

Multichannel K-matrix method

$$|\gamma \mu\rangle = \begin{cases} |\gamma_0 \mu\rangle & \text{(open channel)} \\ |\gamma_c \mu\rangle & \text{(closed channel)} \end{cases}$$

where γ : channel index, μ : configuration series

$$\langle \gamma' \mu' | \gamma \mu \rangle = \delta_{\gamma' \gamma} \delta_{\mu' \mu} \quad \text{(orthogonality)}$$

$$\langle \gamma'_0 \mu' | H | \gamma_0 \mu \rangle = \delta_{\gamma'_0 \gamma_0} \delta_{\mu' \mu} \epsilon_{\gamma_0 \mu} \quad \text{(Prediagonalization)}$$

Define

$$|\Phi_{\gamma_0 E}\rangle = |\gamma_0 E\rangle + \sum_{\gamma} \mathcal{P} \int d\varepsilon \frac{\langle \gamma \varepsilon | K(E) | \gamma_0 E \rangle}{E - \varepsilon} |\gamma \varepsilon\rangle$$

where $\Sigma' = (\Sigma + \int)$, \mathcal{P} denotes the principal integral

$$\Rightarrow \langle \gamma' \varepsilon | H - E | \Phi_{\gamma_0 E} \rangle = 0$$

$$\begin{aligned} \Rightarrow \langle \gamma' \varepsilon' | K(E) | \gamma_0 E \rangle &= \langle \gamma' \varepsilon' | V(E) | \gamma_0 E \rangle \\ &+ \sum_{\gamma} \mathcal{P} \int d\varepsilon \langle \gamma' \varepsilon' | V(E) | \gamma \varepsilon \rangle \\ &\times \frac{\langle \gamma \varepsilon | K(E) | \gamma_0 E \rangle}{E - \varepsilon} \end{aligned}$$

The energy-dependent \mathbf{V} matrix is given by

$$\langle \gamma' \varepsilon' | V(E) | \gamma \varepsilon \rangle = \langle \gamma' \varepsilon' | H - E | \gamma \varepsilon \rangle - (\varepsilon - E) \delta_{\gamma' \gamma} \delta(\varepsilon - \varepsilon')$$

Orthogonality:

$$\begin{aligned} \langle \Phi_{\gamma_0' E'} | \Phi_{\gamma_0 E} \rangle &= \delta(E' - E) \\ &\times \left[\delta_{\gamma_0' \gamma_0} + \pi^2 \sum_{\gamma_0''} \langle \gamma_0'' E | K(E) | \gamma_0' E \rangle \right. \\ &\left. \times \langle \gamma_0'' E | K(E) | \gamma_0 E \rangle \right] \end{aligned}$$

Eigenchannels

$$|\Gamma E\rangle = \sum_{\gamma_0} |\Phi_{\gamma_0 E}\rangle U_{\gamma_0 \Gamma}(E) \cos \eta_{\Gamma}(E) \quad (\text{normalized})$$

where the transformation matrix $U_{\gamma_0 \Gamma}(E)$ and the eigenphase shift η_{Γ} are obtained by diagonalizing the *on the energy shell* K matrix, i.e.

$$\sum_{\gamma_0'} \langle \gamma_0 E | K(E) | \gamma_0' E \rangle U_{\gamma_0' \Gamma}(E) = -\pi^{-1} \tan \eta_{\Gamma} U_{\gamma_0 \Gamma}(E)$$

Total eigenphase shift:

$$\begin{aligned} \eta_{\text{tot}} &= \sum_{\Gamma} \eta_{\Gamma} \\ &= \eta_0 + \tan^{-1} [(\Gamma_w / 2) / (E_{\text{res}} - E)], \quad \eta_0: \text{background} \end{aligned}$$

Channel wave function

$$|\gamma_0 E^-\rangle = \sum_{\Gamma} |\Gamma E\rangle C_{\Gamma \gamma_0}, \quad (\text{normalized})$$

where $C_{\Gamma \gamma} = e^{-i \eta_{\Gamma}} \tilde{U}_{\Gamma \gamma} e^{-i \delta_{\gamma}}$, and δ_{γ} is the scattering phase shift of the outgoing electron in the open channel γ , i.e.,

$$\xi_{\varepsilon l}(r) \xrightarrow{\text{at large } r} [2/(\pi k)]^{1/2} \sin[kr - (\pi l/2) + (q/k) \ln(2kr) + \sigma_l^C + \delta_{\gamma}],$$

where q is the effective nuclear charge, $k = \varepsilon^{1/2}$ is the momentum, and σ_l^C is the Coulomb phase shift.

Photoionization cross sections

Total and partial cross sections:

$$\sigma_{\text{tot}} = \sum_{\gamma_0} \sigma_{\gamma_0}, \quad \sigma_{\gamma_0} = 4 \pi^2 \alpha f_{\gamma_0 I} \quad (\text{in unit of } a_0^2)$$

Where

$$f_{\gamma_0 I}^{(L)}(E) = \frac{\Delta E}{3g_I} \sum_{\text{all M's}} \left| \langle \Psi_I | \Sigma \vec{r}_i | \gamma_0 E^- \rangle \right|^2$$

$$f_{\gamma_0 I}^{(V)}(E) = \frac{4}{3g_I (\Delta E)} \sum_{\text{all M's}} \left| \langle \Psi_I | \Sigma \vec{\nabla}_i | \gamma_0 E^- \rangle \right|^2$$

$\Delta E = E - E_I$ is the excitation energy in Ry,

$g_I = (2S_I + 1)(2L_I + 1)$, and Ψ_I is the initial state.

Discretized K matrix

If γ represents a closed channel,

$$\begin{aligned} K^{(1)} &\equiv \sum_{\gamma} \sum_{v_{\gamma}} \langle \gamma' \varepsilon' | V(E) | \gamma \varepsilon_{v_{\gamma}} \rangle \frac{\langle \gamma \varepsilon_{v_{\gamma}} | K(E) | \gamma_0 E \rangle}{E - \varepsilon_{v_{\gamma}}} \\ &= \langle \gamma' \varepsilon' | V(E) X(E) | K(E) | \gamma_0 E \rangle \end{aligned}$$

where

$$\langle \gamma' \varepsilon_{v'} | X(E) | \gamma \varepsilon_v \rangle = \delta_{\gamma' \gamma} \delta_{v' v} \frac{1}{E - \varepsilon_v}$$

Discretized K matrix (Conti.)

If γ represents an open channel,

$$K^{(2)} \equiv \sum_{\gamma} \sum_i \mathcal{P} \int_{\varepsilon_i}^{\varepsilon_{i+1}} d\varepsilon \langle \gamma' \varepsilon' | V(E) | \gamma \varepsilon \rangle \frac{\langle \gamma \varepsilon | K(E) | \gamma_0 E \rangle}{E - \varepsilon}$$

Use N -point Lagrange interpolation:

$$f(x) = \sum_{m'} \sum_m \alpha_{mm'} f(x_m) x^{m'-1}, \quad m, m' = 1, \dots, N$$

$$\langle \gamma' \varepsilon' | V(E) | \gamma \varepsilon \rangle \sim \sum_{m'} \sum_m \alpha_{mm'} \langle \gamma' \varepsilon' | V(E) | \gamma \varepsilon_m \rangle \varepsilon^{m'-1}$$

$$\langle \gamma \varepsilon | K(E) | \gamma_0 E \rangle \sim \sum_{n'} \sum_n \alpha_{nn'} \langle \gamma \varepsilon_n | K(E) | \gamma_0 E \rangle \varepsilon^{n'-1}$$

Discretized K matrix (Conti.)

$$K^{(2)} = \langle \gamma' \varepsilon' | V(E) Y(E) K(E) | \gamma_0 E \rangle$$

where

$$\langle \gamma' \varepsilon_m | Y(E) | \gamma \varepsilon_n \rangle = \delta_{\gamma' \gamma} \sum_i Y_{mn}^{(i)}$$

and

$$Y_{mn}^{(i)} = \sum_{m'} \sum_{n'} \alpha_{mm'} \alpha_{nn'} \int_{\varepsilon_i}^{\varepsilon_{i+1}} d\varepsilon \frac{\mathcal{P} \varepsilon^{m'+n'-2}}{E - \varepsilon}$$

Discretized K matrix (Conti.)

$$\langle \gamma' \varepsilon' | [1 - V(E)P(E)] K^d(E) | \gamma_0 E \rangle = \langle \gamma' \varepsilon' | V(E) | \gamma_0 E \rangle$$

