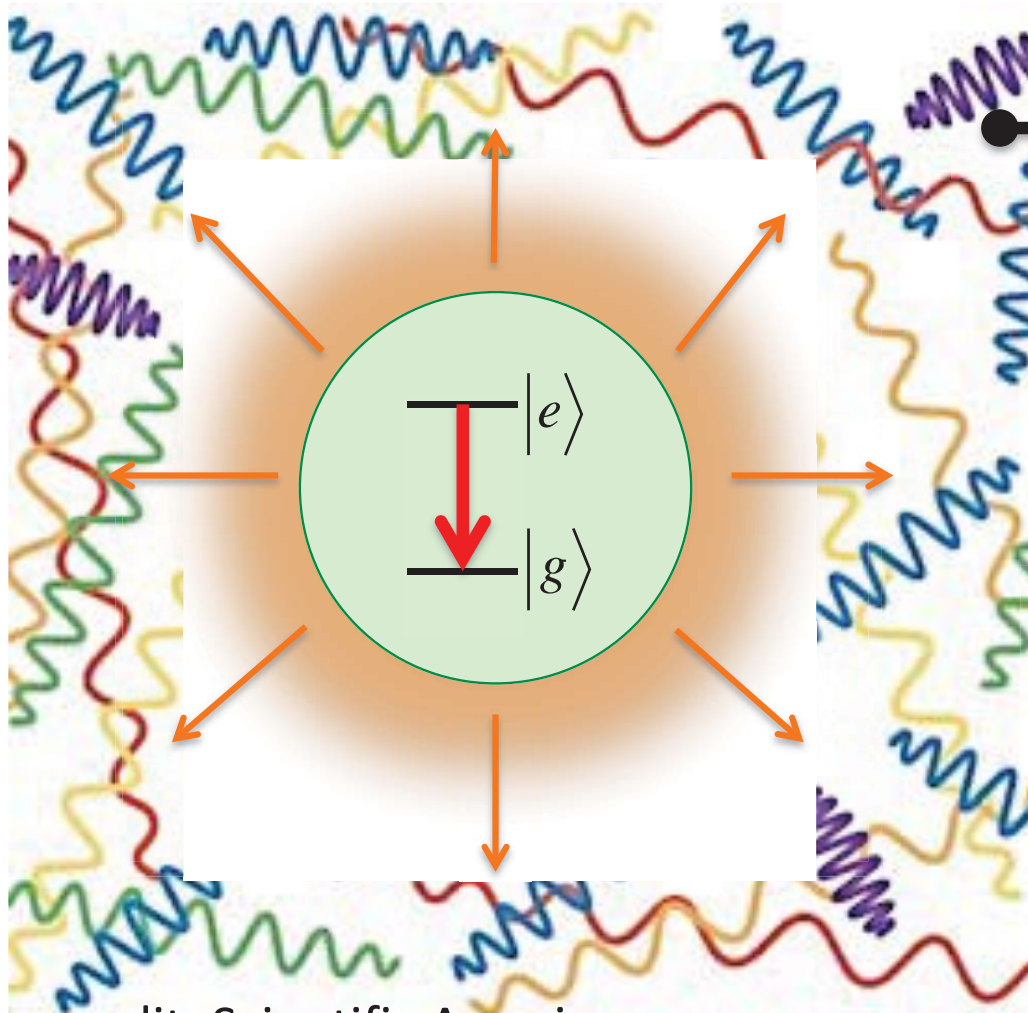
an atom

Schrödinger's Equation

$$i\hbar \frac{\partial}{\partial t} \psi(\mathbf{r}, t) = -\frac{\hbar^2}{2m} \nabla^2 \psi(\mathbf{r}, t) + V(\mathbf{r}, t) \psi(\mathbf{r}, t)$$

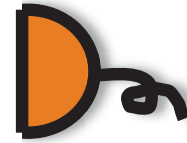
eigenstate???

something in vacuum

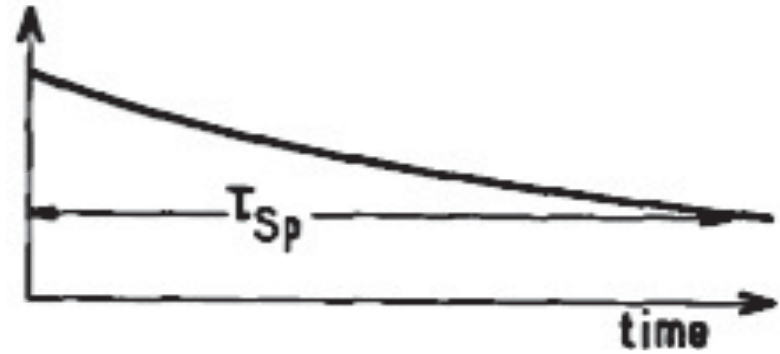


credit: Scientific American

Vacuum

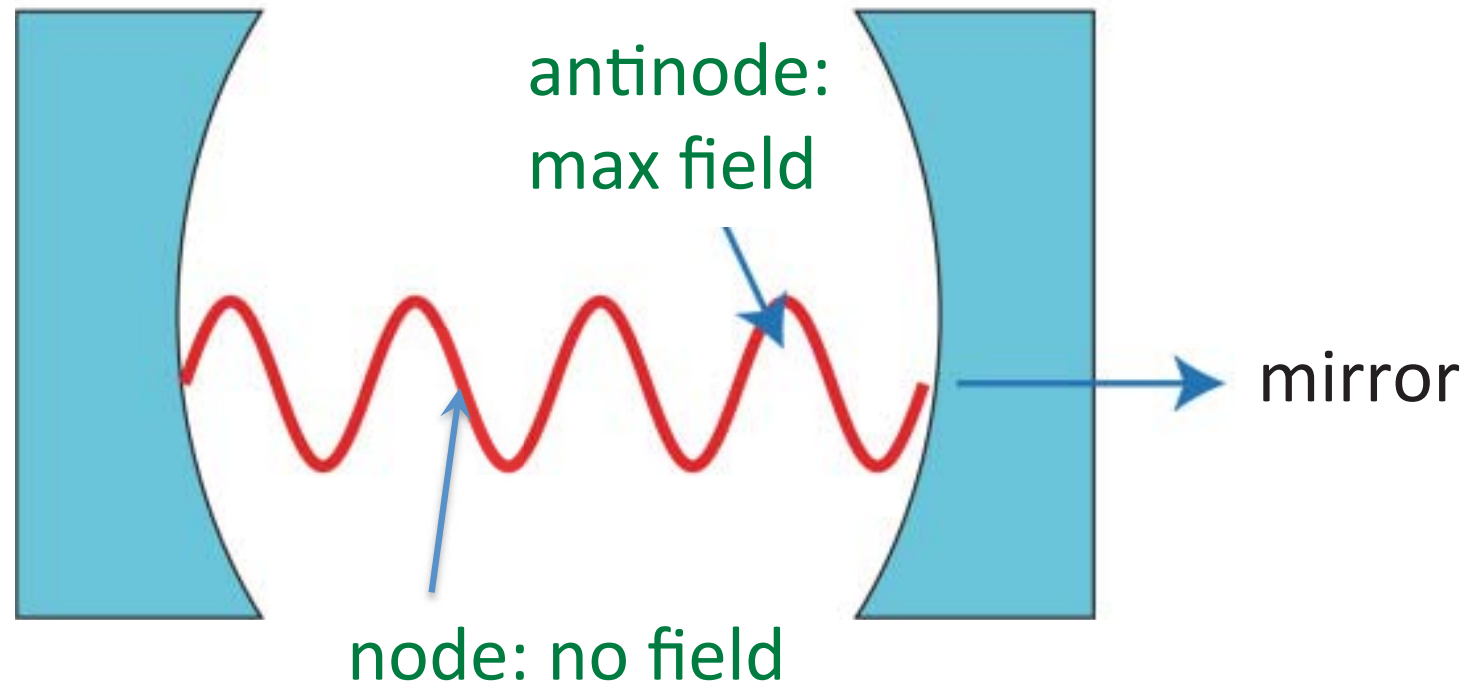


Detector

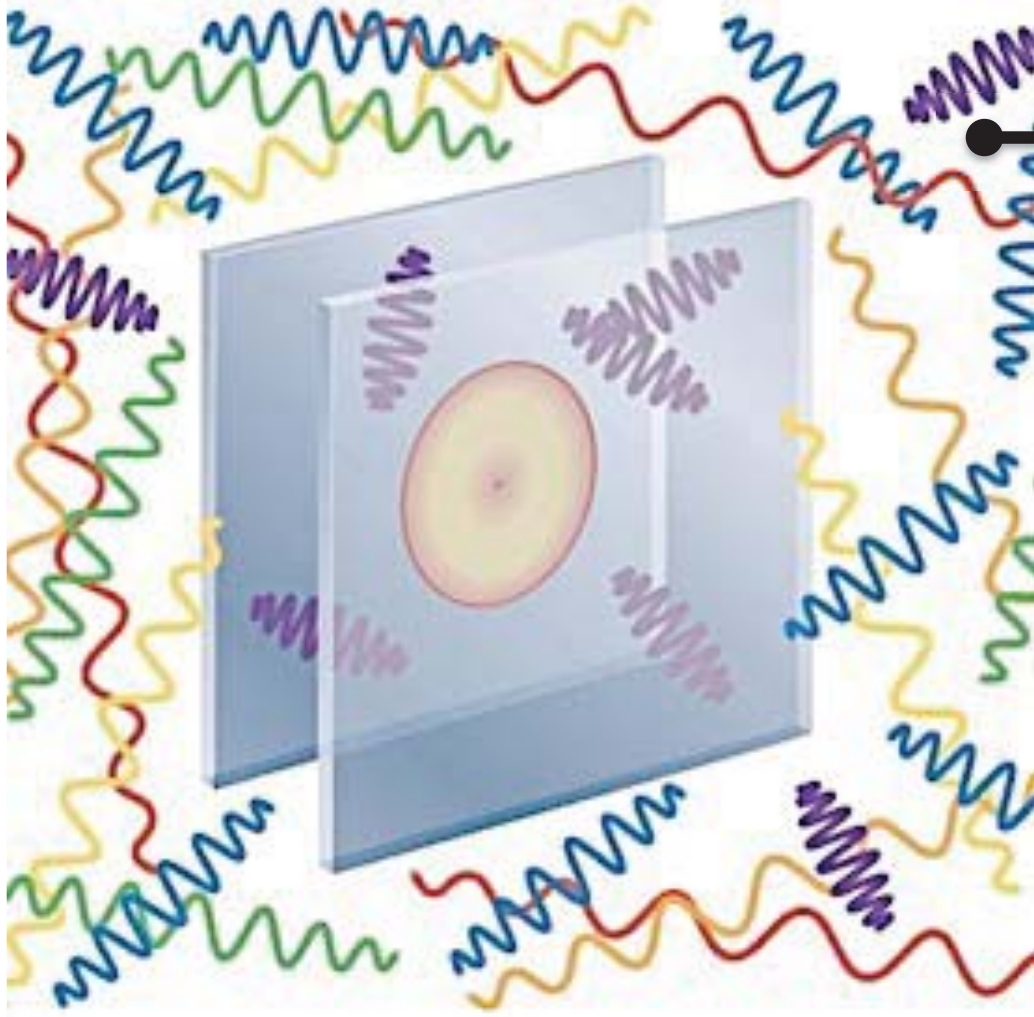


spontaneous decay

Standing wave in a cavity



Interaction in a cavity



Vacuum outside

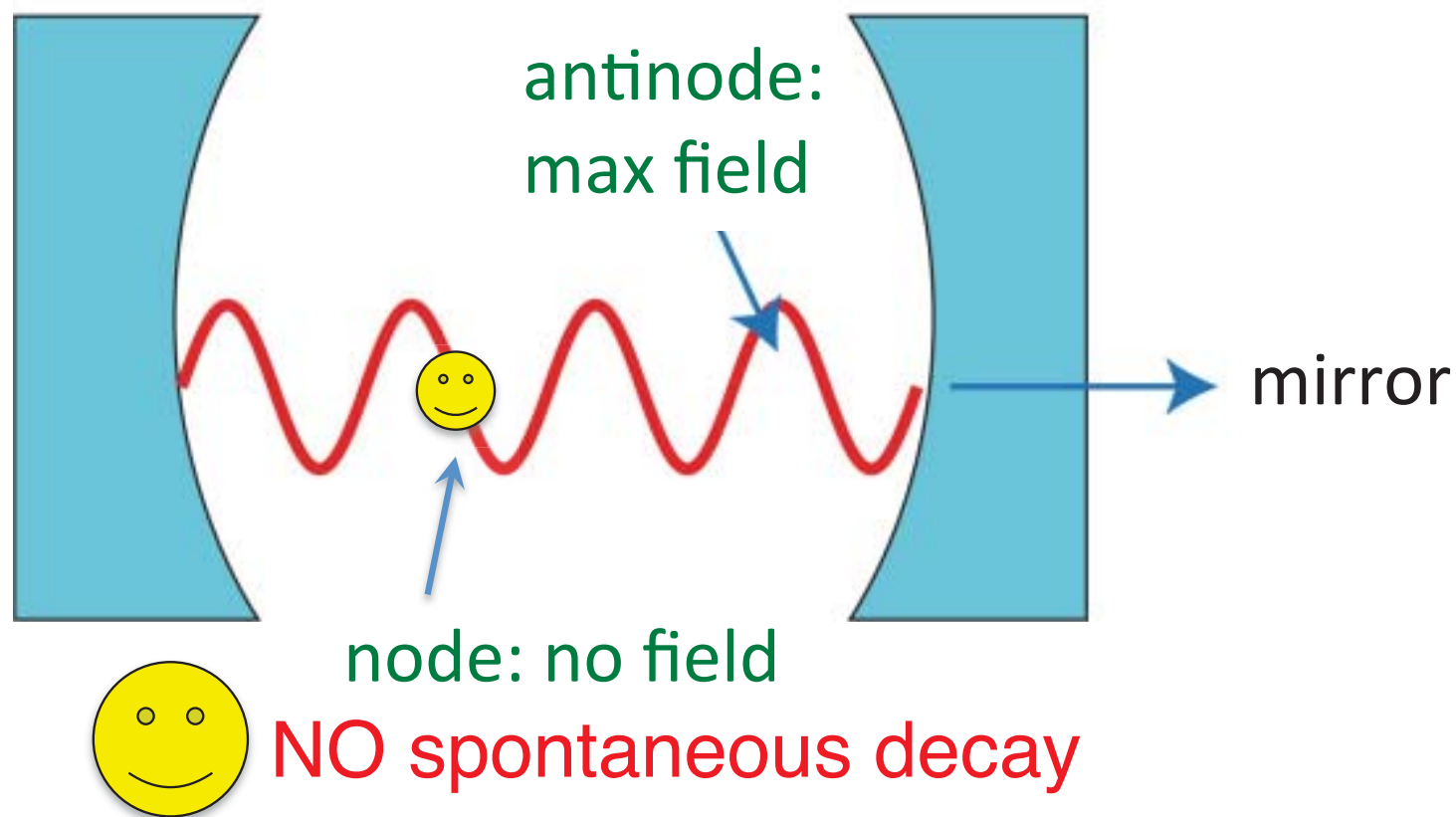
$$\hat{H}_{\text{JC}} = \hbar\omega_c \hat{a}^\dagger \hat{a} + \hbar\omega_a \frac{\hat{\sigma}_z}{2} + \frac{\hbar\Omega}{2} (\hat{a}\hat{\sigma}_+ + \hat{a}^\dagger\hat{\sigma}_-)$$

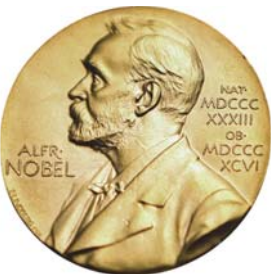
Jaynes–Cummings model

single mode photons v.s.
a single atom

credit: Scientific American

Cavity modifies atom's behavior

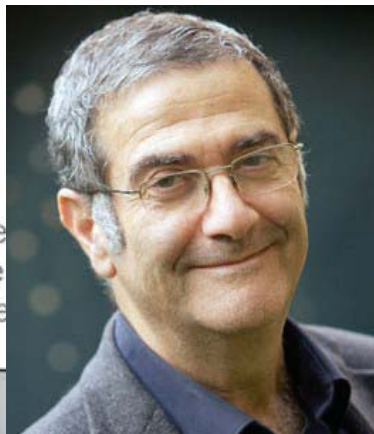
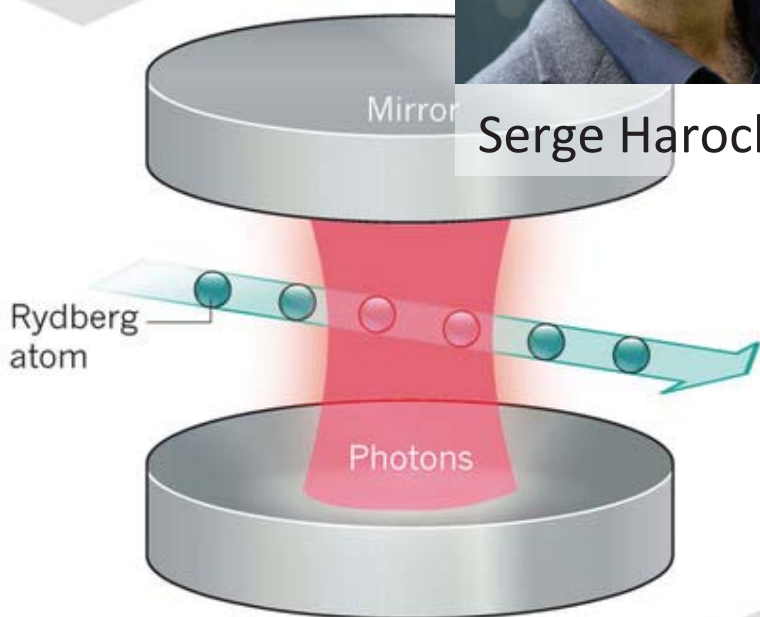




The Nobel Prize in Physics 2012

HAROUCHE METHOD

Microwave photons are placed between two reflective mirrors that enable them to bounce back and forth between the mirrors.

