

Attosecond Physics

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2015 AMO summer school

Aug. 27, Hsinchu



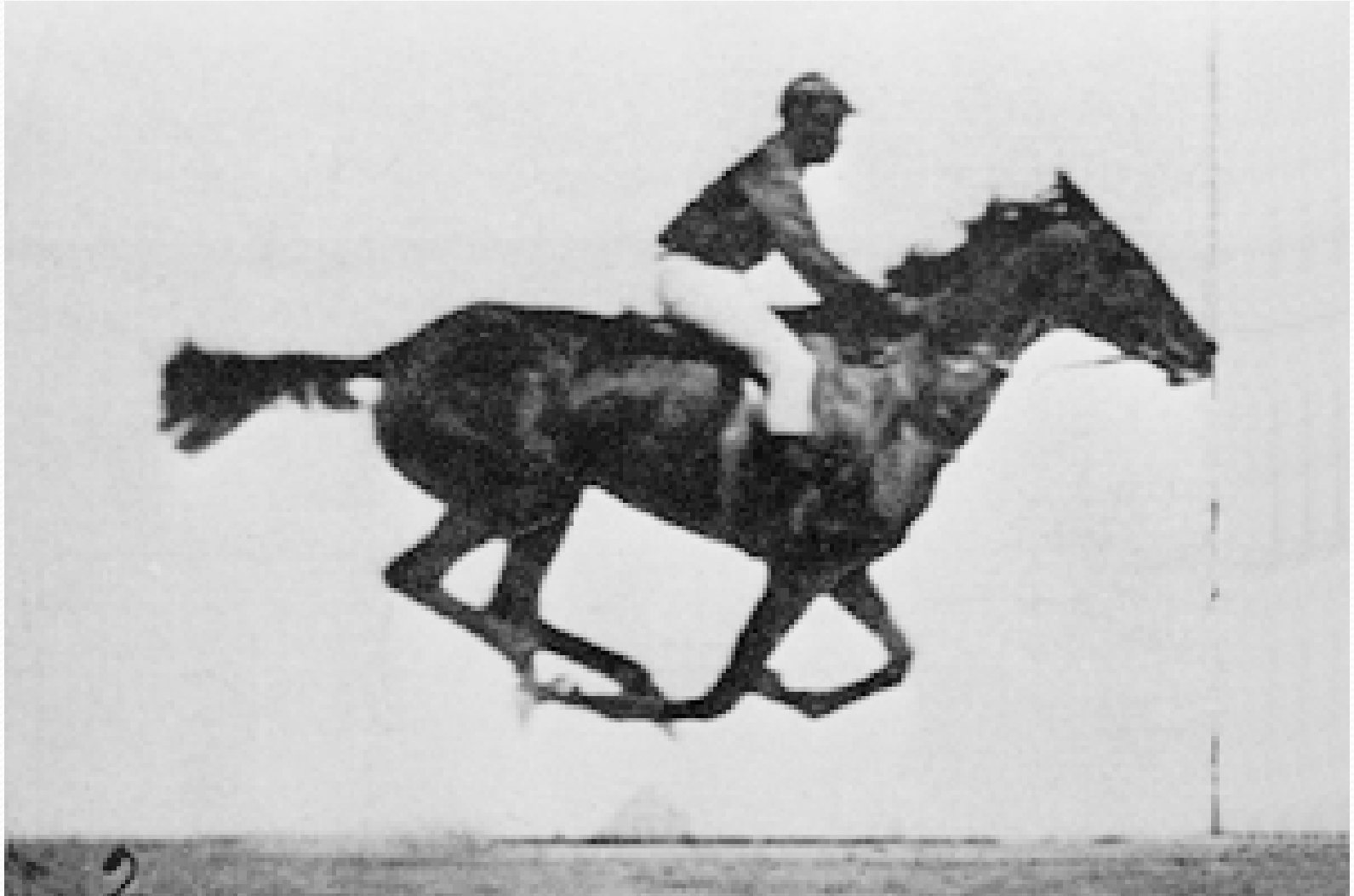
Nobel Prize in Chemistry

- **1999: Ahmed Zewail**

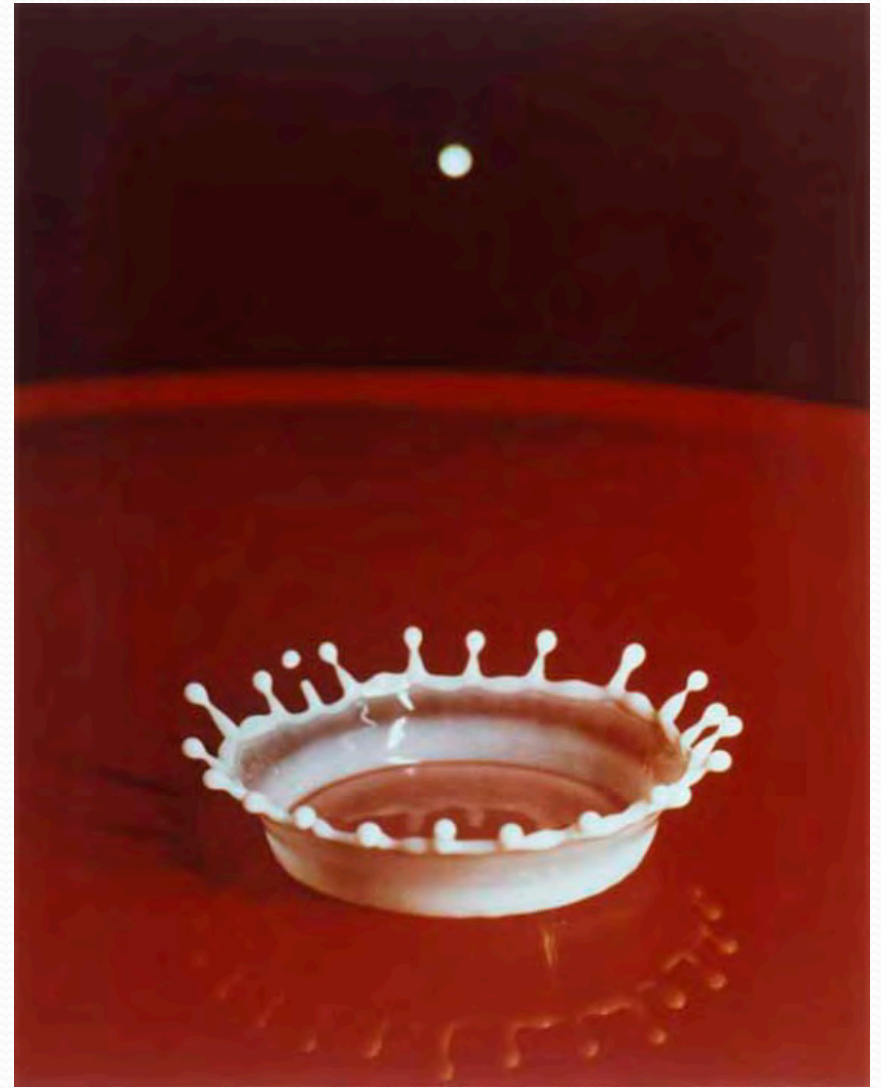
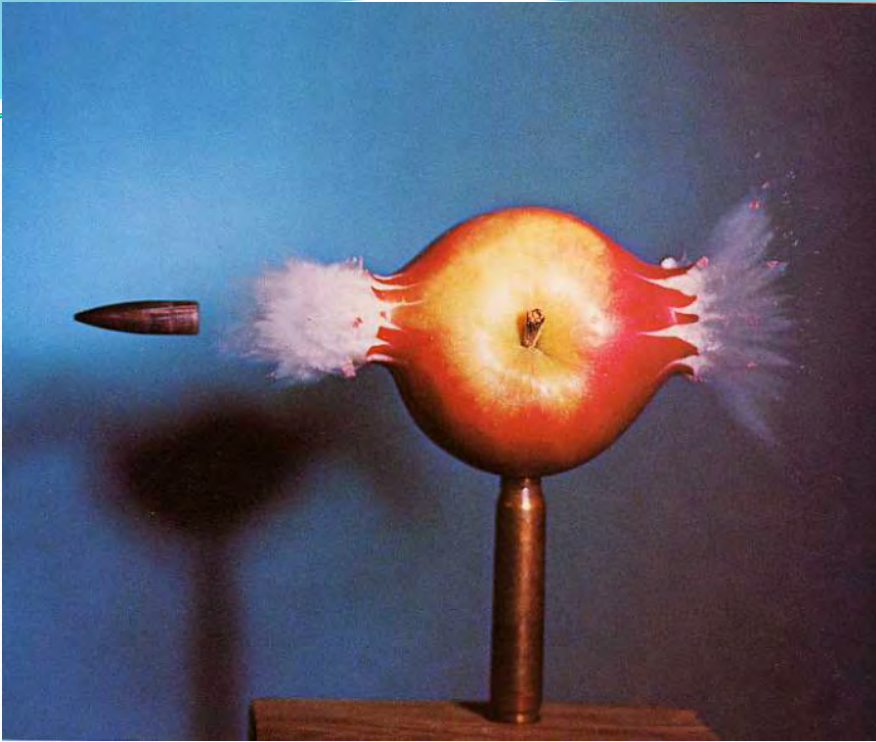
Zewail successfully used a rapid laser technique called **femtosecond spectroscopy** to observe how atoms in a molecule move during a chemical reaction. In femtosecond spectroscopy, a pump-probe experiment "photographs" chemical reactions as they happen, using an ultrafast laser as "flash".

Femtochemistry





Eadweard Muybridge



Harold Edgerton

What is Attosecond Physics

- Ultrafast optics
- Strong field physics

One Atomic unit of length $a_0 = 0.0529$ nm

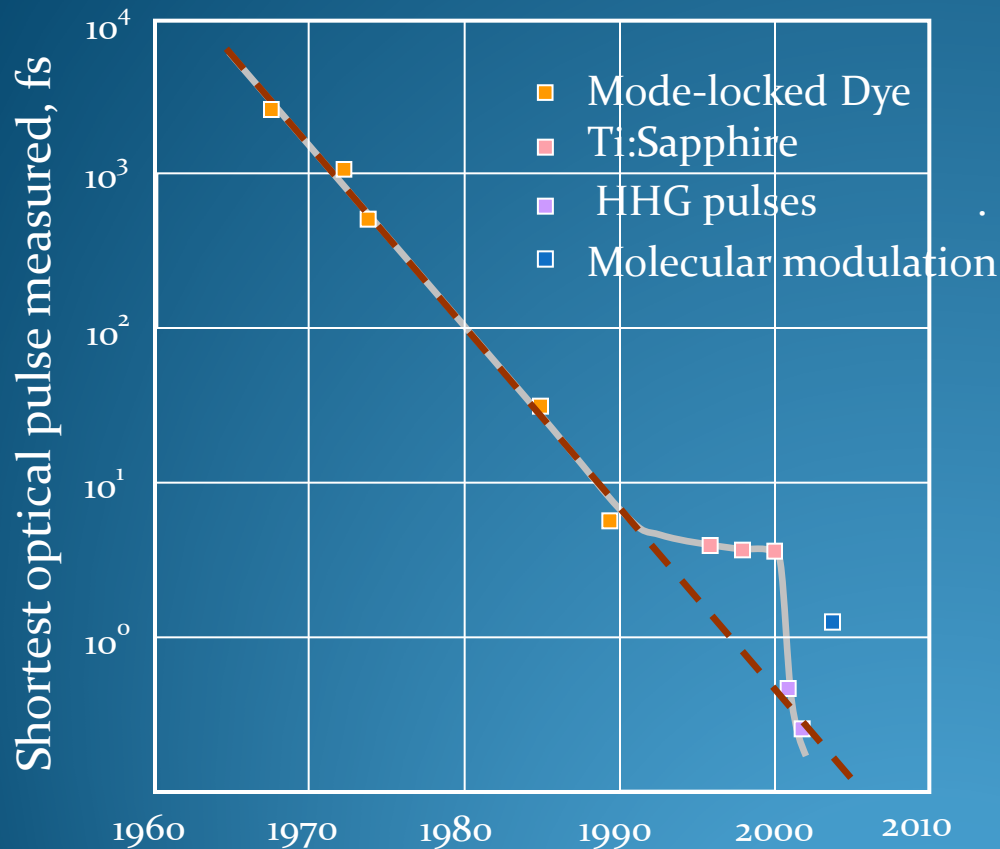
One Atomic unit of electric field $E_H = 5.142 \times 10^9$ V/cm

One Atomic unit of intensity $E_o = 3.55 \times 10^{16}$ W/cm²

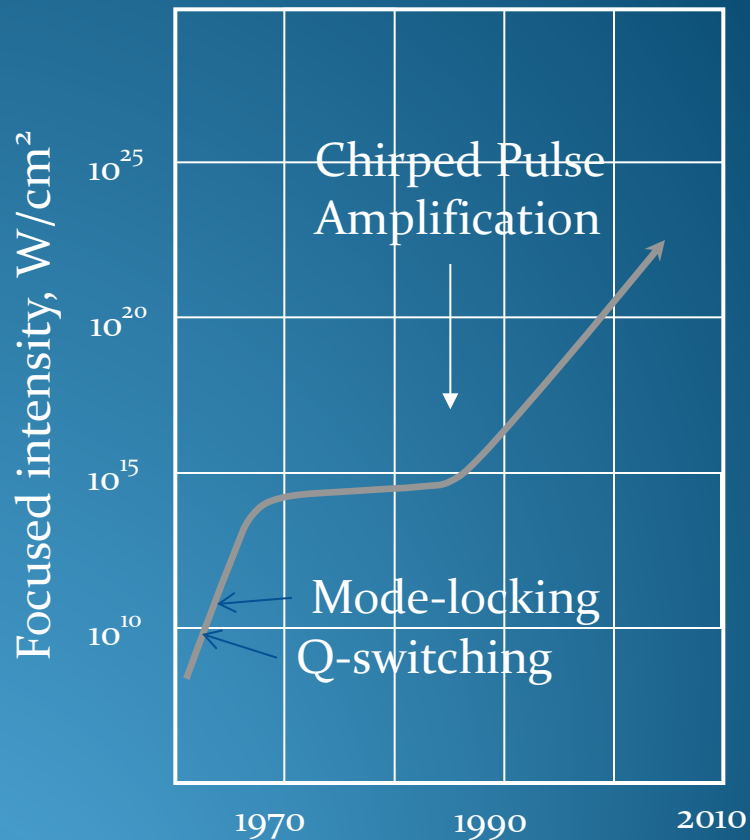
One Atomic unit of time $T_a = 24.2$ as

Laser pulses got shorter
over the years

Peak intensity
increased



Ultrafast science

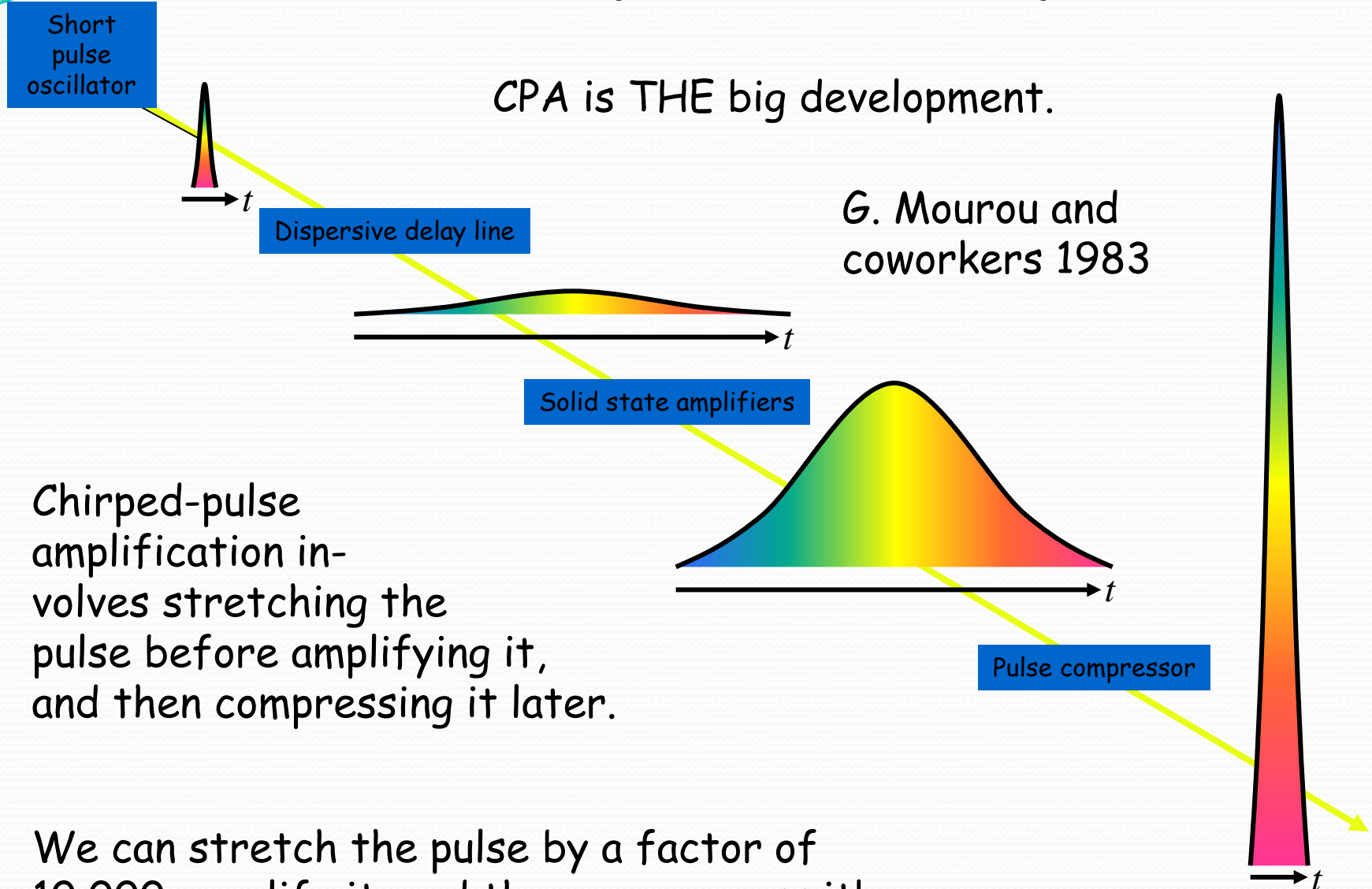


High field physics

Chirped-Pulse Amplification

CPA is THE big development.

G. Mourou and coworkers 1983

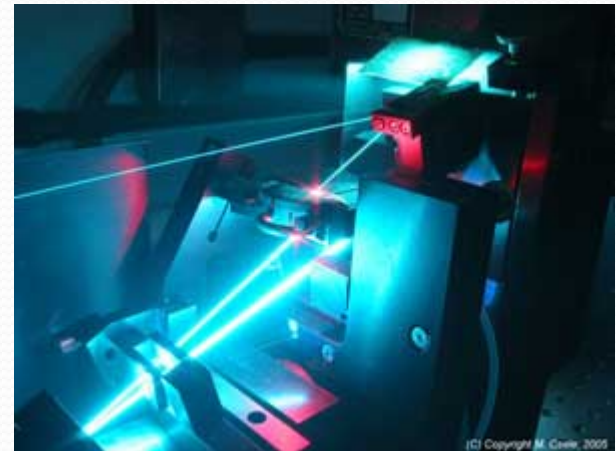


Chirped-pulse amplification involves stretching the pulse before amplifying it, and then compressing it later.

We can stretch the pulse by a factor of 10,000, amplify it, and then recompress it!

Femtosecond Laser

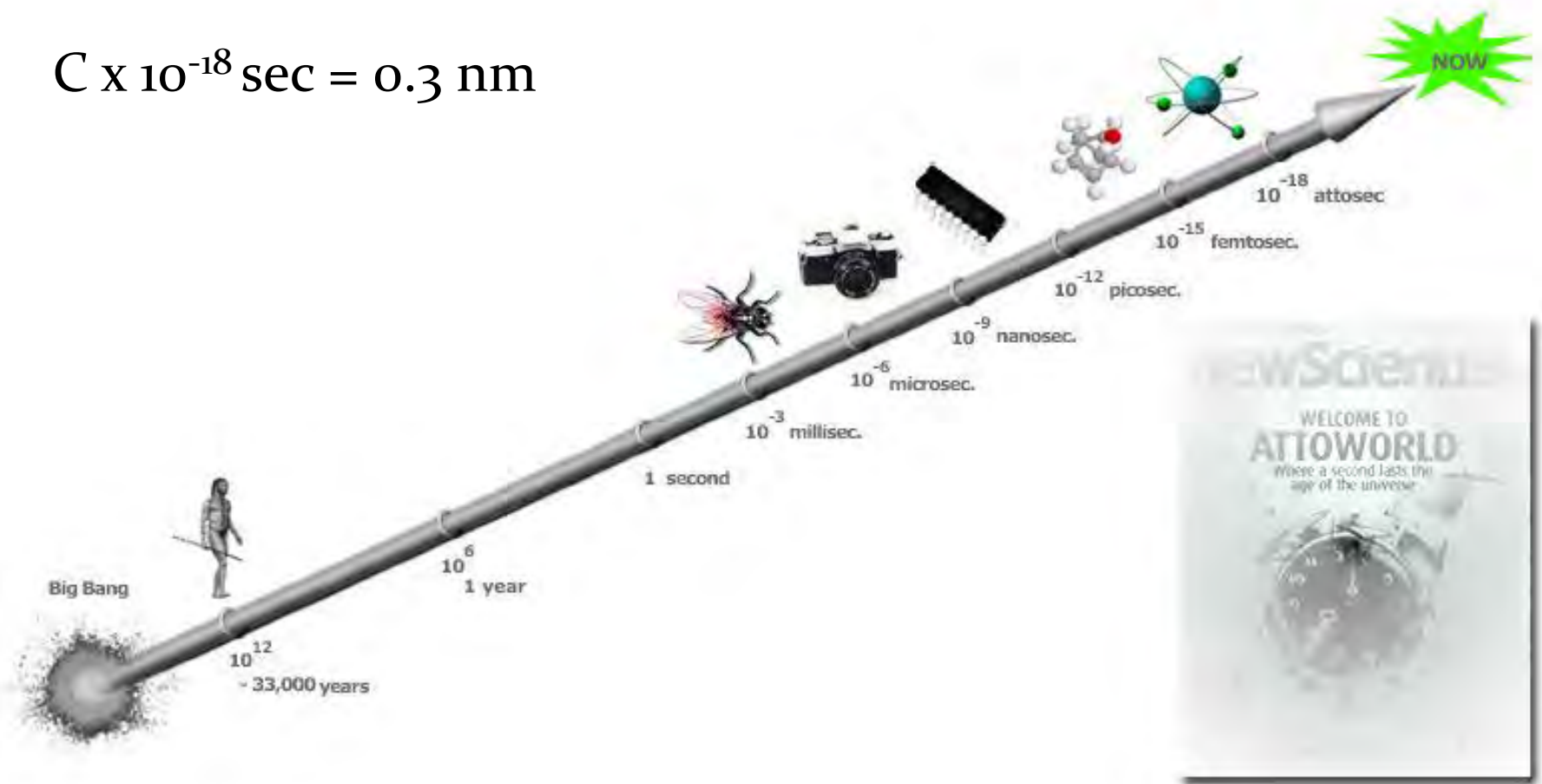
- 1974: E. P. Ippen and C. V. Shank develop the sub-picosecond mode-locked CW **dye laser**, establishing ultrafast optical science.
- 1982: P. F. Moulton develops **titanium -sapphire laser**. The titanium -sapphire laser replaces the dye laser for tunable and ultrafast laser applications.