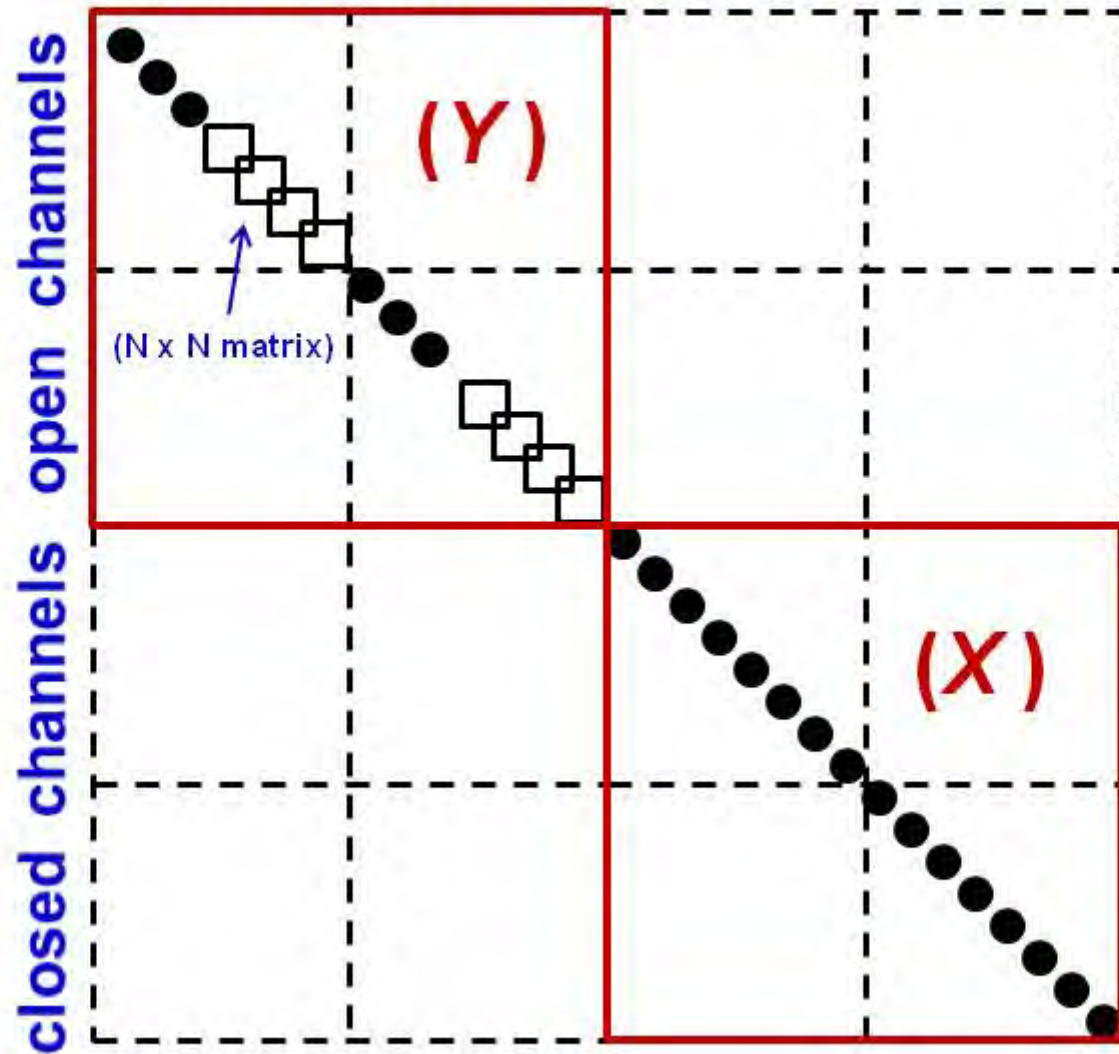
where

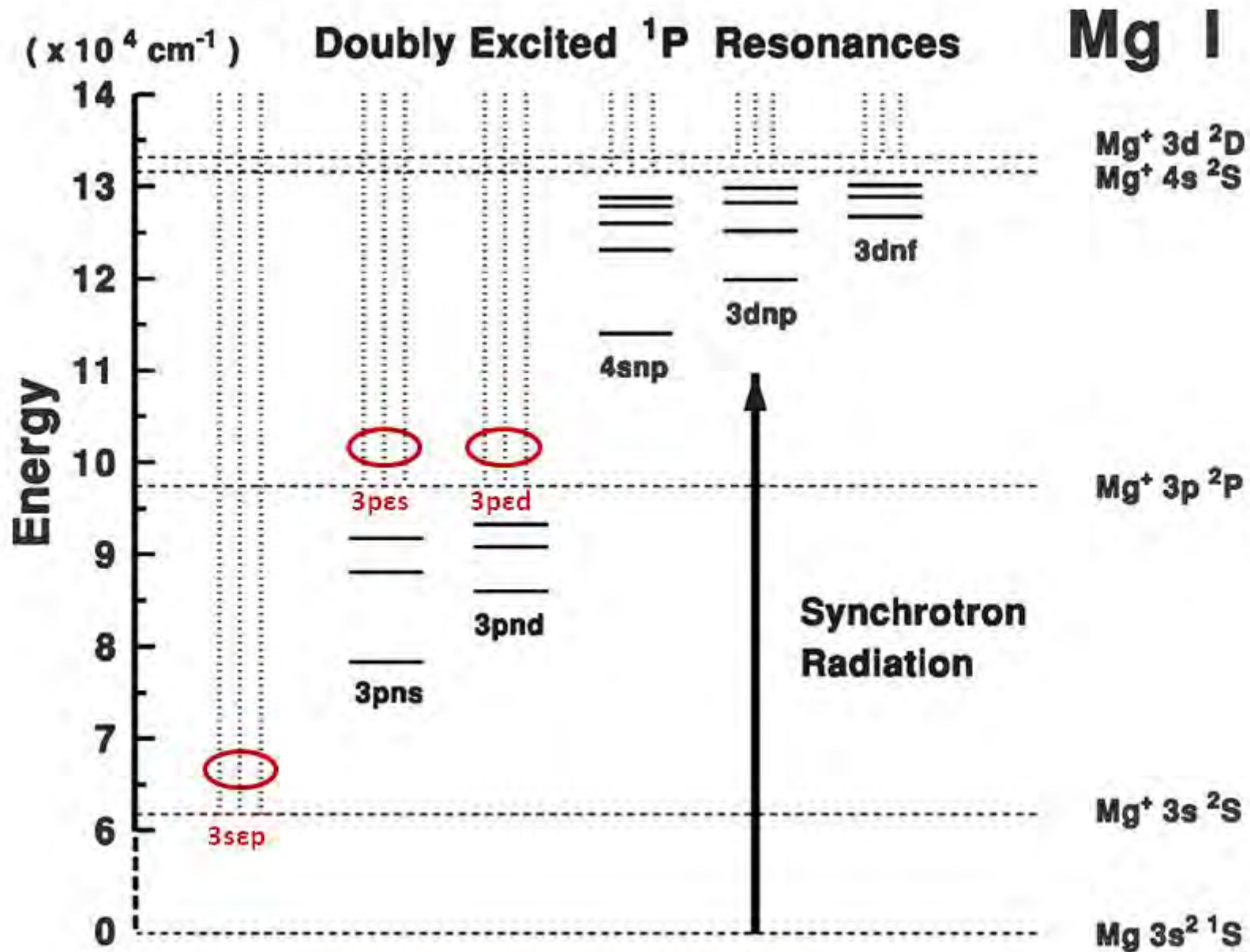
$$P(E) = X(E) + Y(E)$$

$$\Rightarrow K^d(E) = \frac{1}{[1 - V(E)P(E)]} V(E)$$

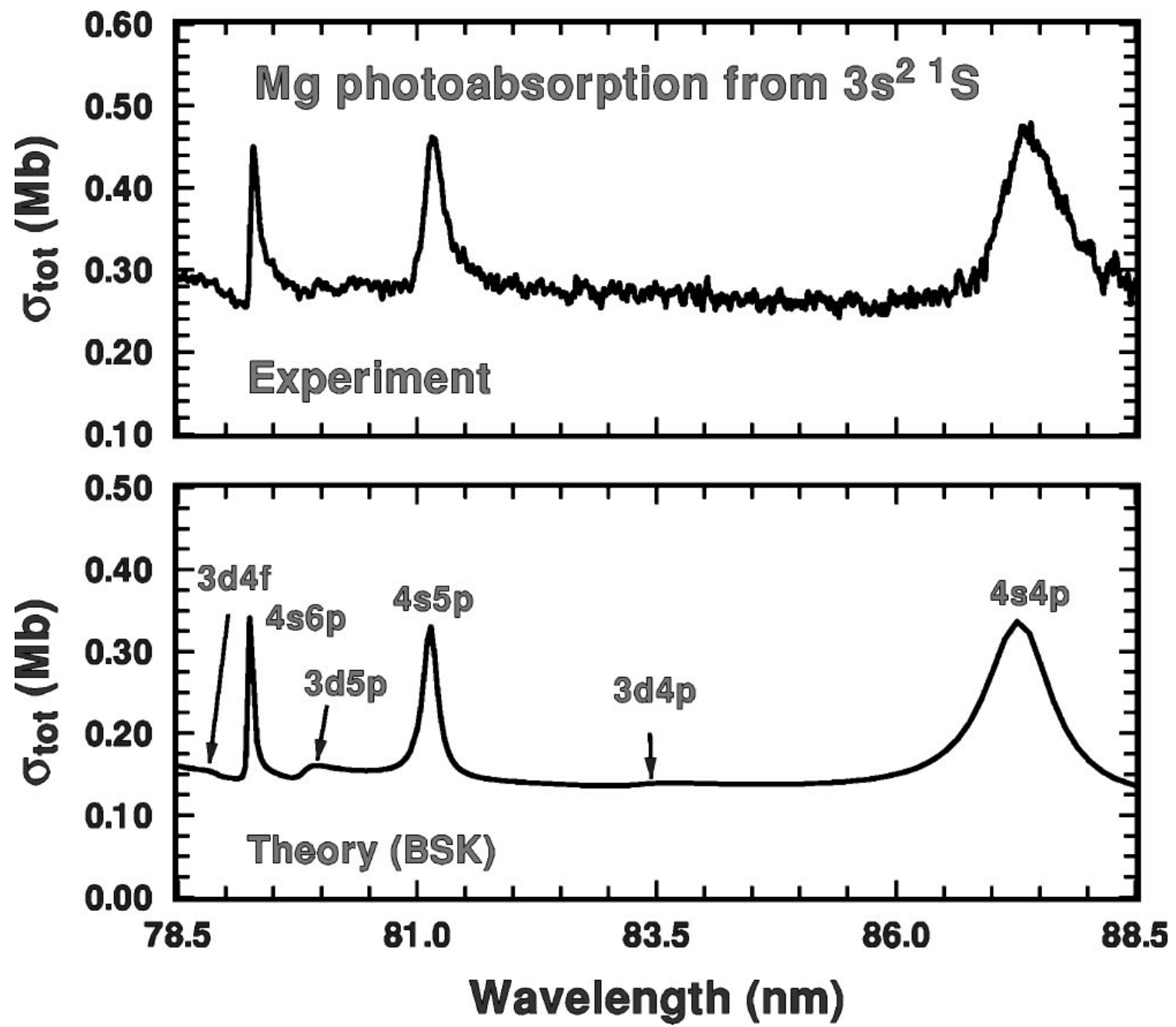
P matrix ($= X + Y$)