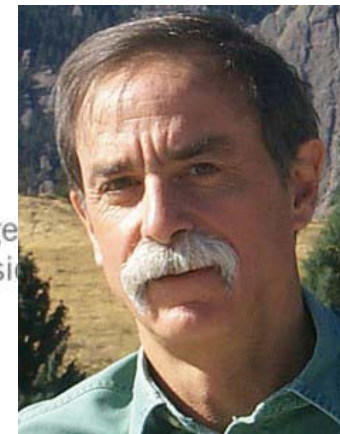
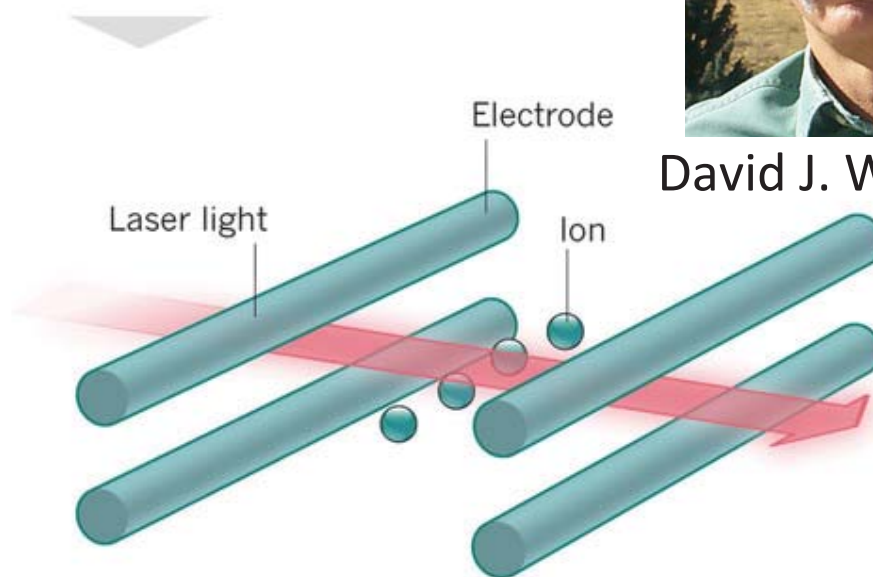


Serge Haroche

Rydberg atoms, which have one electron in a high-energy level, are sent through the system to measure and manipulate the photon's quantum state.

WINELAND METHOD

An electric field produced by an arrangement of electrodes holds one or several ions inside a Paul trap.



David J. Wineland

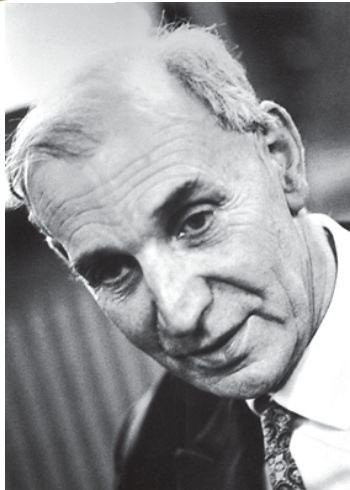
Ed Hinds & Rainer Blatt *Nature* 492, 55 (2012)

Laser light is shone on the ion, suppressing its thermal vibration and allowing its quantum state to be measured and controlled.

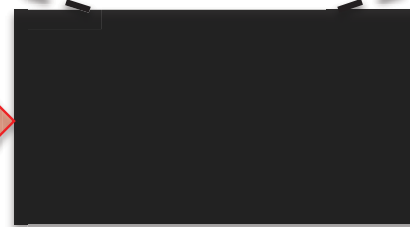
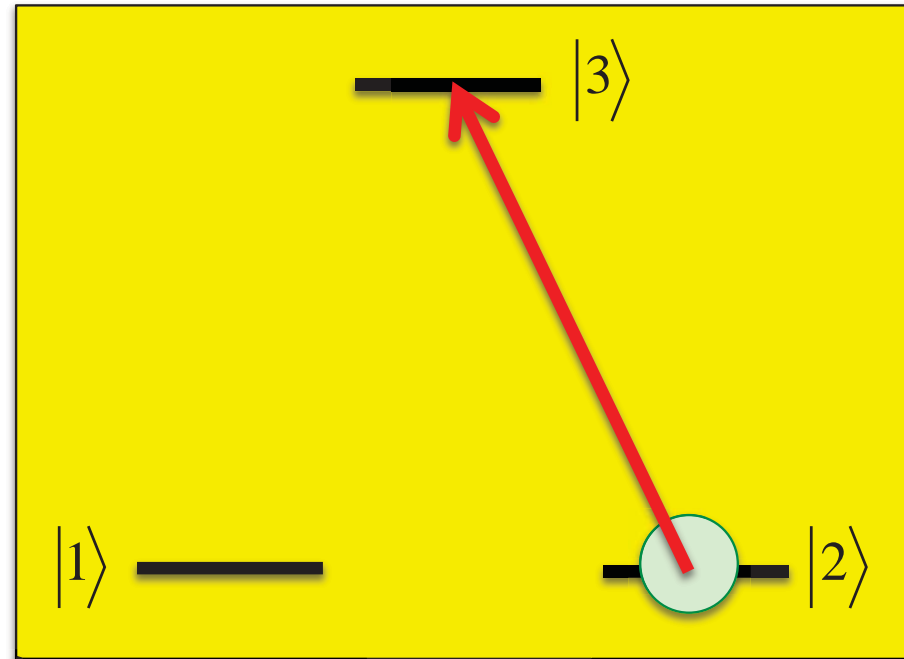
Control Matter with light



Optical pumping (1966 Nobel Prize)



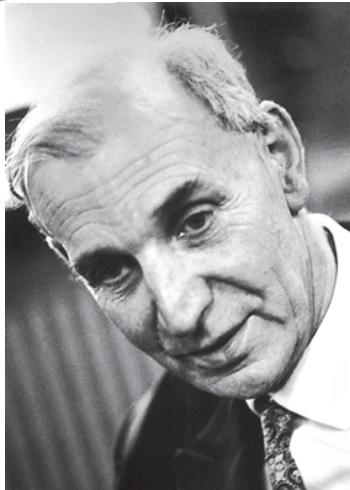
Alfred Kastler



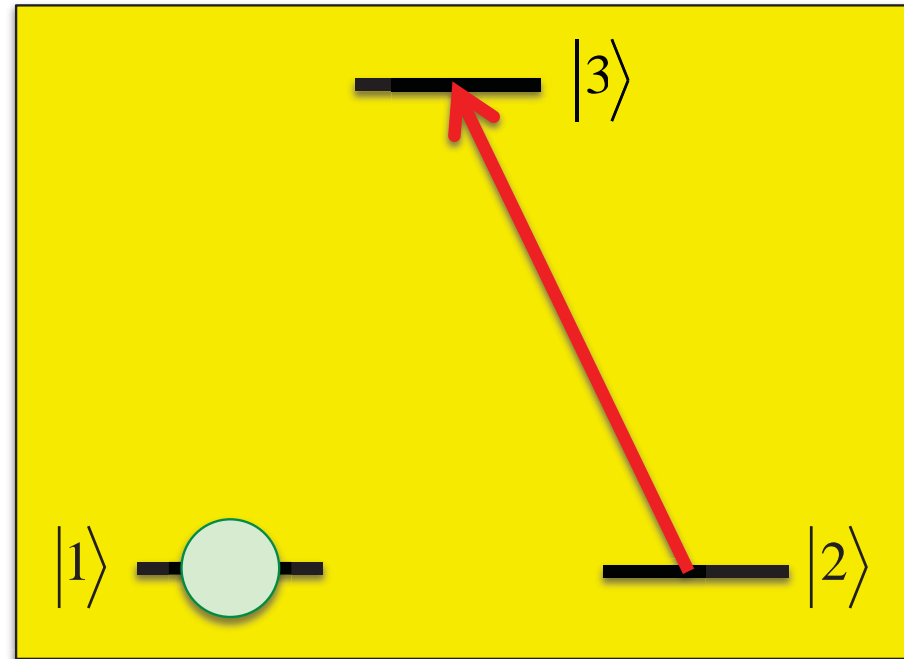
opaque



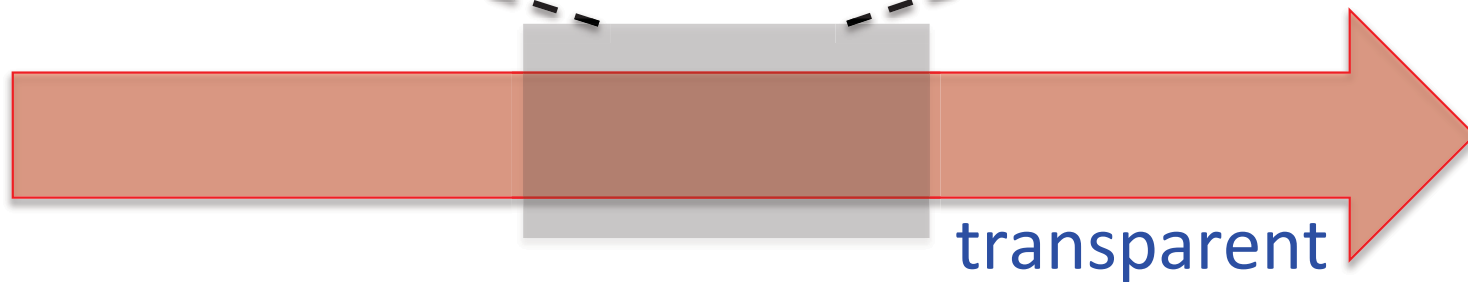
Optical pumping (1966 Nobel Prize)



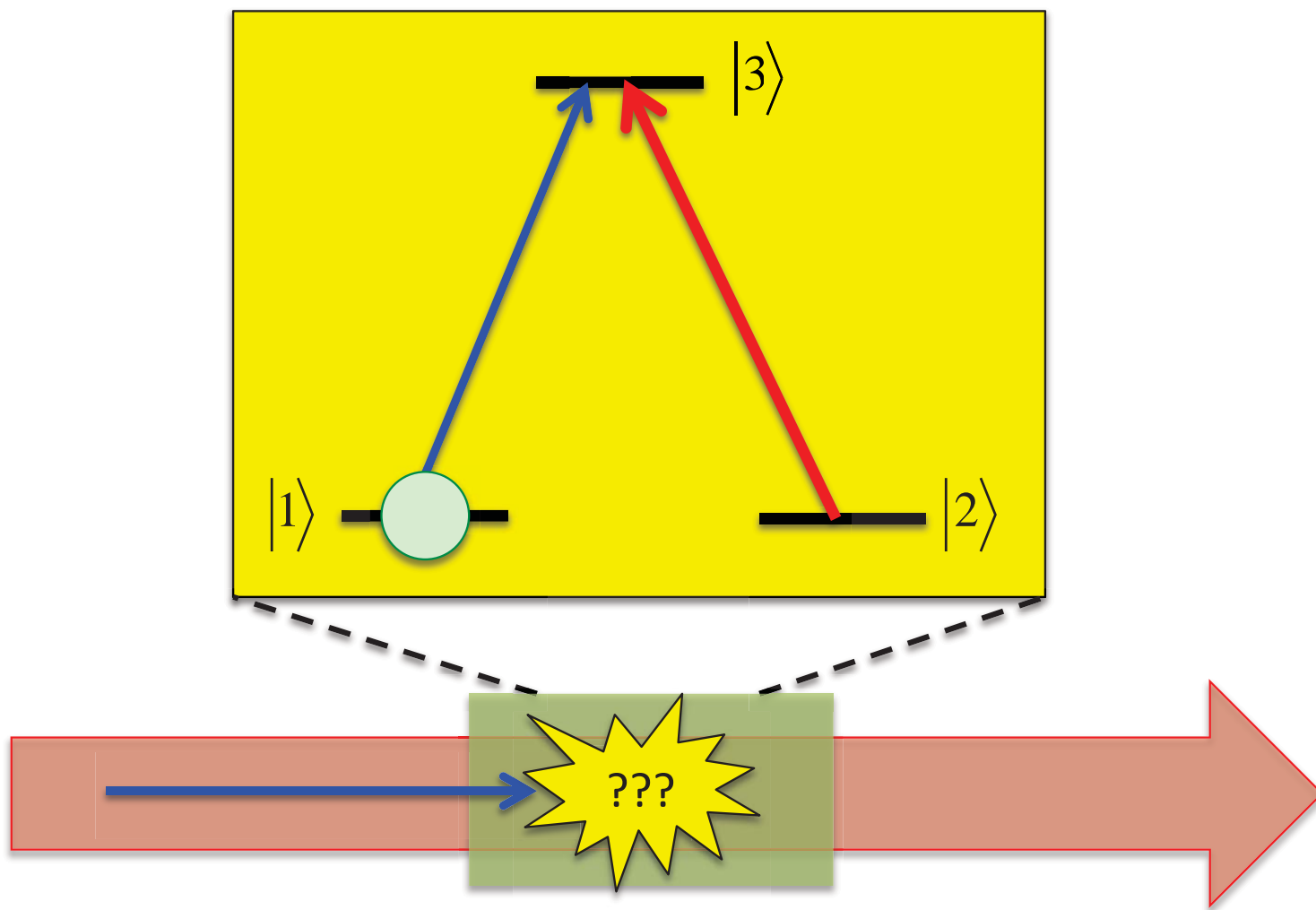
Alfred Kastler



$|1\rangle$ dark state



How about applying 2nd light?