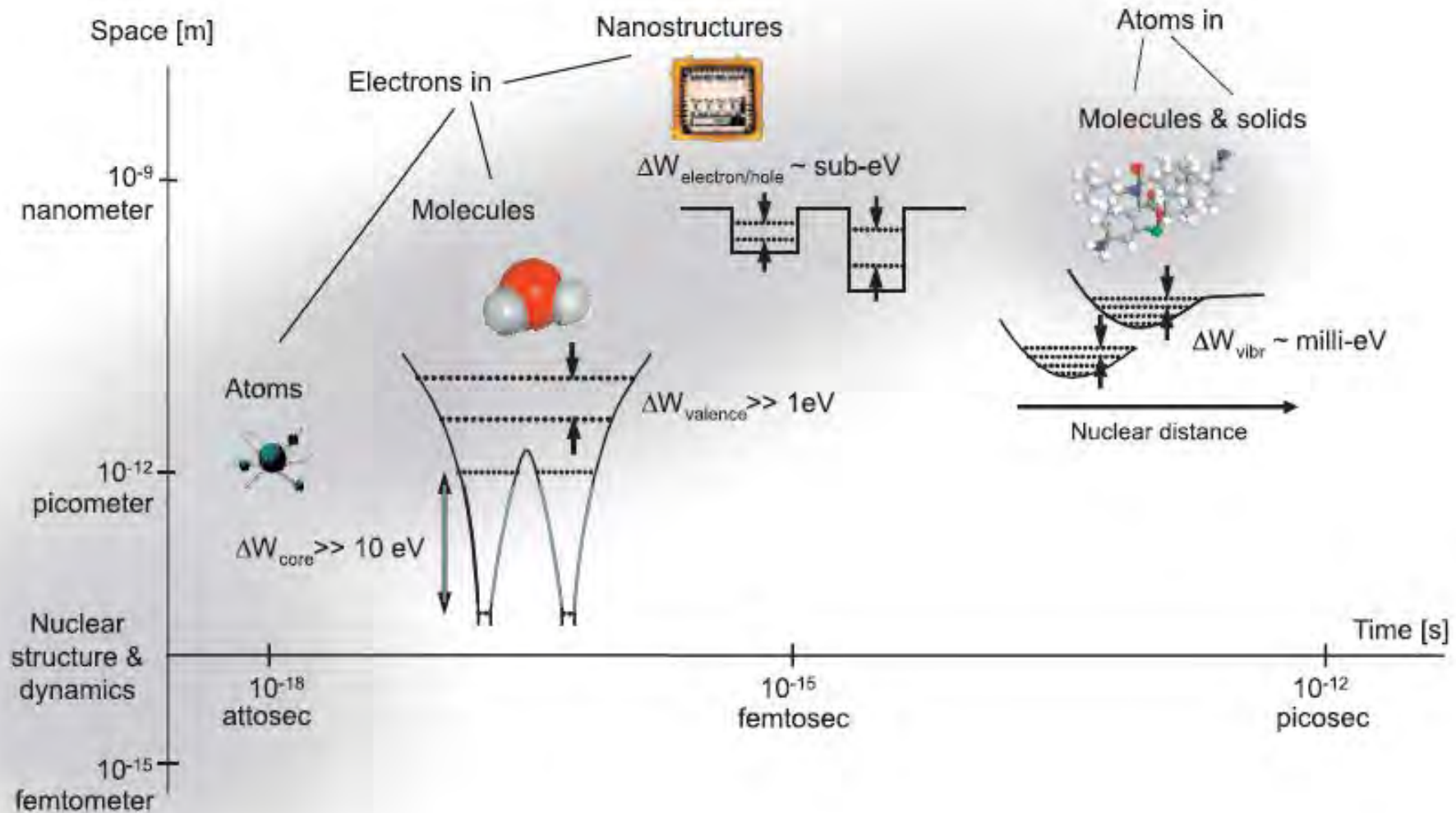
Attoworld

Attosecond: 10^{-18} second

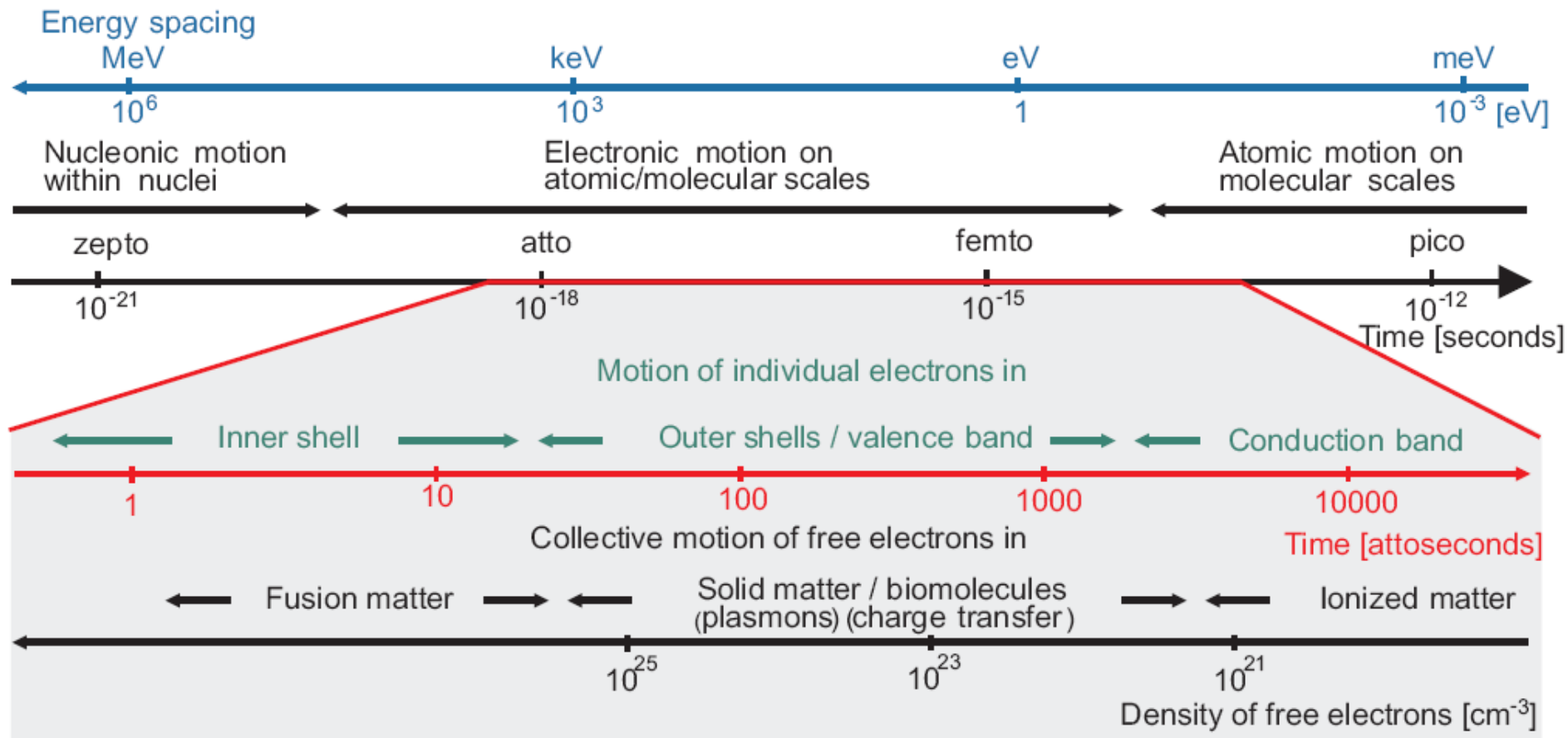
$C \times 10^{-18} \text{ sec} = 0.3 \text{ nm}$



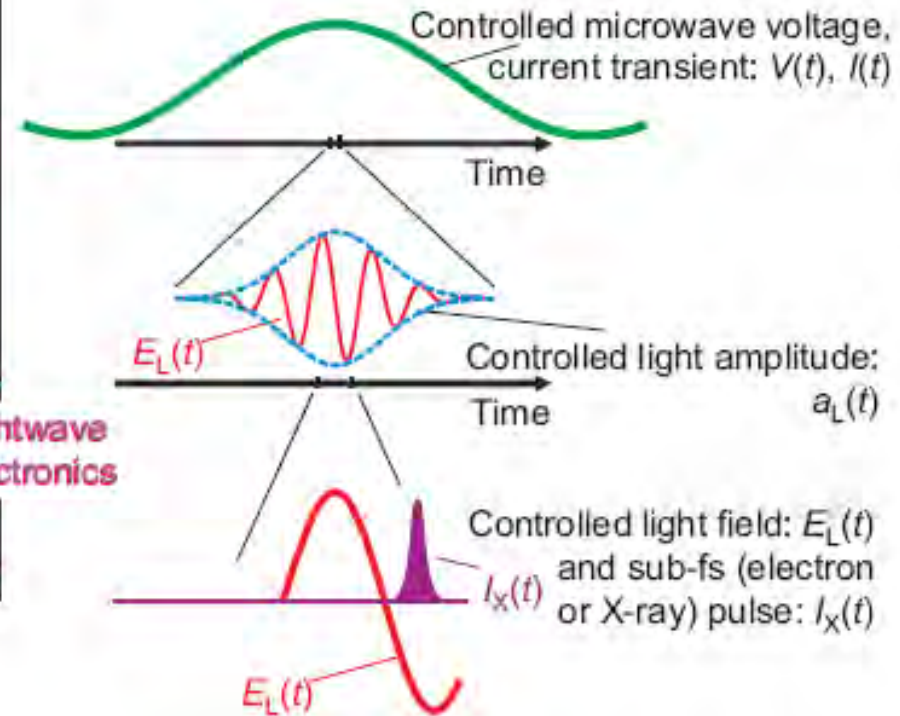
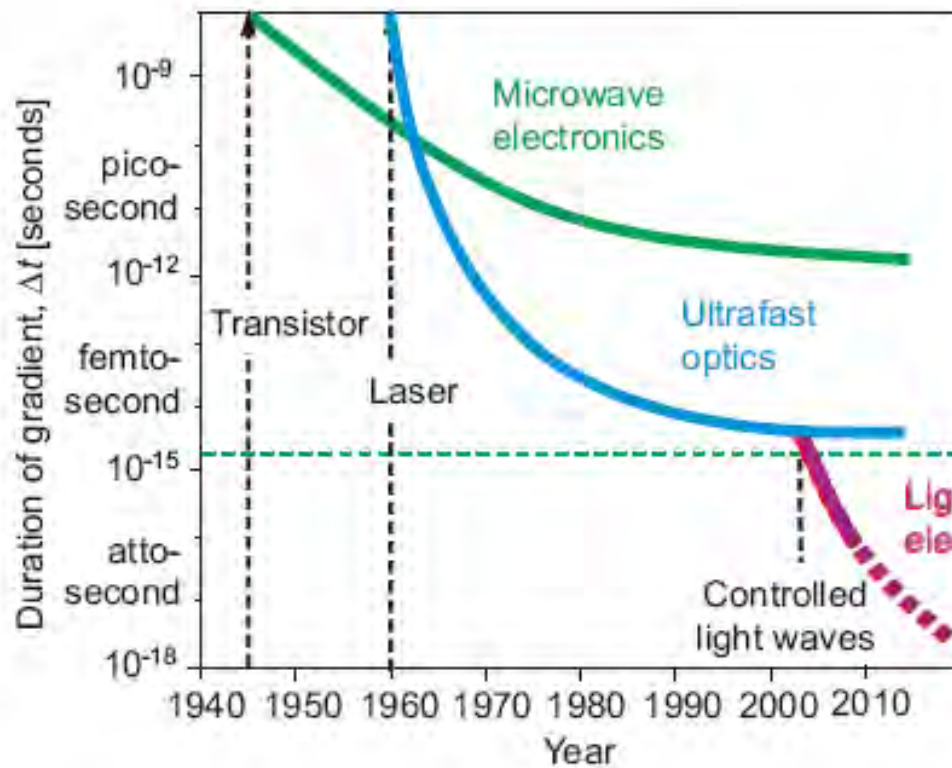
Characteristic length and time scales in the microcosm



Electronic Motion



Evolution of Ultrafast Science



Attosecond physics

- Attosecond Pulse Generation and Characterization
- Broadband High-Harmonics Generation
- High Harmonic Spectroscopy
- Ultrafast Phenomena
- Strong Field Electronic and Nuclear Dynamics
- New Ultrafast Sources and Applications
- ...

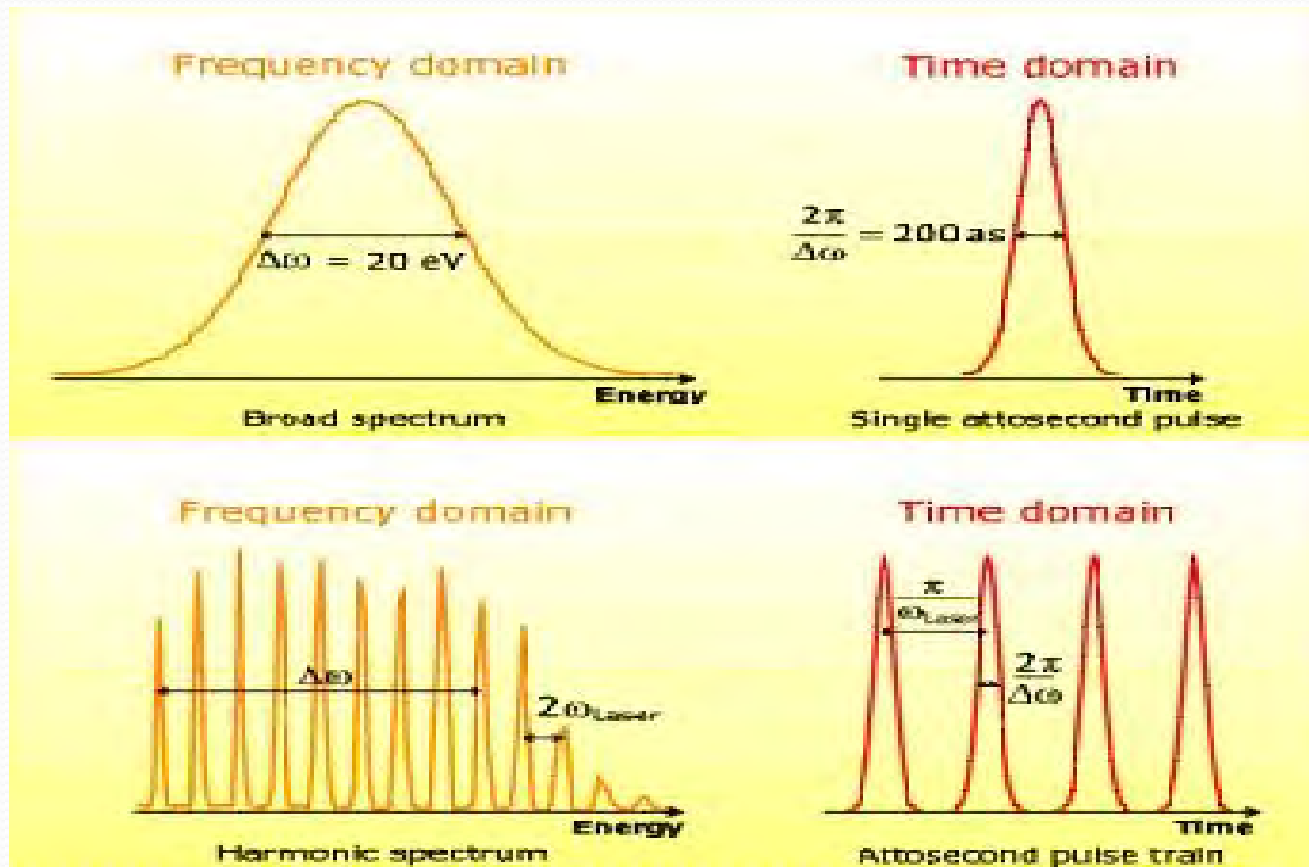
The experimental tools and techniques for electronic dynamics

- Few-cycle pulse with controllable CEP in NIR-VIS optical range
- Attosecond XUV pulse (isolated pulse, pulse train)
- Attosecond electron pulse
- Detectors for ion, electron, or photon

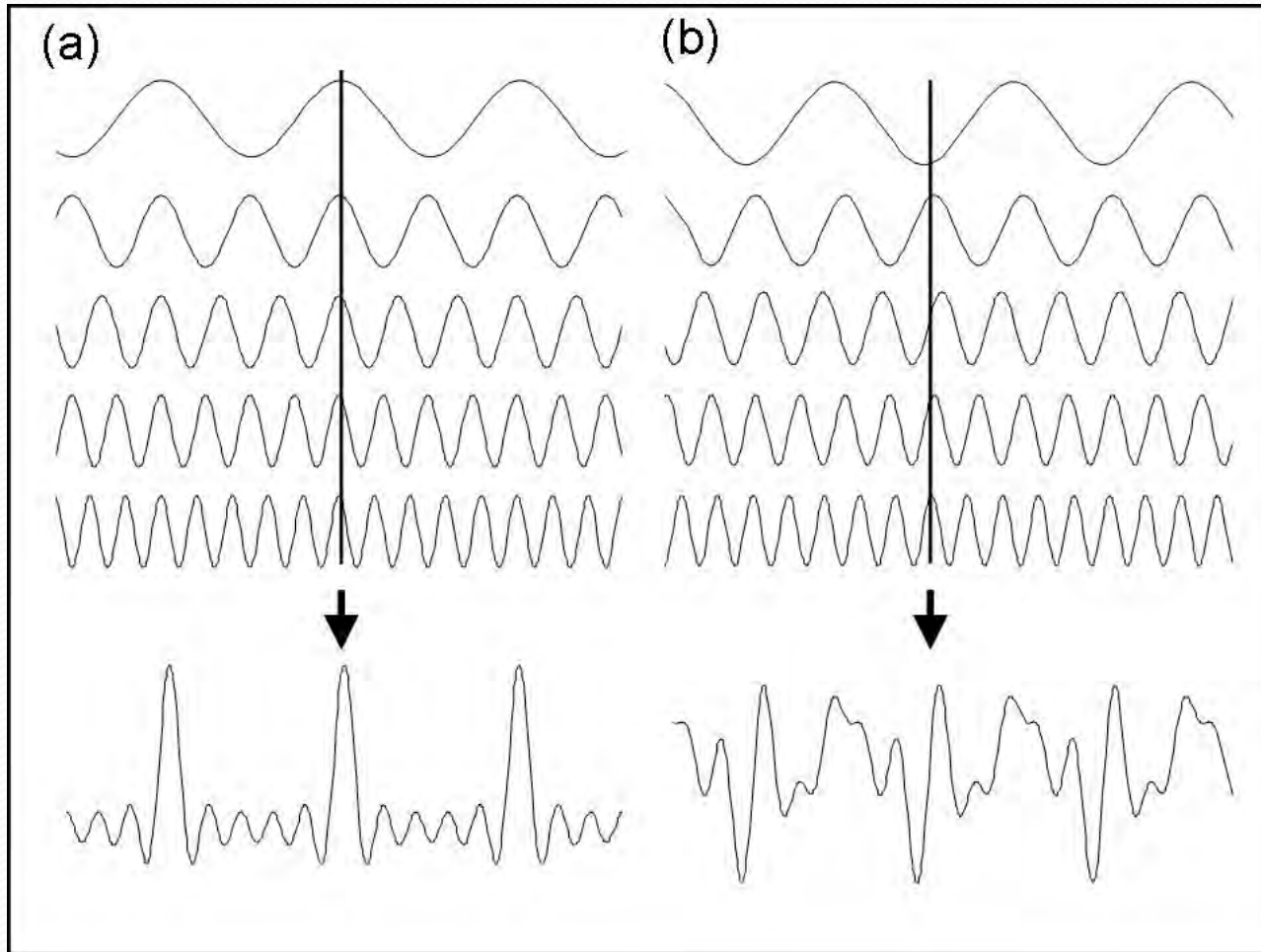
Correlation Between Time and Frequency

$$x(t - t_0) \xleftrightarrow{FT} e^{-j\omega t_0} X(\omega)$$

Fourier transform:
$$X(\omega) = \int_{-\infty}^{\infty} x(t) e^{-i\omega t} d\omega$$



Principle of optical interference of coherent light fields



In phase

Random phase

Optical cycle

$$E(t) = \tilde{E}(t) + c.c.$$

$$\tilde{E}(t) = A(t)e^{i(\omega_0 t + \phi)}$$

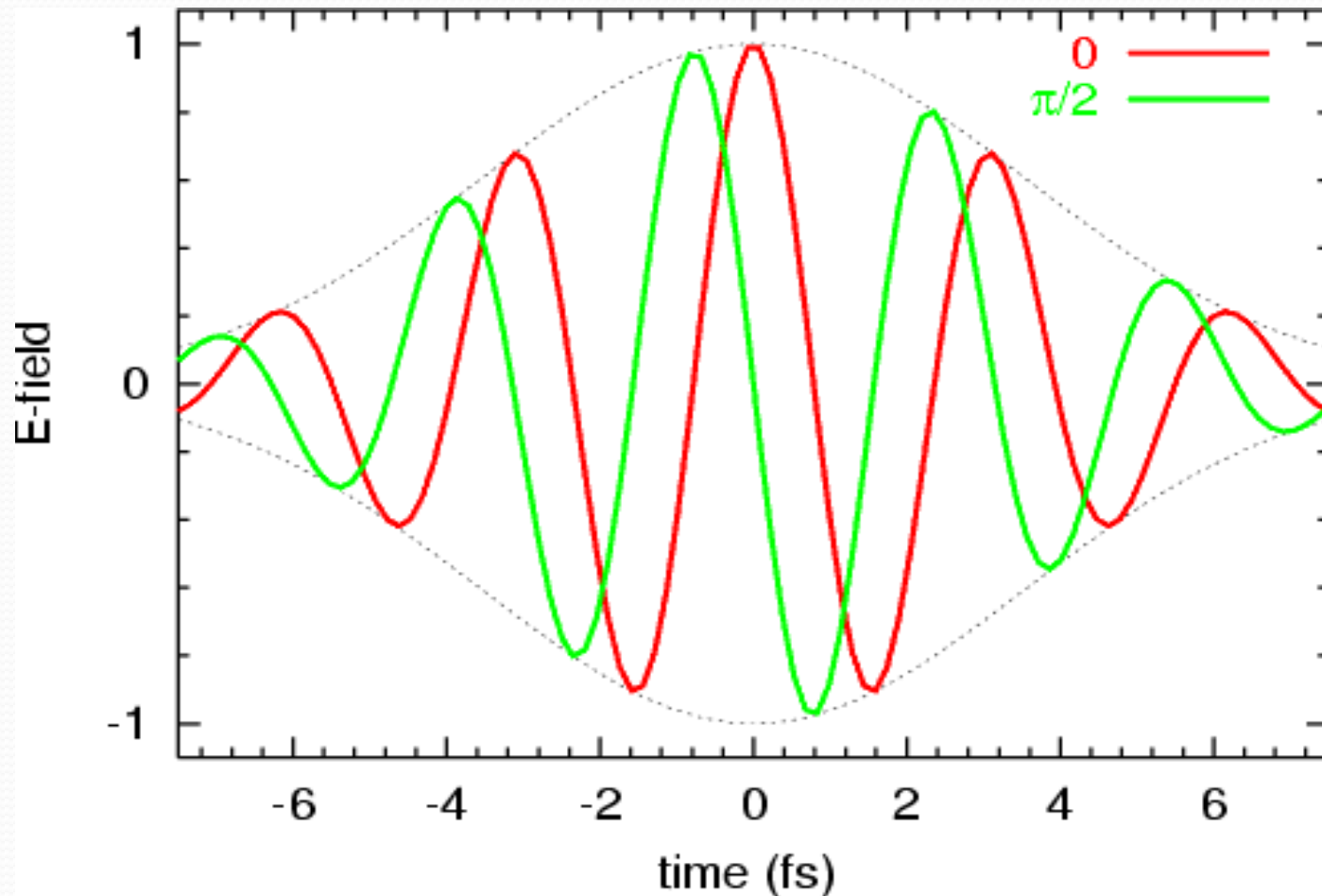
$$\omega_0 = \frac{\int_0^\infty \omega |E(\omega)|^2 d\omega}{\int_0^\infty |E(\omega)|^2 d\omega}$$

Carrier frequency

$E(\omega)$: Fourier transform of $E(t)$

Carrier envelope phase

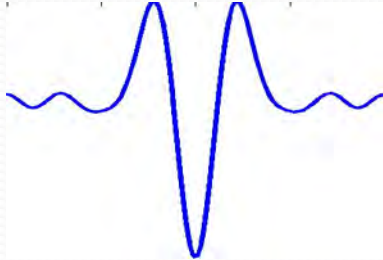
$$E(t) = E_0(t) \cos(\omega_0 t + \underline{\phi})$$



Single cycle waveforms

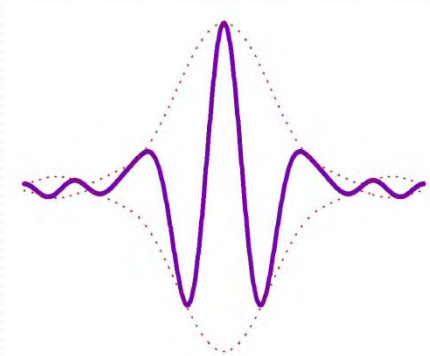
Inverted cosine

$$\phi_n = \pi$$



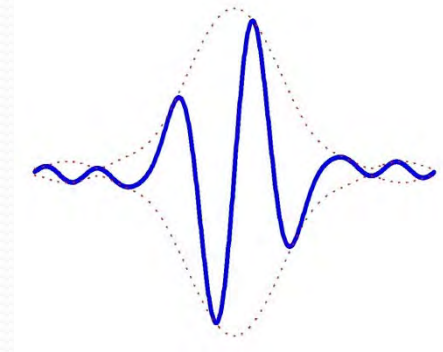
cosine pulse

$$\phi_n = 0$$



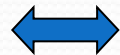
sine pulse

$$\phi_n = \pi/2$$



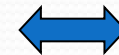
780 nm

200 nm



12,820 cm⁻¹

50,000 cm⁻¹



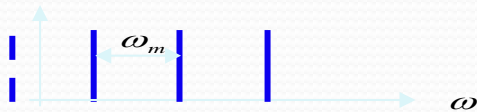
2.6 fs

684 as

Constant carrier envelope phase

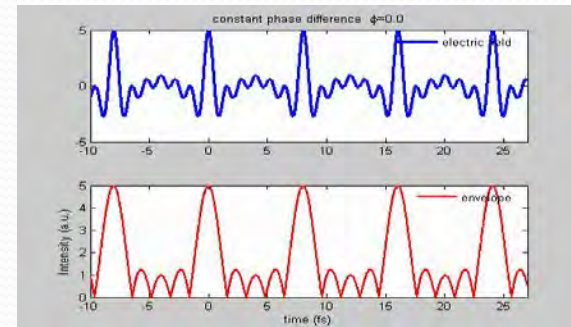
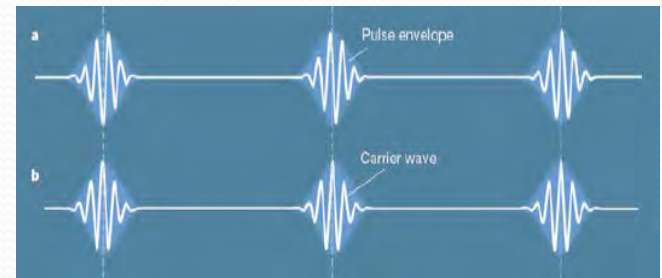
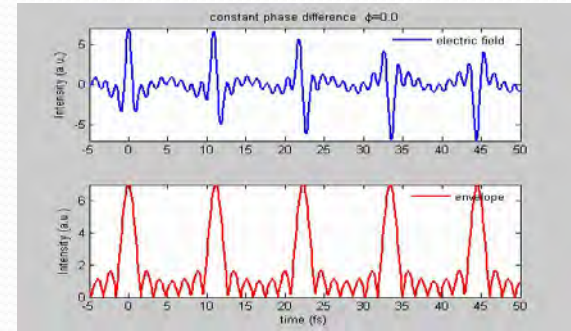
$$E(t) = \sum_n E_n(t) \cos(\omega_n t + \phi_n)$$

incommensurate $\omega_n = n\omega_m + \omega_{ceo}$ $\phi_n'(t) = \omega_{ceo}t + \phi_n$



commensurate $\omega_q = n\omega_m$ $\phi_n = \phi_{CEP} + n\phi_m$

$$E(t) = \sum_n A_n(t) \cos(n\omega_m(t + \phi_m/\omega_m) + \phi_{CEP})$$

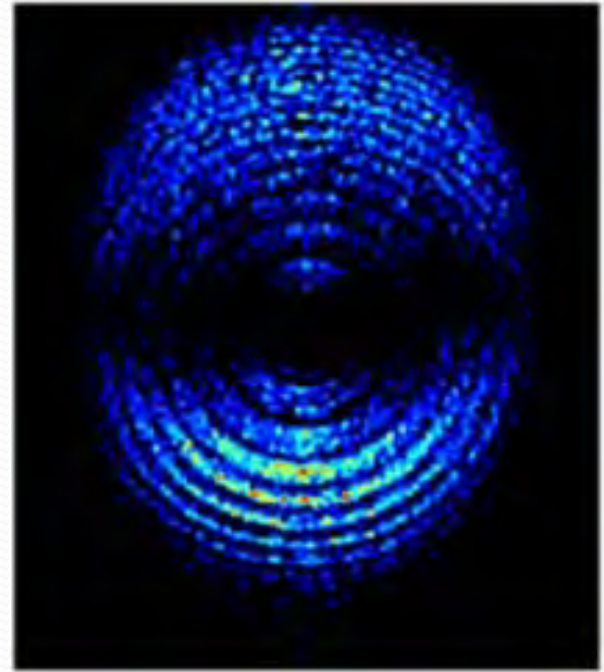


Constant CEP requires that the frequencies are commensurate and the relative phases form an arithmetic series