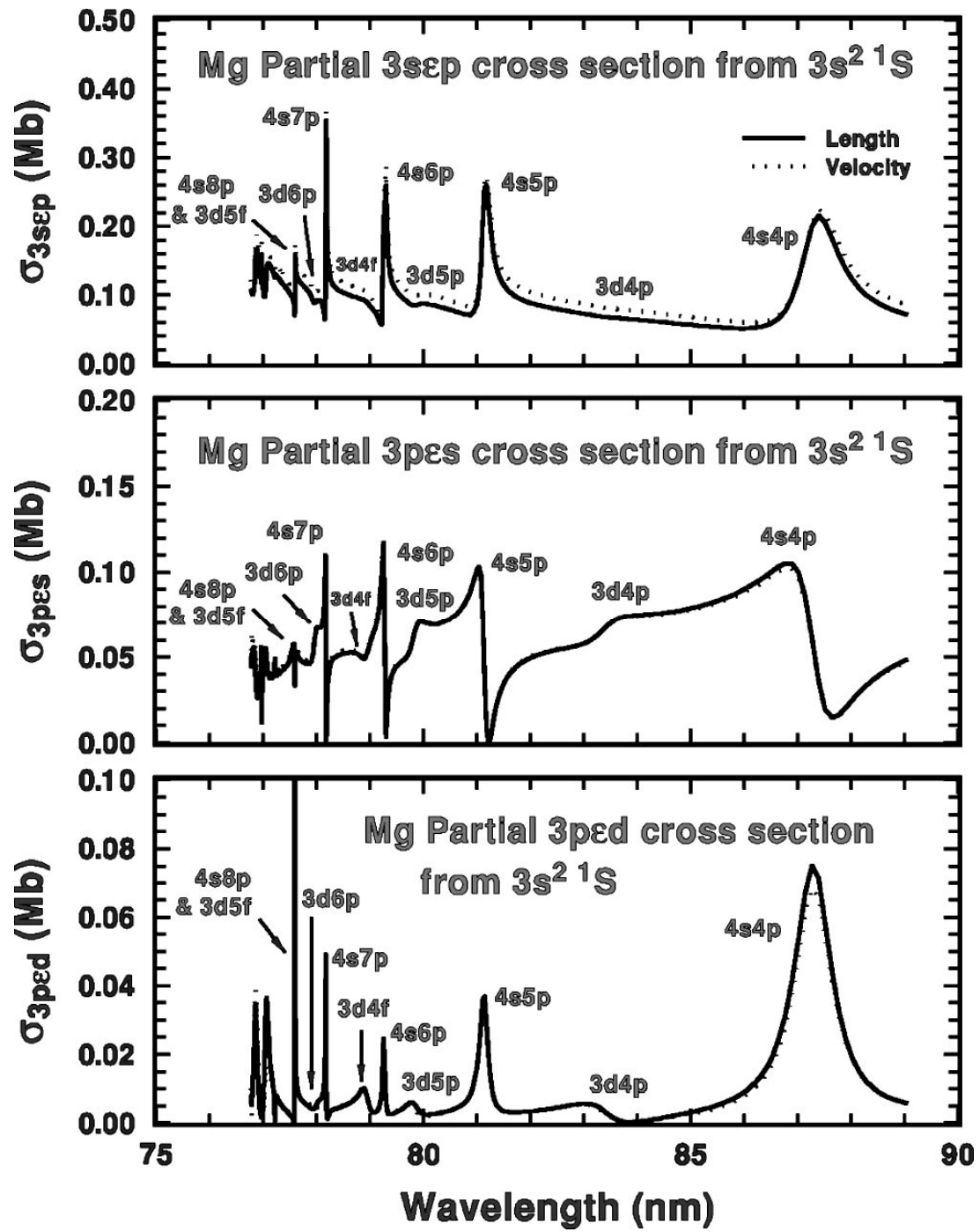
open channels closed channels

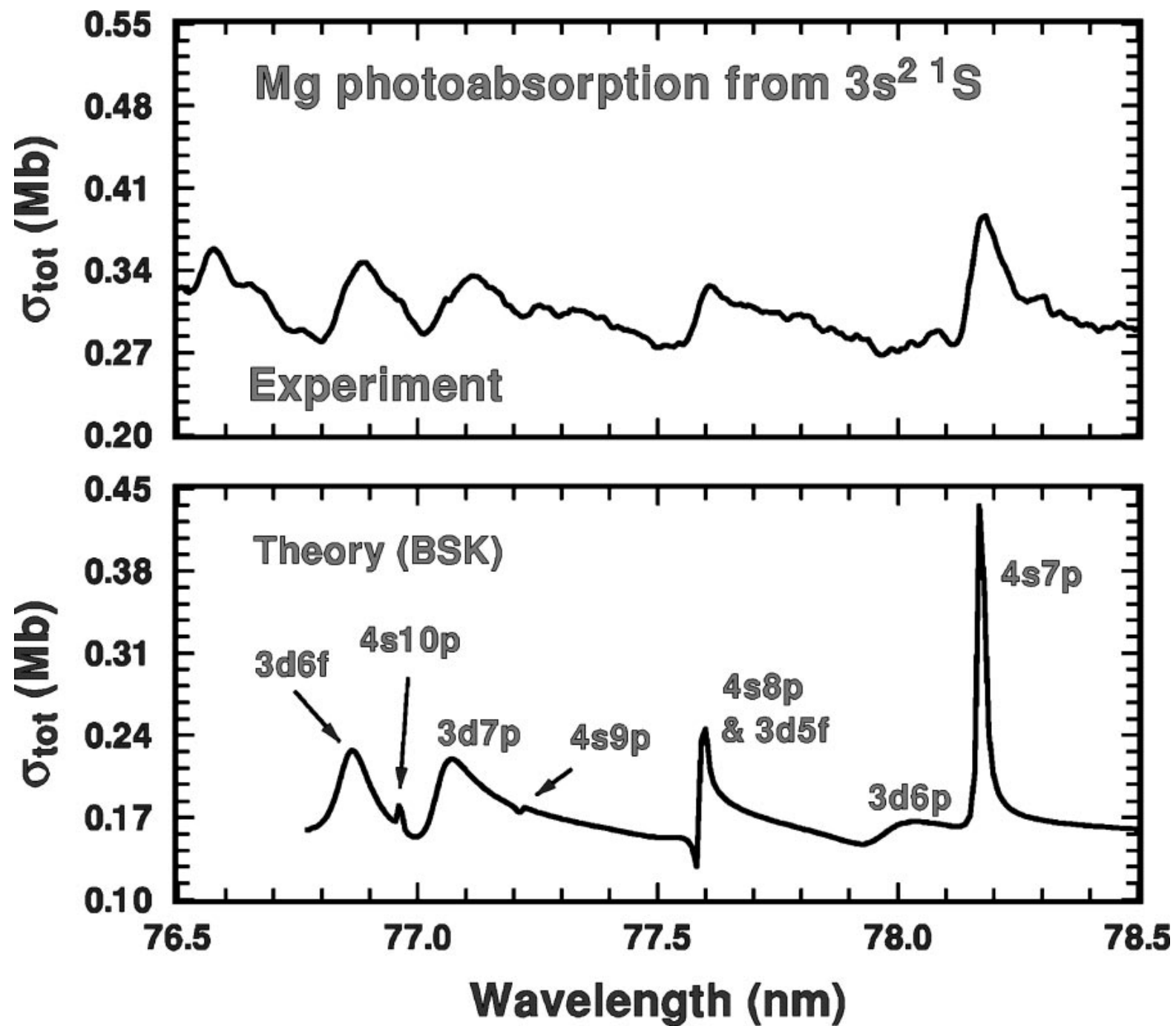


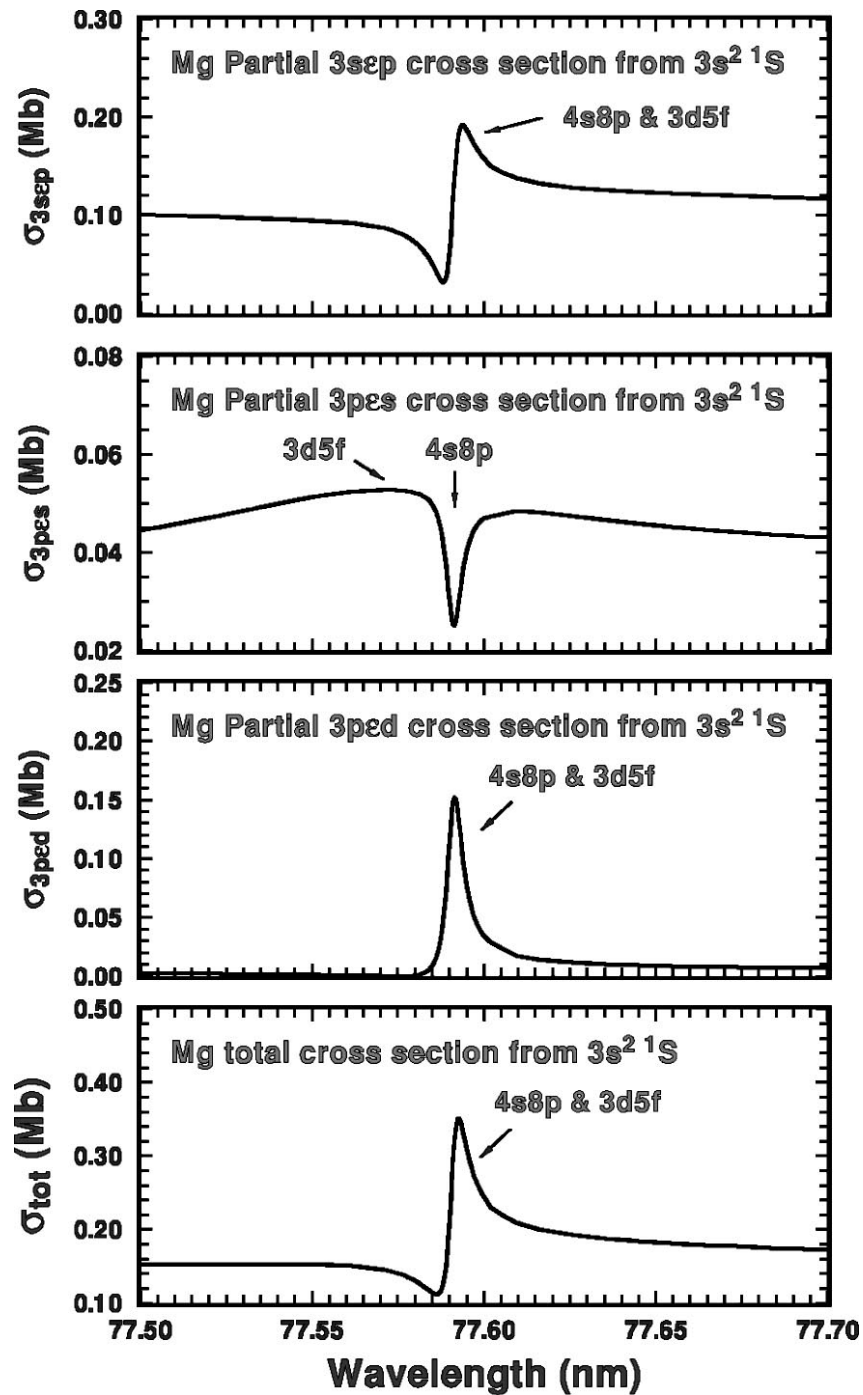


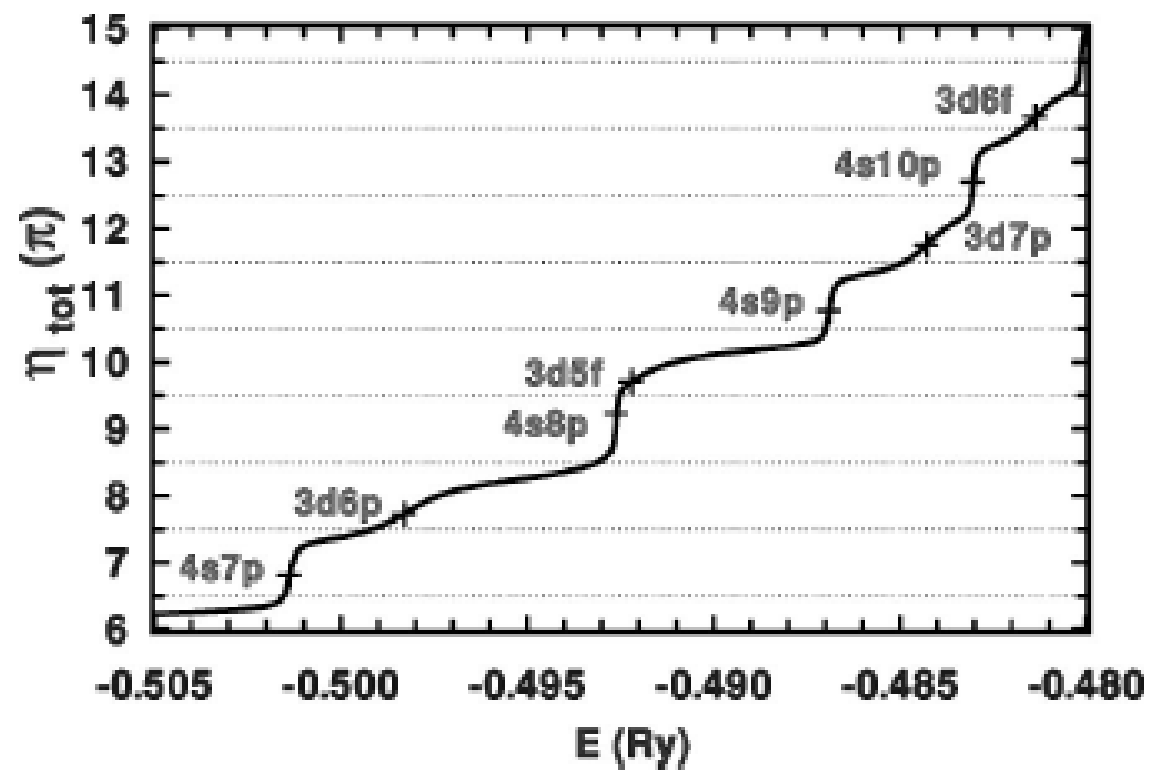
Phys. Rev. A 64, 052716 (2001).