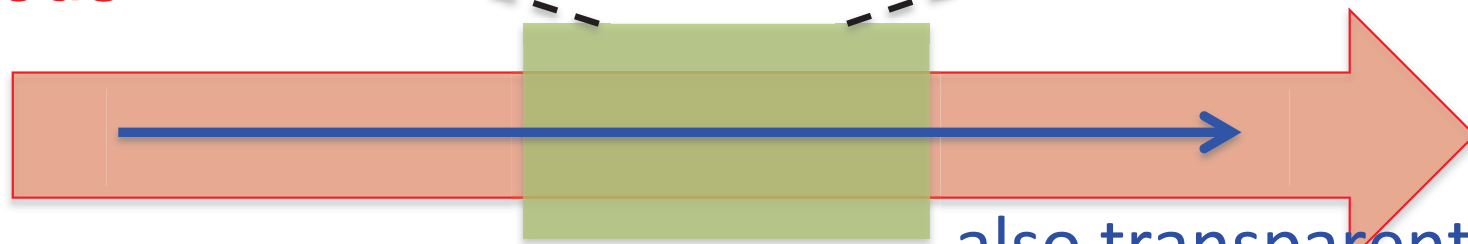
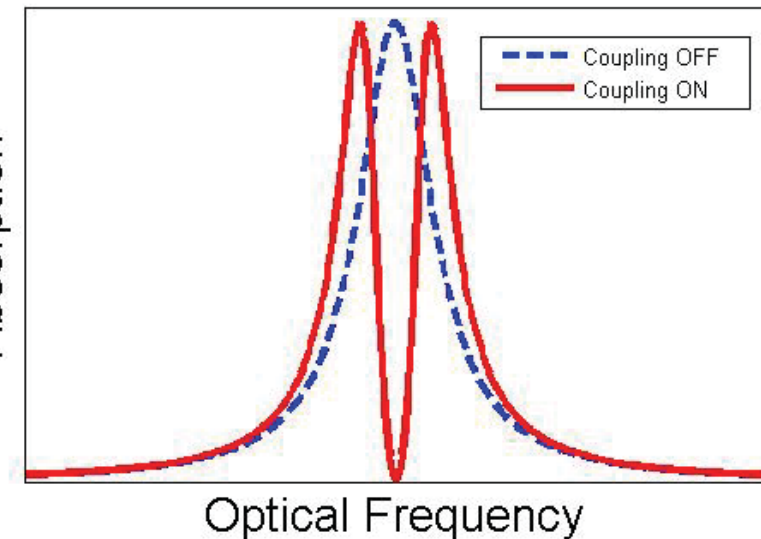
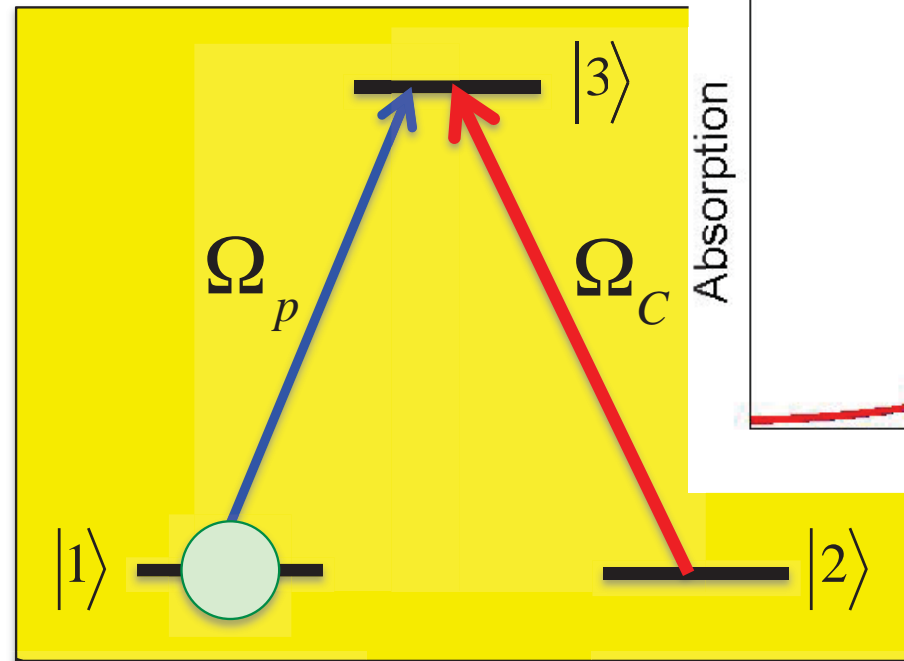
Electromagnetically induced transparency

$$|a^0\rangle = \cos \theta |1\rangle - \sin \theta |2\rangle$$

new dark state

$$\tan \theta = \frac{\Omega_p}{\Omega_c}$$

NO spontaneous decay



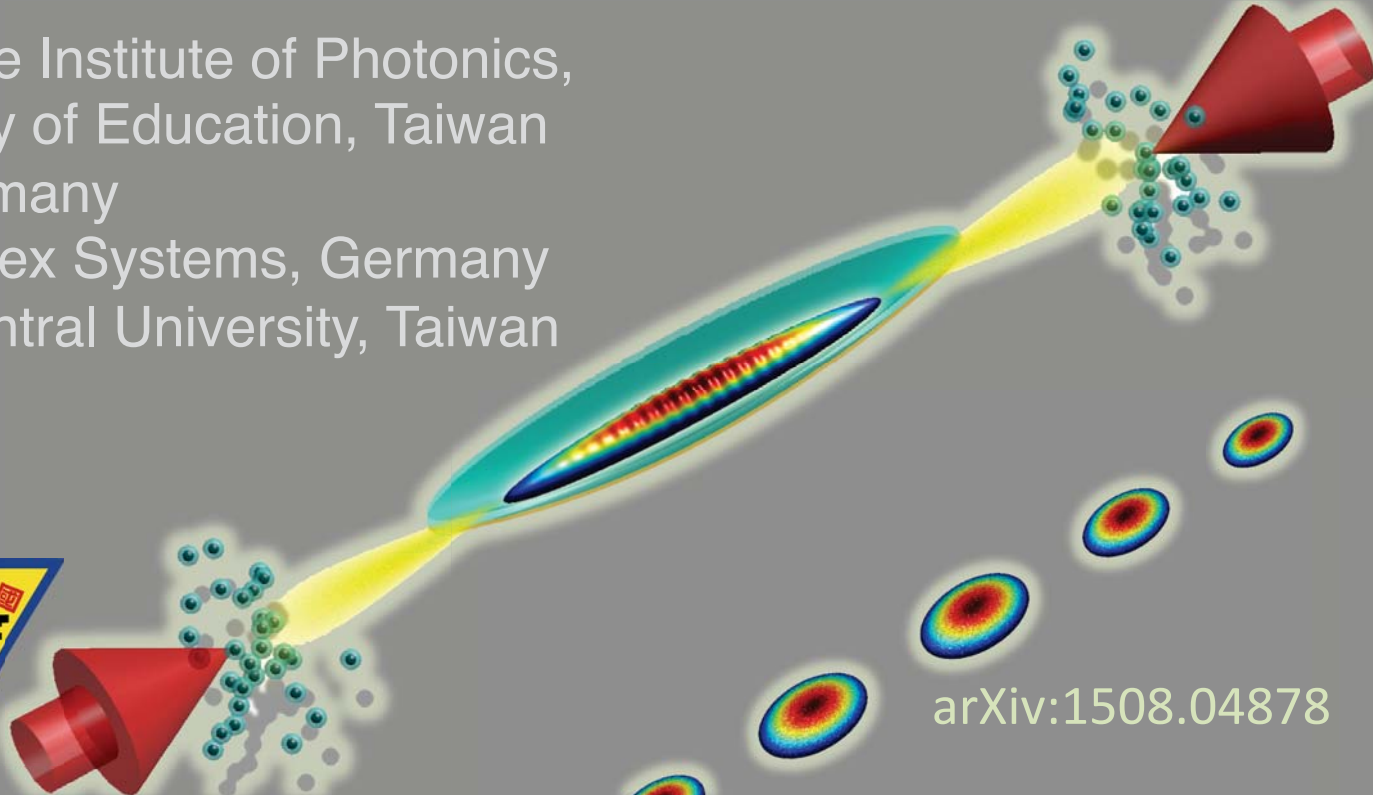
also transparent!!

All-optical cavity using atomic mirrors

Shih-Wei Su(蘇士煒)¹, Zhen-Kai Lu(呂振凱)²,
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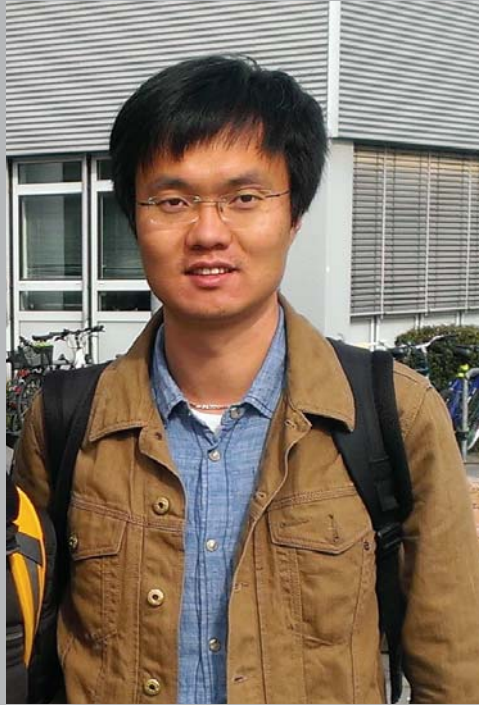


arXiv:1508.04878

Acknowledgement



Shih-Wei Su
(蘇士煒)



Zhen-Kai Lu
(呂振凱)

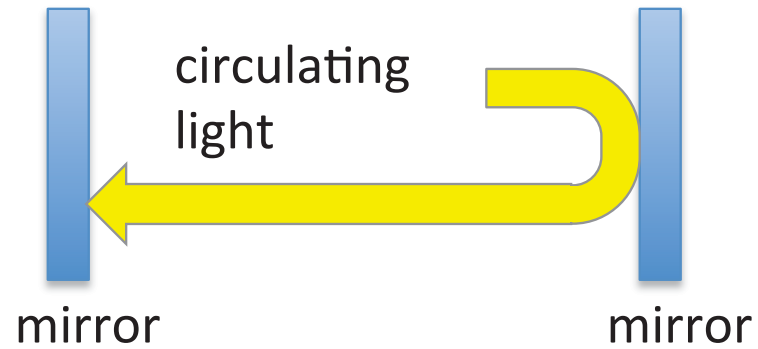


I-Kang Liu
(劉翼綱)



Shih-Chuan Gou
(郭西川)

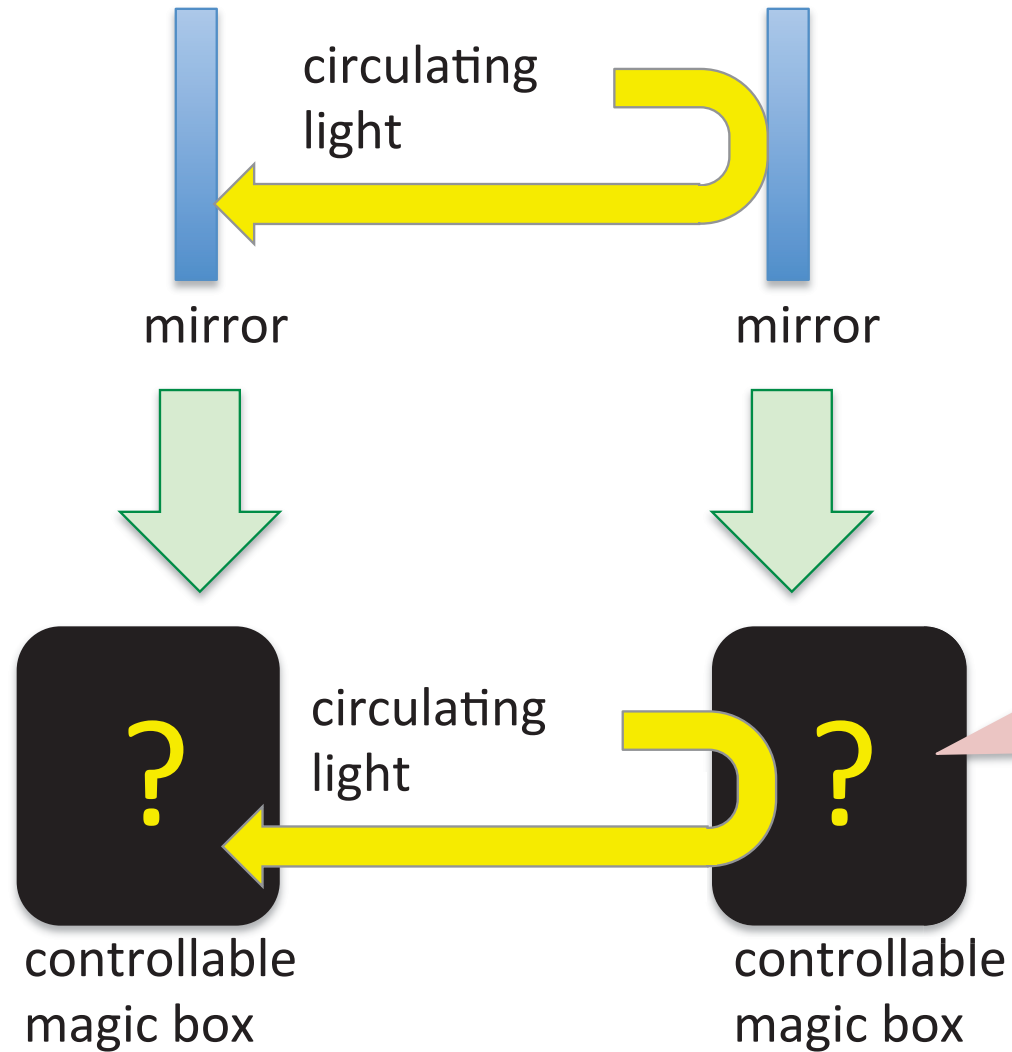
What we want to do?



conventional cavity

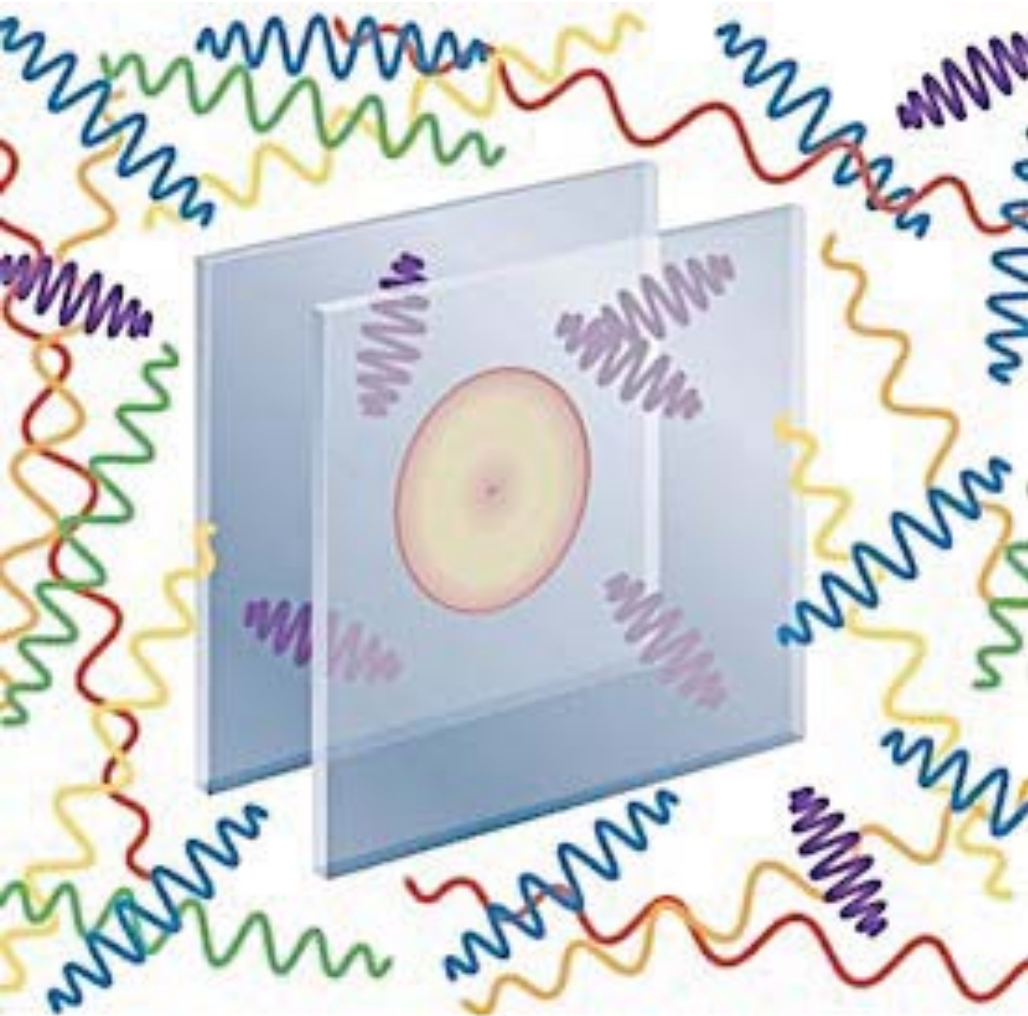


What I want to do



reflection on/ off
whenever you
want

Motivation – different cavity QED



credit: Scientific American

conventional mirror is not controllable.