Attosecond pulse train for quantum stroboscope



Hummingbird wing



Electron velocity mapping

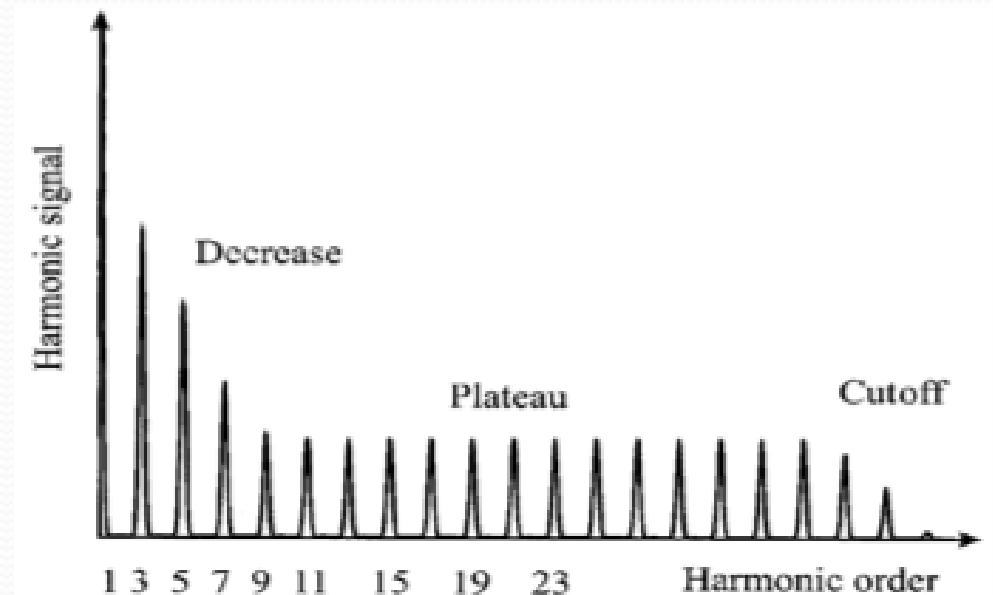
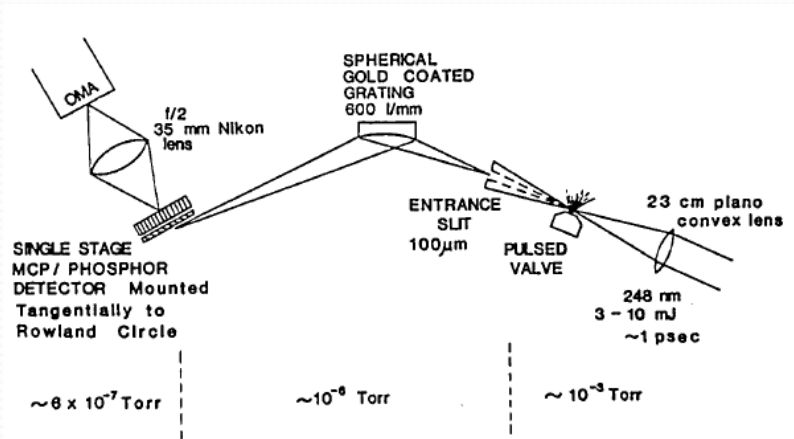
Ingredients of an attosecond single-cycle optical pulse:

1. Broad spectrum - 2 or more octaves
2. In phase condition
3. Constant carrier envelope phase:
 - Commensurate frequencies
 - Constant phase difference between adjacent spectral components
4. Stable and controllable carrier envelope phase

Lightwave control

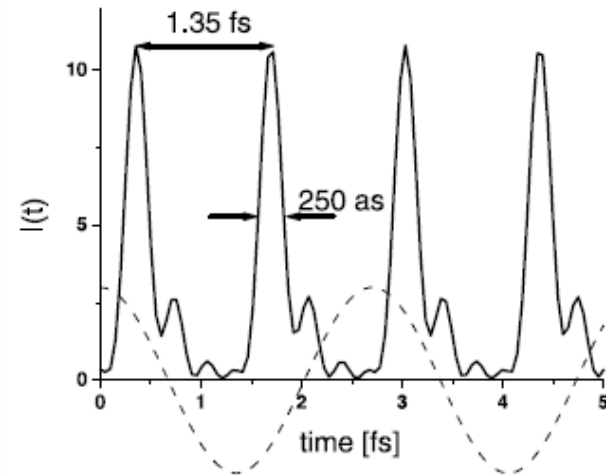
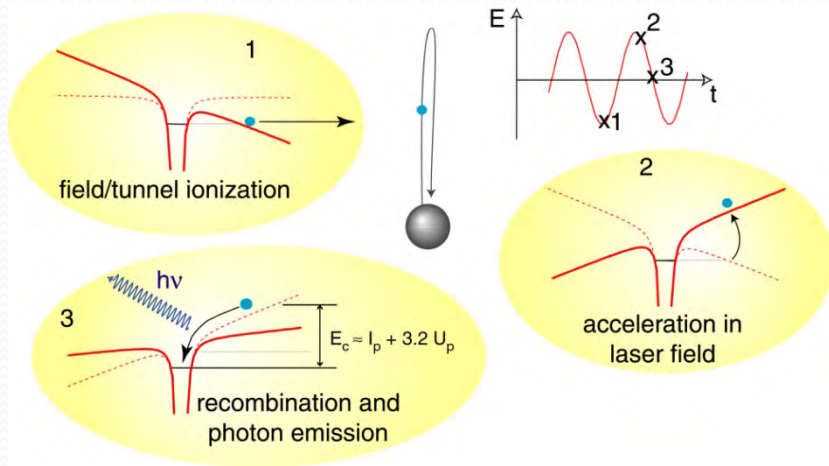
High Harmonic Generation

High-order Harmonic Generation has been a regular process to generate attosecond pulse and coherent soft X-ray.



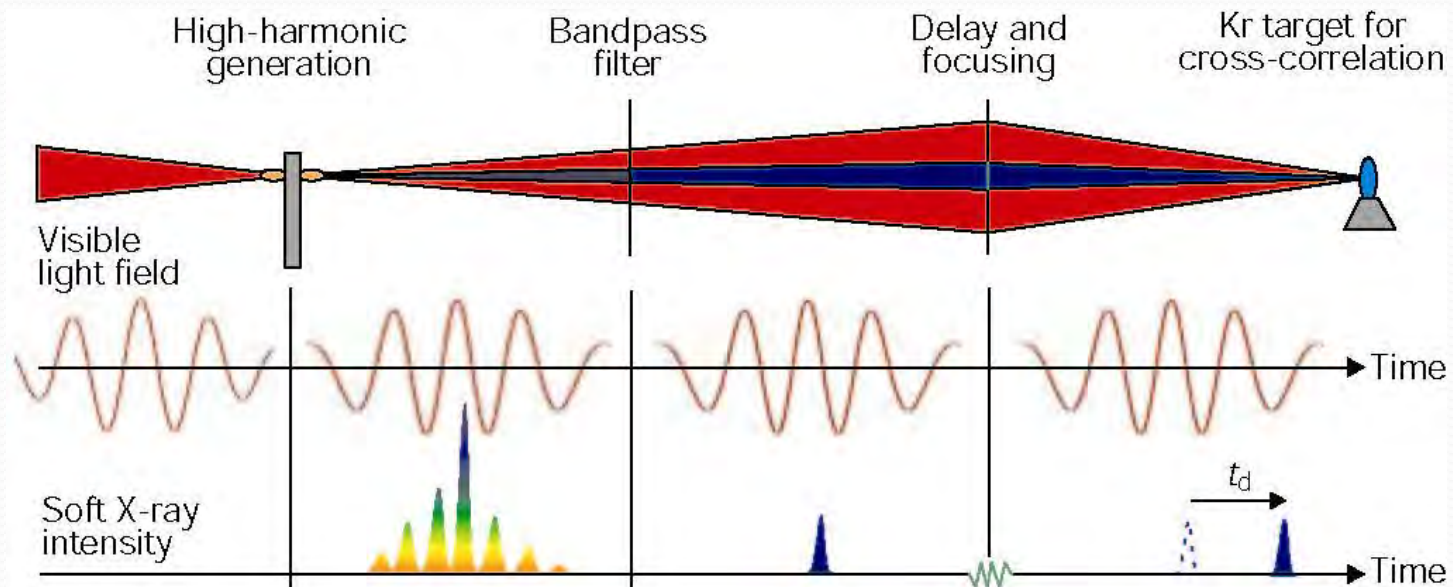
J. Opt. Soc. Am. B 4, 595 (1987)

Attosecond Pulse Generation



P. Corkum, Phys. Rev. Lett. **71**, 1994 (1993)

P. M. Paul, *et al.*, Science **292**, 1689 (2001)



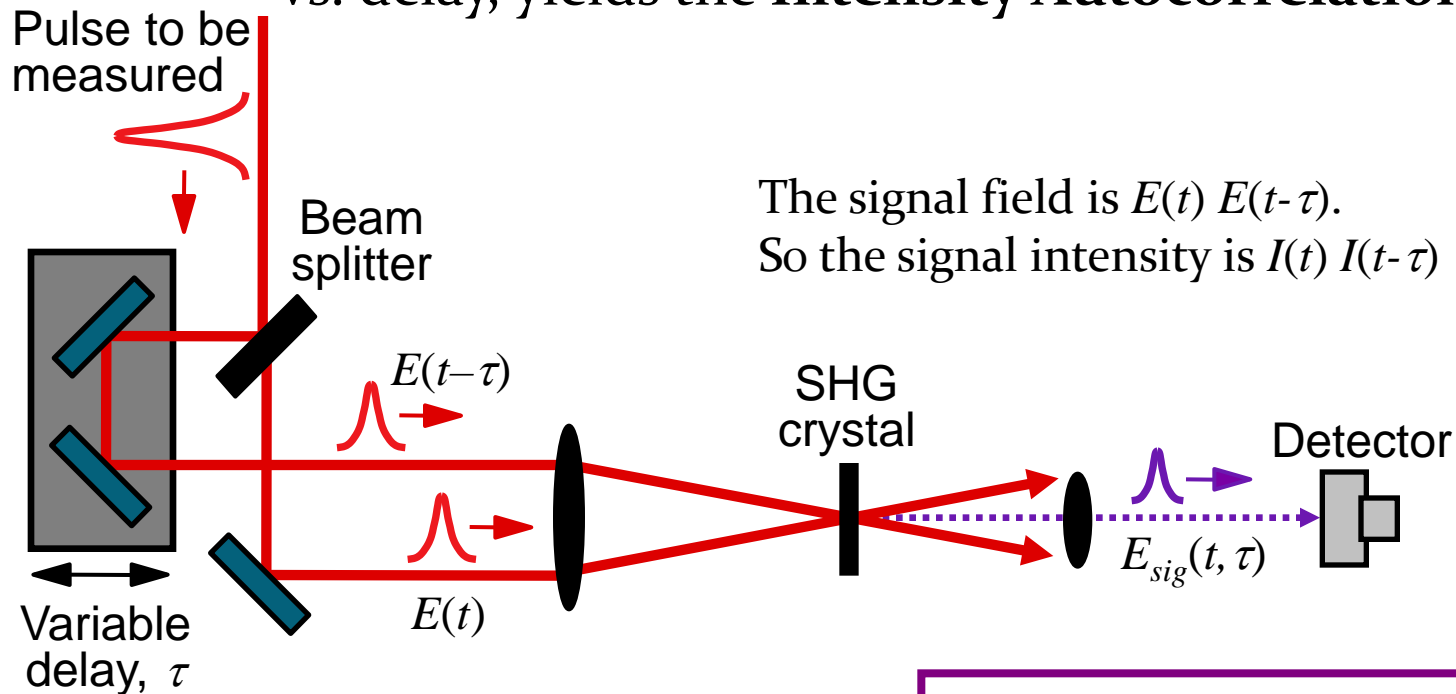
M. Hentschel *et al.*, Nature **414**, 509 (2001)

How to measure pulses

- Autocorrelation
- Frequency-Resolved Optical Gating (FROG)
- Spectral Phase Interferometry for Direct Electric-field Reconstruction (SPIDER)
- Reconstruction of Attosecond Beating by Interference of Two-Photon Transition (RABITT)
- Complete Reconstruction of Attosecond Bursts (CRAB)

Pulse Measurement in the Time Domain: *The Intensity Autocorrelator*

Crossing beams in a nonlinear-optical crystal, varying the delay between them, and measuring the signal pulse energy vs. delay, yields the **Intensity Autocorrelation**, $A^{(2)}(\tau)$.

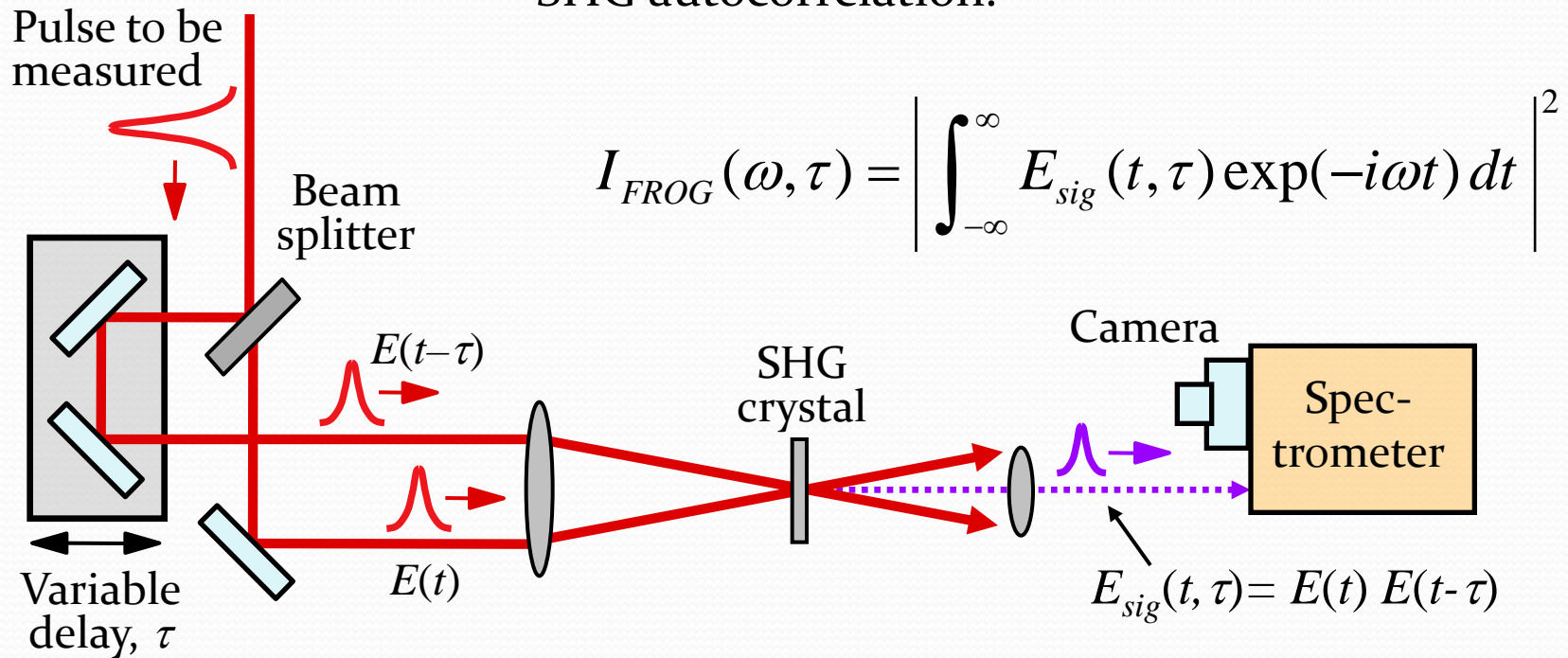


The Intensity
Autocorrelation:

$$A^{(2)}(\tau) \equiv \int_{-\infty}^{\infty} I(t) I(t-\tau) dt$$

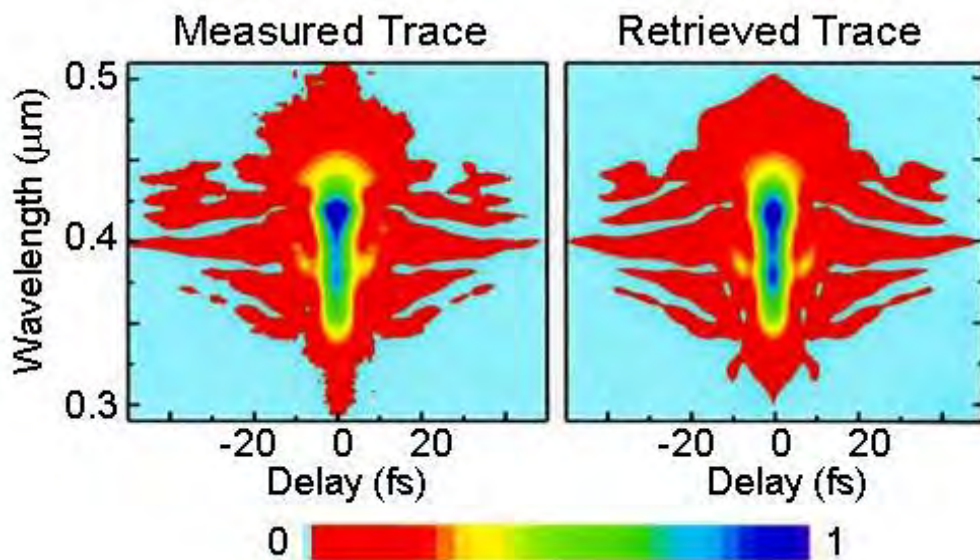
Second-harmonic-generation FROG

SHG FROG is simply a spectrally resolved SHG autocorrelation.

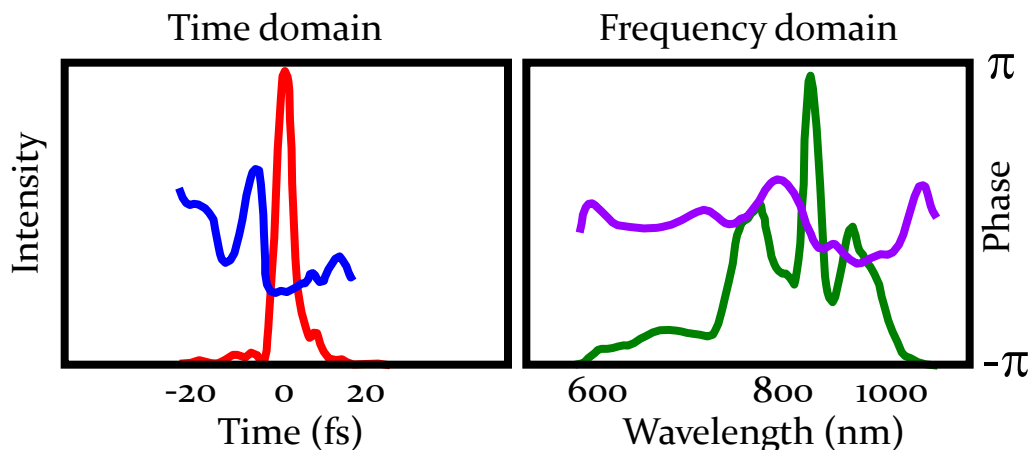


SHG FROG is the most sensitive version of FROG.

SHG FROG measurements of a 4.5-fs pulse!



Agreement between the experimental and reconstructed FROG traces provides a nice check on the measurement.



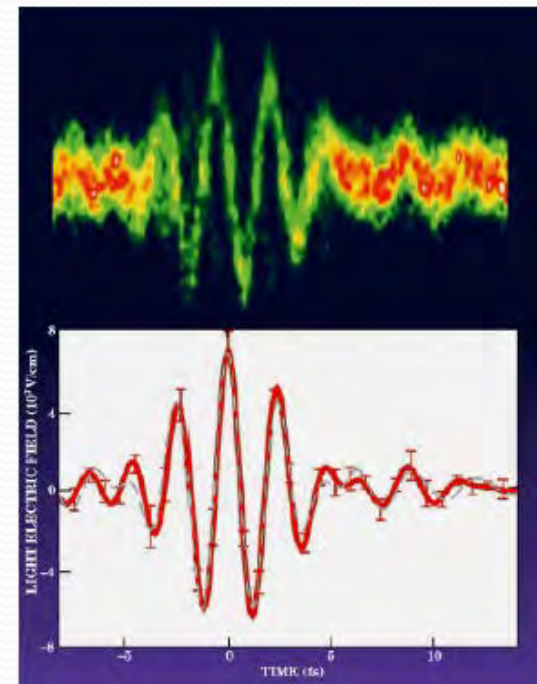
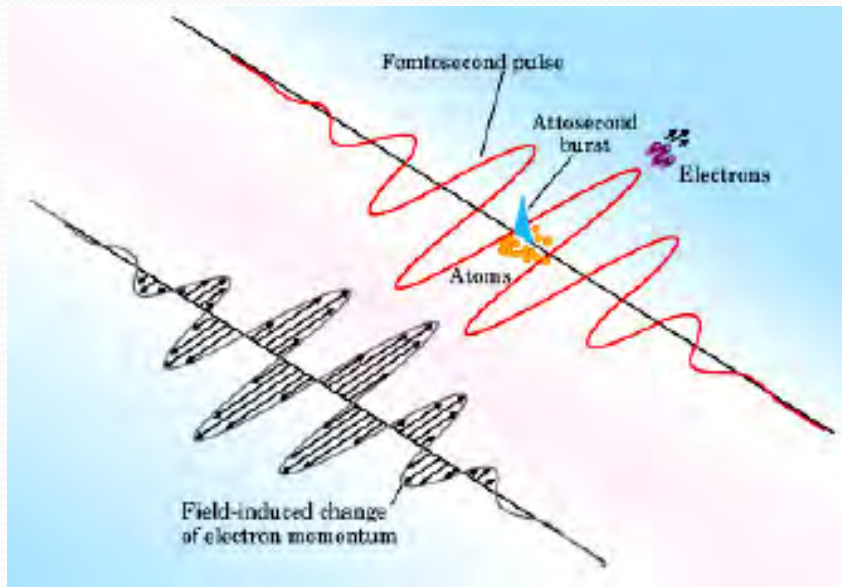
Baltuska, Pshenichnikov, and Weirisma, *J. Quant. Electron.*, 35, 459 (1999).

PHYSICS TODAY October 2004

Search and Discovery

Attosecond Bursts Trace the Electric Field of Optical Laser Pulses

The familiar textbook sketch of light's oscillating electric field can now be drawn directly from measurements.