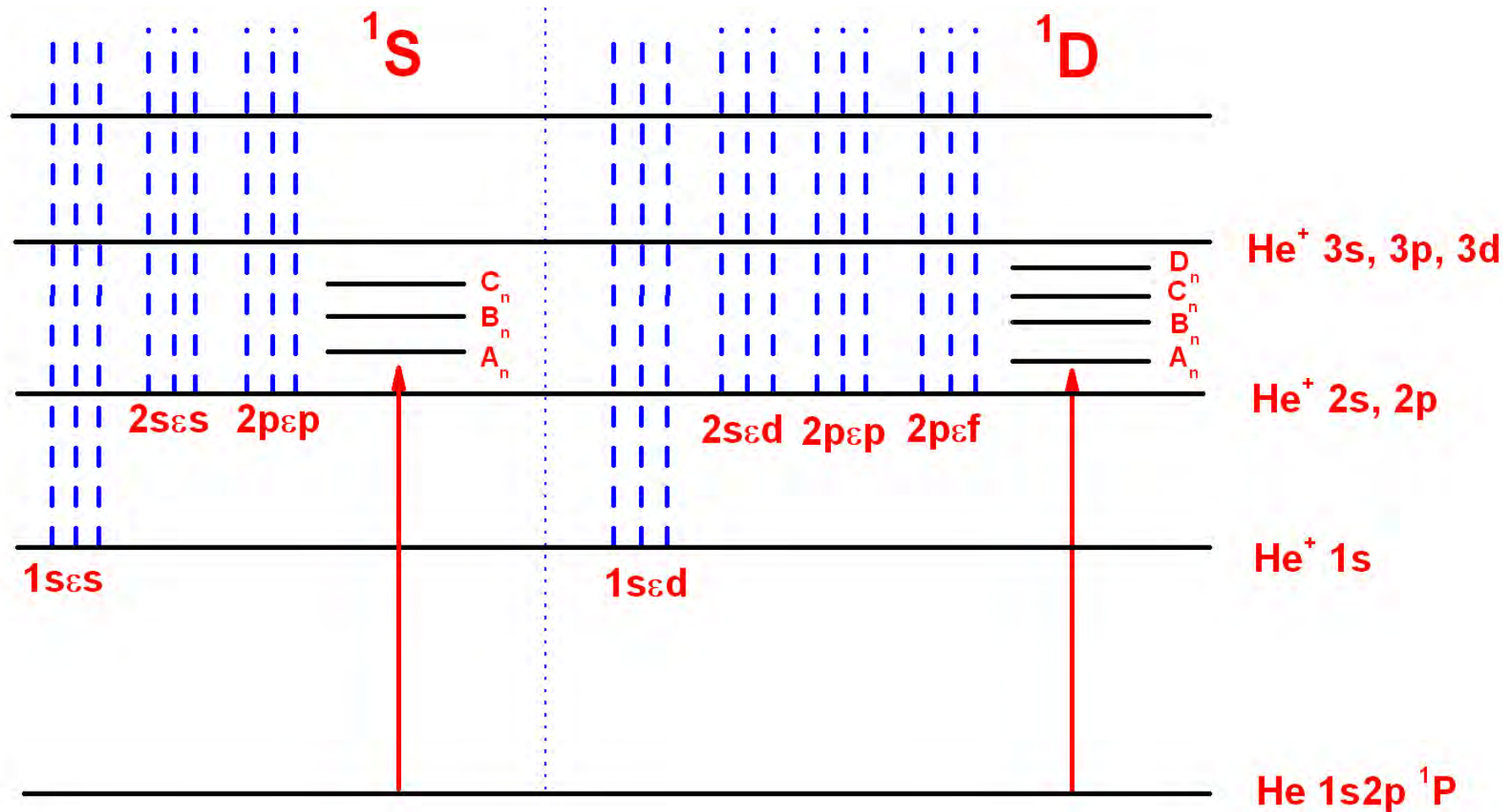


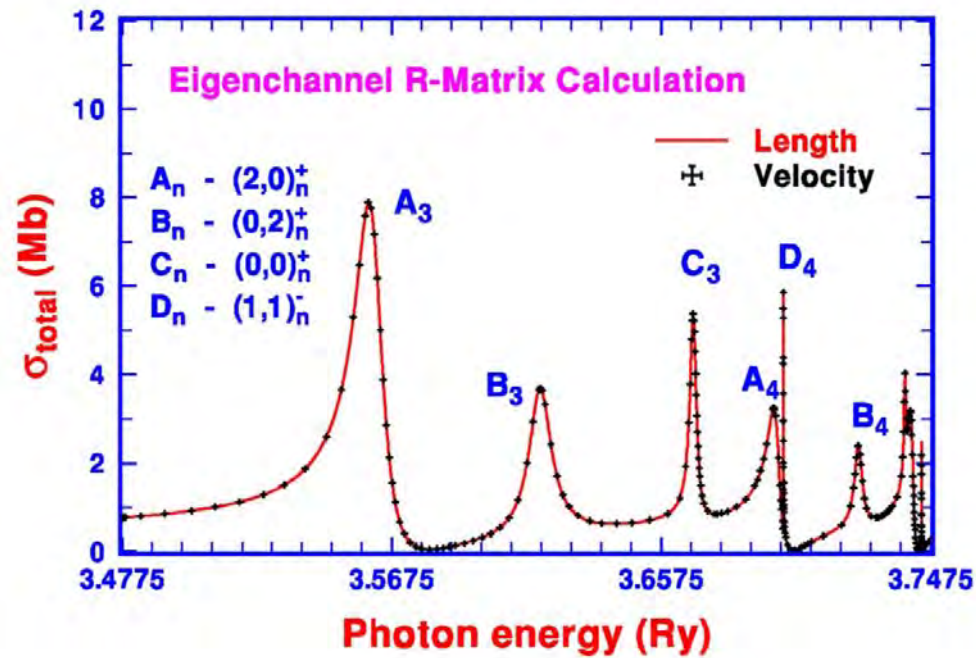
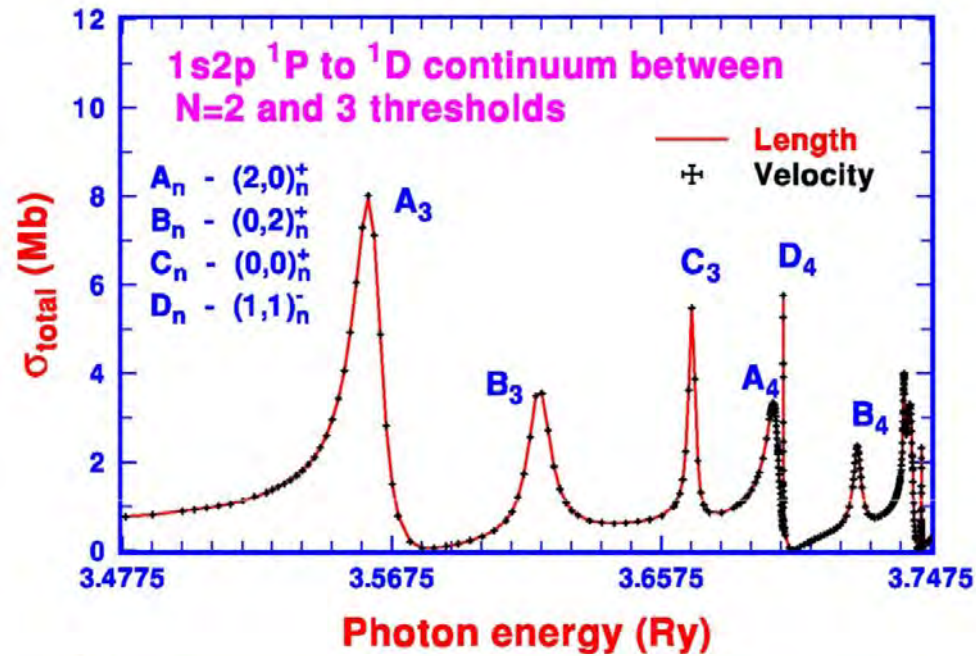


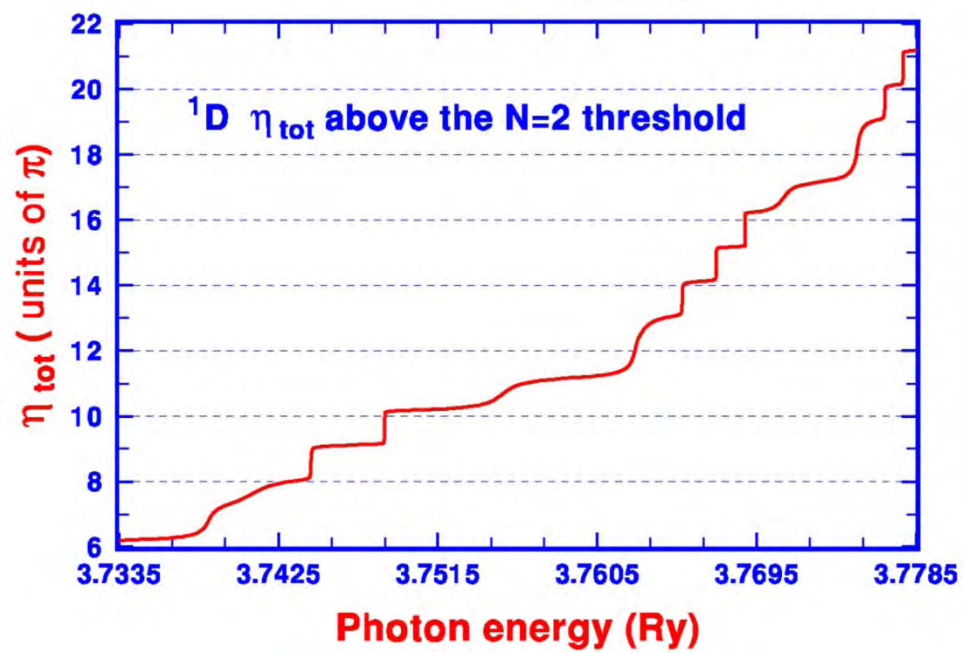
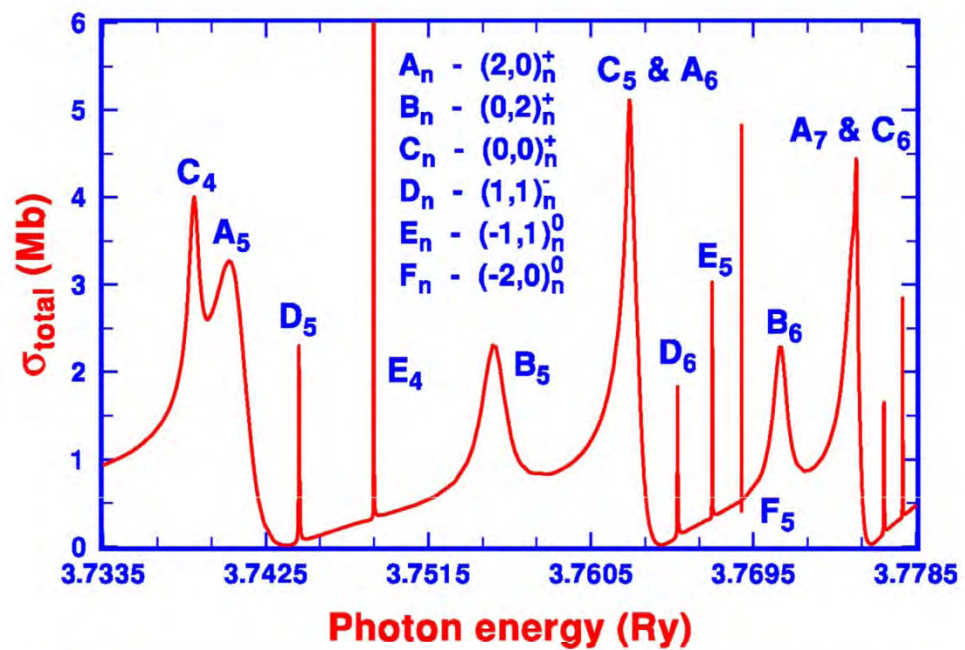