? Can fast control happen?

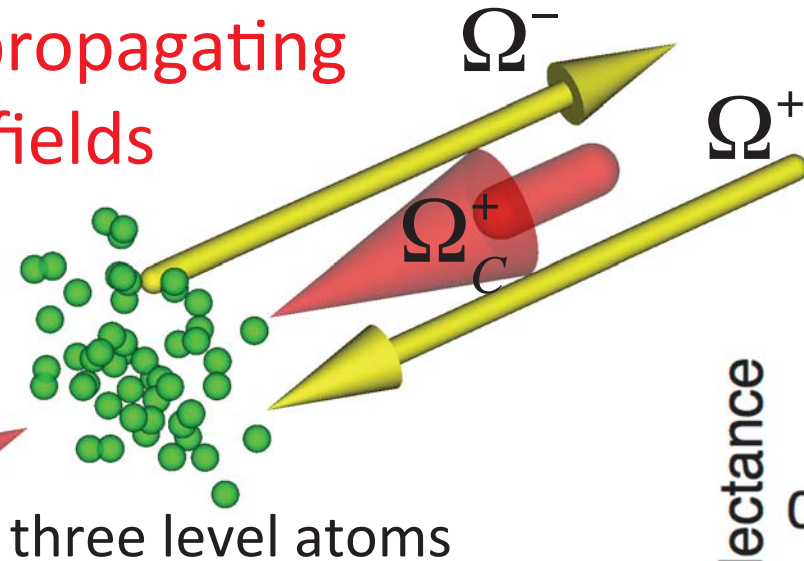
? can the mirror be as tiny as atom clouds in a cavity?

? any cavity effect different from vacuum Rabi oscillation?

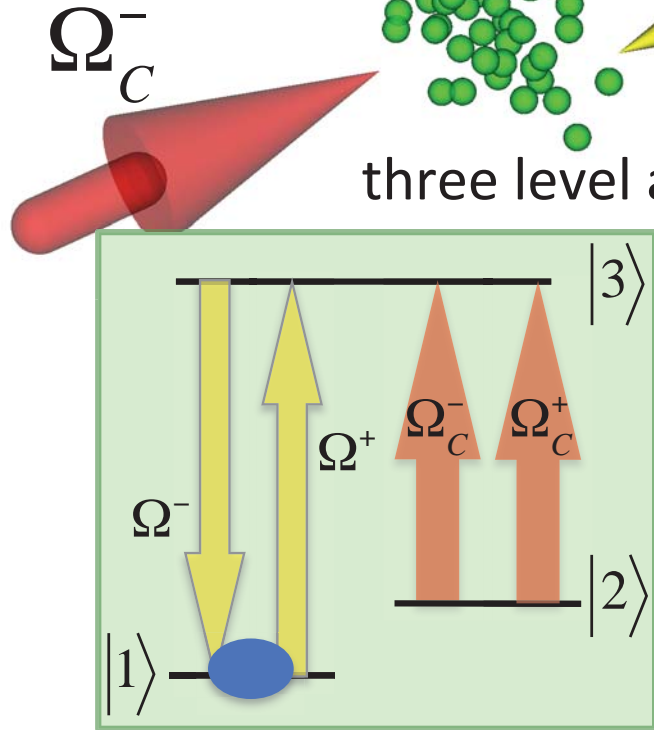
? any novel way of controlling matter?

Dispersive Mirror

counter propagating
coupling fields



three level atoms



$$\partial_t \hat{\rho} = \frac{1}{i\hbar} [\hat{H}, \hat{\rho}] + \hat{\rho}_{dec}$$

$$\left(\frac{1}{c} \partial_t \pm \partial_z\right) \Omega^\pm = i\eta_{EIT} \rho_{31}^\pm$$

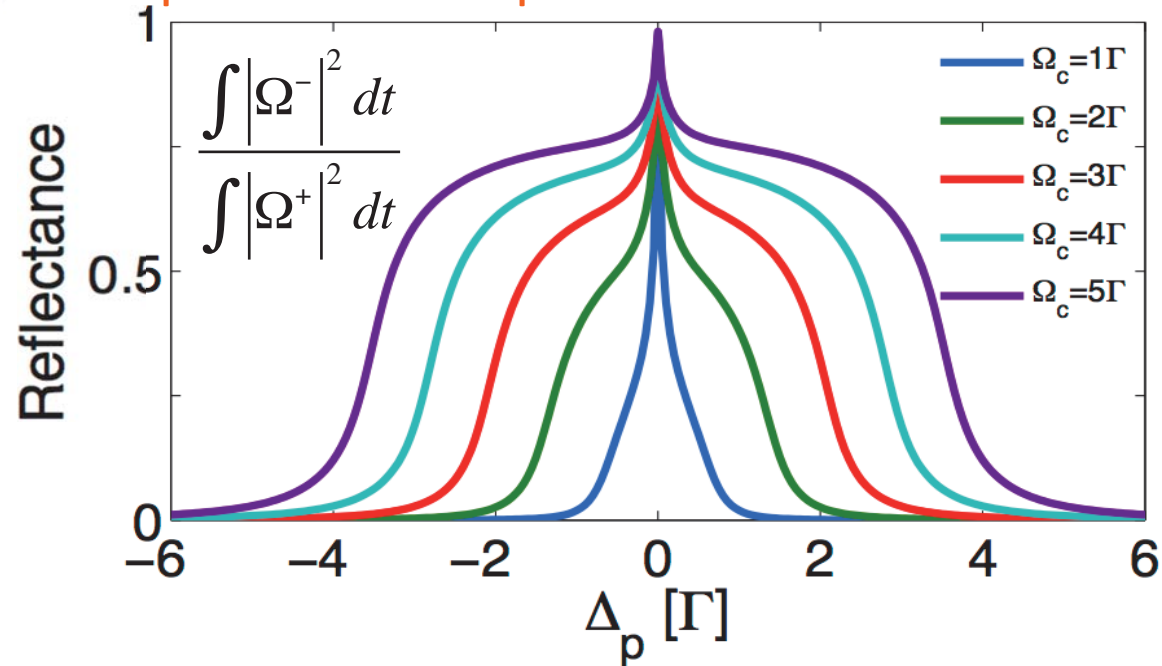
Optical Bloch equation

optical depth

$$\eta_{EIT} = \Gamma d^{opt} / 2L$$

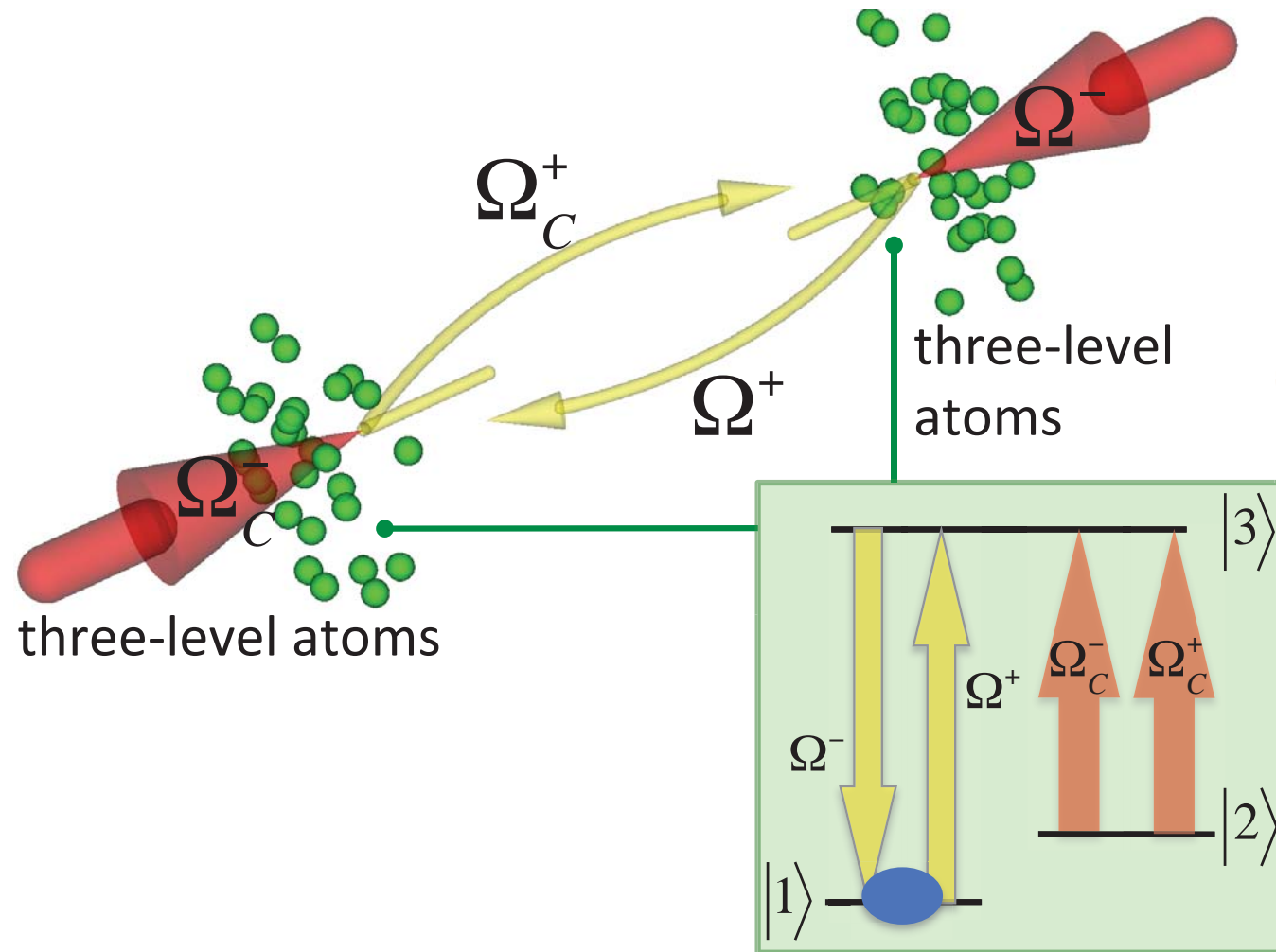
linewidth

length

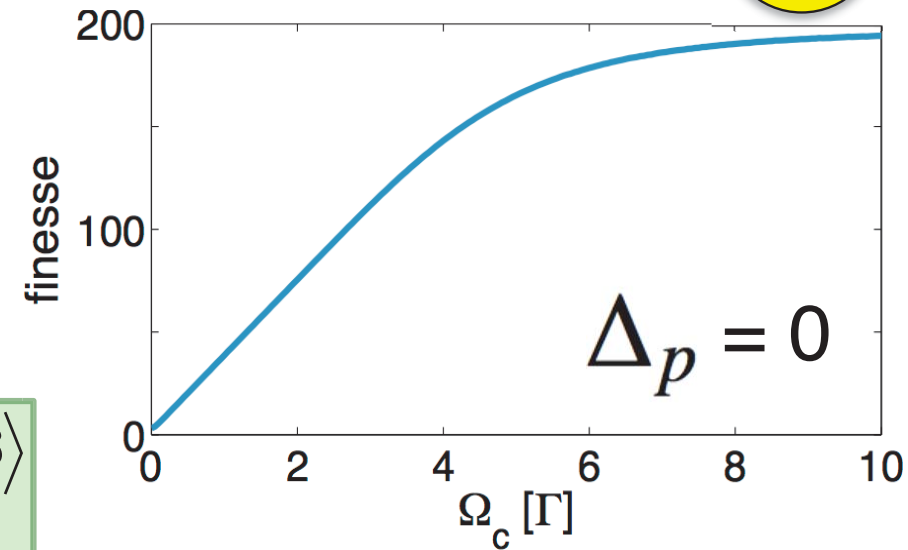


O. Kocharovskaya and Y. I. Khanin, Sov. Phys. JETP 63, 945 (1986).
 M. Bajcsy, A. S. Zibrov, and M. D. Lukin, Nature 426, 638 (2003).
 Y.-W. Lin et. al., Phys. Rev. Lett. 102, 213601 (2009).

Dispersive Cavity & its Finesse



all-optical control

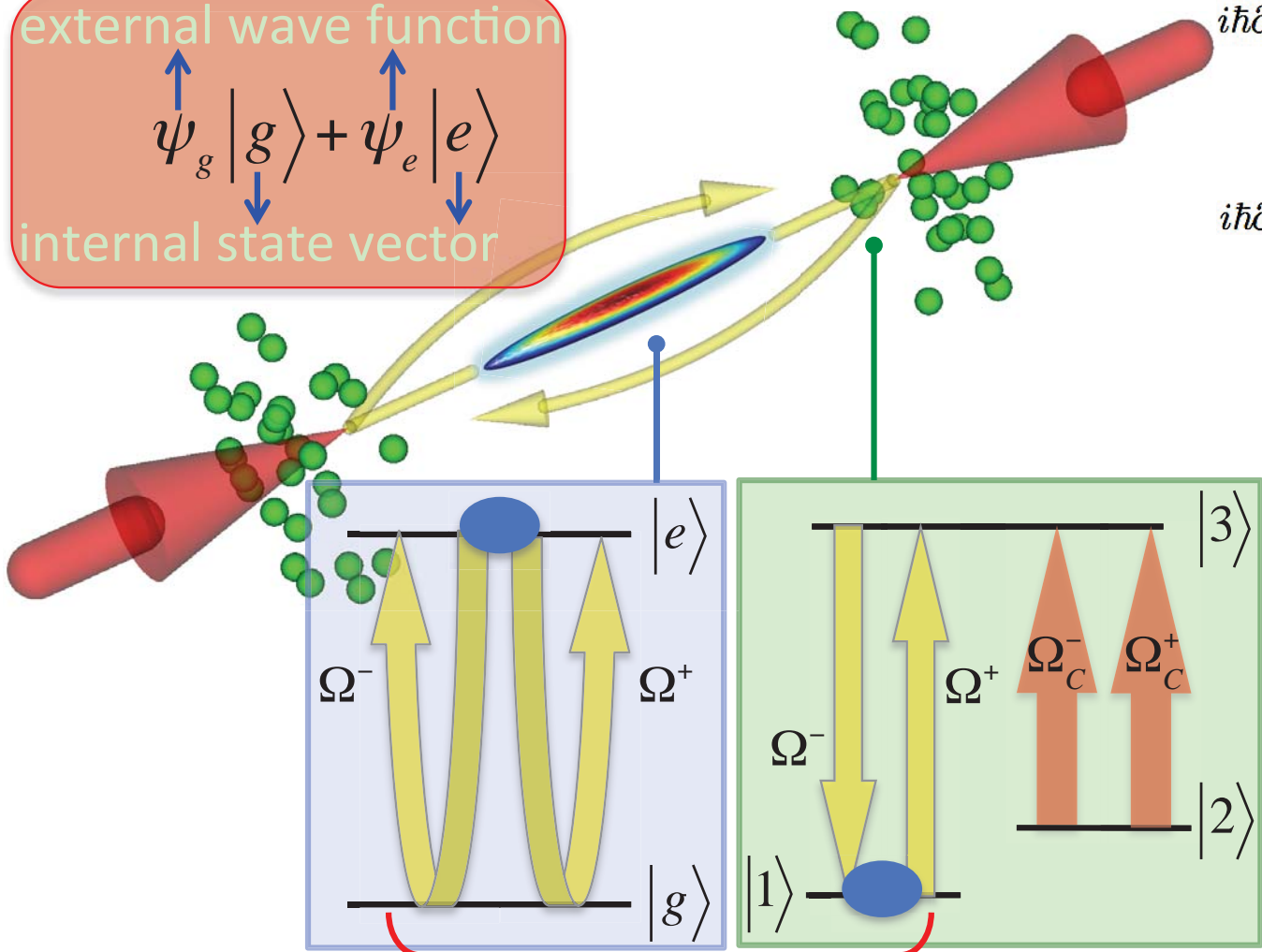


$$f(\Delta_p) = \frac{\pi}{1-R(\Delta_p)}$$

$R(\Delta_p)$ reflectance

Loading a 1D Bose-Einstein Condensate

external wave function
 $\psi_g |g\rangle + \psi_e |e\rangle$
 internal state vector



polarization & energy matched

$$i\hbar\partial_t\psi_g = \left(-\frac{\hbar^2}{2m}\partial_z^2 + V_{trap} + g_{gg}|\psi_g|^2 + g_{ge}|\psi_e|^2\right)\psi_g - \frac{\hbar}{2}(\Omega^{+*}e^{-ik_pz} + \Omega^{-*}e^{ik_pz})\psi_e,$$

$$i\hbar\partial_t\psi_e = \left(-\frac{\hbar^2}{2m}\partial_z^2 + V_{trap} + g_{ee}|\psi_e|^2 + g_{ge}|\psi_g|^2\right)\psi_e - \frac{\hbar}{2}(\Omega^+e^{ik_pz} + \Omega^-e^{-ik_pz})\psi_g,$$