A. Baltuska et al., *Nature* **421**, 611 (2003)

E. Goulielmakis et al., *Science* **305**, 1267 (2004)

Attosecond spectroscopy in condensed matter

A. L. Cavalieri¹, N. Müller², Th. Uphues^{1,2}, V. S. Yakovlev³, A. Baltuška^{1,4}, B. Horvath¹, B. Schmidt⁵, L. Blümel⁵, R. Holzwarth⁵, S. Hendel², M. Drescher⁶, U. Kleineberg³, P. M. Echenique⁷, R. Kienberger¹, F. Krausz^{1,3} & U. Heinzmann²

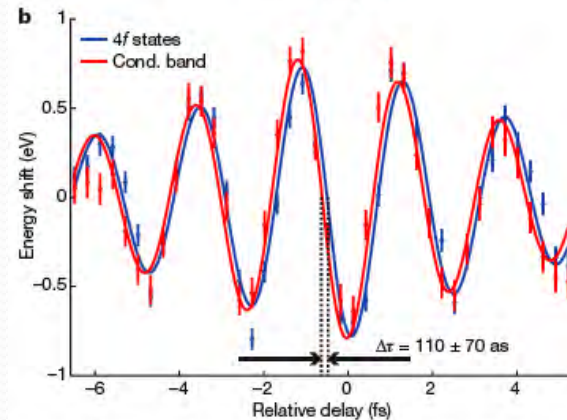
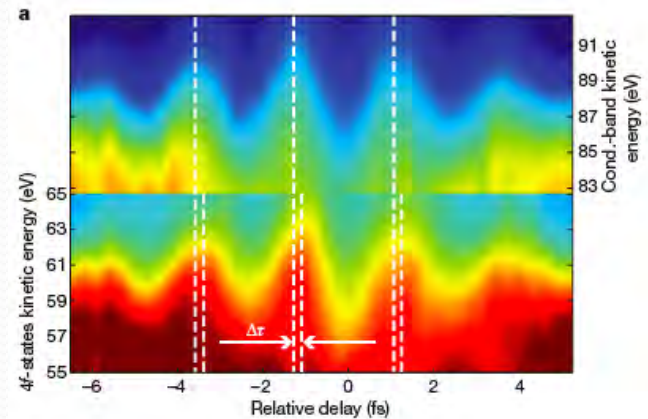
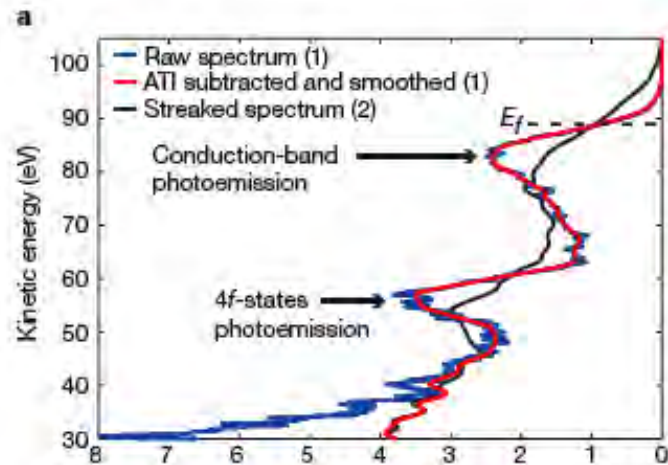
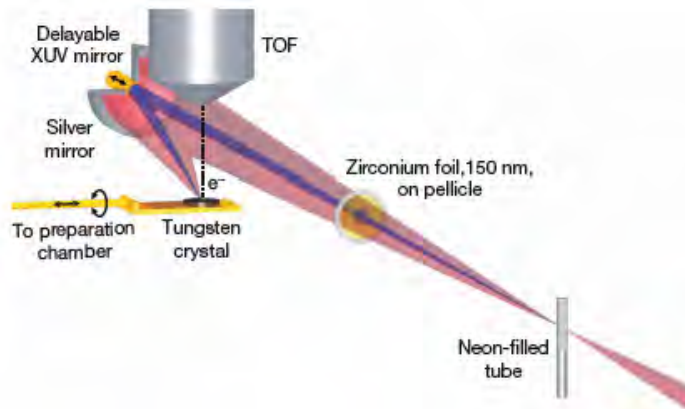
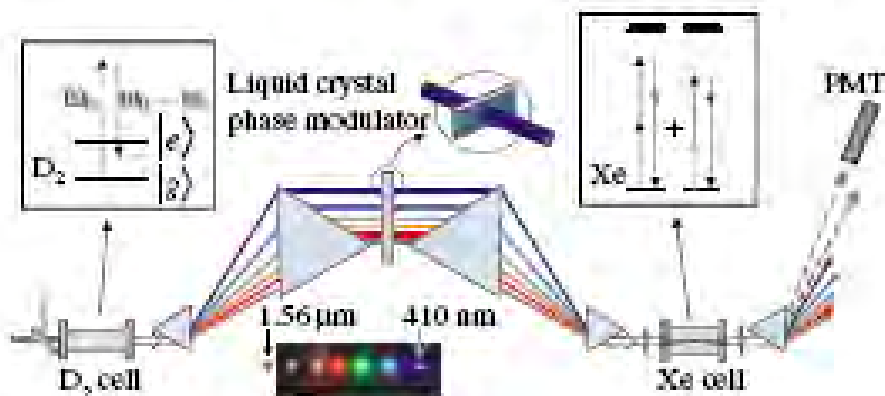
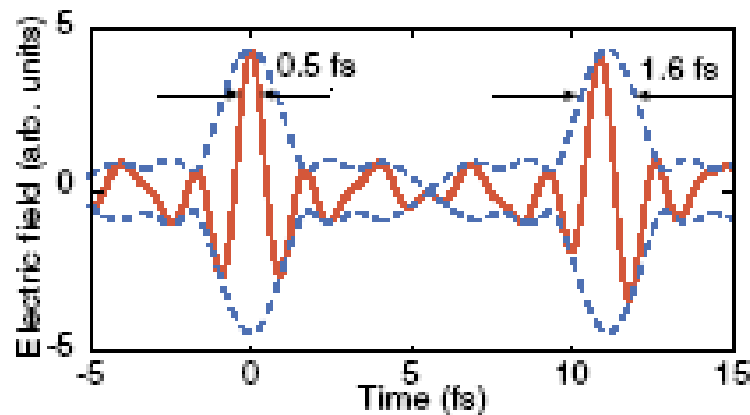
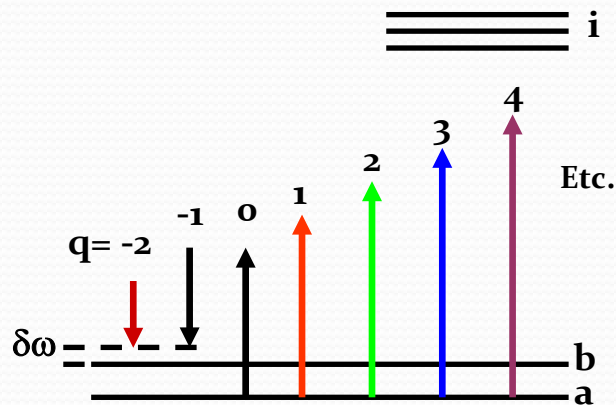


Figure 3 | Evidence of delayed photoemission. a, The 4f and conduction-band spectrograms, following cubic-spline interpolation of the measured data

Nature 449, 1029 (2007)

Methods of generating attosecond pulses

High-order stimulated Raman scattering using molecular modulation



Advantages: IR-UV region

good power

single-cycle

Disadvantages: complex setup

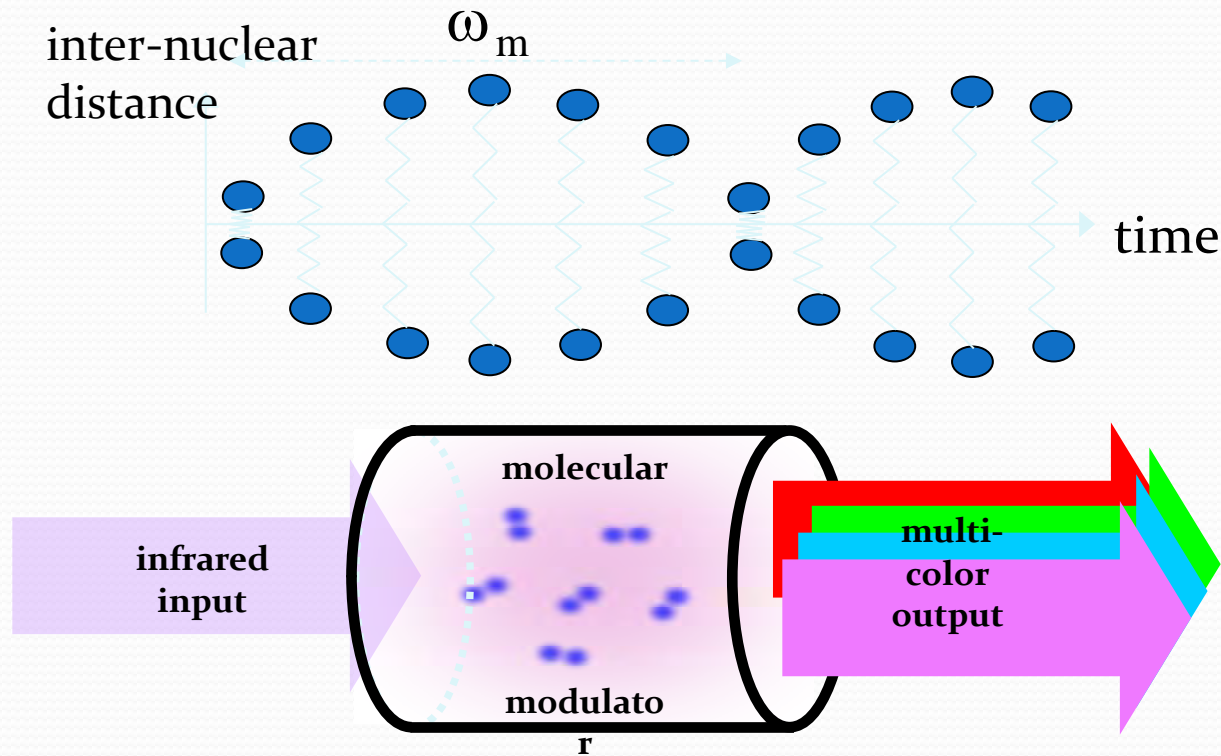
8-50 fs pulse spacing

limited to ~ 300-500 as

M.Y. Shverdin et.al., PRL 94, 033904 (2005)

Molecular Modulation

Molecular modulation is analogous to electro-optic modulation

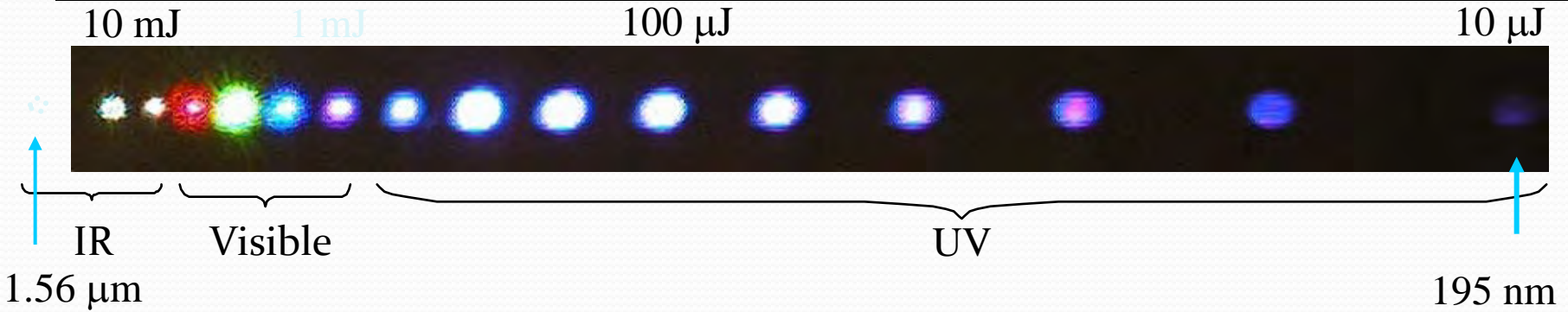


Refractive Index $n = n_0 + \delta \cos \omega_m t$

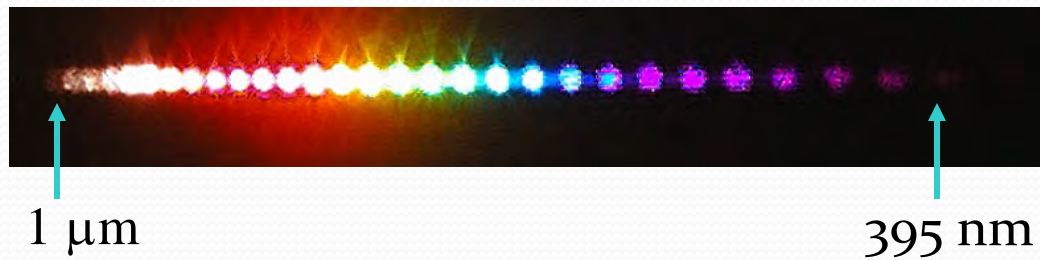
$$\omega_q = \omega_0 + q\omega_m \quad q = -2, -1, 0, 1, 2, 3, \dots$$

Alexei Sokolov
Steve Harris

D₂ Vibration Spectra: 16 sidebands, spaced by 2994 cm⁻¹

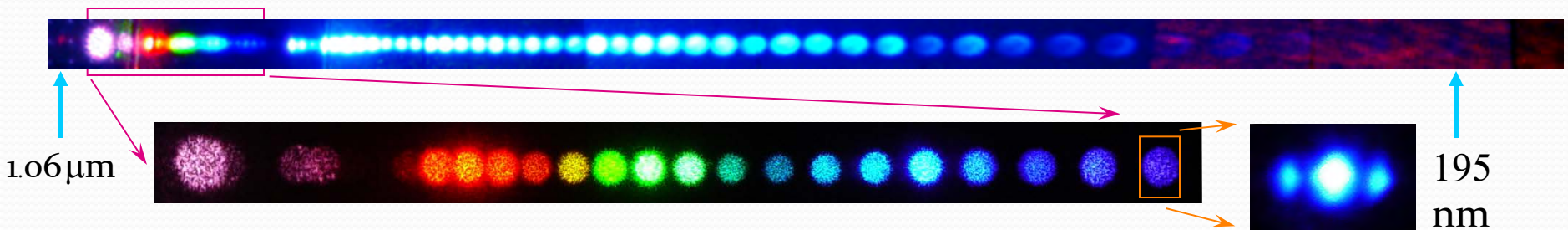


H₂ Rotation Spectra: 29 sidebands, spaced by 587 cm⁻¹

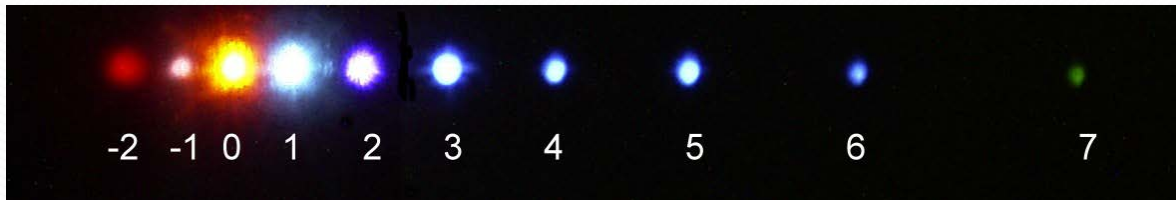
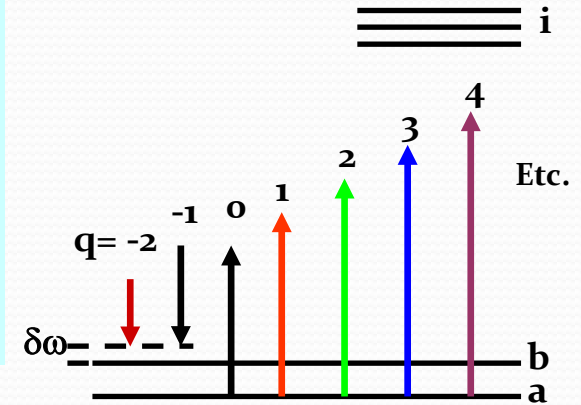
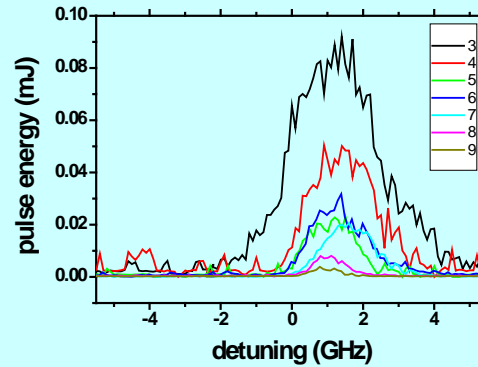
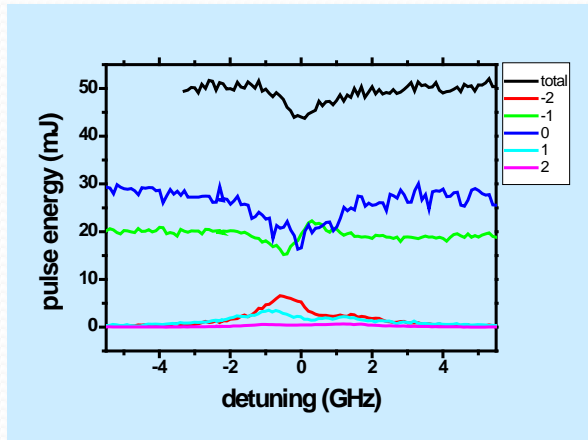


Phys. Rev. A (R)(1997)
Phys. Rev. Lett. 81 (1998)
Opt. Lett. 24 (1999)
Phys. Rev. Lett. 84 (2000)
Phys. Rev. Lett. 85 (2000)
Phys. Rev. A 63 (2001)
Phys. Rev. Lett. 91 (2003)
Phys. Rev. Lett. 93 (2005)

Multiplicative Spectra: ~ 200 sidebands, spaced by < 587 cm⁻¹



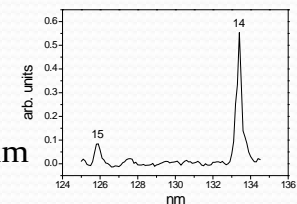
Raman sidebands generated



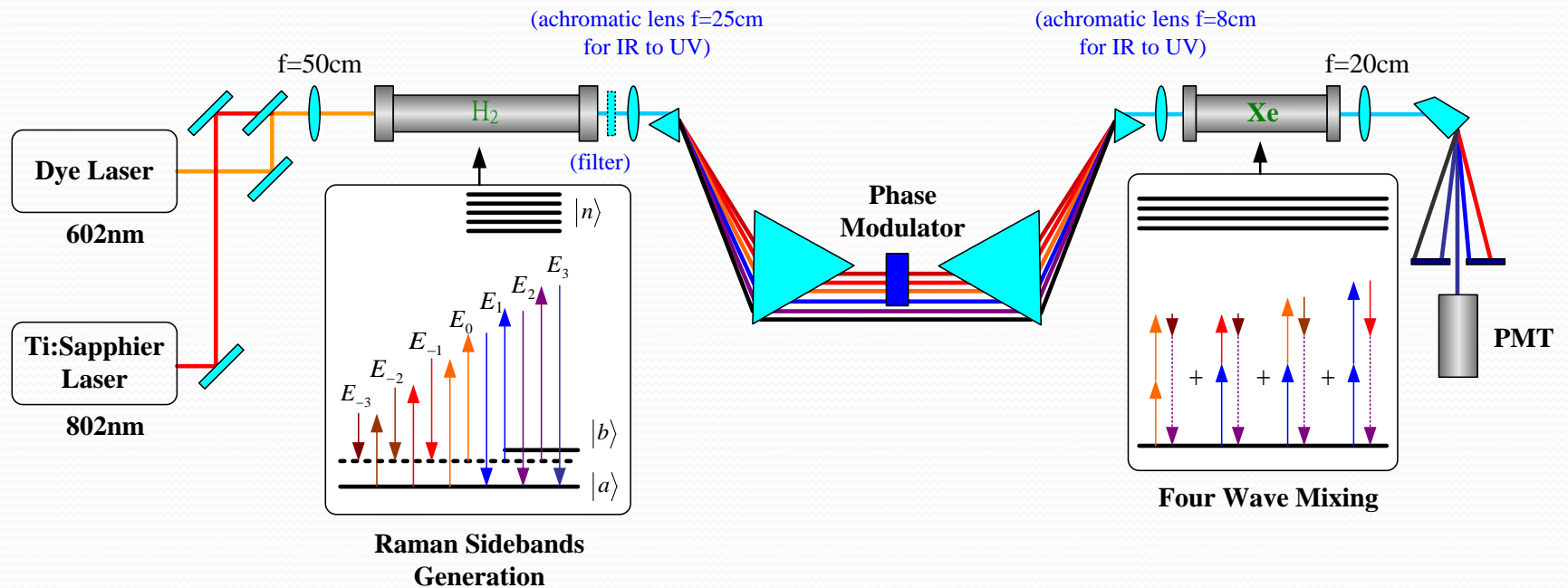
$q = 3$	$\lambda = 339.6$ nm
6	238.6 nm
9	183.9 nm
12	149.6 nm
14	133.0 nm

Total spectral span $>70,000 \text{ cm}^{-1}$
($\sim 500 \text{ as}$)

15th order at 126 nm observed



Experiment Setup



Status of sub-cycle optical pulse generation by molecular modulation

IAMS sub-cycle source

0.833 cycle per pulse

1.4 fs envelope

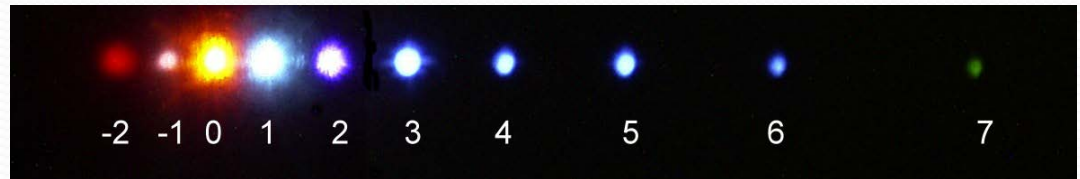
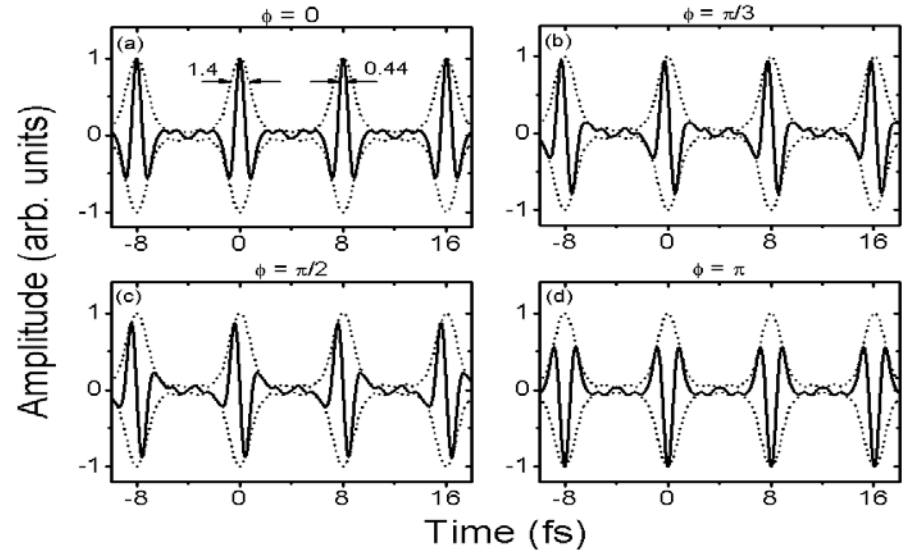
440 as cycle width

constant carrier envelope phase

2 ns pulse train duration

8.0 fs pulse spacing

~1 MW peak power



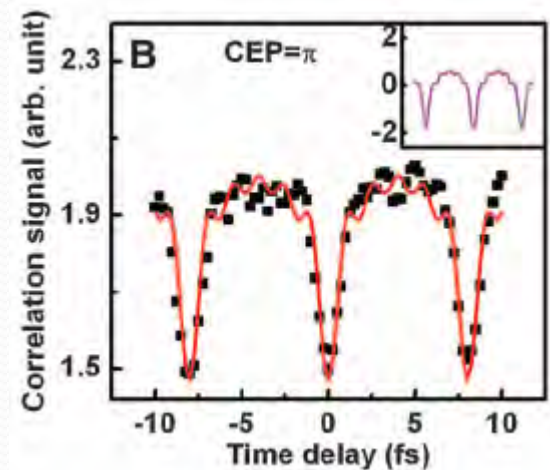
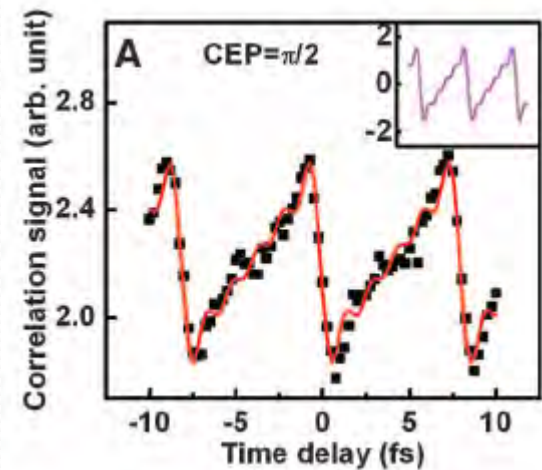
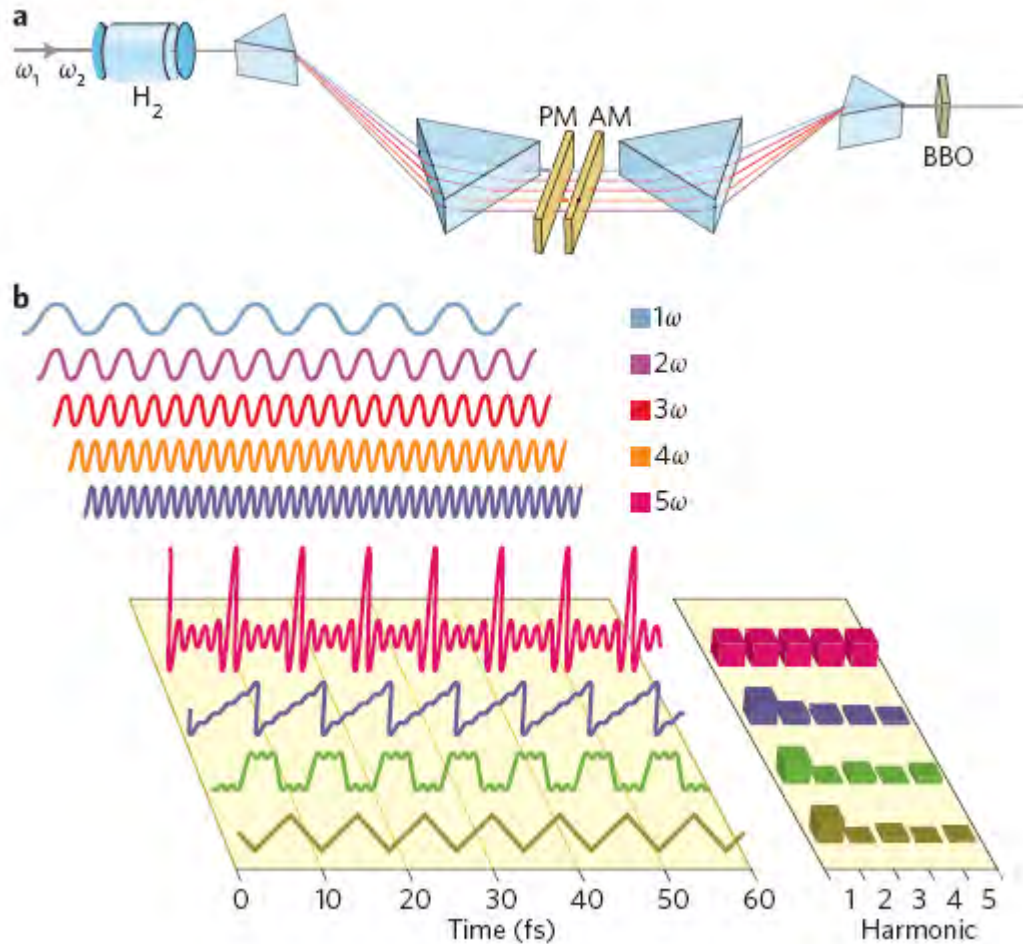
Phys. Rev. Lett. 100, 163906