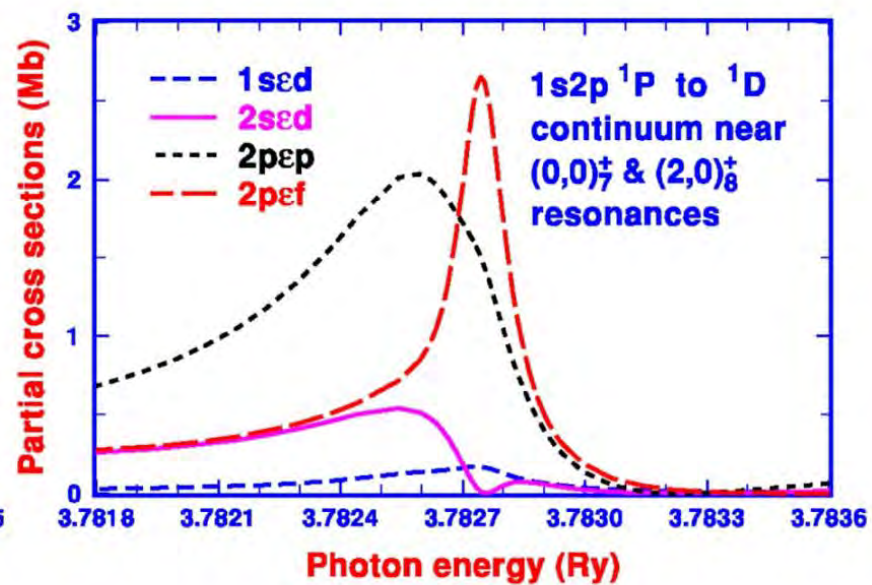
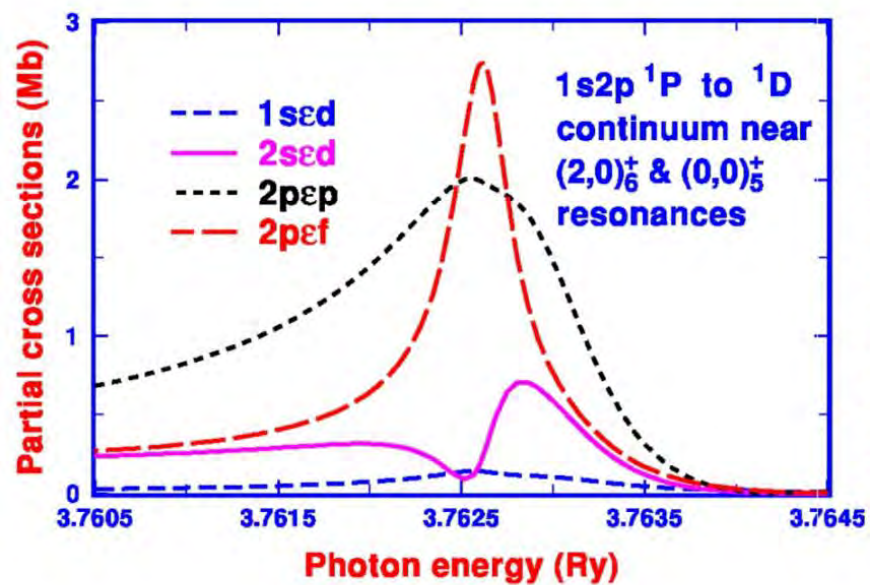
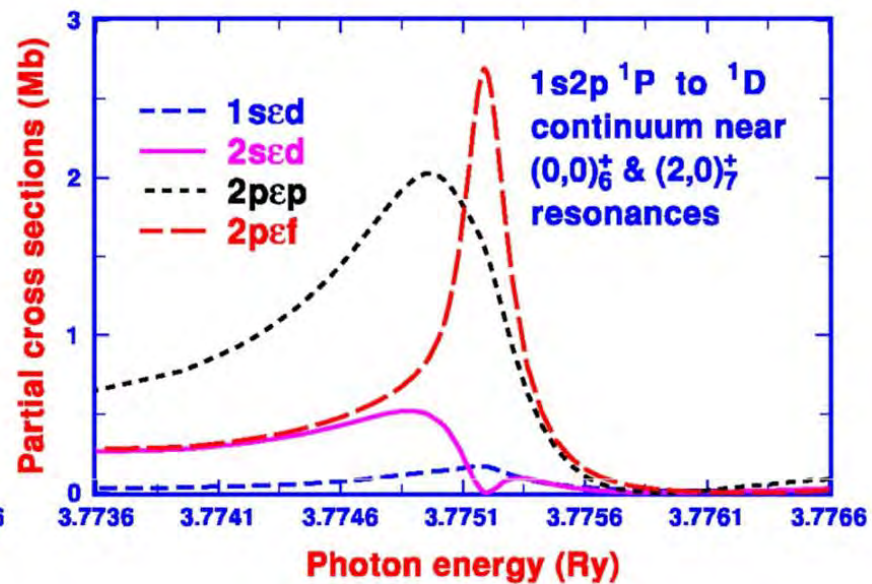
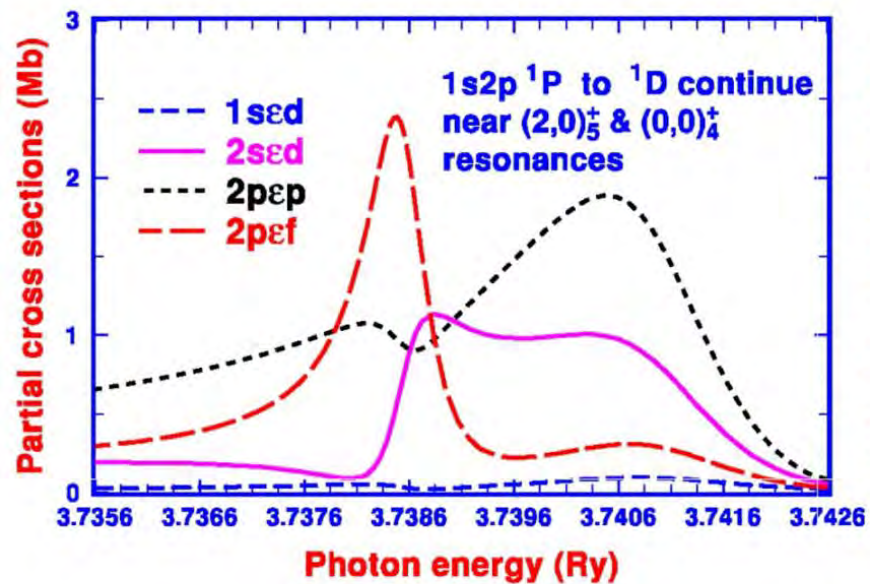


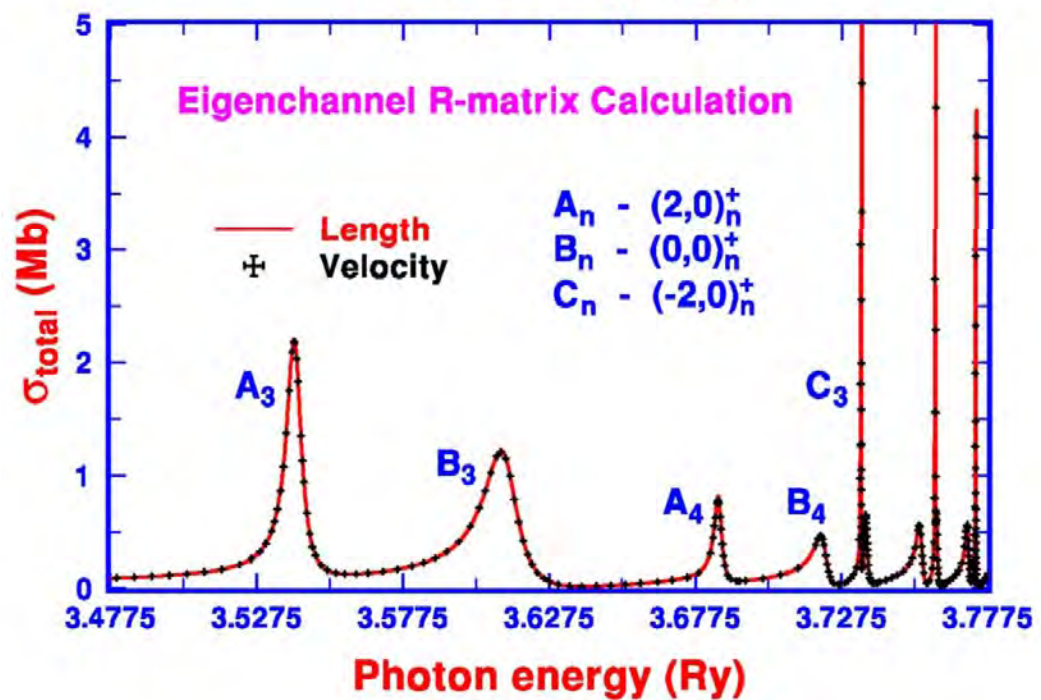
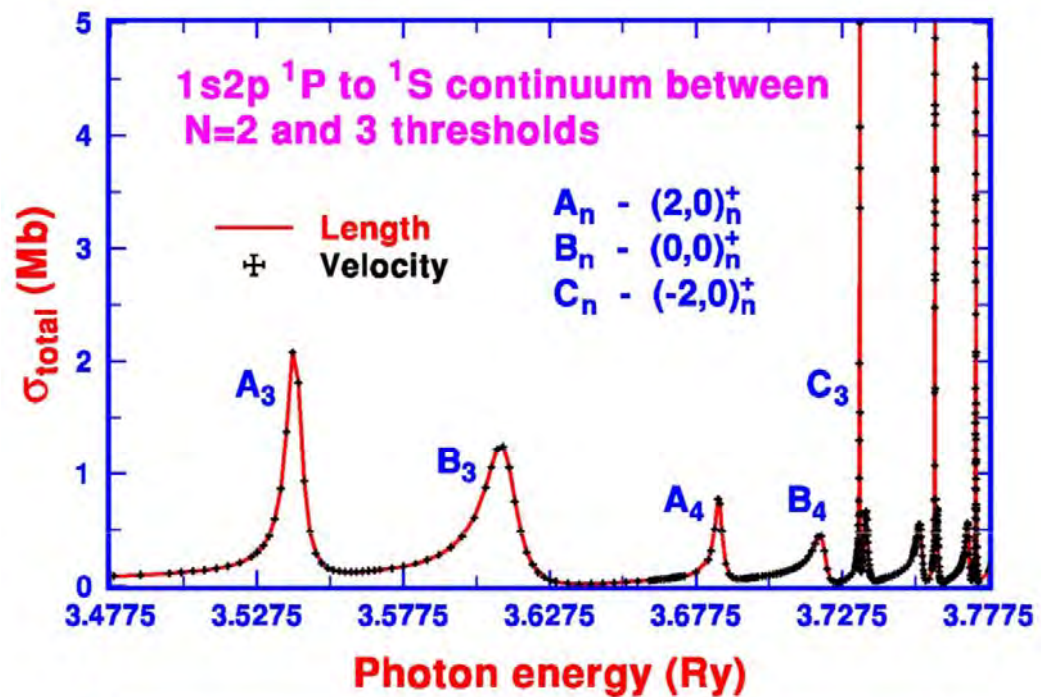
Overlapping doubly excited resonances