Gross-Pitaevskii equation

$$\left(\frac{1}{c}\partial_t \pm \partial_z\right)\Omega^\pm = i\eta_{BEC}\sigma_{eg}^{(\pm 1)}$$

wave equation in a cavity

$$\sigma_{eg} = \psi_e\psi_g^*$$

BEC coherence

particle number

$$\eta_{BEC} = 3\Gamma N_{BEC}^{\uparrow} \lambda^2 / 4\pi A$$

BEC's cross section

BEC-photon coupling constant

Bose-Einstein Condensate

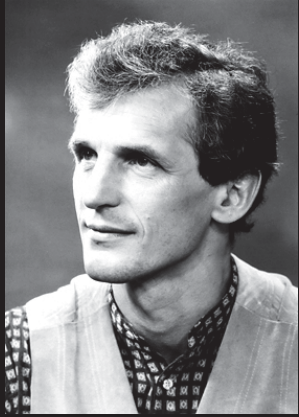
video link

<http://toutestquantique.fr/en/bose-einstein-condensate/>

Bose-Einstein Condensate (2001 Nobel Prize)



Eric Cornell



Wolfgang Ketterle

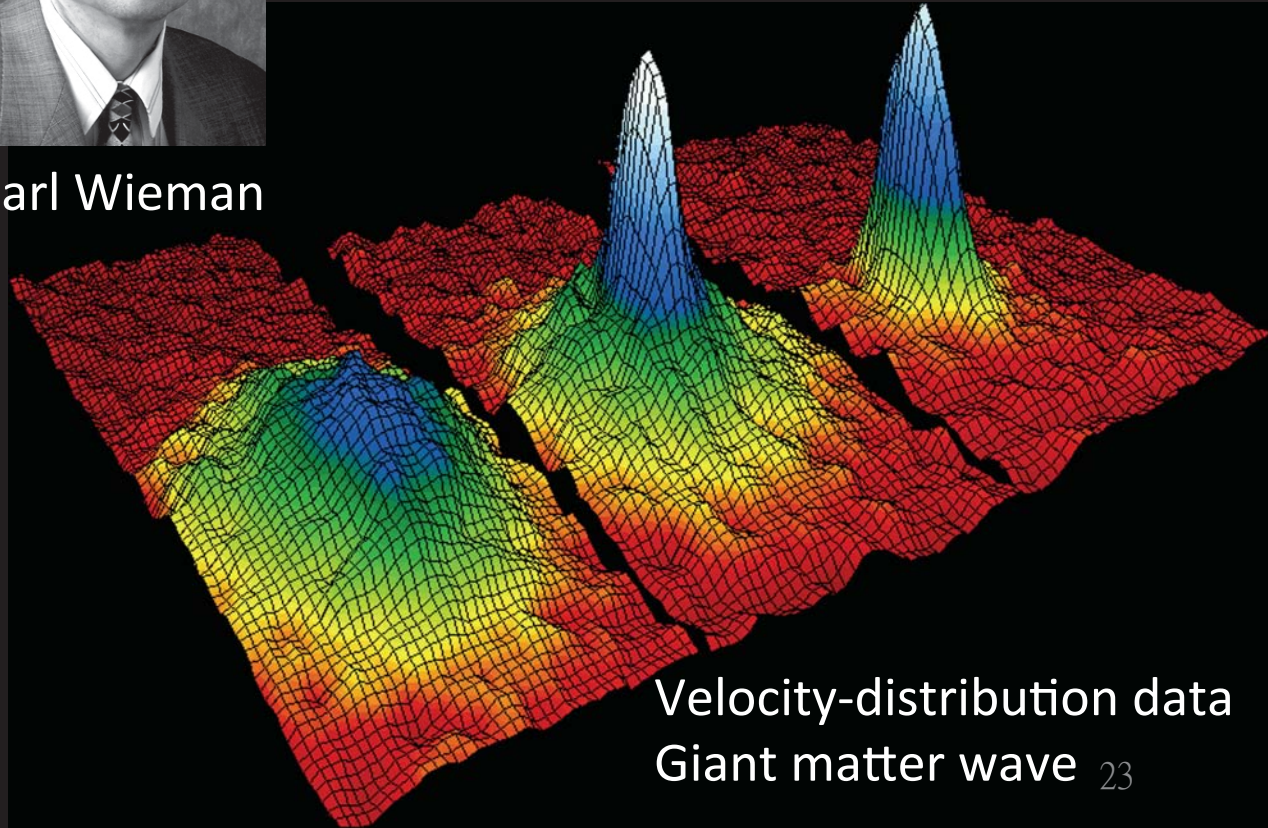
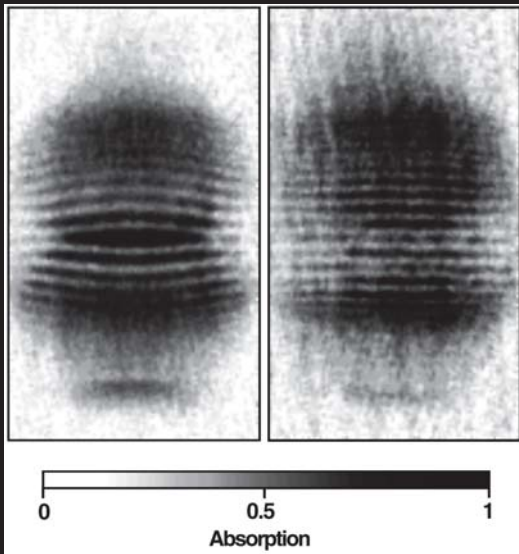


Carl Wieman



interference
between
2 BECs

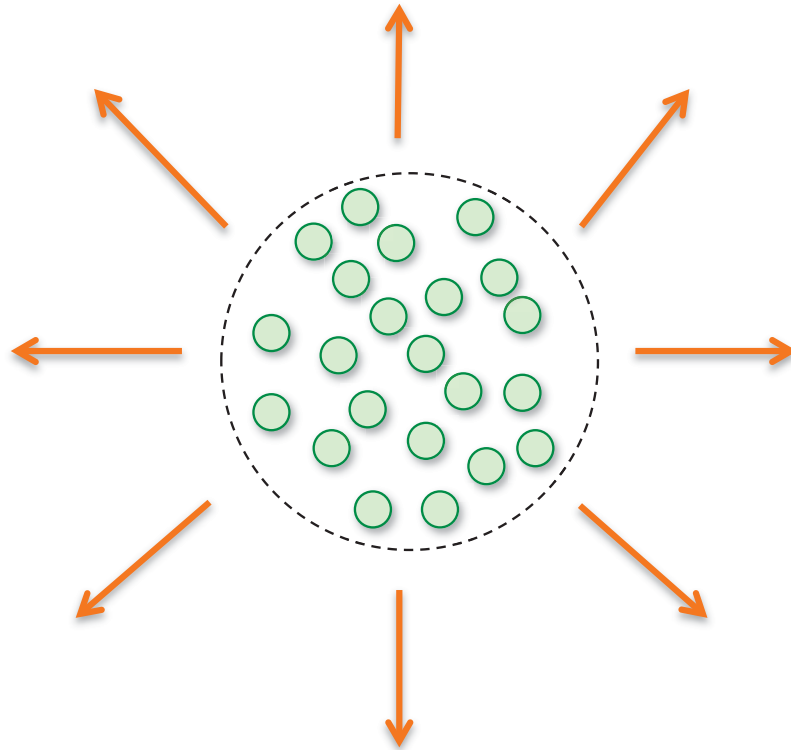
Coherent !!!



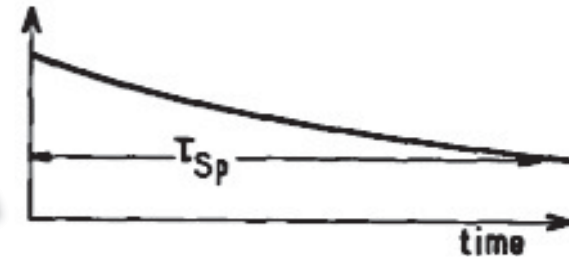
Velocity-distribution data
Giant matter wave ²³

Superradiance from 1D gas

round

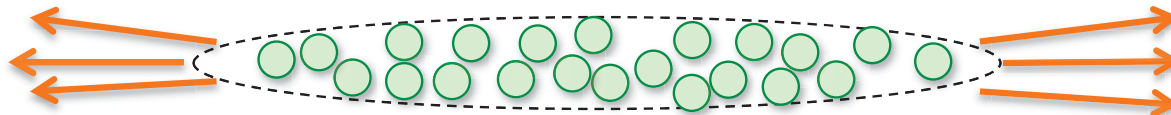


Detector

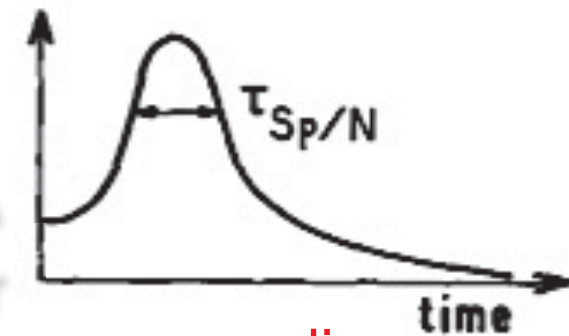


spontaneous decay

cigar



Detector



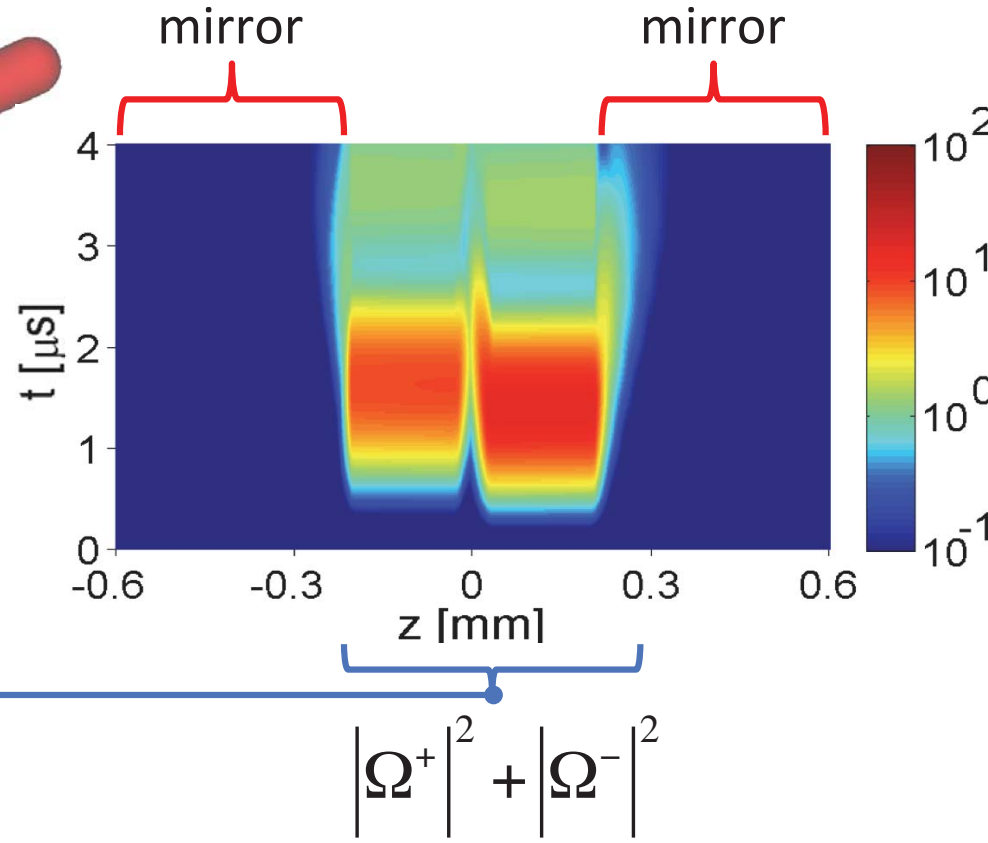
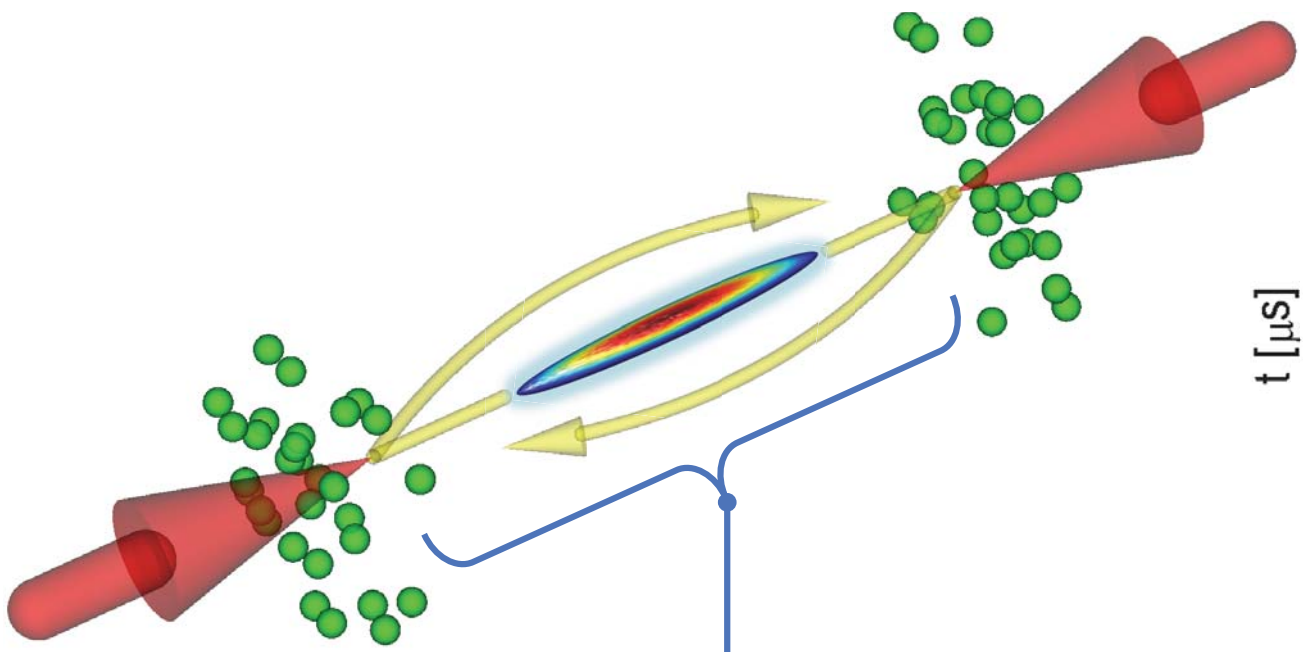
superradiance