(2008)

Phys. Rev. Lett. 102, 213902 (2009)

Total spectral span $>70,000 \text{ cm}^{-1}$

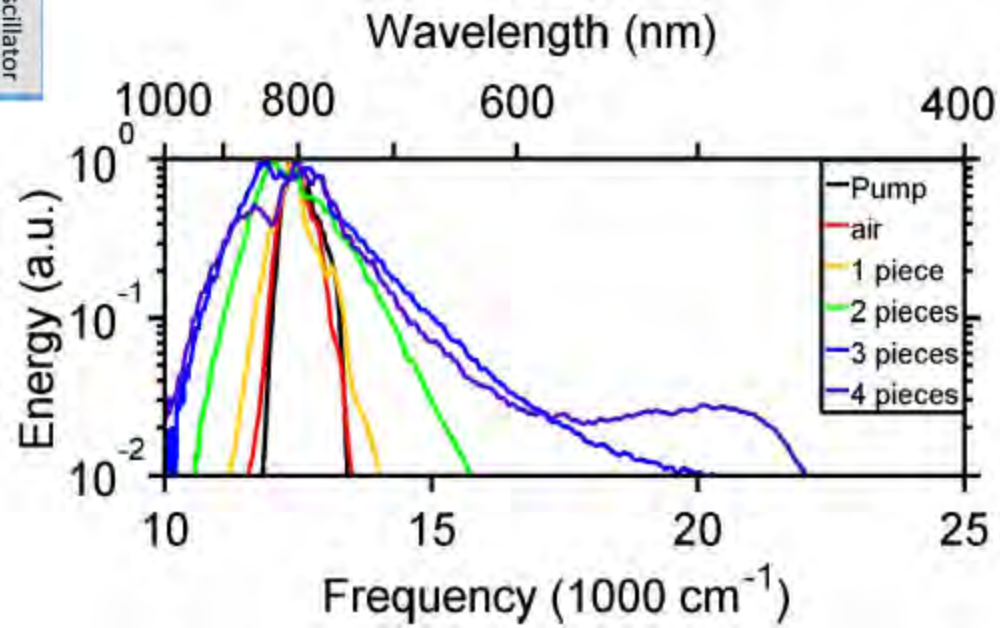
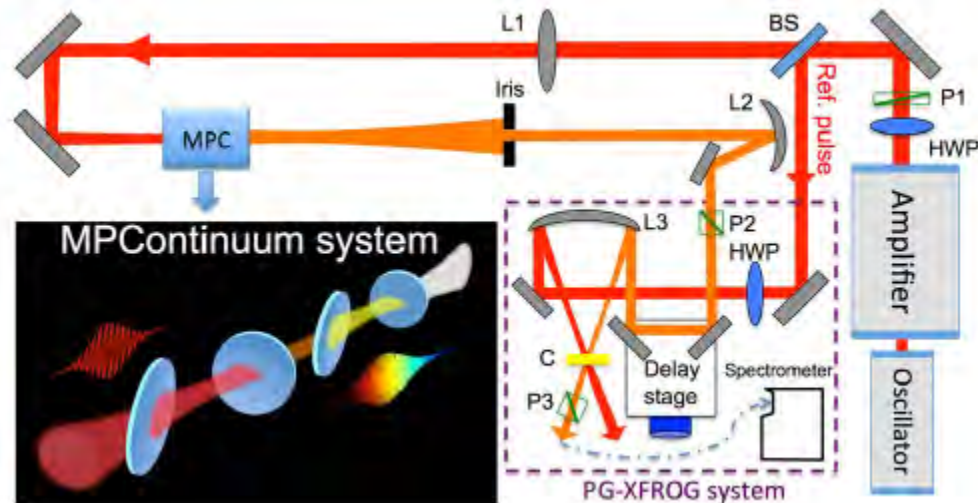
Synthesis and measurement of ultrafast waveform



Han-Sung Chan, et al.,
Science 331, 1165 (2011)

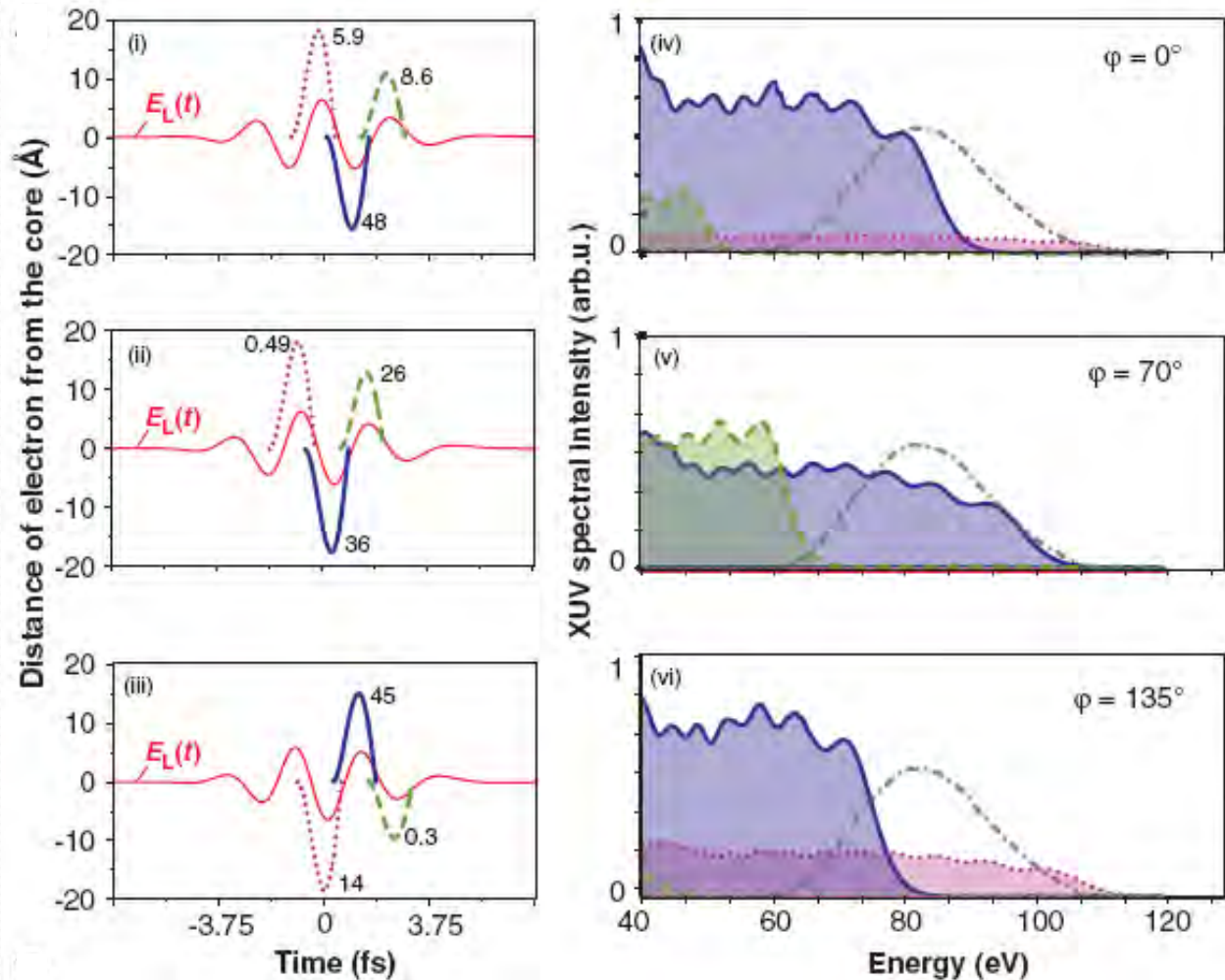
Generation of intense supercontinuum in condensed media

CHIH-HSUAN LU,¹ YU-JUNG TSOU,¹ HONG-YU CHEN,¹ BO-HAN CHEN,¹ YU-CHEN CHENG,² SHANG-DA YANG,¹ MING-CHANG CHEN,¹ CHIA-CHEN HSU,³ AND A. H. KUNG^{1,2,*}



Single-cycle Nonlinear Optics

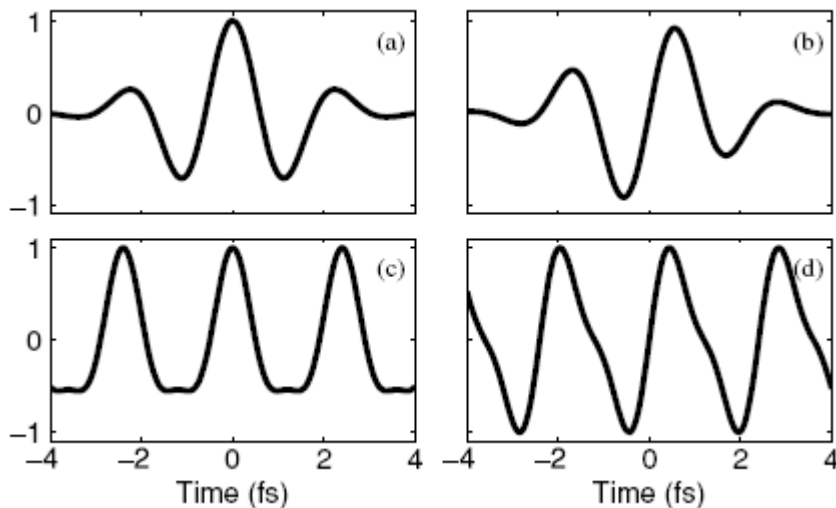
Simulation of sub-femtosecond XUV emission from neon atoms ionized by a linearly polarized, sub-1.5-cycle, 720 nm laser field.



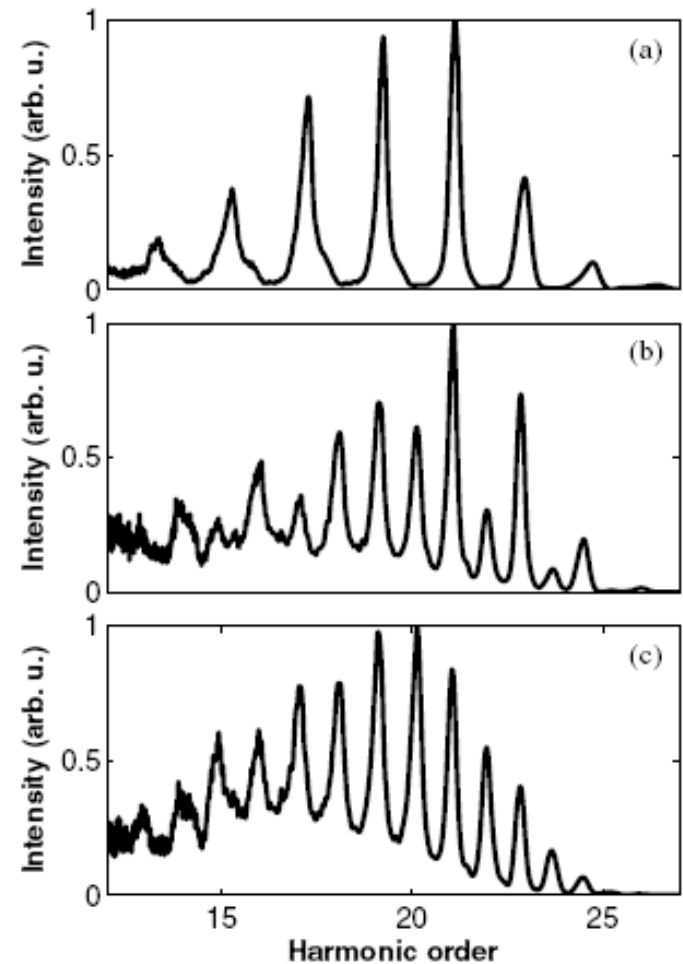
E. Goulielmakis, et al., Science 320, 1614 (2008)

Two-color multi-cycle field

Few-cycle pulse



Two-color pulse train



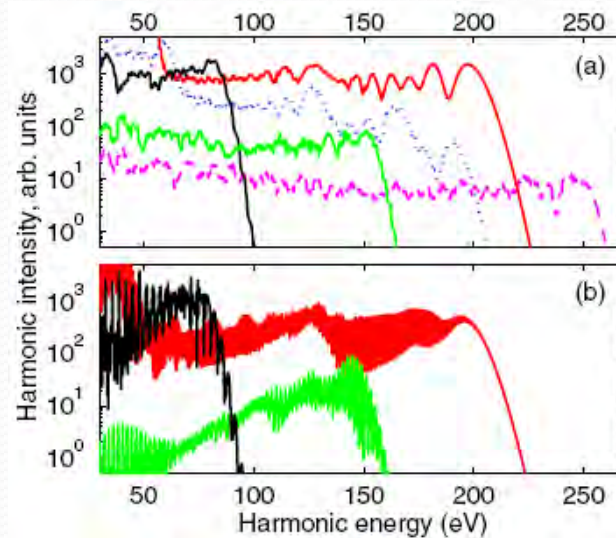
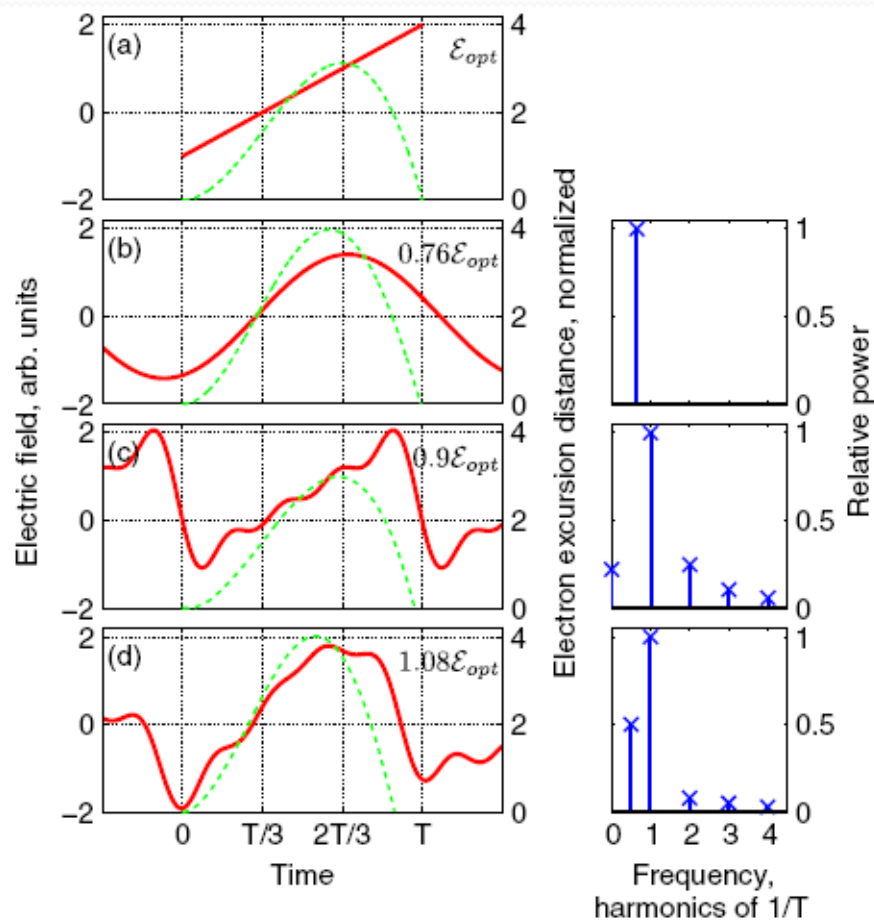
J Mauritsson et al., J. Phys. B: At. Mol. Opt. Phys. 42, 134003 (2009)

Ideal Waveform to Generate the Maximum Possible Electron Recollision Energy for Any Given Oscillation Period

L. E. Chipperfield,* J. S. Robinson, J. W. G. Tisch, and J. P. Marangos

Imperial College London, London SW7 2BW

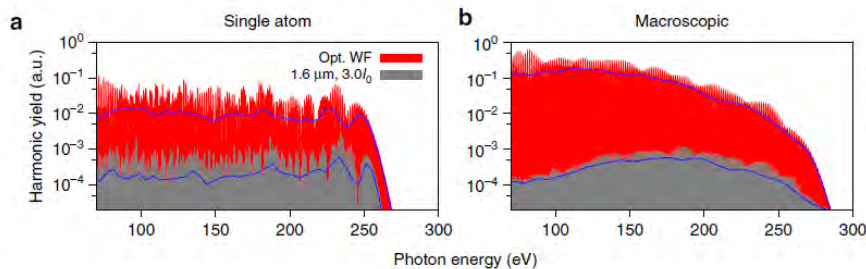
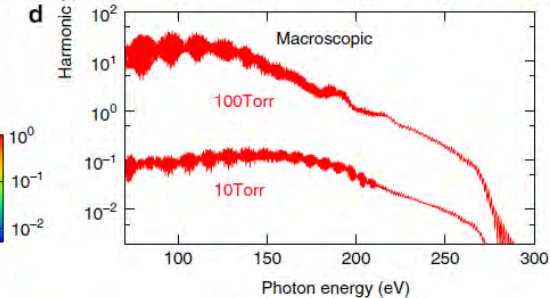
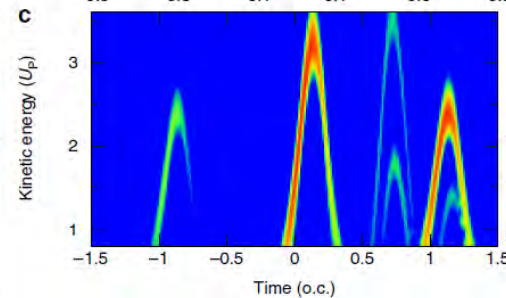
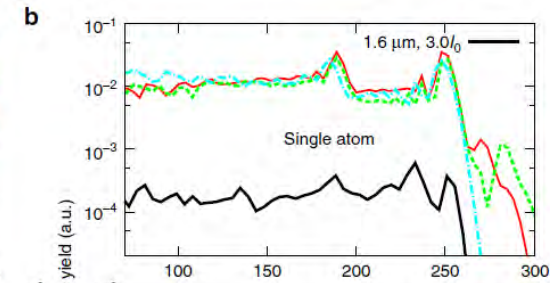
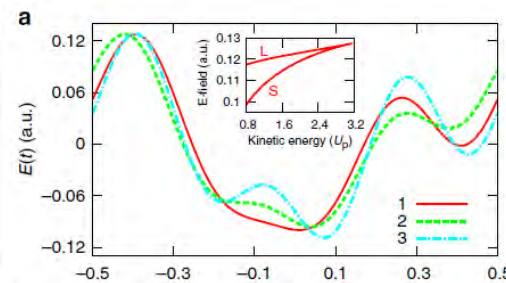
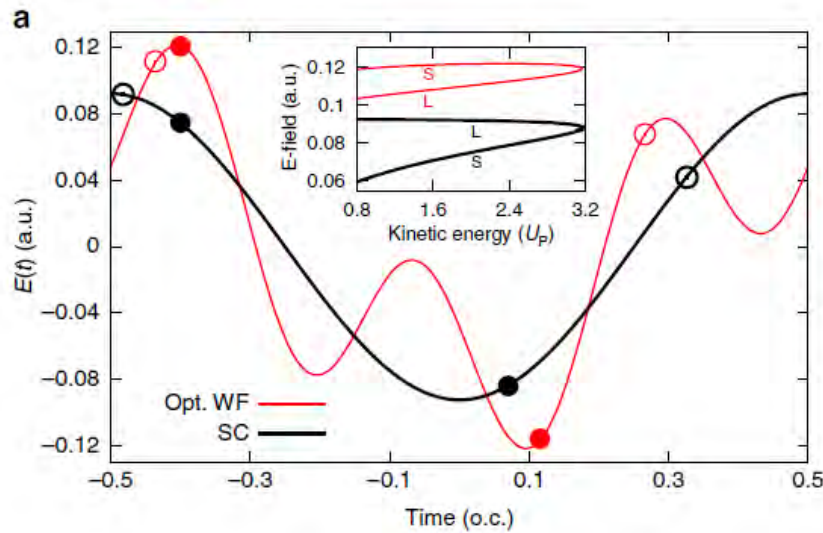
(Received 25 April 2008; published 12 February 2009)



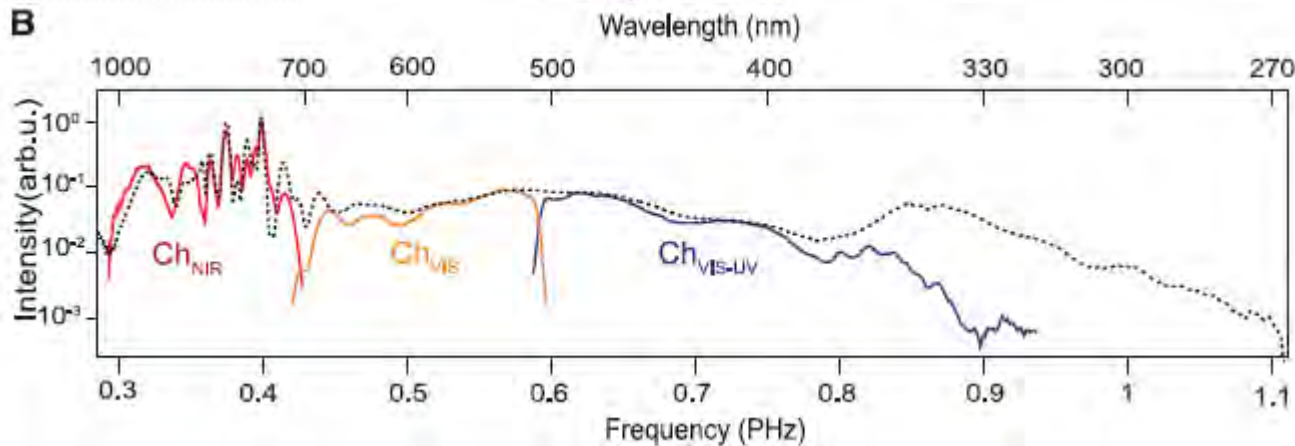
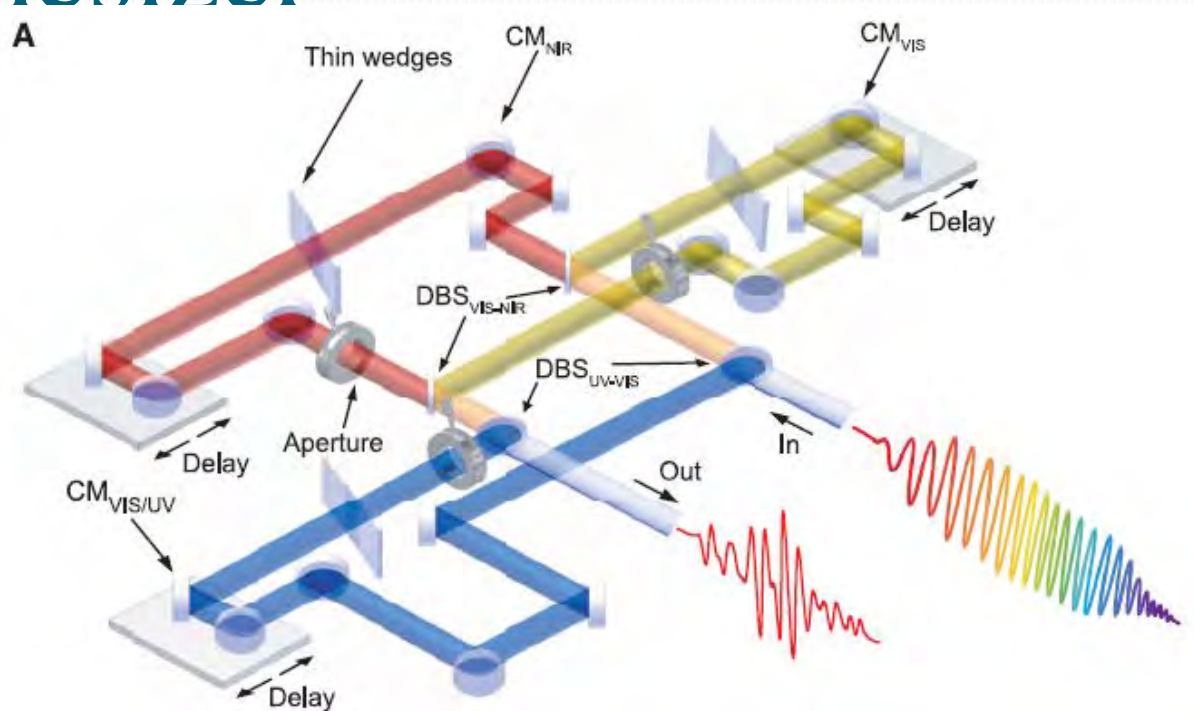
L. E. Chipperfield, et al.,
Phys. Rev. Lett. 102, 063003 (2009)

Waveforms for optimal sub-keV high-order harmonics with synthesized two- or three-color laser fields

Cheng Jin, et al.,
Nat. Commun. 5, 4003 (2014)