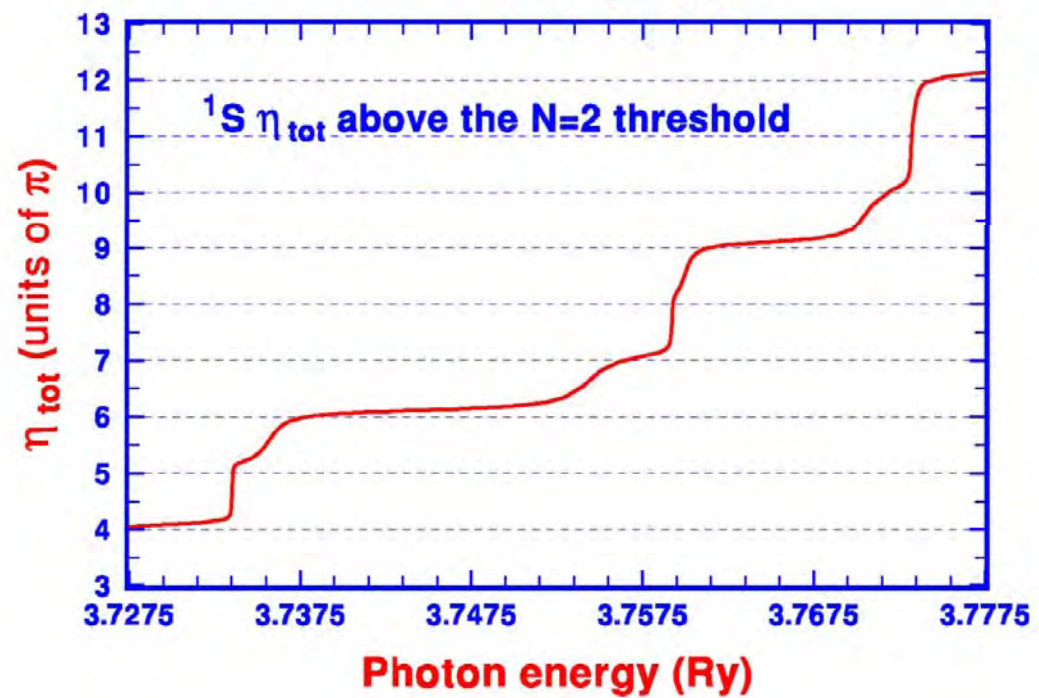
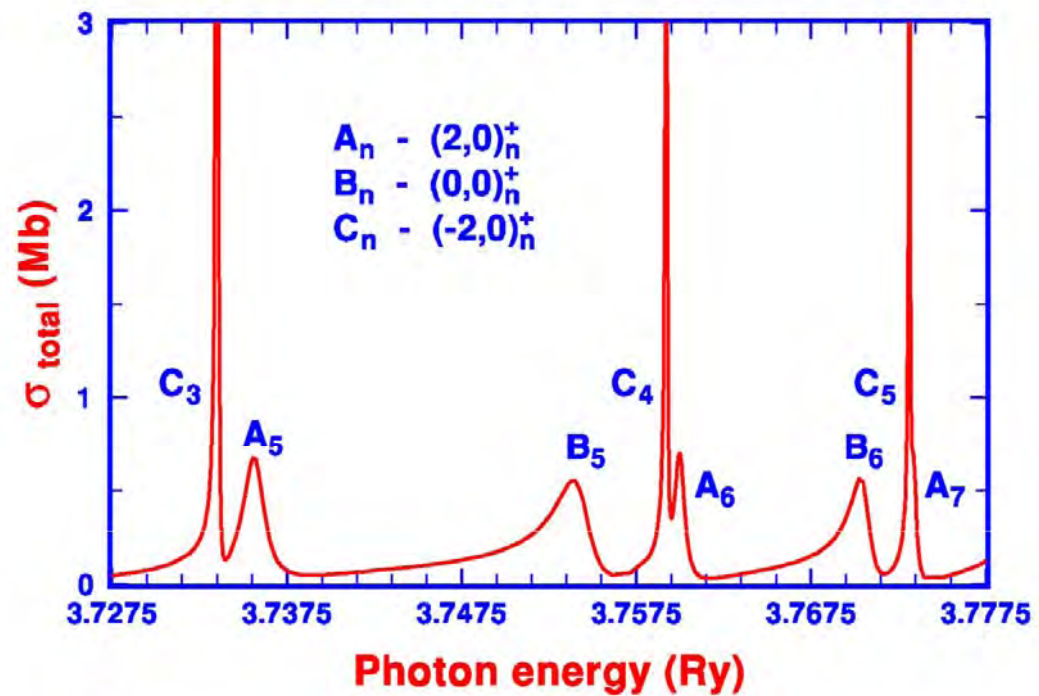


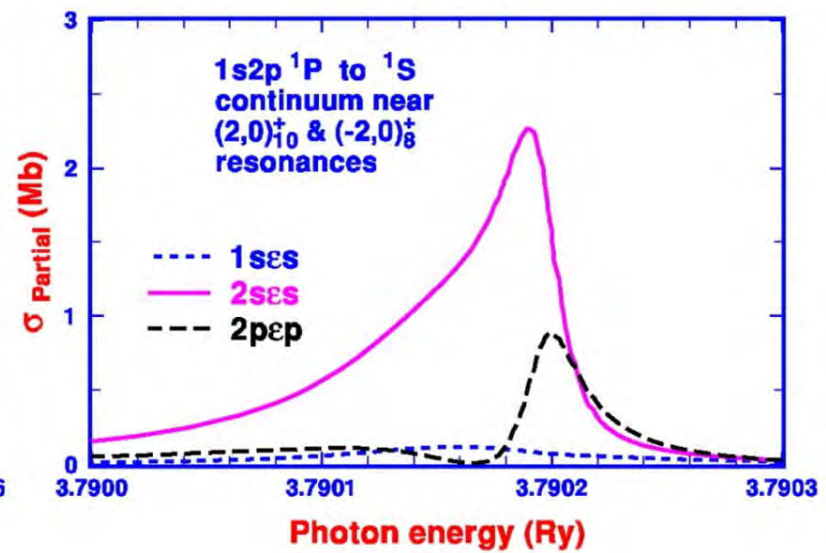
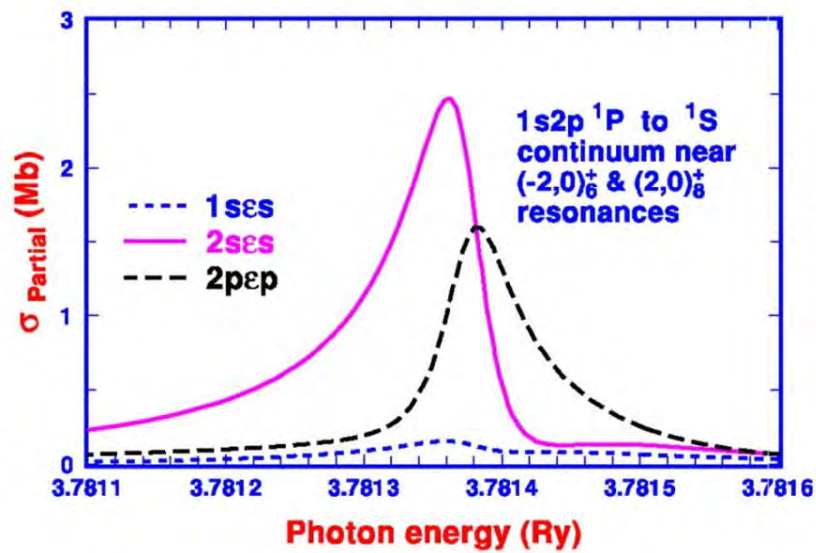
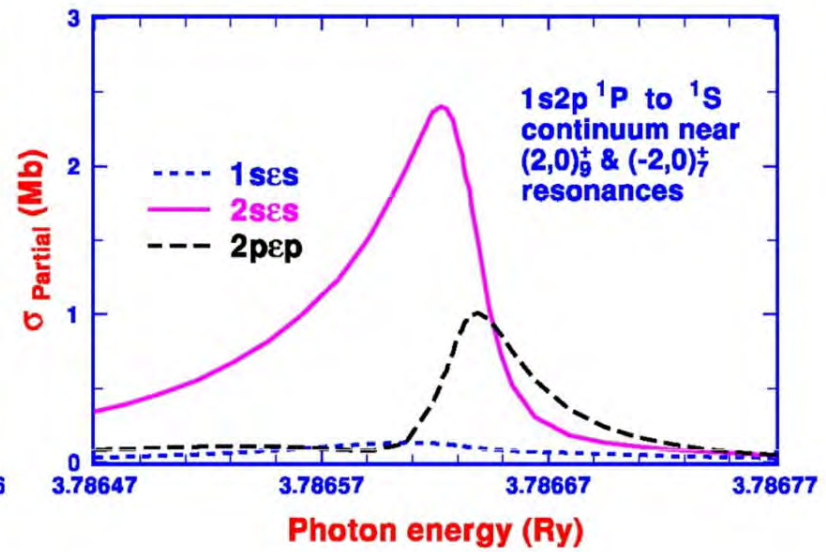
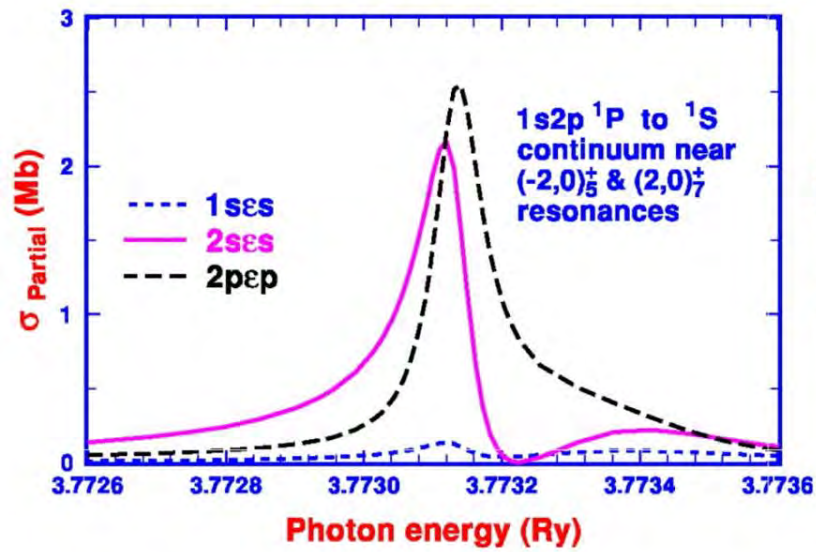






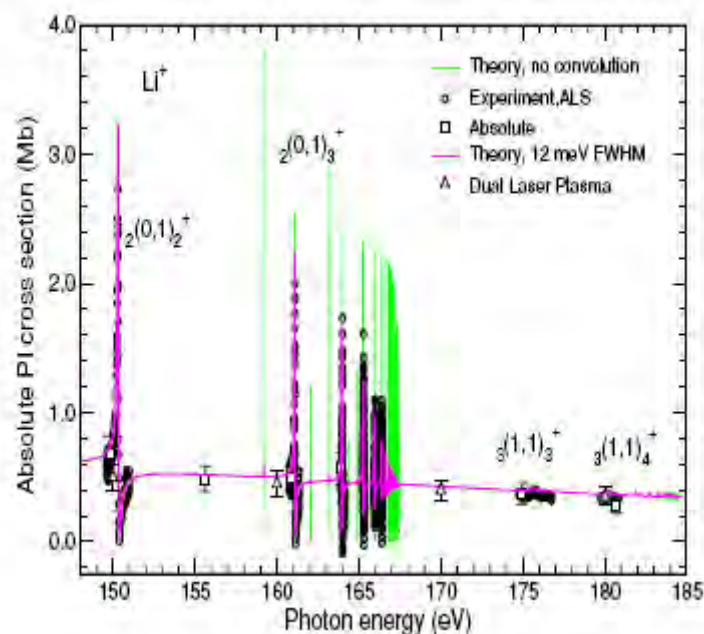


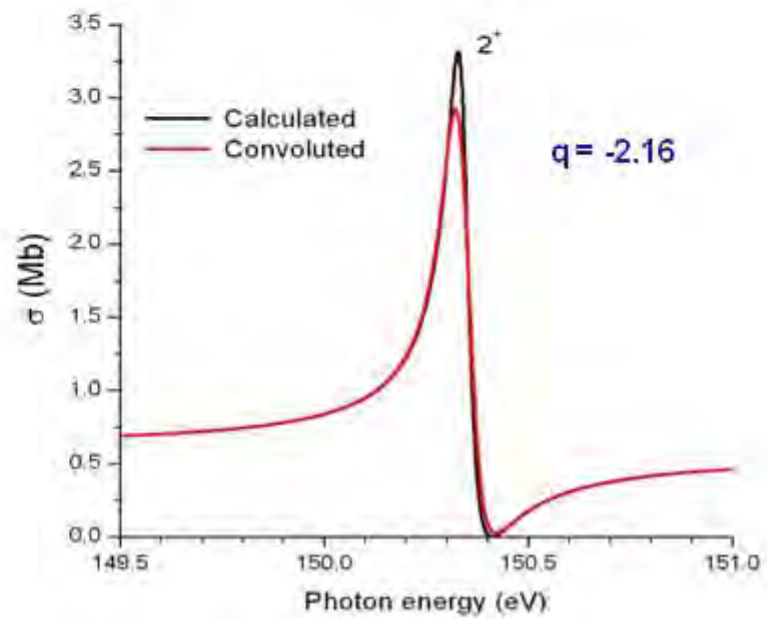
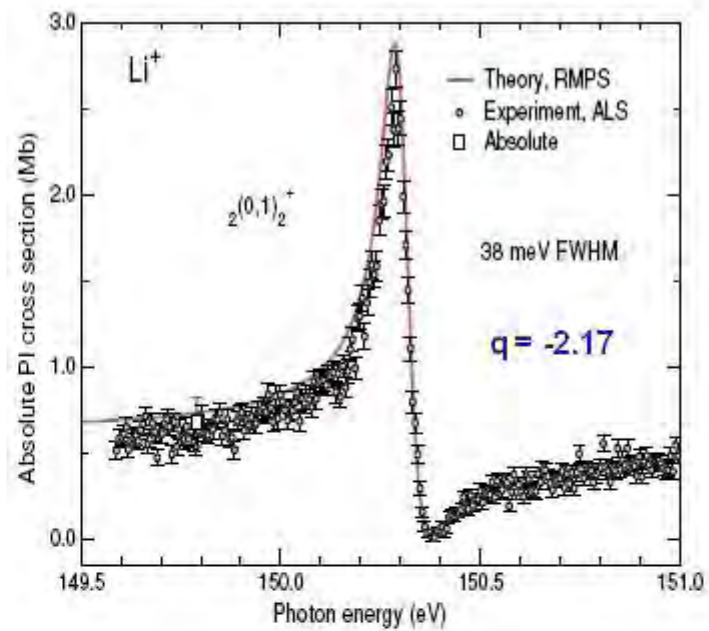