M. Gross, S. Haroche, Phys. Rep. 93, 301 (1982).

F. Haake et al., Phys. Rev. A 20,2047 (1979).

S. Inouye et al. Science 285, 571 (1999).

Confined superradiance from BEC

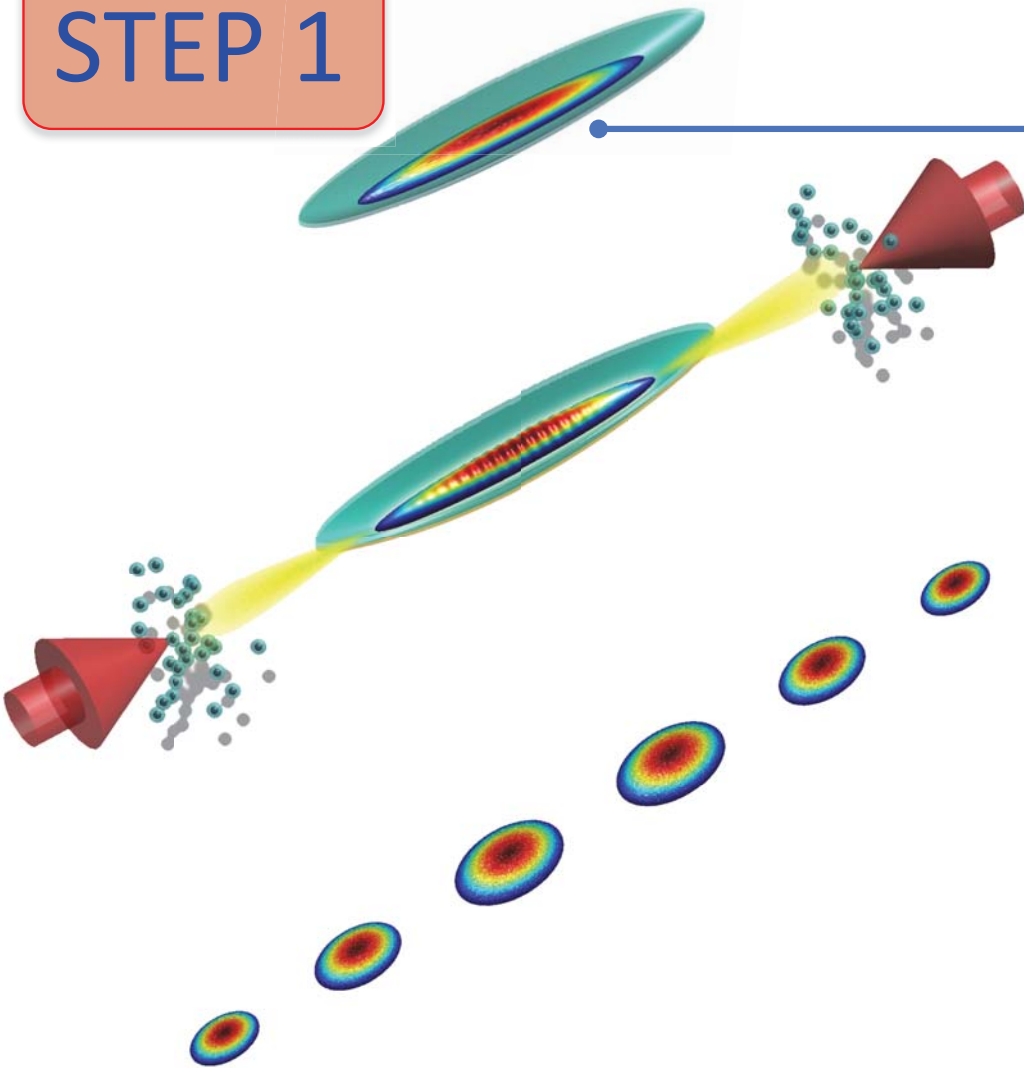


$A = 9\pi \mu\text{m}^2$
 $N_{BEC} = 3 \times 10^5 \text{ } ^7\text{Li}$
 $d^{opt} = 500$
 $\Omega_c = 2\Gamma$

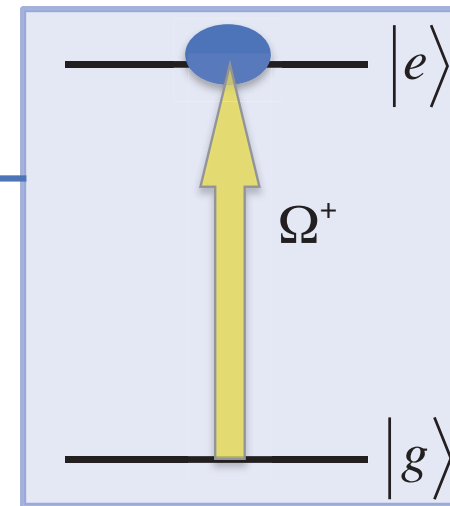
SR confined between two atomic mirrors.

Vacuum induced atomic diffraction

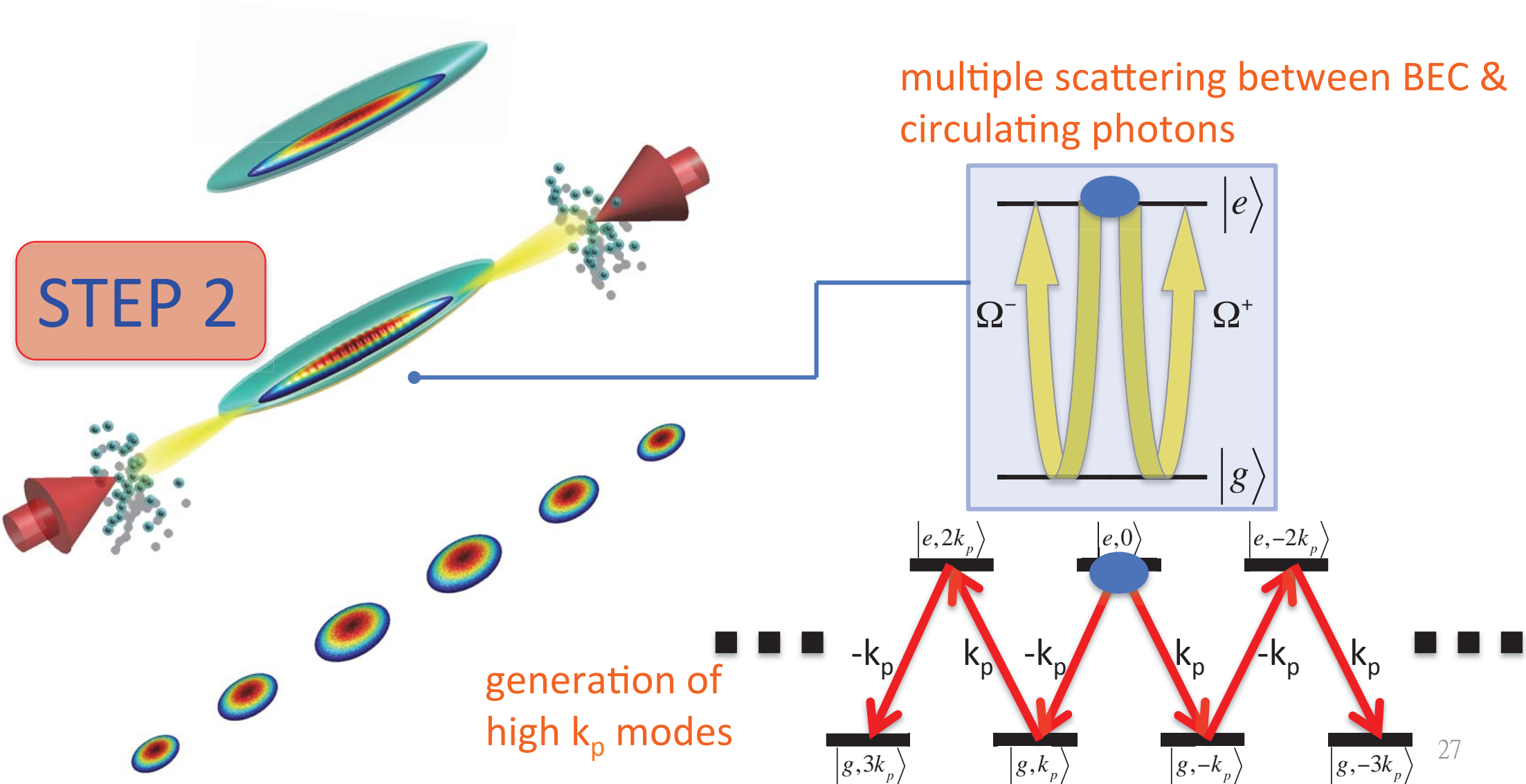
STEP 1



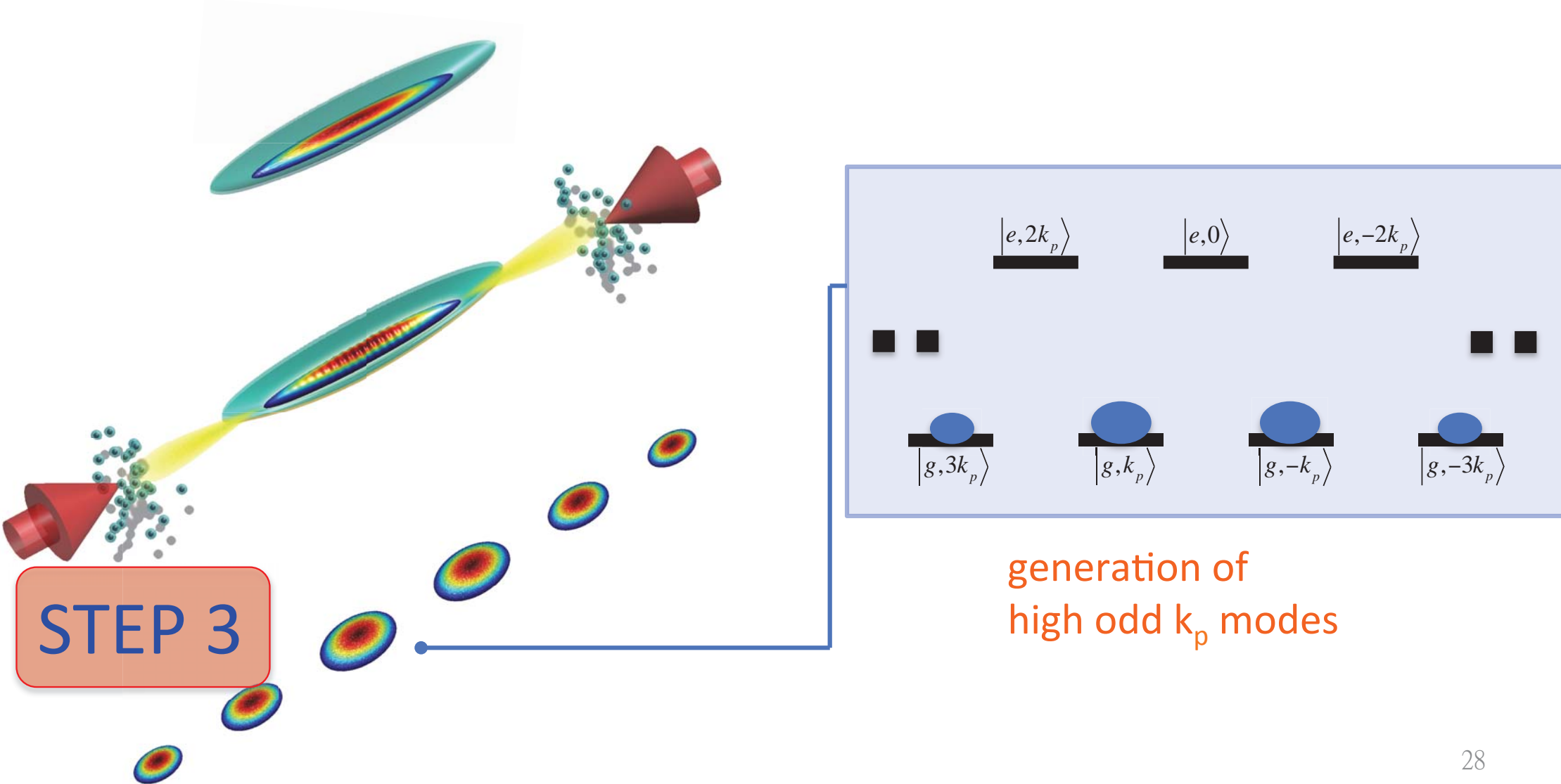
prepare and promote a binary BEC



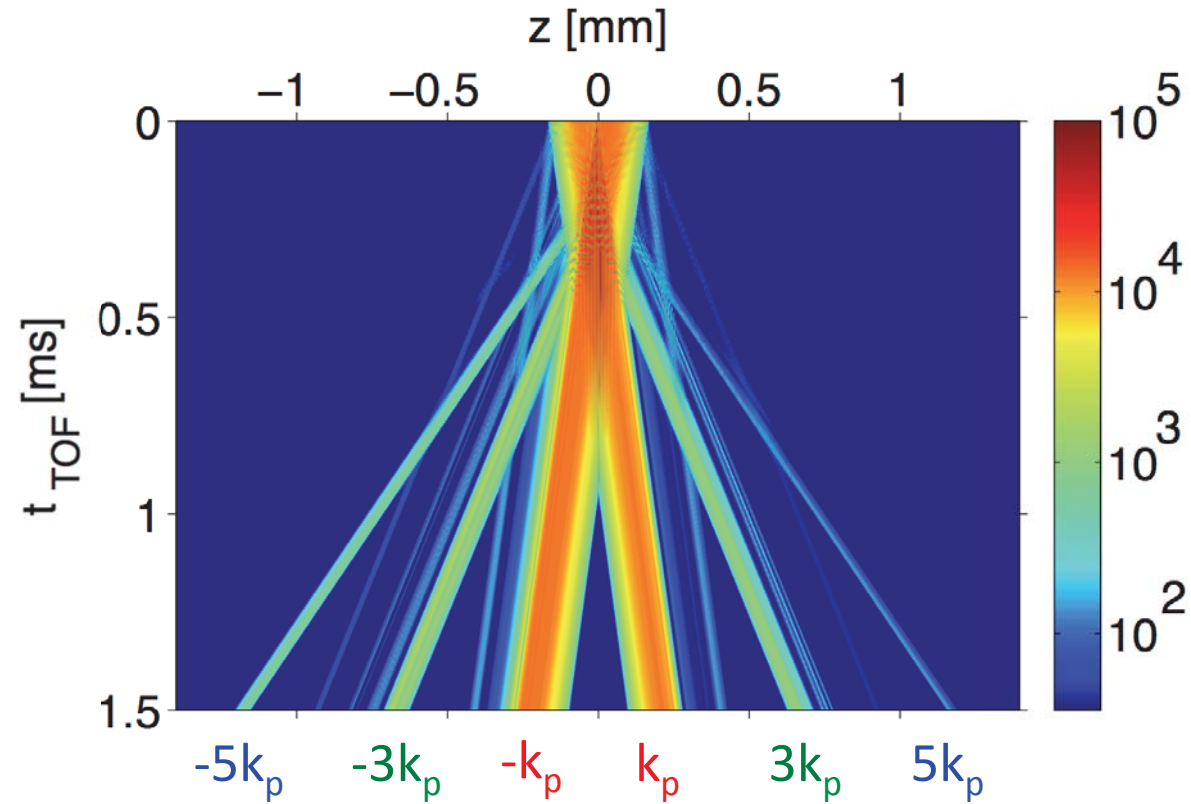
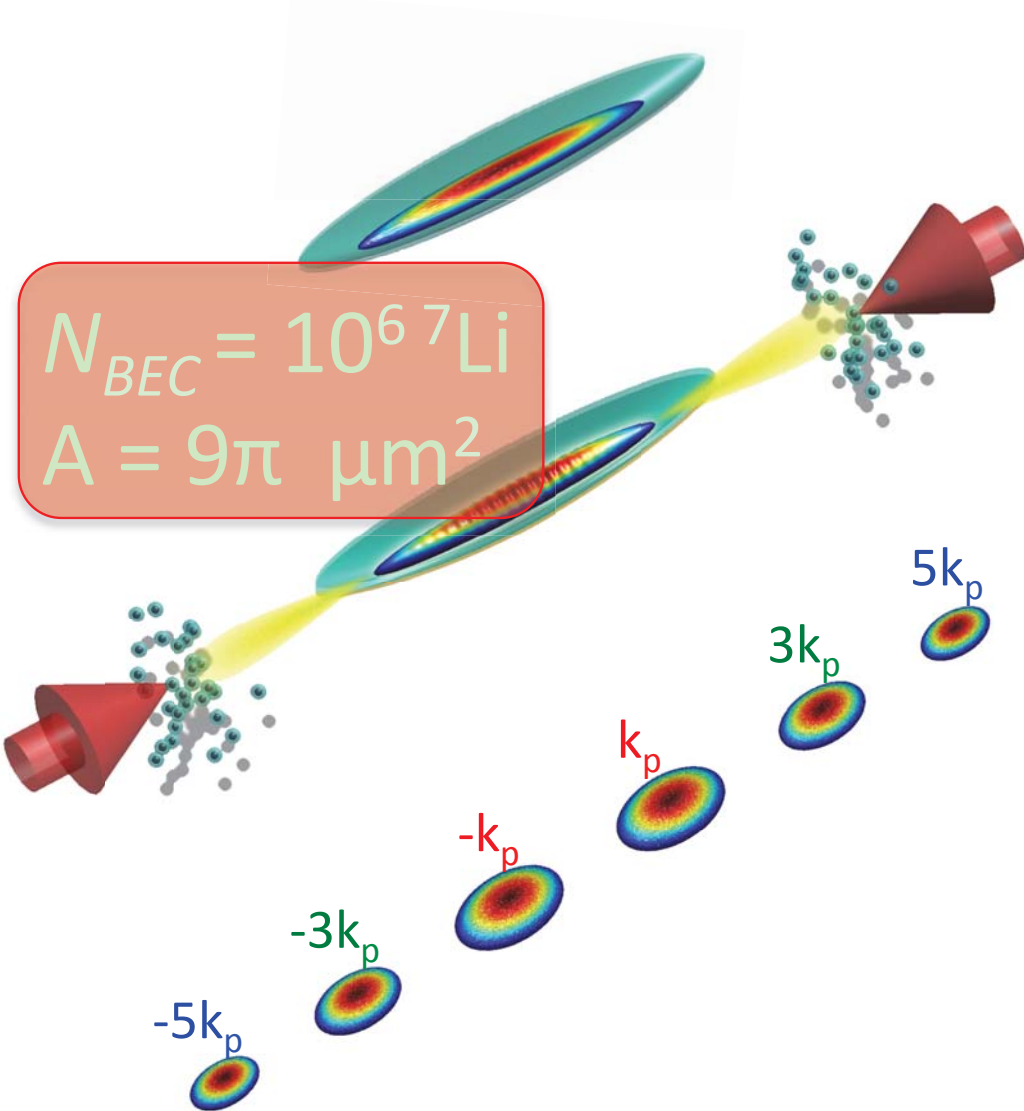
Vacuum induced atomic diffraction