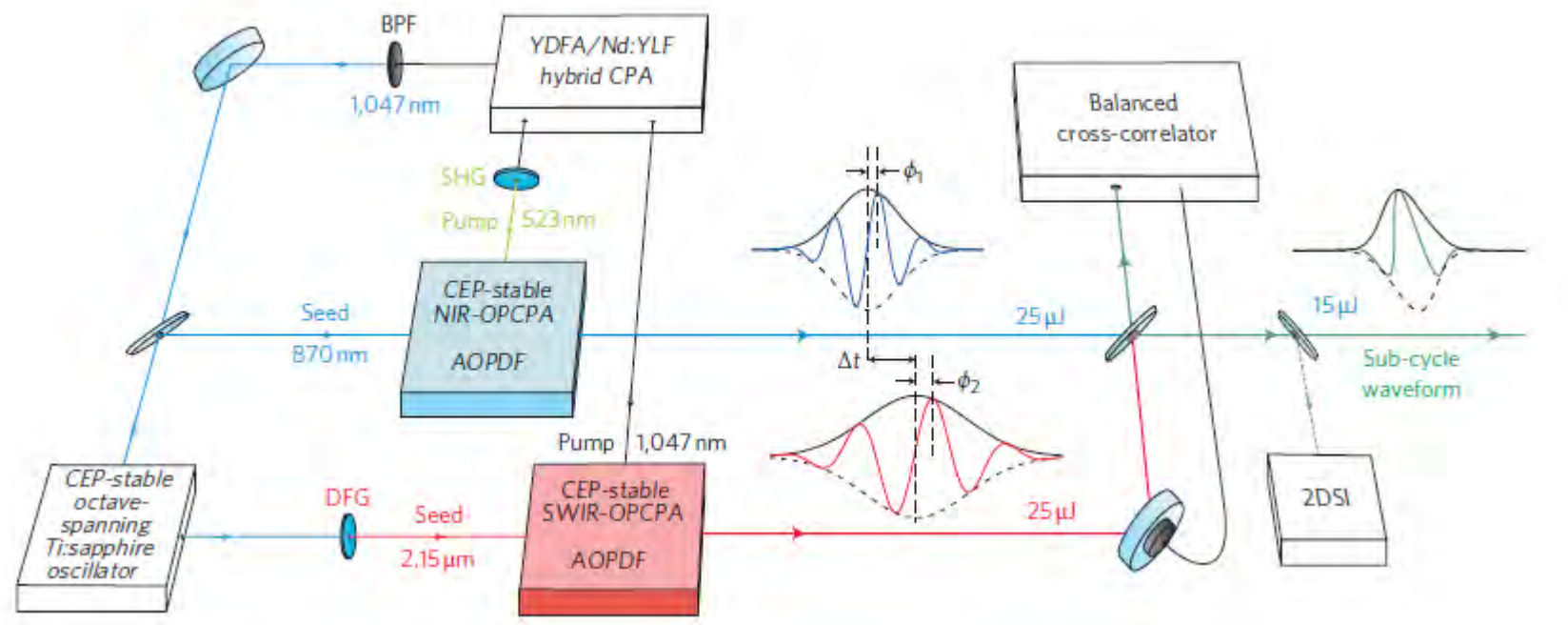
Three-channel Optical Field Synthesizer



A. Wirth et al.
Science 334, 195 (2011)

High-energy pulse synthesis with sub-cycle waveform control for strong-field physics

Shu-Wei Huang¹, Giovanni Cirimi¹, Jeffrey Moses¹, Kyung-Han Hong¹, Siddharth Bhardwaj¹, Jonathan R. Birge¹, Li-Jin Chen¹, Enbang Li², Benjamin J. Eggleton², Giulio Cerullo³ and Franz X. Kärtner^{1,4*}



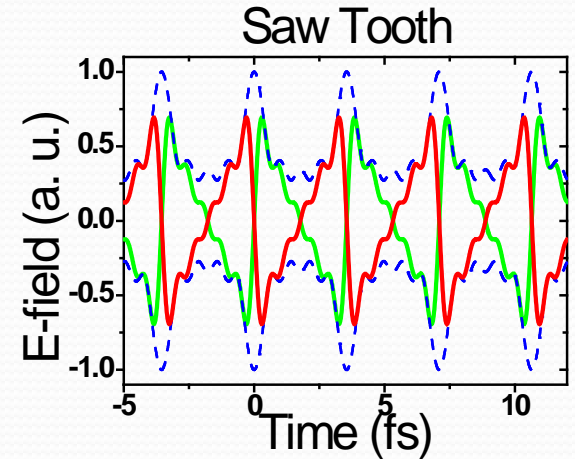
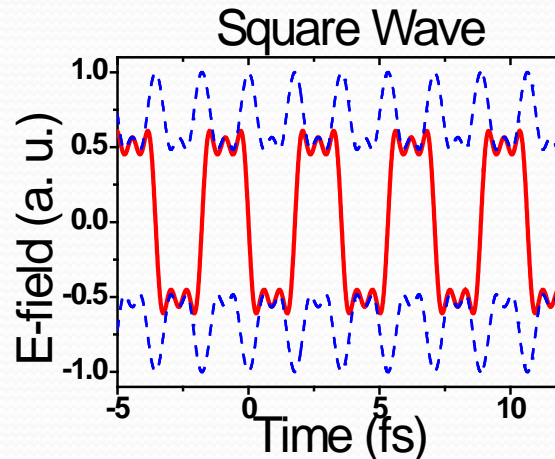
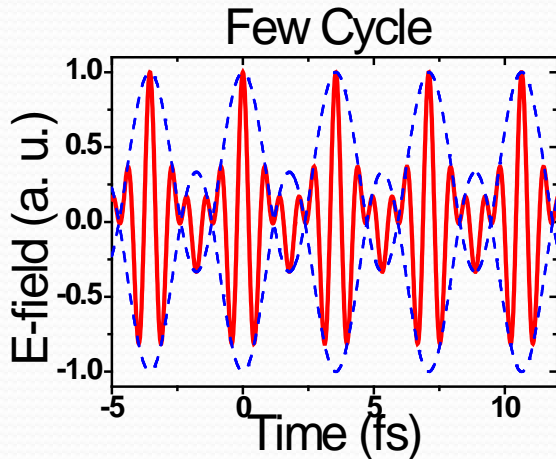
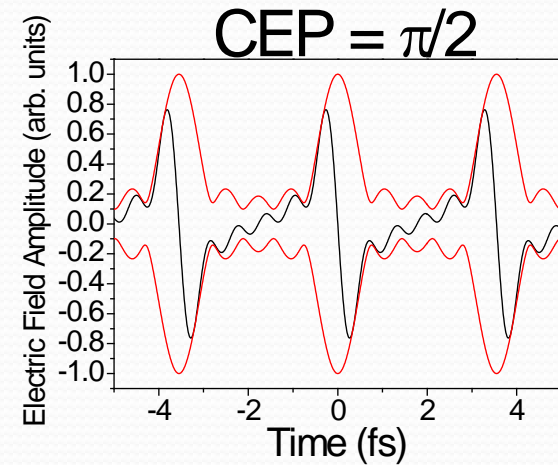
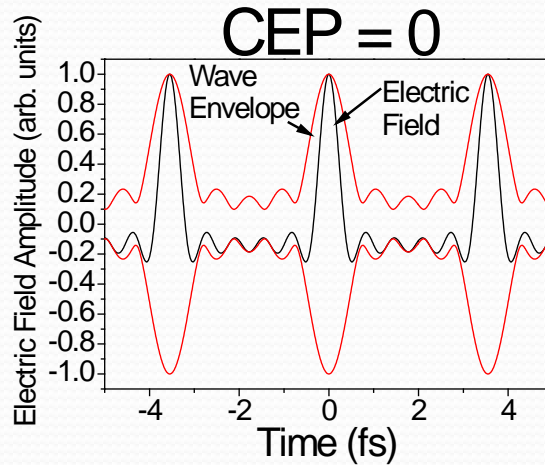
Multi-color laser field

- Broadband source: larger than 2 octaves
coherent and commensurate
- High peak power enough
 10^{13} - 10^{14} W/cm²
- Simple experiment setup
- Light waveform controllable

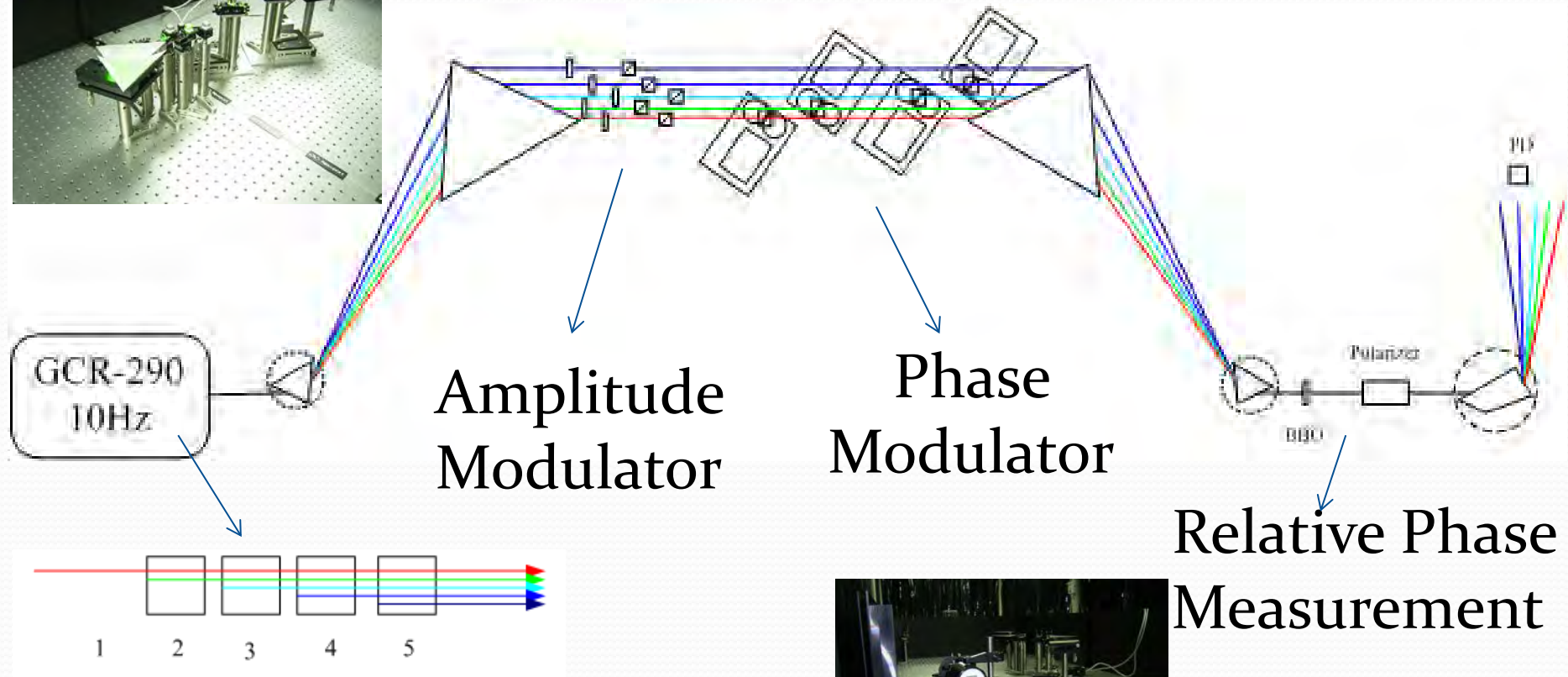
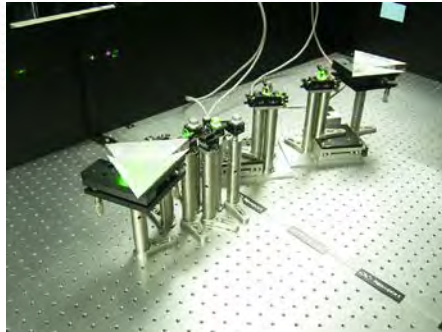
Waveform by Harmonics

1064
532
355
266
213

$\sim 37,600\text{cm}^{-1}$

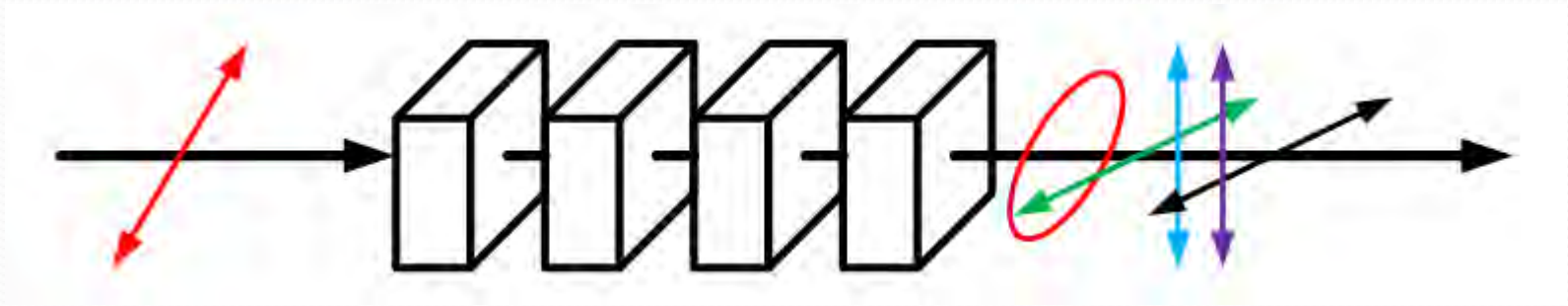





Experimental Setup



Harmonics generation
(collinearly)

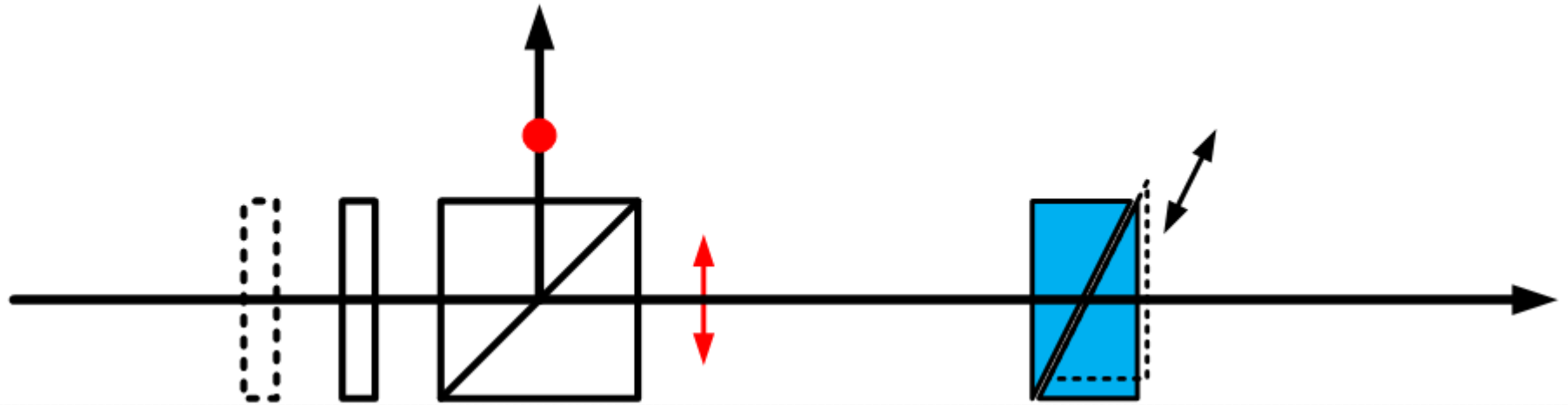
Harmonic Generation



	1064 nm
	532 nm
	355 nm
	266 nm
	213 nm

			1064: 380 mJ
KD*P (II)	$1064+1064 \rightarrow$	532: 178 mJ	
KD*P (I)	$1064+532 \rightarrow$	355: 70 mJ	
BBO (I)	$532+532 \rightarrow$	266: 41 mJ	
BBO (I)	$1064+266 \rightarrow$	213: 22 mJ	

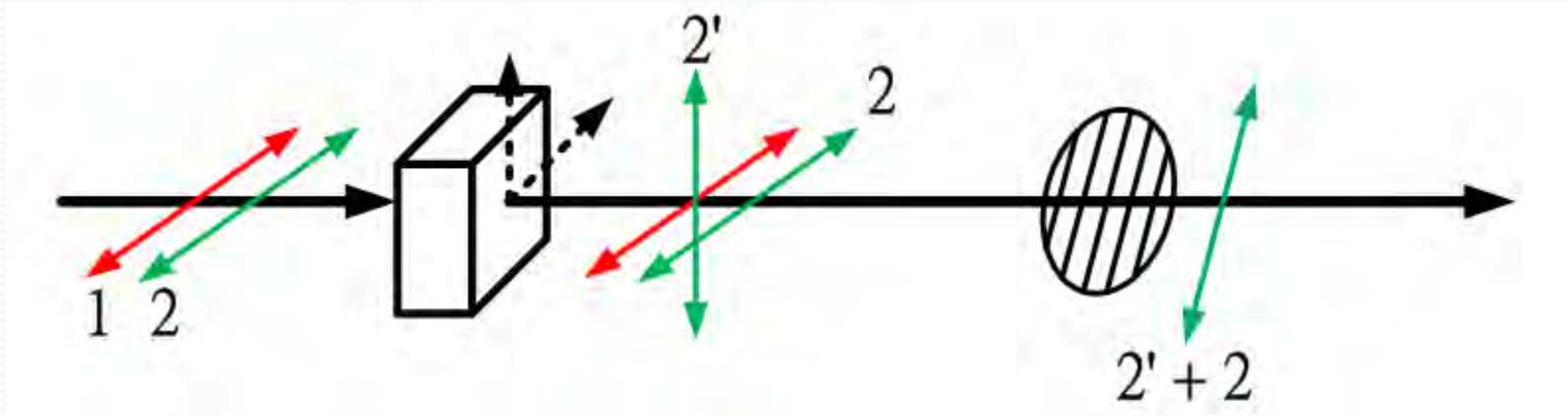
Amplitude and Phase modulator



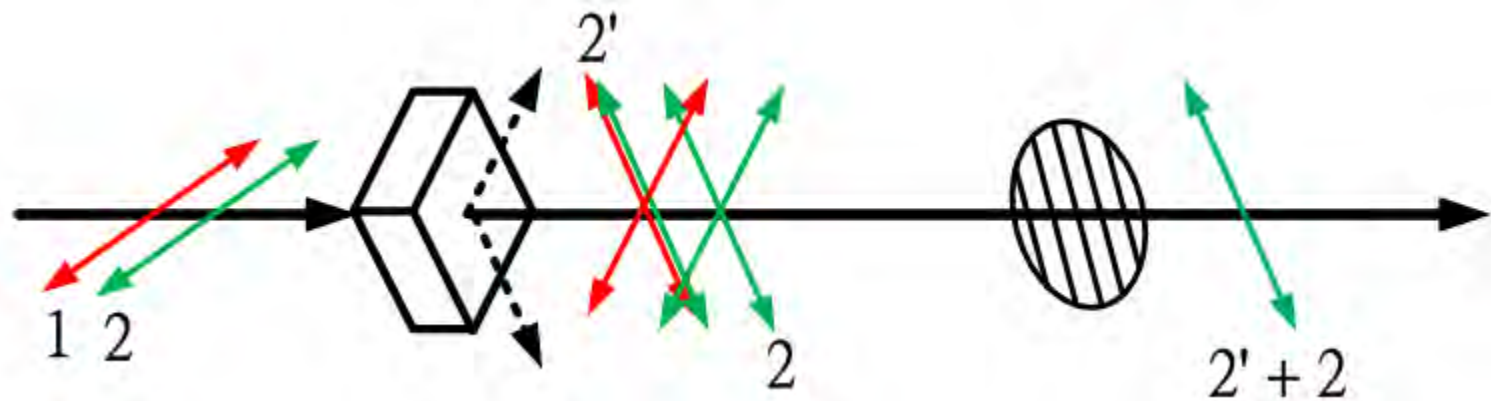
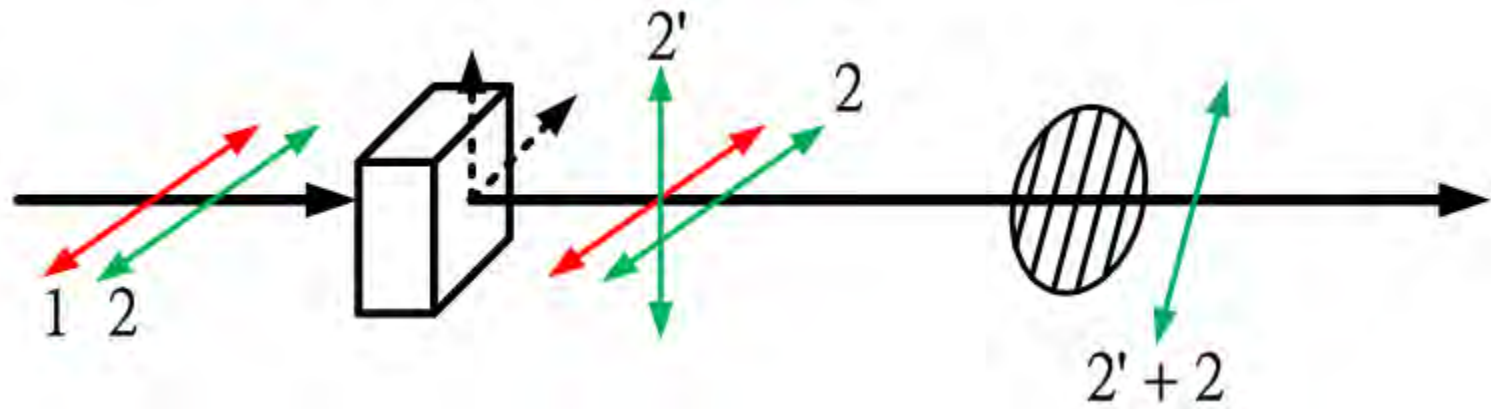
$\lambda/4$ $\lambda/2$ PBC
Amplitude
Modulator
For 1064
532
355
266

Prism Pair
Phase Modulator
For 1064
532
355
266

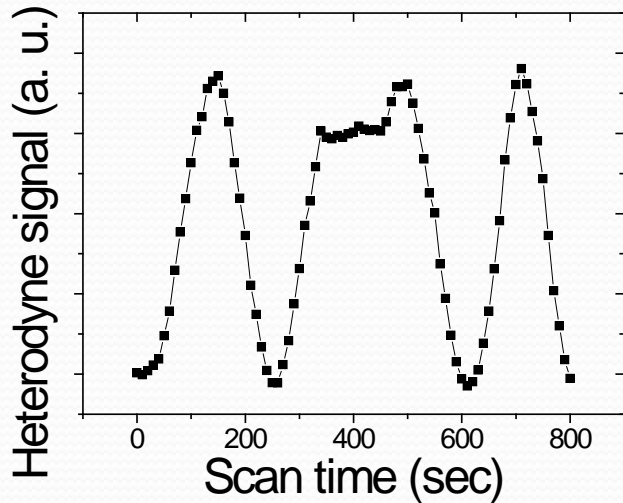
Relative Phase Measurement



Relative Phase Measurement



Relative Phase Measurement



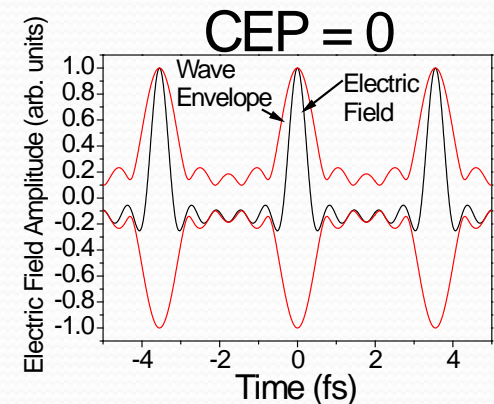
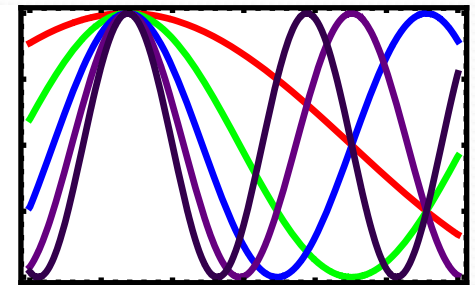
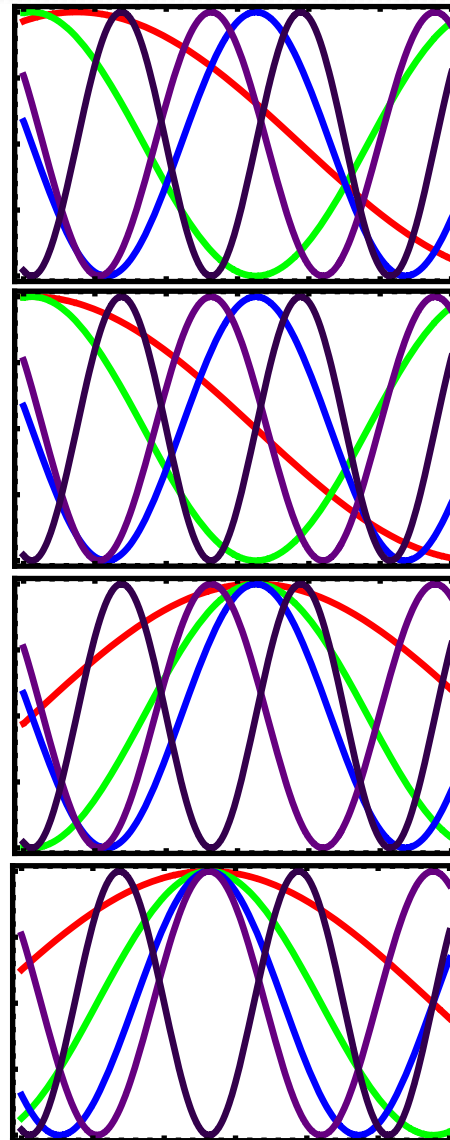
The relative phase between harmonics