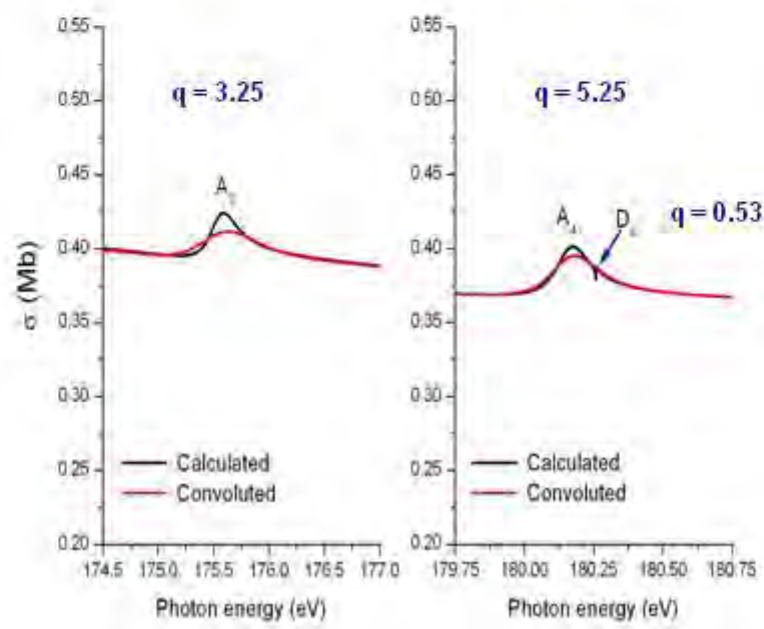
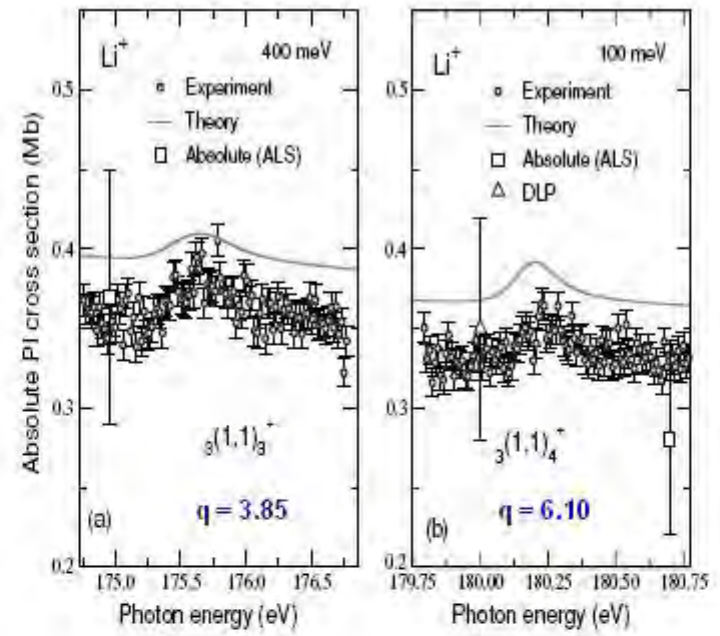


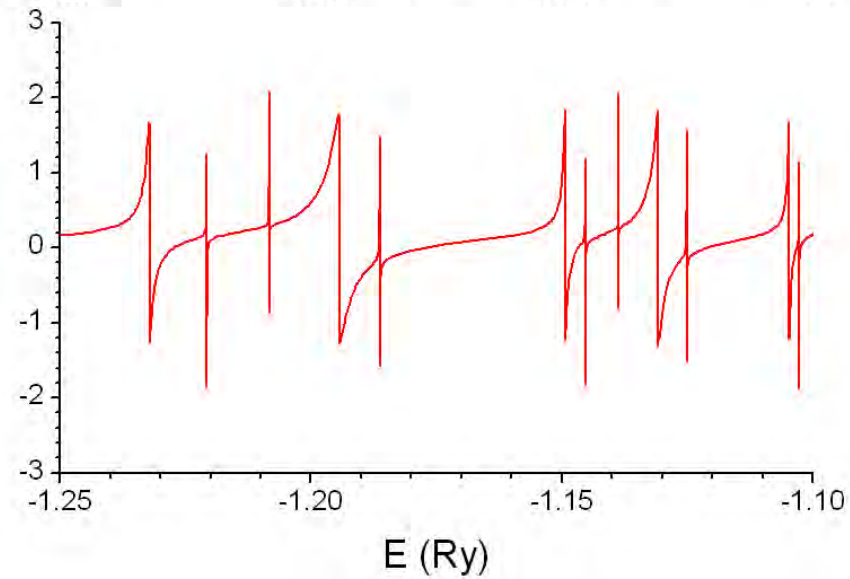
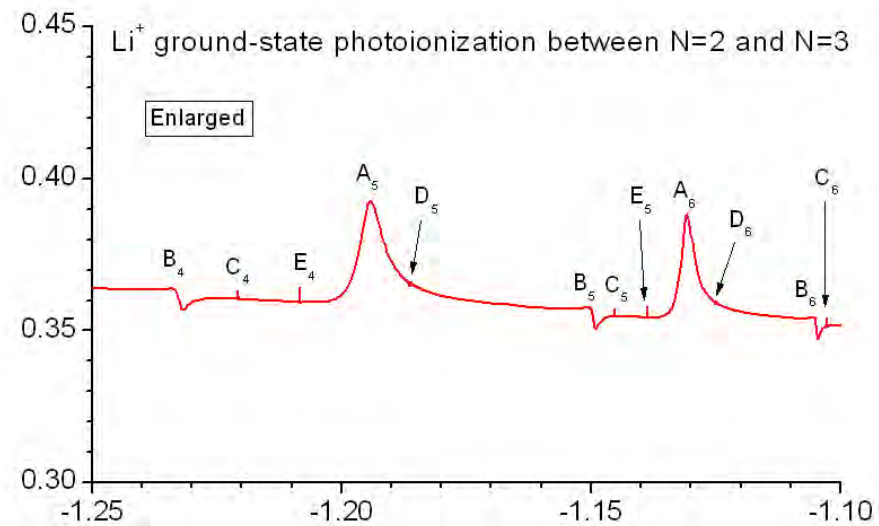
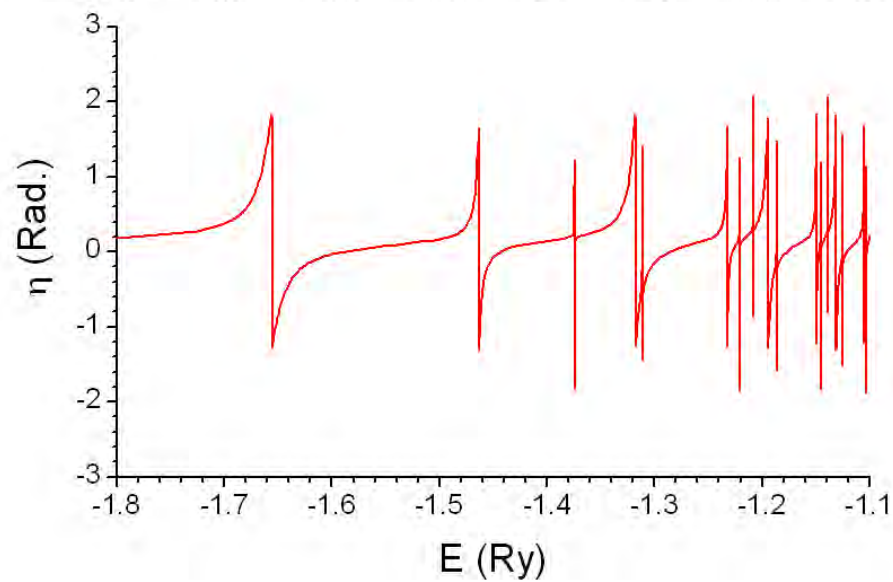
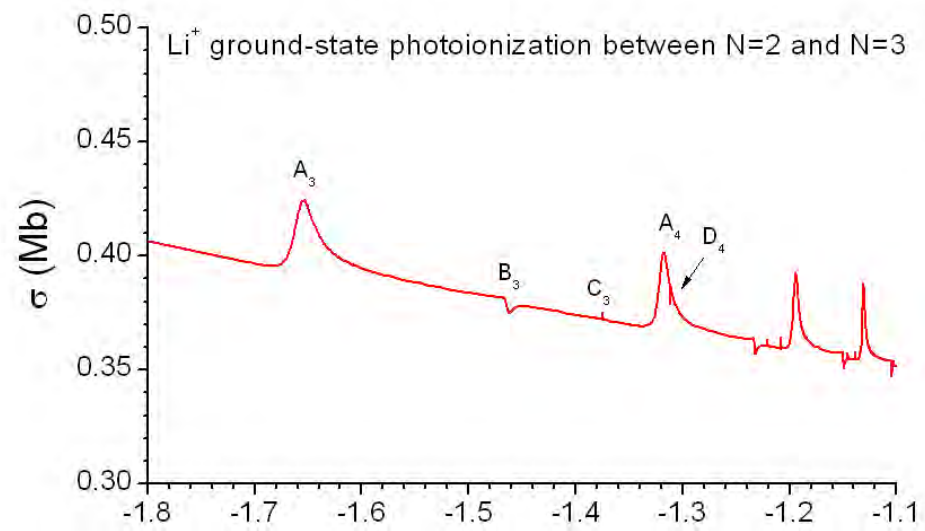
Doubly excited resonances in the photoionization spectrum of Li^+ : experiment and theory