Vacuum induced atomic diffraction

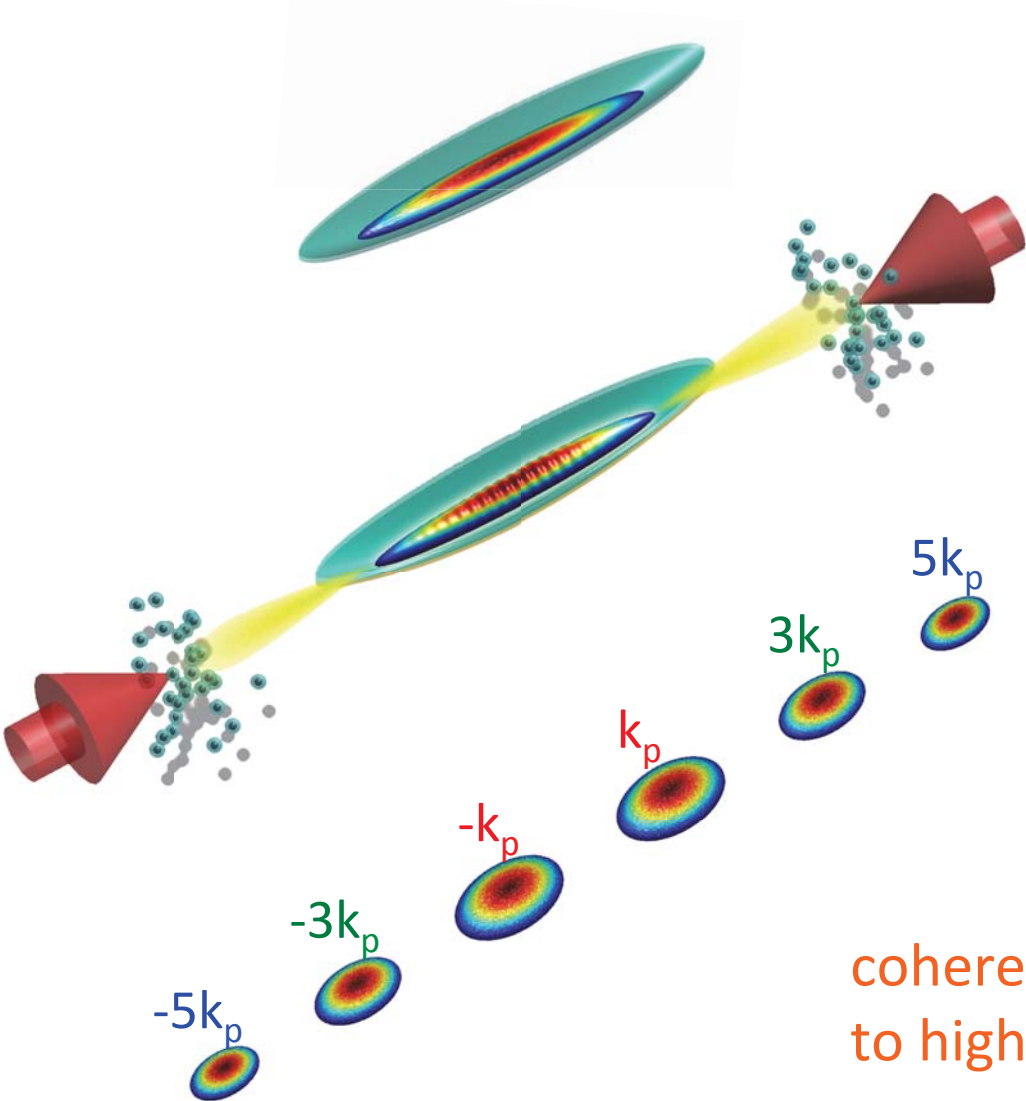


Time of flight simulation

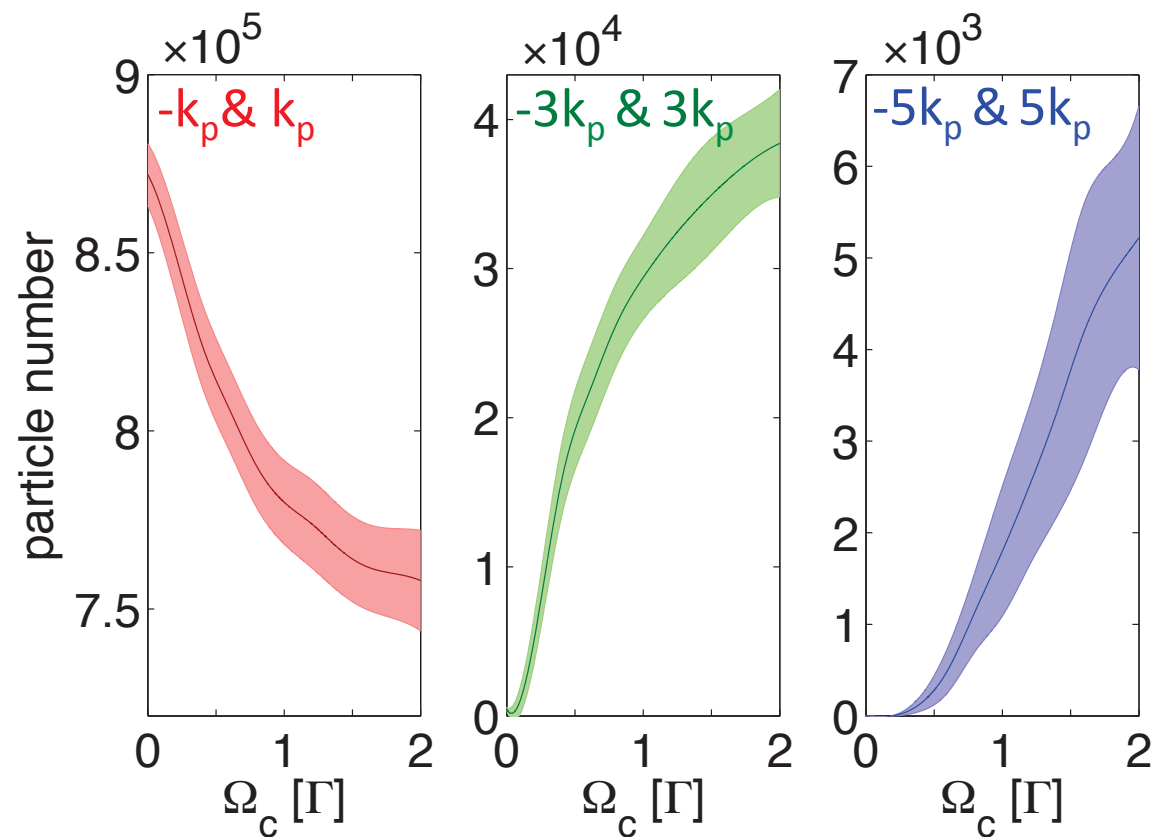


generation of
high odd k_p modes

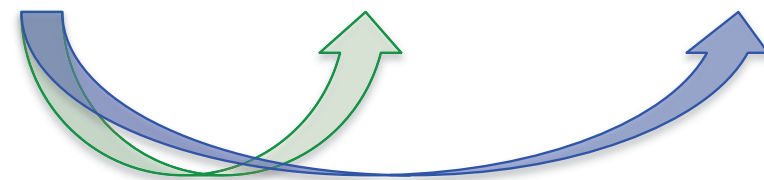
Coupling-strength-dependent diffraction



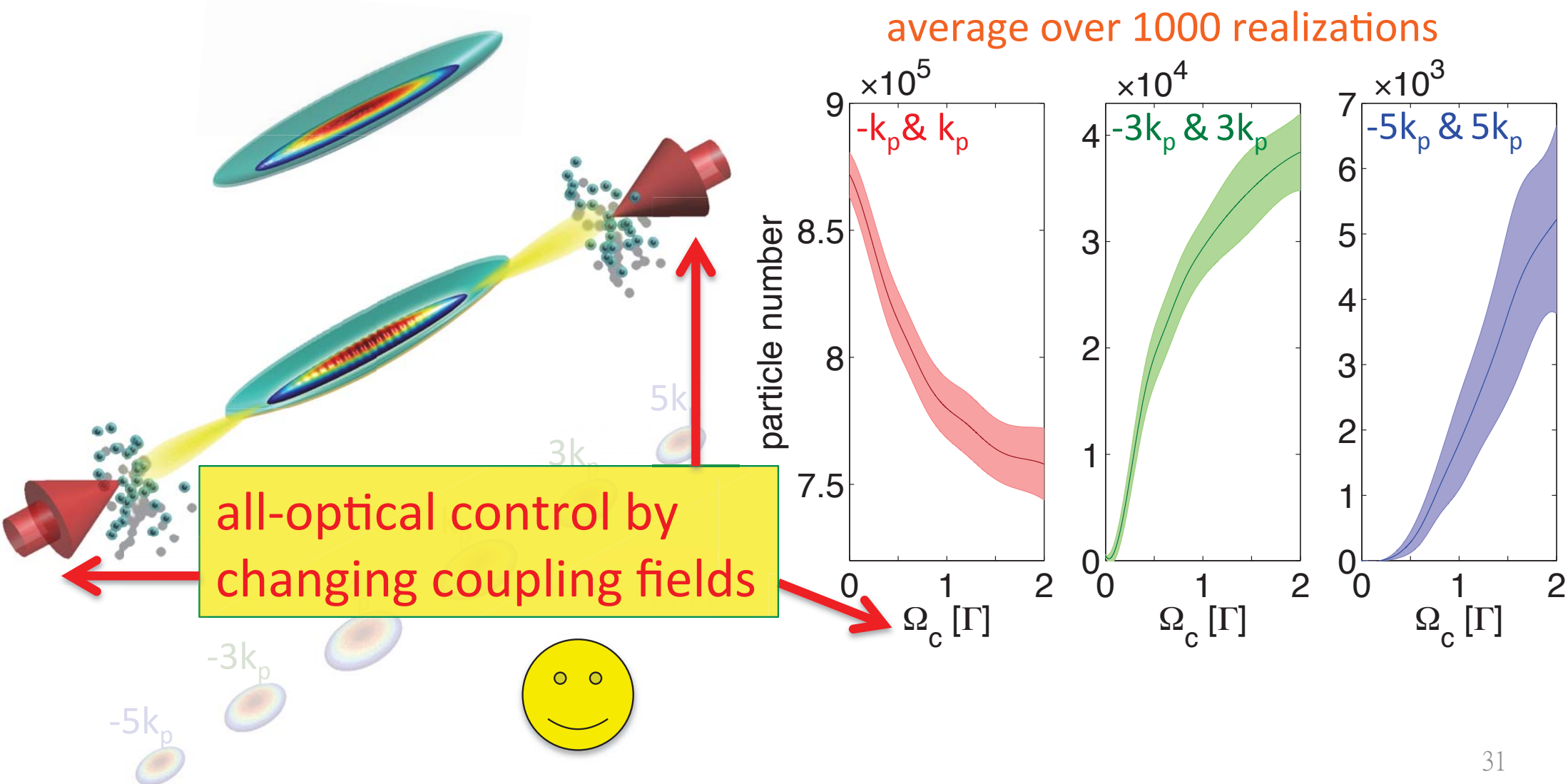
average over 1000 realizations



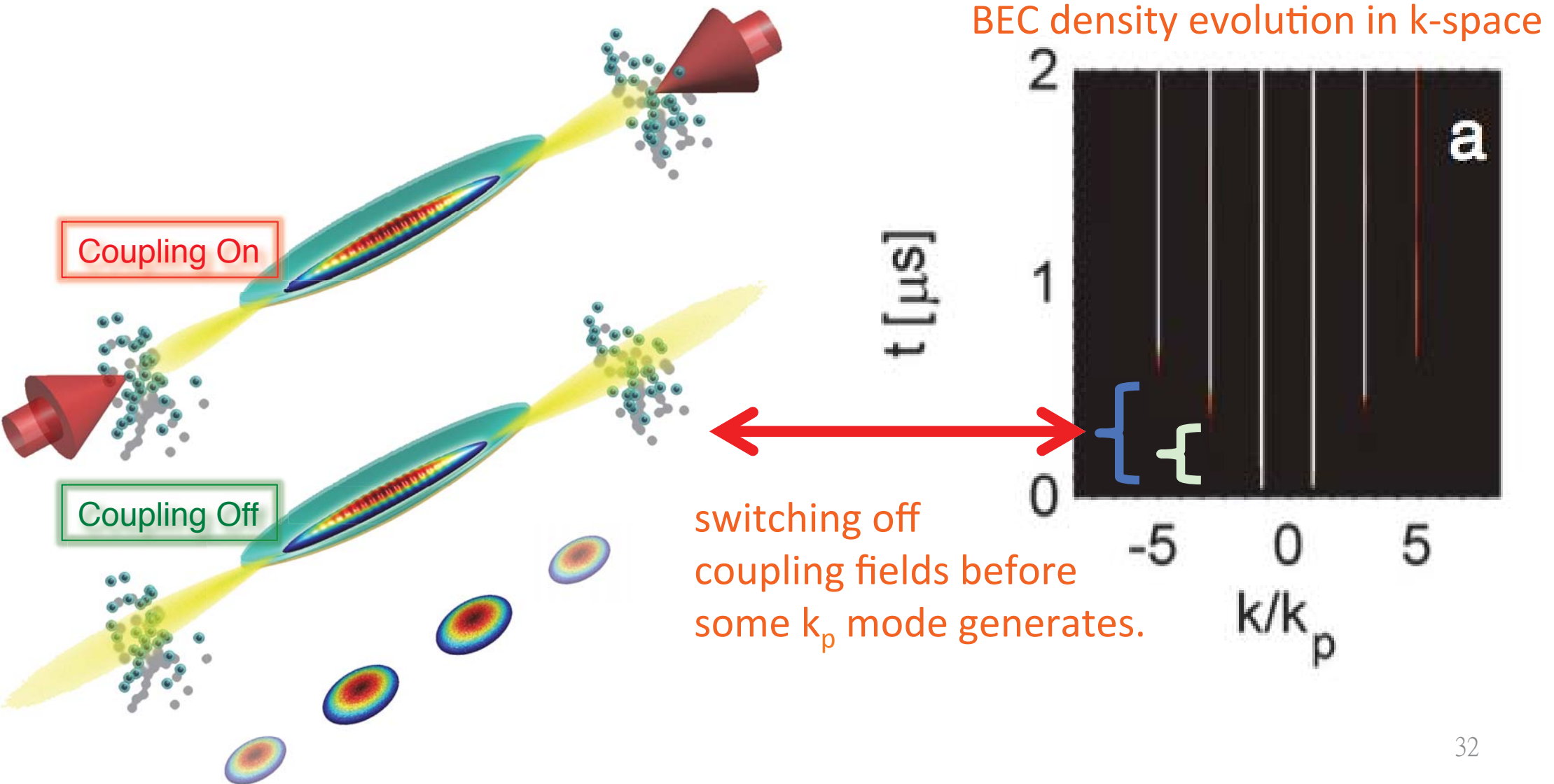
coherent transfer
to high k_p modes



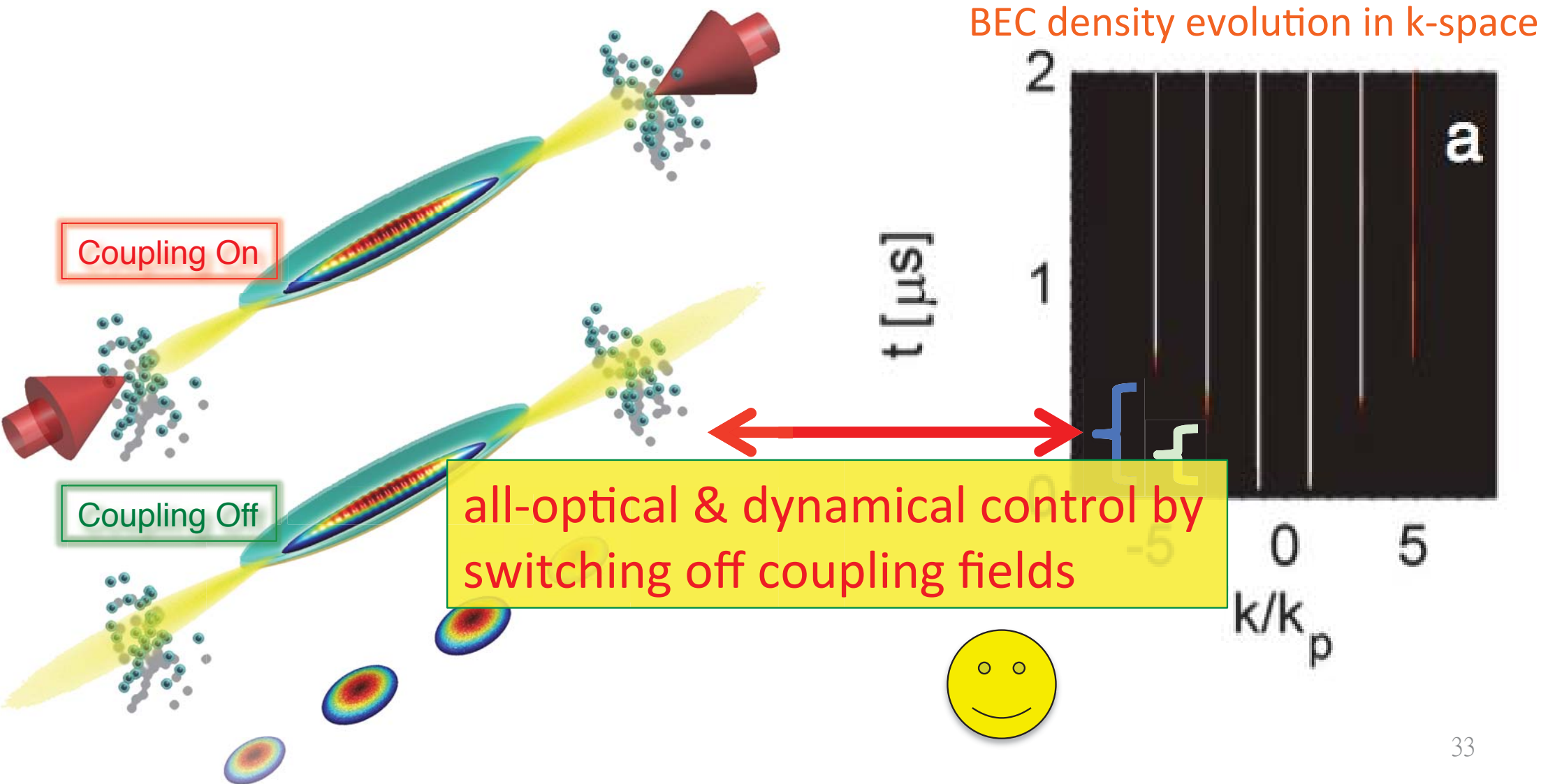
Coupling-strength-dependent diffraction



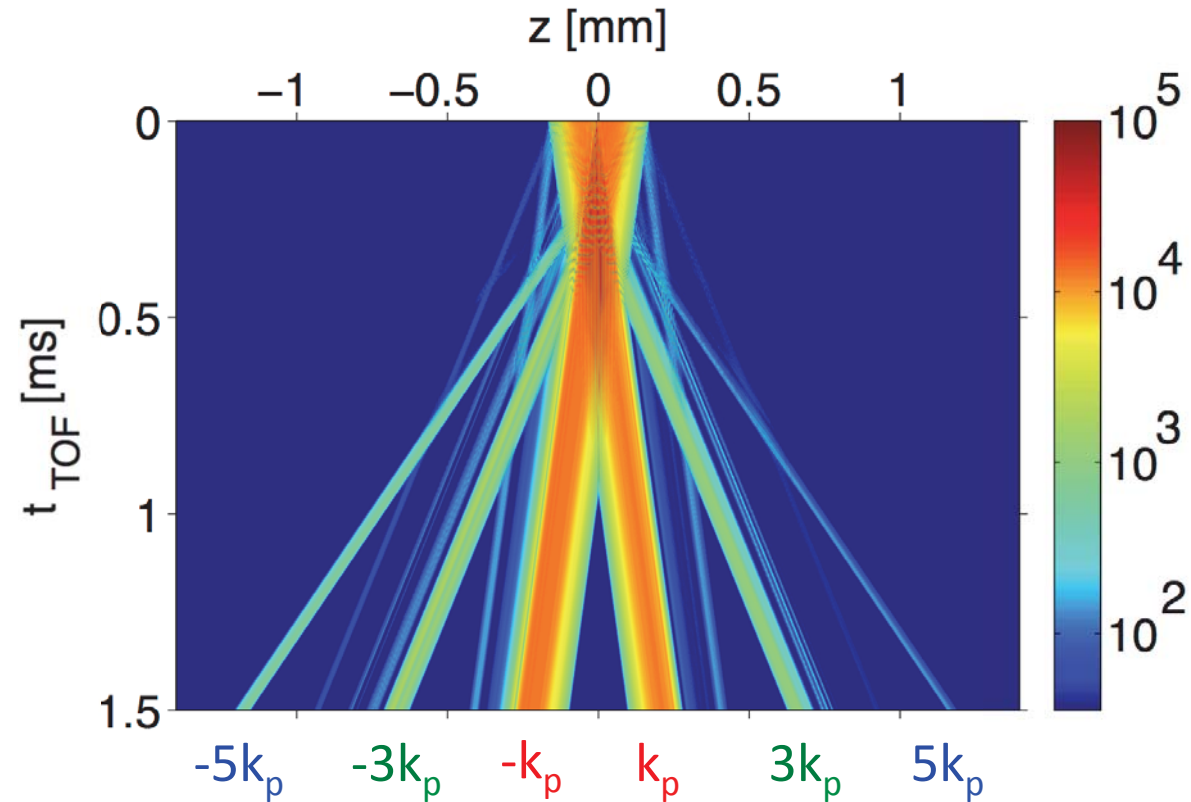
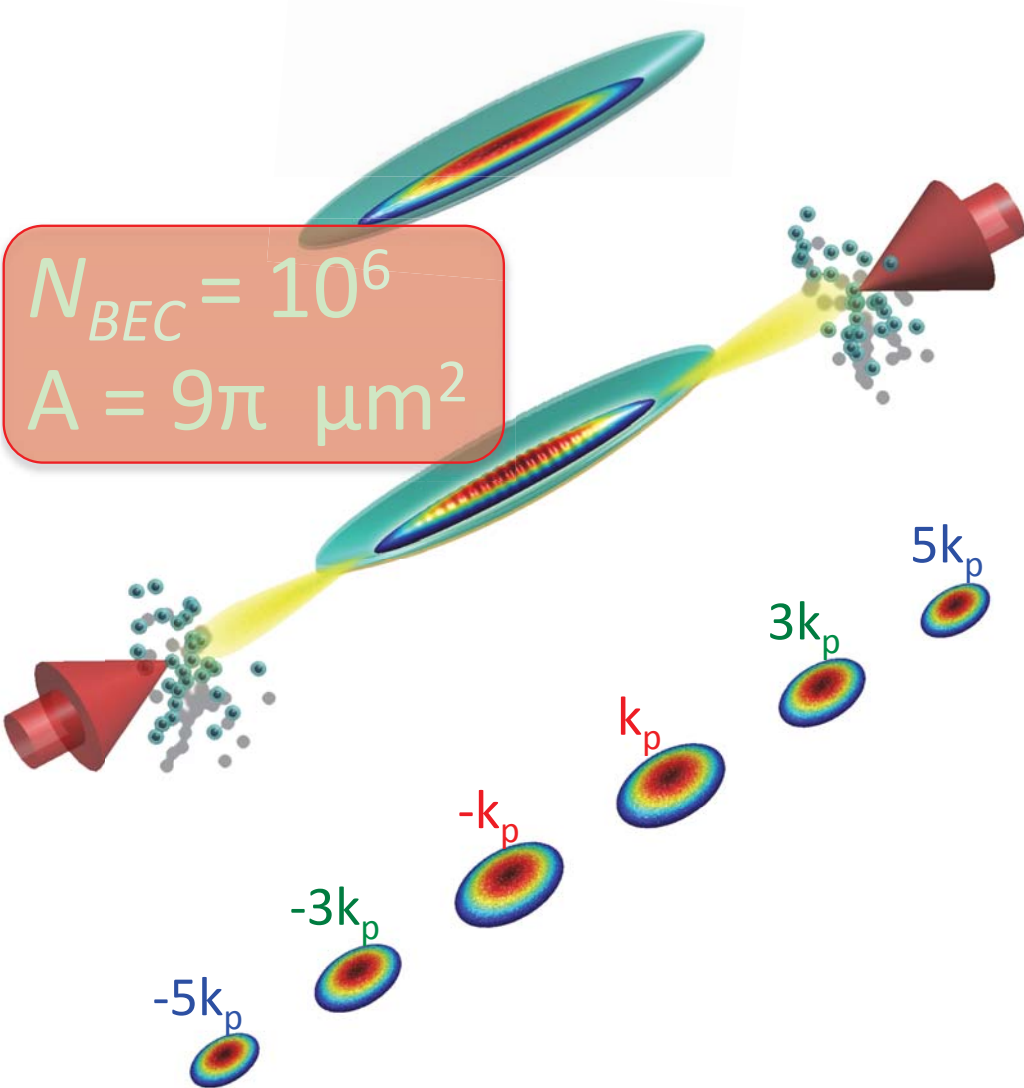
Dynamical control of diffraction



Dynamical control of diffraction