$$\text{SHG: } \Phi_{532} = \Phi_{1064} + \pi/2$$

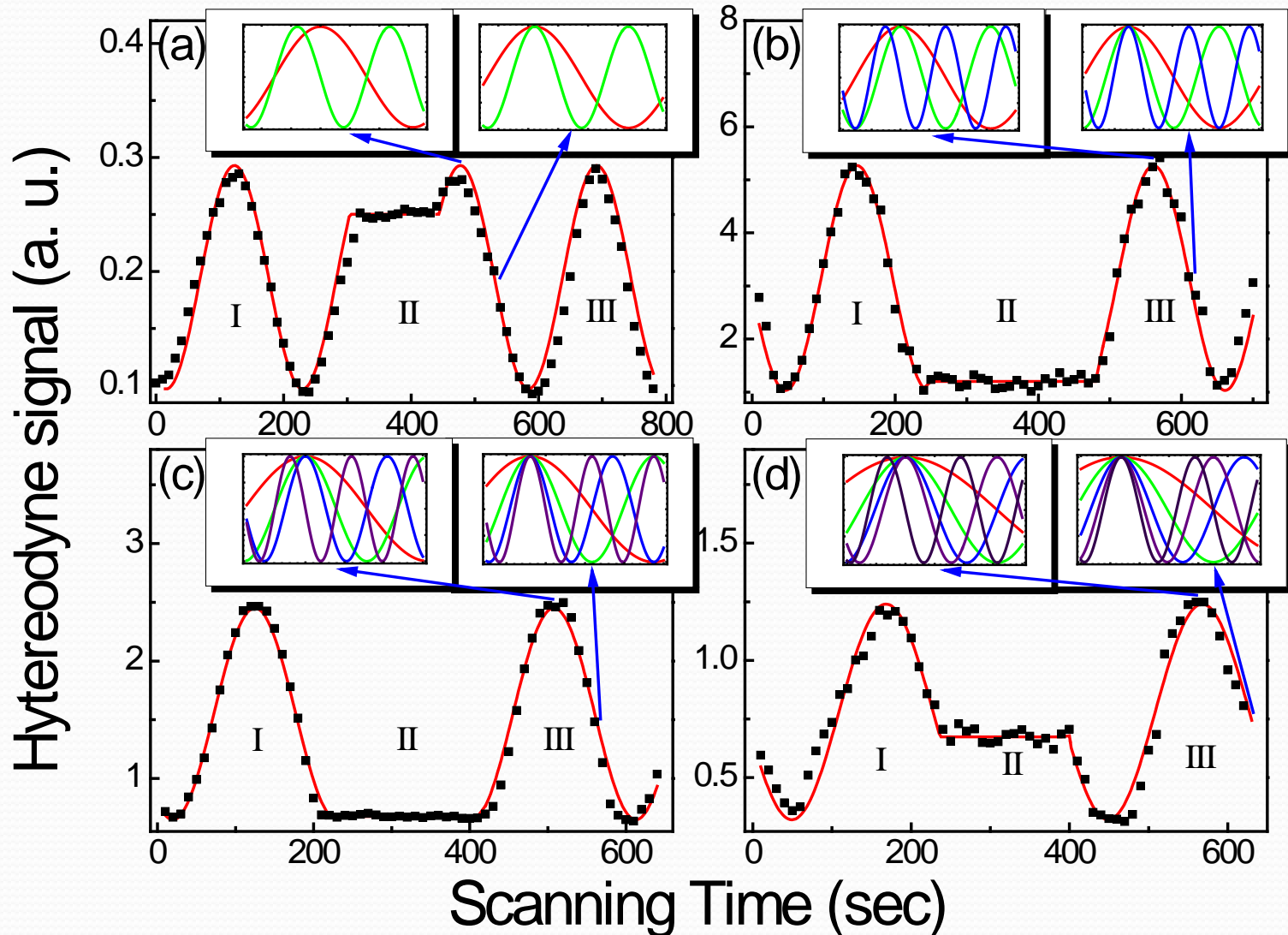
$$\text{SFG: } \Phi_{355} = \Phi_{1064} + \Phi_{532} + \pi/2$$

$$\text{SHG: } \Phi_{266} = \Phi_{532} + \pi/2$$

$$\text{SFG: } \Phi_{213} = \Phi_{1064} + \Phi_{266} + \pi/2$$

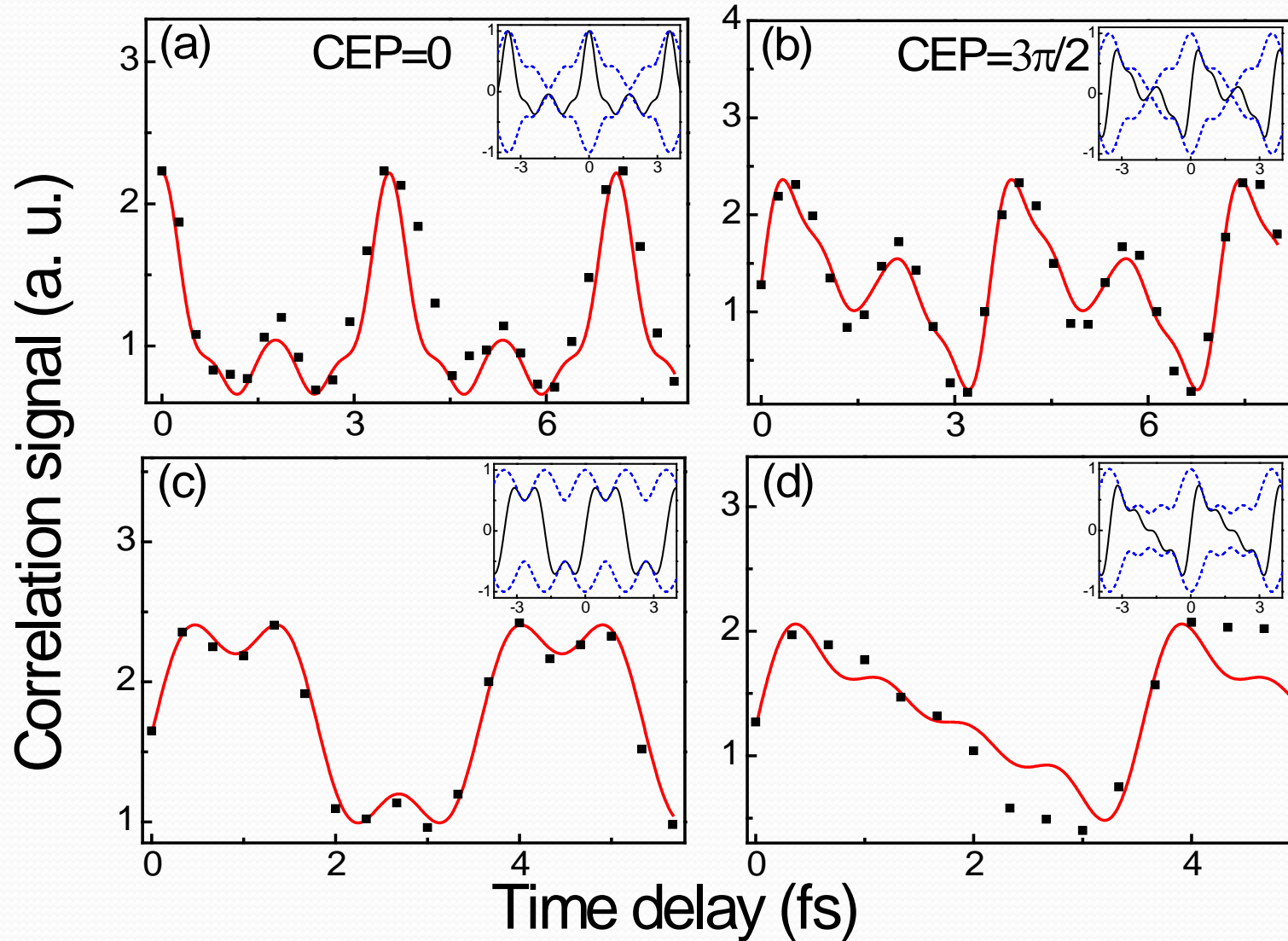


Relative Phase Measurement

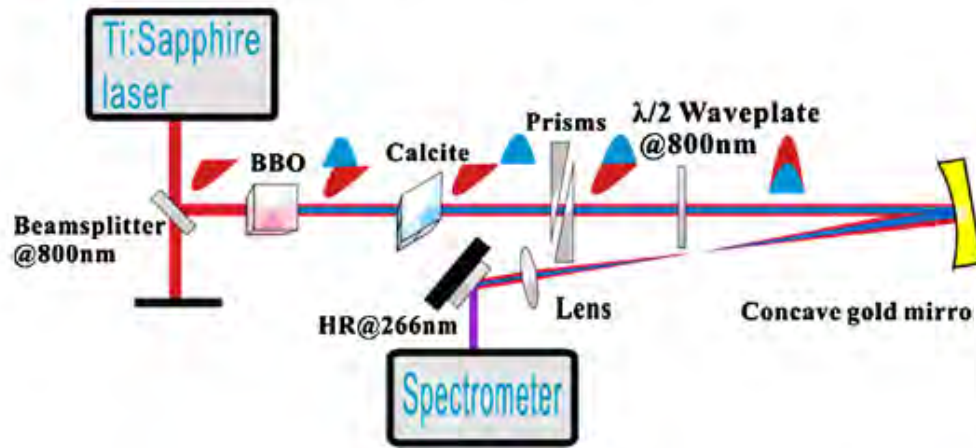


The heterodyne signal: (a) 532 nm, (b) 355 nm, (c) 266 nm, (d) 213 nm

Waveform Synthesized by five harmonics

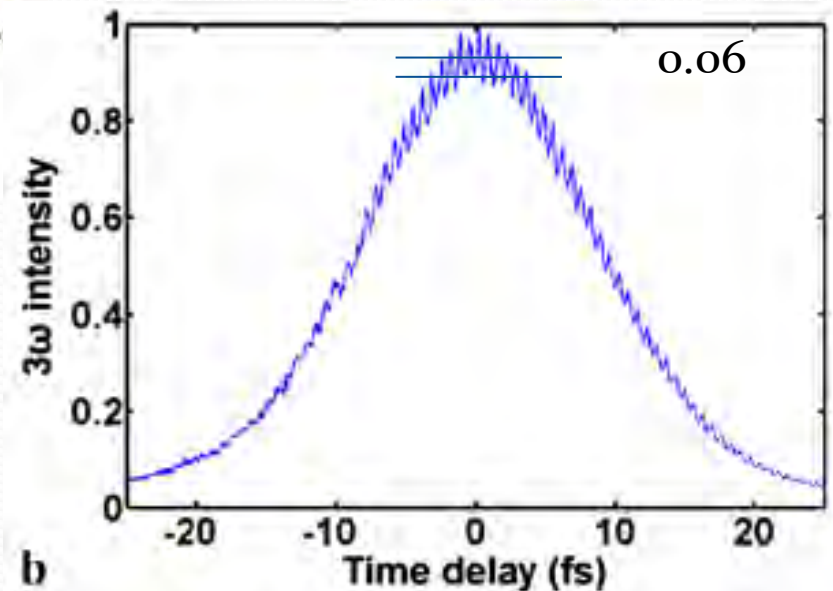


Third-harmonic generation in relative-phase-controlled two-color laser field



H. Xu, et al.,
Appl. Phys. B 104, 909 (2011)

$$S_{3\omega}(\tau) \propto a^2 A_{\omega}^6 \sqrt{\frac{\pi}{6}} + 9b^2 A_{2\omega}^4 A_{\omega}^2 \sqrt{\frac{\pi}{6}} \cdot \exp\left(-\frac{4\tau^2}{3\tau_{\omega}^2}\right) + 6ab A_{\omega}^4 A_{2\omega}^2 \sqrt{\frac{\pi}{6}} \cos 4\omega\tau \cdot \exp\left(-\frac{4\tau^2}{3\tau_{\omega}^2}\right)$$



Third-Harmonic Generation by two-color

$$\tilde{E}_i(z, t) = E_i e^{-i\omega_i t} + c.c. \quad E_i = A_i e^{i(k_i z + \phi_i)}$$

$$\tilde{P}^{(3)}(z, t) = \varepsilon_0 \chi^{(3)} \tilde{E}^3(z, t)$$

$$\tilde{E}(z, t) = \tilde{E}_1(z, t) + \tilde{E}_2(z, t) \quad \phi_1 = 0, \Delta\phi = \phi_2 - \phi_1 = \phi_2$$

$$\tilde{P}_3^{(3)}(z, t) = \varepsilon_0 (a \tilde{E}_1^3(z, t) + 3b \tilde{E}_2^2(z, t) \tilde{E}_1^*(z, t) + c.c.)$$

$$a = \chi^{(3)}(\omega_3; \omega_1, \omega_1, \omega_1) \quad b = \chi^{(3)}(\omega_3; \omega_2, \omega_2, -\omega_1)$$

Third-Harmonic Generation by two-color

$$\begin{aligned}\tilde{E}_3(t) &= \varepsilon_0 (a e^{-i3\omega_1 t} \int_{-L/2}^{L/2} A_1^3 e^{-i\Delta k_{13} z} dz + 3b e^{-i3\omega_1 t} \int_{-L/2}^{L/2} A_2^2 A_1^* e^{-i(\Delta k_{123} z + 2\Delta\phi)} dz) + c.c. \\ &= \varepsilon_0 (a A_1^3 \sin c(\frac{\Delta k_{13} L}{2}) \sin(3\omega_1 t) + 3b A_2^2 A_1^* \sin c(\frac{\Delta k_{123} L}{2}) \sin(3\omega_1 t + 2\Delta\phi))\end{aligned}$$

$$I_3(\Delta\phi) \propto \int \tilde{E}_3^2(t) dt$$

$$\begin{aligned}&= a^2 A_1^6 \sin^2 c(\frac{\Delta k_{13} L}{2}) + 9b^2 A_2^4 A_1^2 \sin^2 c(\frac{\Delta k_{123} L}{2}) \\ &\quad + 6ab A_1^4 A_2^2 \sin c(\frac{\Delta k_{13} L}{2}) \sin c(\frac{\Delta k_{123} L}{2}) \cos(2\Delta\phi)\end{aligned}$$

Relative phase measurement for multi-color waveform synthesis

$$E(t) = \sum_n A_n \cos(n\omega t + \phi_n)$$

$$E(t) = \sum_n A_n \cos(n\omega(t - \phi_1/\omega) + \phi_n) = \sum_n A_n \cos(n\omega t + \phi_n - n\phi_1)$$

$$\phi'_1 \rightarrow 0, \phi'_n \rightarrow \phi_n - n\phi_1.$$

$$\phi_n - \phi_{n-1} = \phi_{n-1} - \phi_{n-2} = \dots = \phi_2 - \phi_1$$

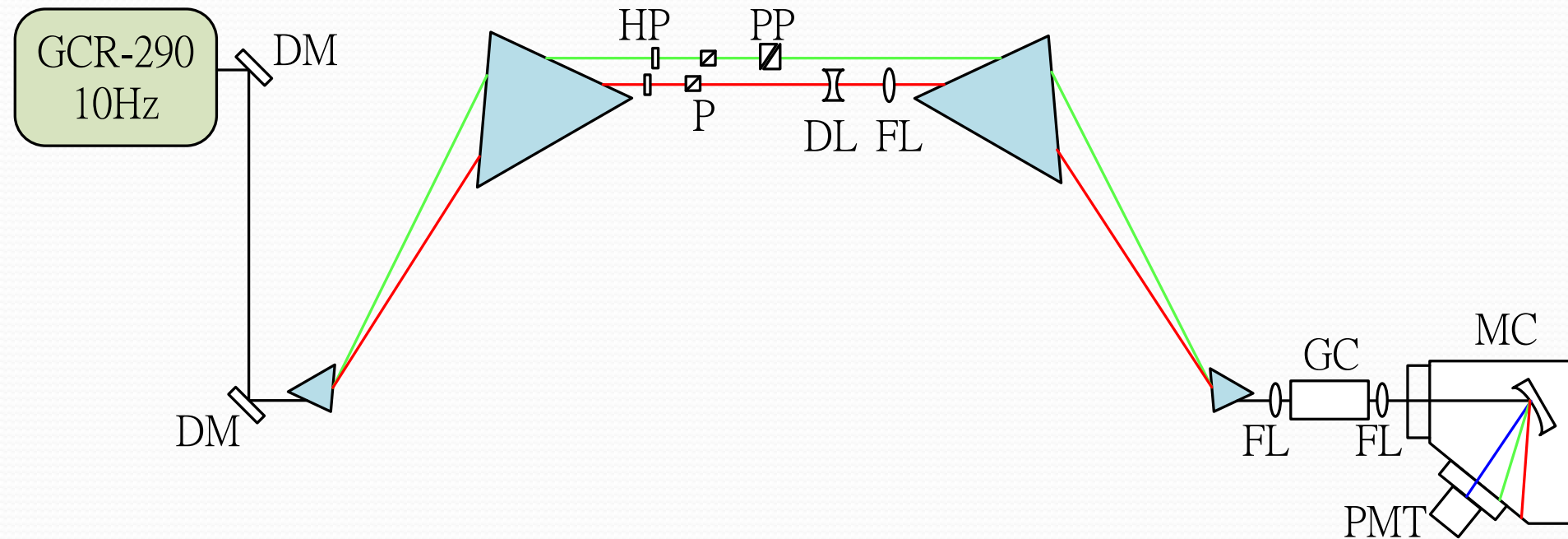
Interference of FWMs

Phys. Rev. Lett. 100, 163906 (2008)

$$\Delta\phi = \phi_2 - \phi_1$$

Interference of FWM and THG

Experiment Setup for two-color THG



DM: dichroic mirror; HP: half-wave plate; P: polarizer;
PP: prism pair; DL: defocus lens; FL: focus lens; GC: gas
cell;
MC: monochromator; PMT: photomultiplier tube

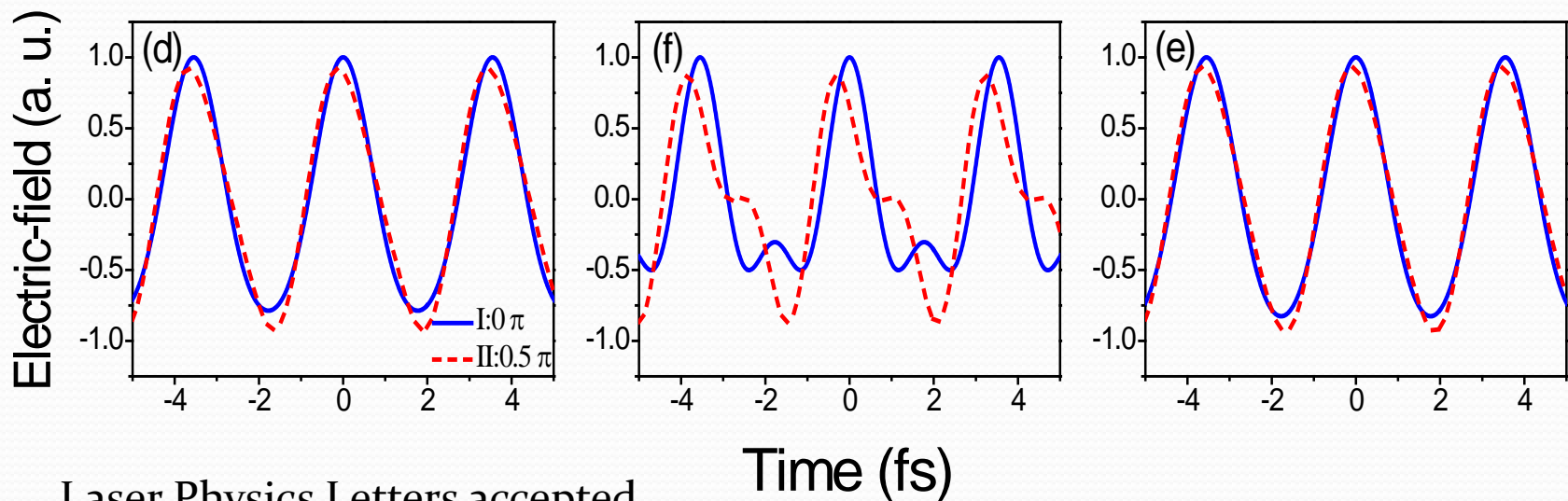
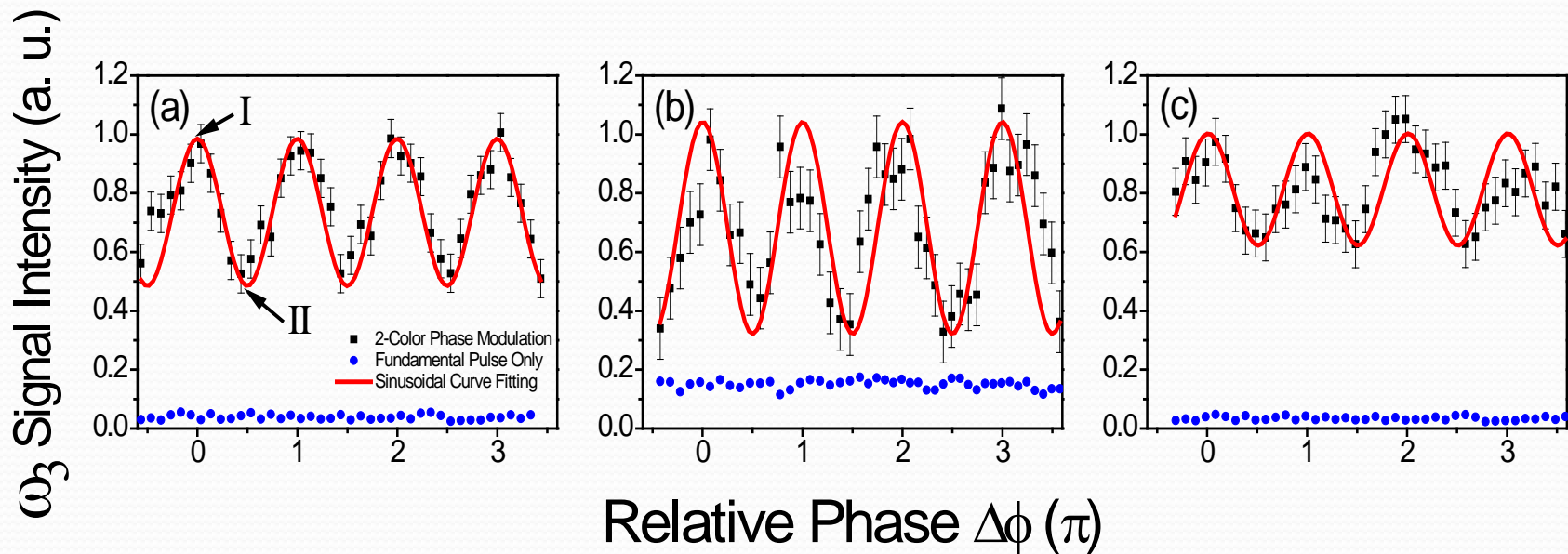
Third-Harmonic Generation by 2-color

$$I_3(\Delta\phi) \propto A_1^6 \sin^2 c\left(\frac{\Delta k_{13}L}{2}\right) + 9A_2^4 A_1^2 \sin^2 c\left(\frac{\Delta k_{123}L}{2}\right) + 6A_1^4 A_2^2 \sin c\left(\frac{\Delta k_{13}L}{2}\right) \sin c\left(\frac{\Delta k_{123}L}{2}\right) \cos(2\Delta\phi) \quad (6)$$

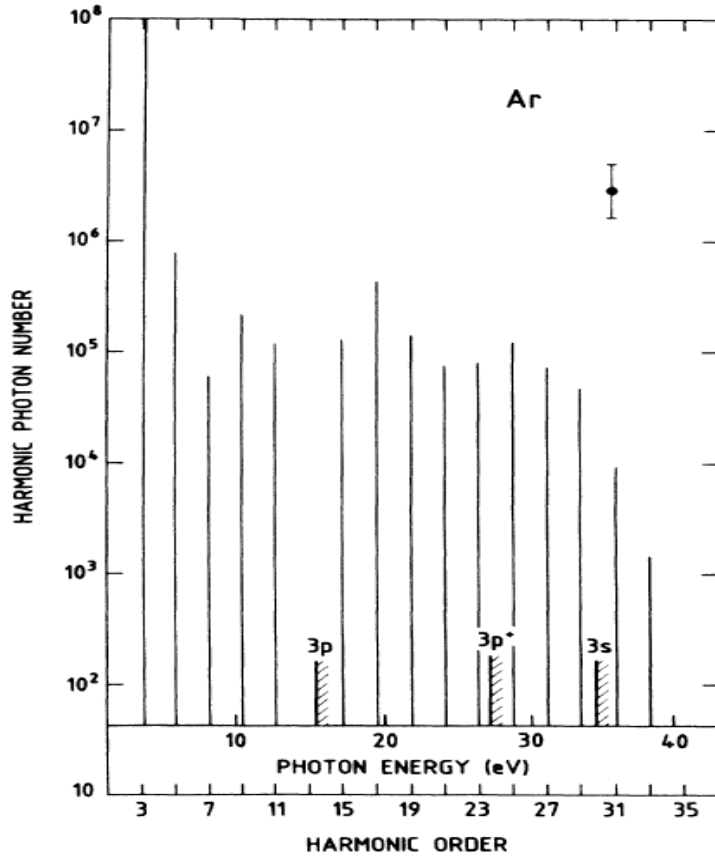
$$= \alpha^2 A_1^6 + 9\beta^2 A_2^4 A_1^2 + 6\alpha\beta A_1^4 A_2^2 \cos(2\Delta\phi)$$

where $\alpha = \sin c\left(\frac{\Delta k_{13}L}{2}\right)$, $\beta = \sin c\left(\frac{\Delta k_{123}L}{2}\right)$

Third-Harmonic Signal



HHG by 1064 nm

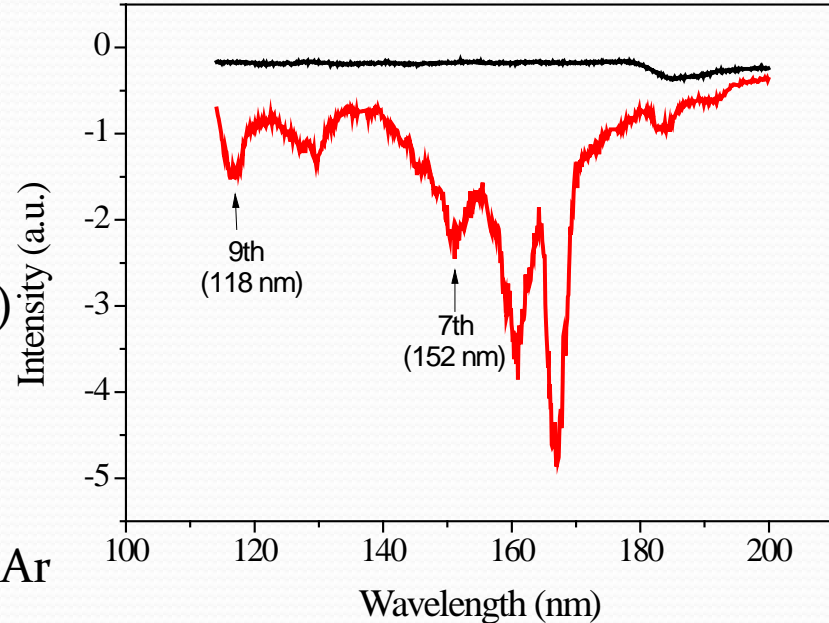
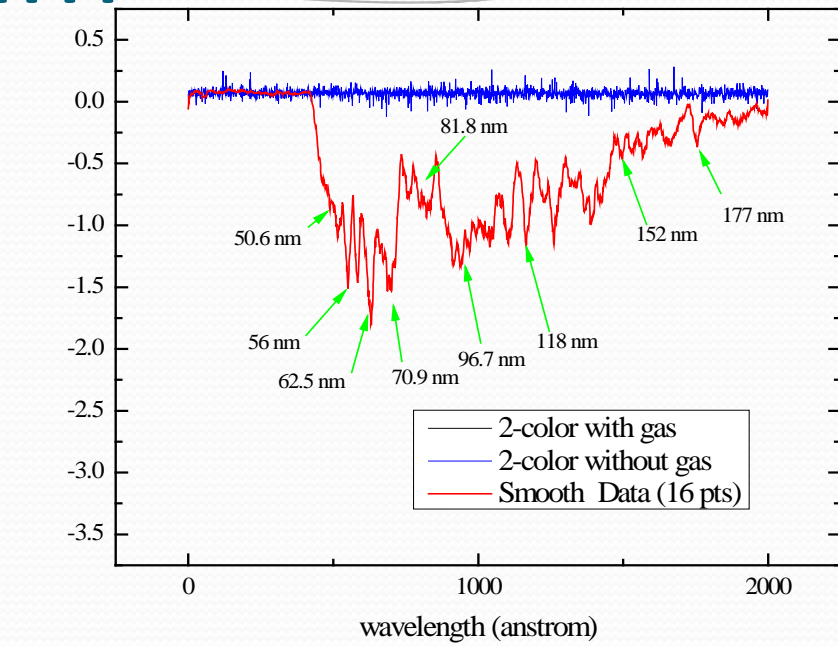


X. F. Li, et al. Phys. Rev. A 39, 5751 (1989)

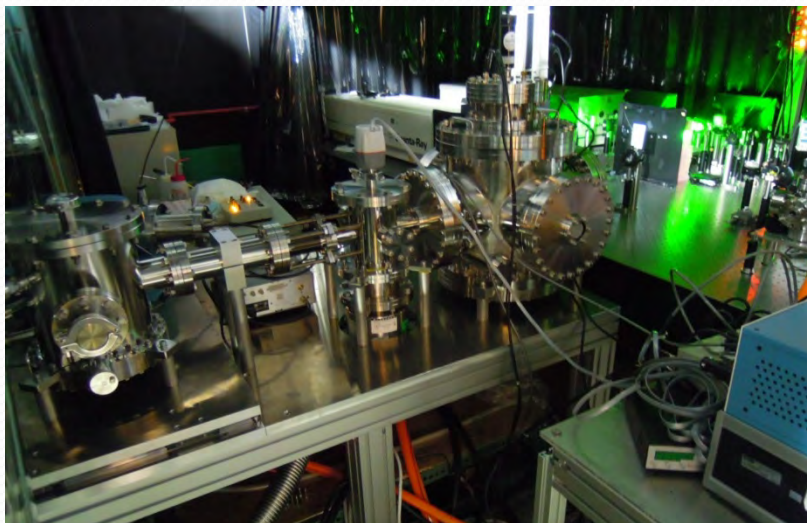
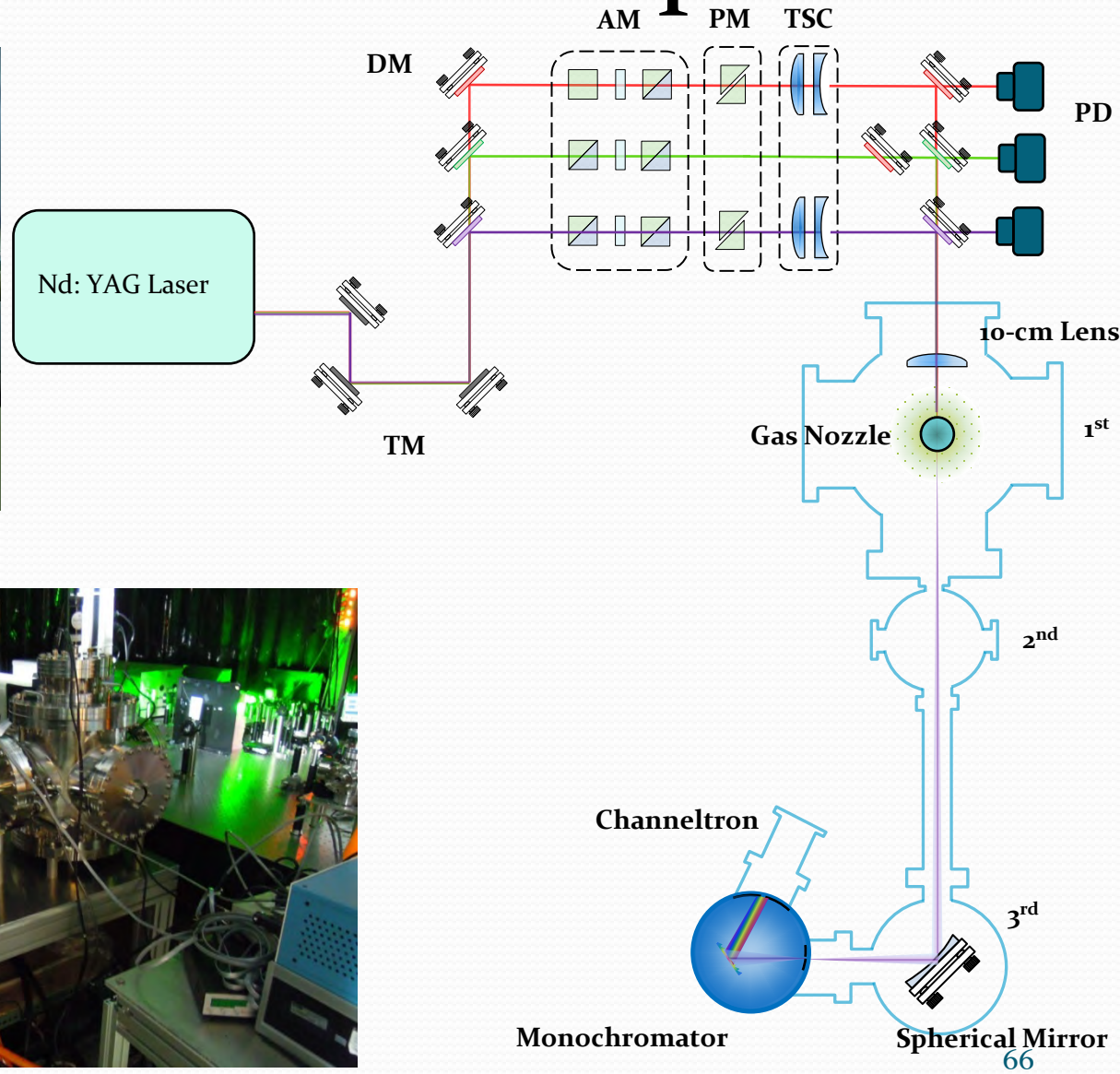
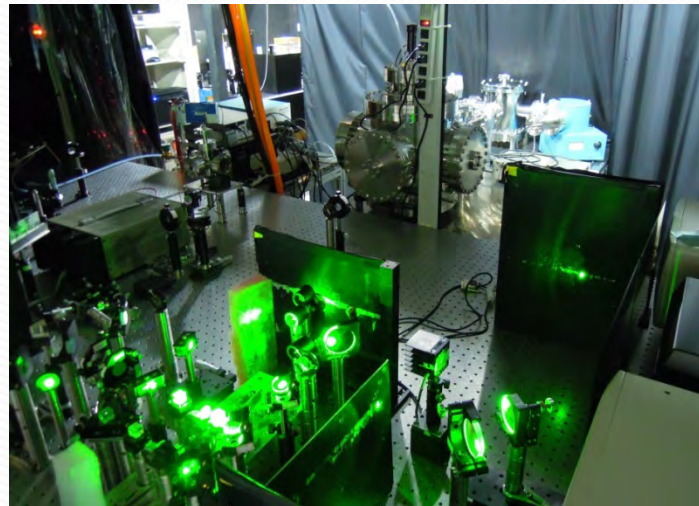
3×10^{13} W/cm², 15 Torr Ar

3.66×10^{13} W/cm², 12 Torr Ar

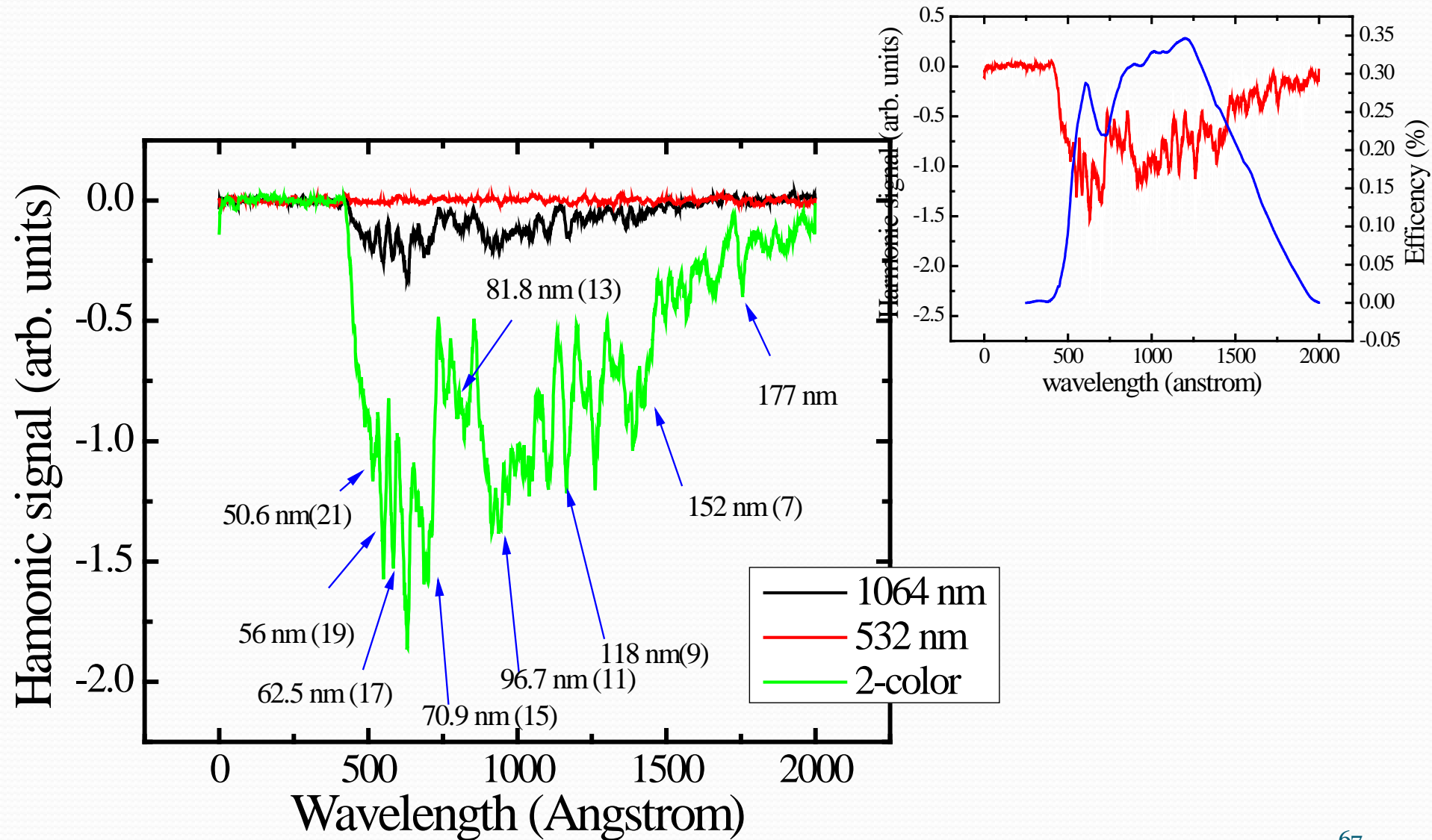
Spectrum pumped by 1064-nm and 532-nm pulses



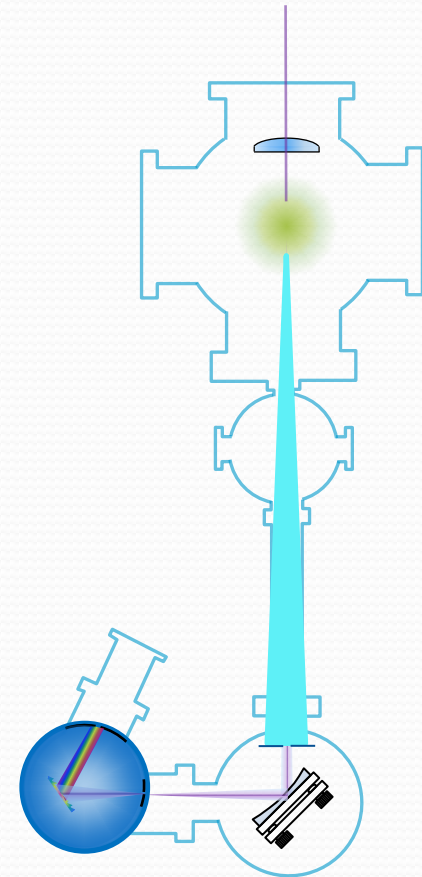
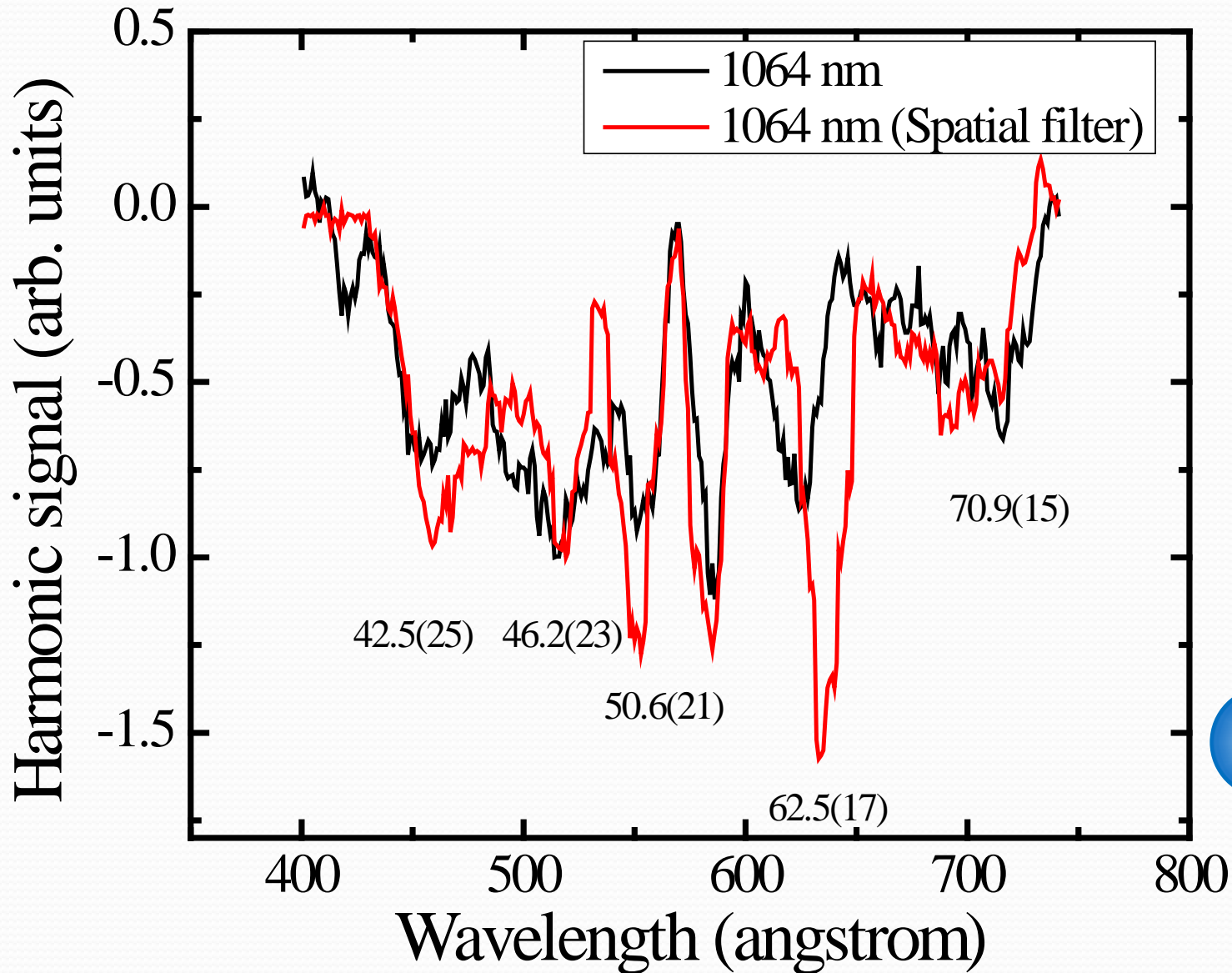
Experiment setup



Two-color Harmonic Generation



Harmonics and Fluorescence

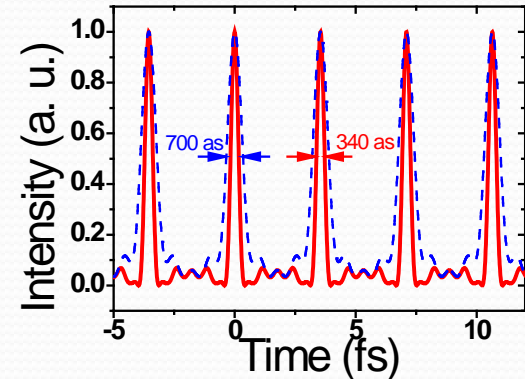


Summary

When CEP=0

=>The pulse width ~ 340 as

Focusing to a $\Phi 20\mu\text{m}$ spot, the intensity will reach 10^{14} W/cm².



We have investigated the generation of TH on the influence of relative phases and amplitudes of the two-color fields. A modulation depth as high as 0.35 has been observed.

We present a promising way for *in situ* determination of the relative phase in multi-color laser field.

Waveform synthesis in the VUV spectral range by higher harmonics generation using waveform-controlled multi-colour quasi-single-frequency laser fields is feasible.

The plasma induced by the two-color laser field is shown to have a significant effect on generation of the harmonic signal.

Photoelectron and/or ion measurement

Possible consequences of recollision (i-iv)

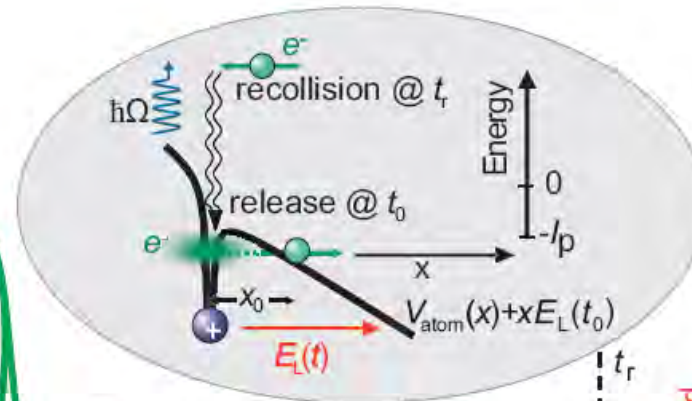
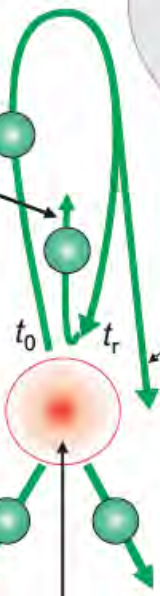
(i) energetic electron emission by elastic backscattering of the electron

(ii) energetic photon emission upon the electron recombining into its ground state

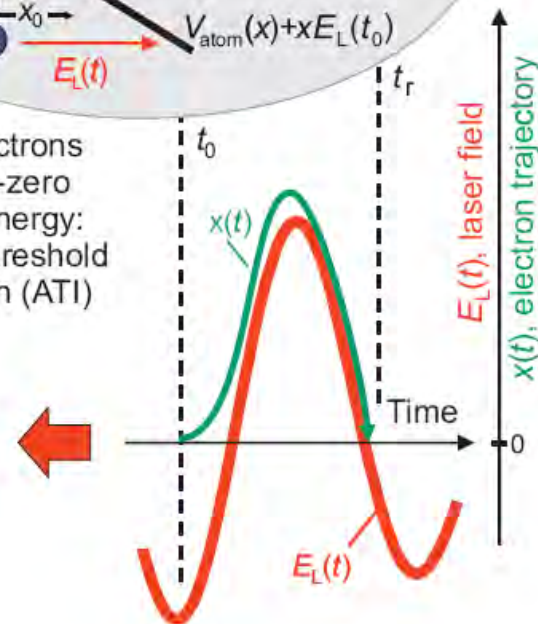
(iii) detachment of another electron: non-sequential double ionization (NSDI)

(iv) excitation of bound electrons upon inelastic collision

$\hbar\Omega$



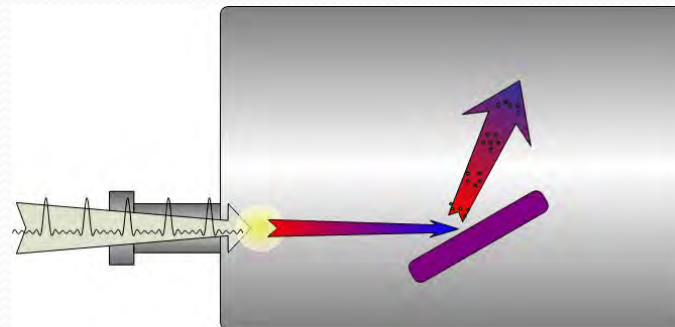
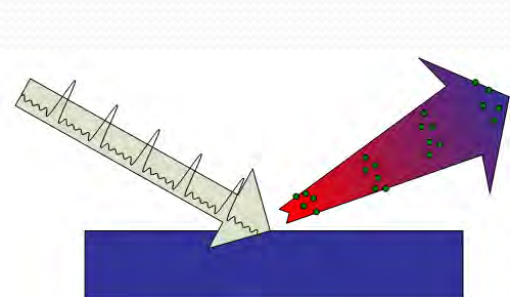
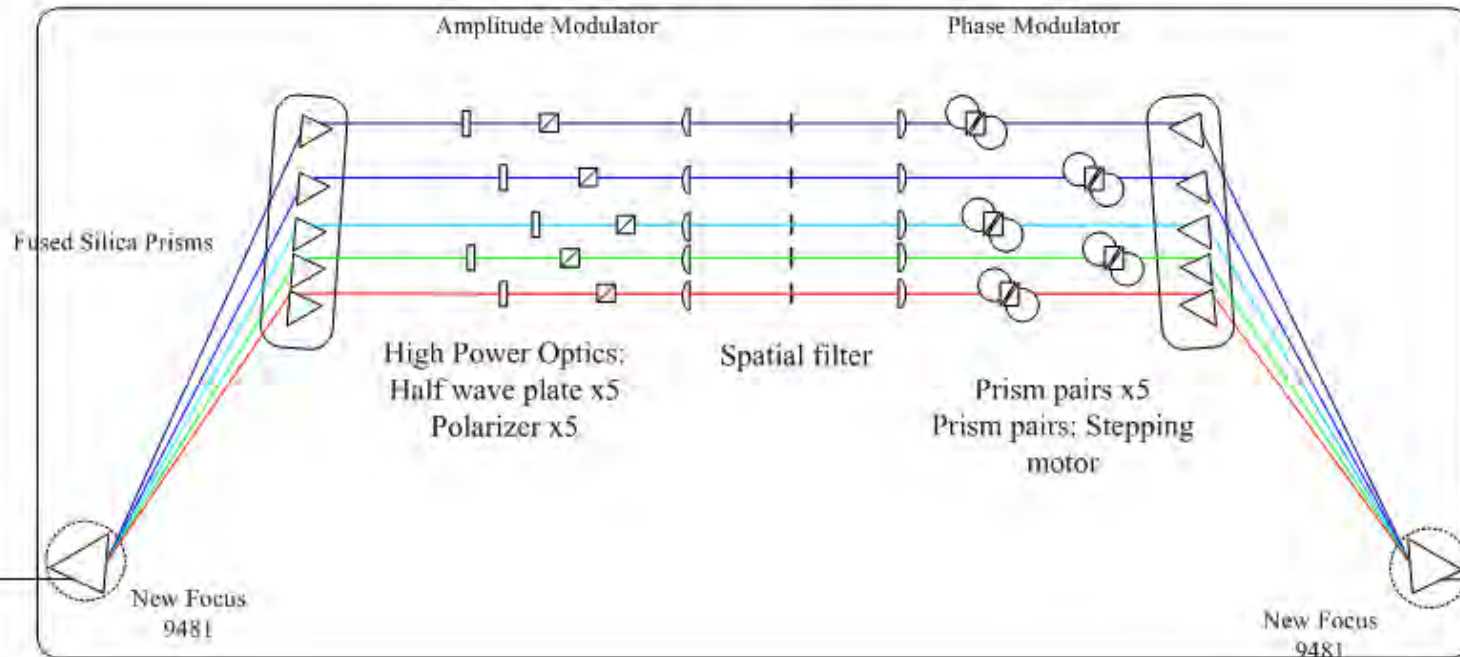
Free electrons with non-zero kinetic energy: above-threshold ionization (ATI)



F. Krausz & M. Ivanov, Rev. Mod. Phys. 81 163 (2009)

Quantum Stroboscope for Electron Motion

AM & PM



Photoelectron emission &
Nonlinear optics

Thanks for your attention



ATTO 2015 - St. Sauveur, Canada