Dynamical control of diffraction



coupling always on

Specific mode can be suppressed

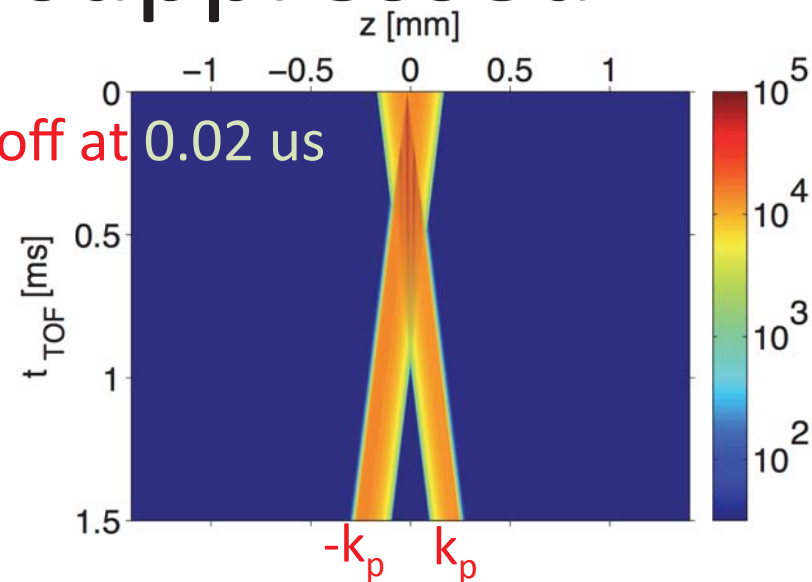
$$N_{BEC} = 10^6$$

$$A = 9\pi \mu\text{m}^2$$

Coupling On

$$N_k = 8.5 \times 10^5$$

coupling off at 0.02 μs

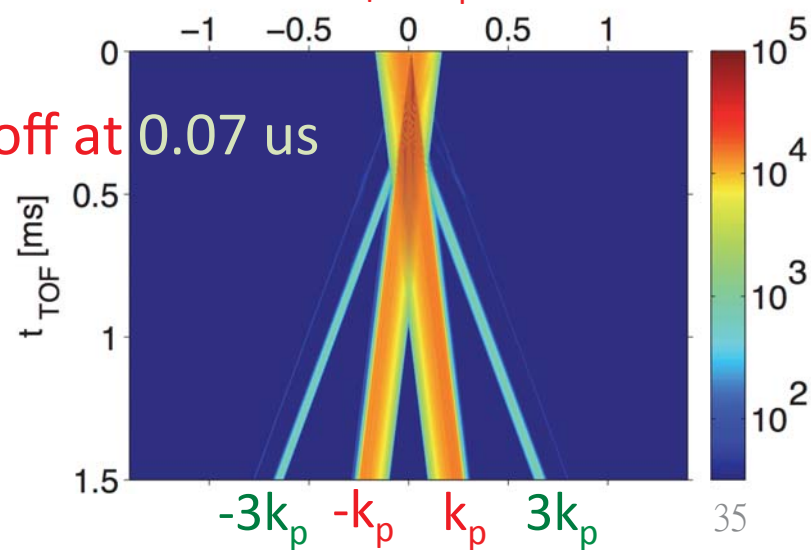


Coupling Off

$$N_k = 8.2 \times 10^5$$

$$N_{3k} = 1.6 \times 10^4$$

coupling off at 0.07 μs



Summary

an all-optical and dynamically controllable cavity is demonstrated.

a dispersive cavity can be used to manipulate atomic diffraction by

- 1) changing the strength of coupling fields
- 2) dynamically switching off the coupling fields

arXiv:1508.04878