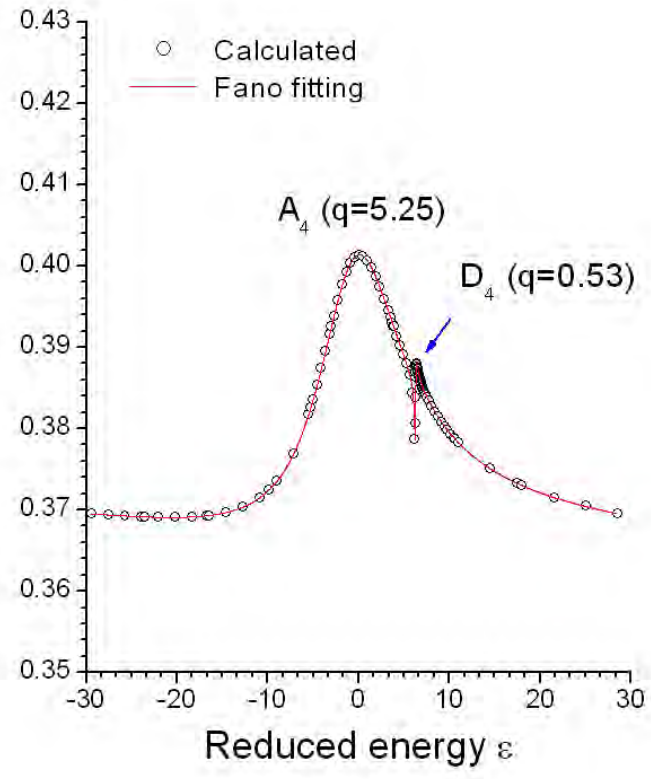
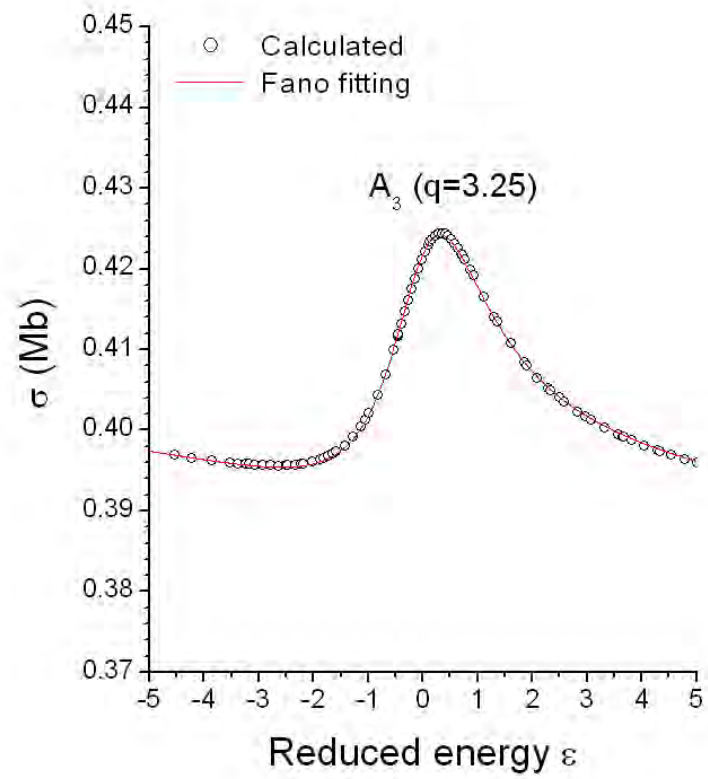
S W J Scully^{1,9}, I Álvarez², C Cisneros², E D Emmons¹, M F Gharaibeh¹,
D Leitner³, M S Lubell⁴, A Müller⁵, R A Phaneuf¹, R Püttner⁶,
A S Schlachter³, S Schippers⁵, W Shi^{5,10}, C P Ballance⁷
and B M McLaughlin^{8,9}

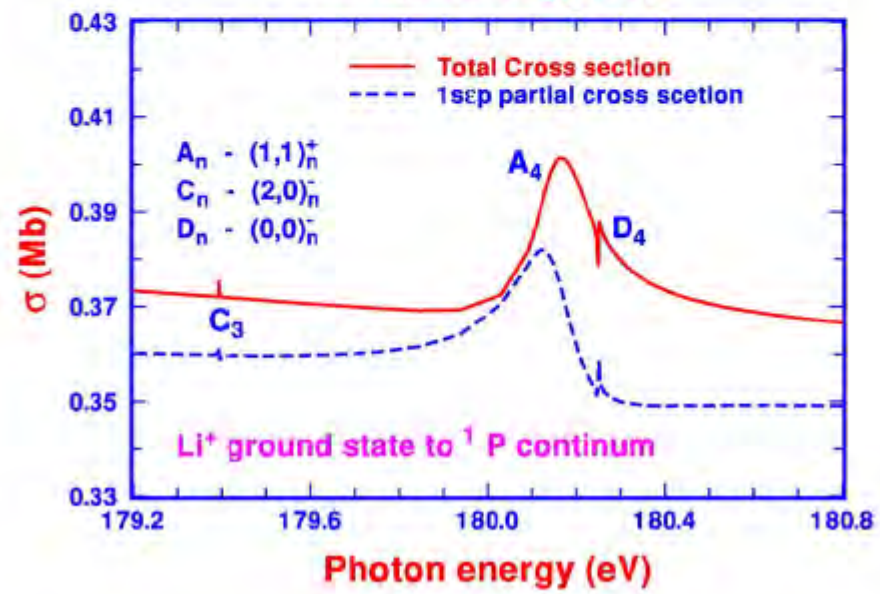
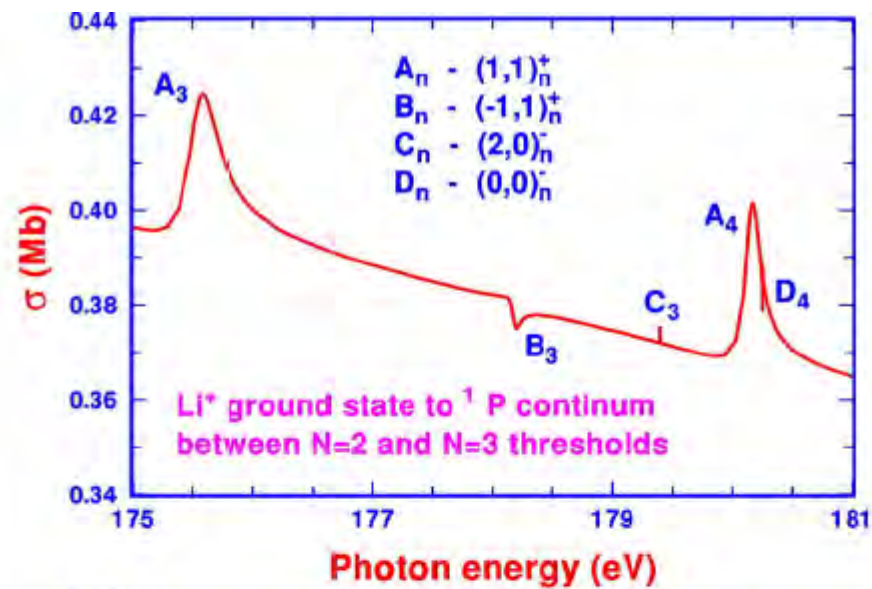


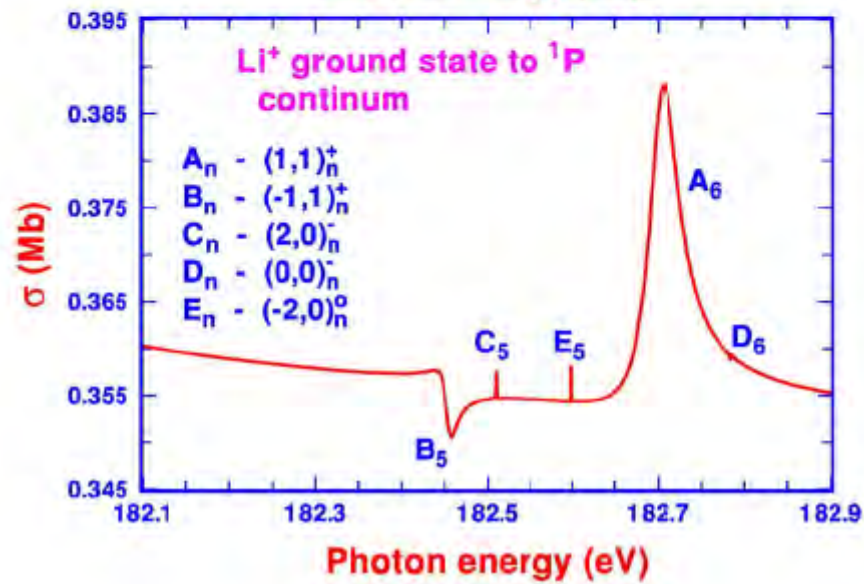
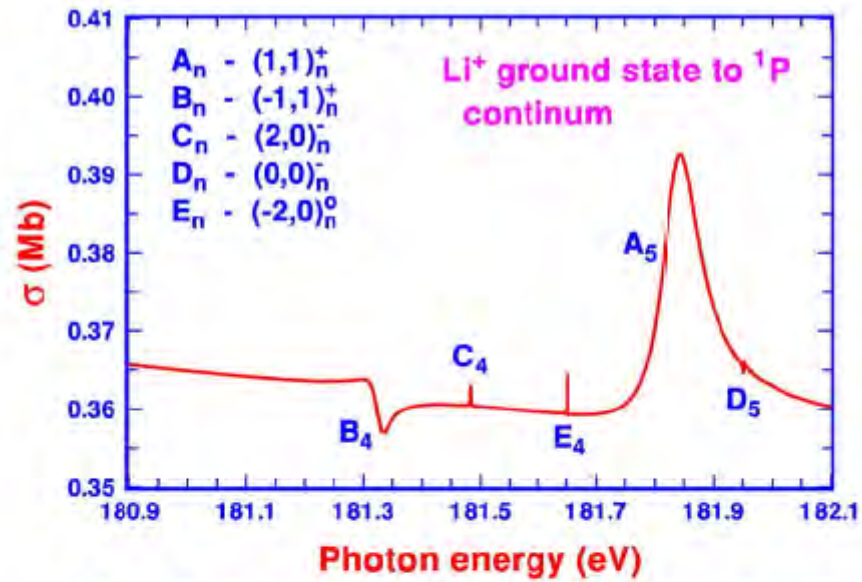


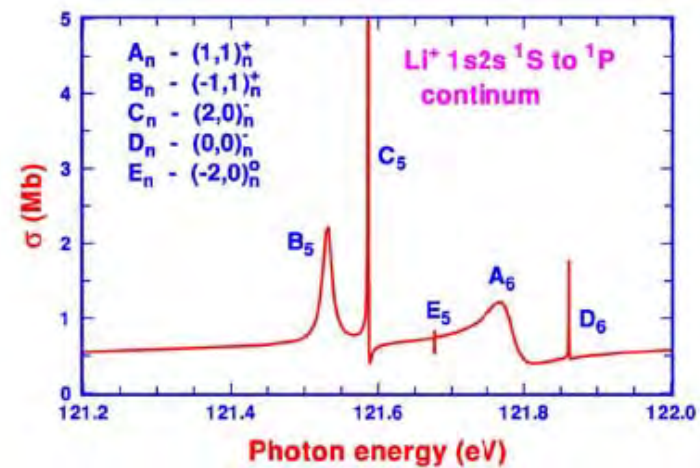
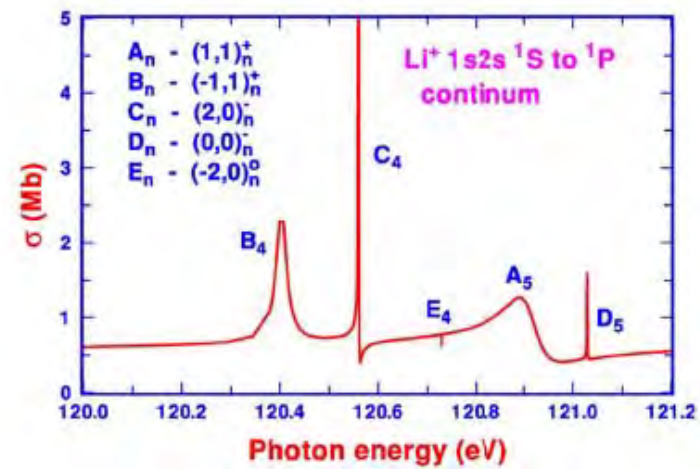
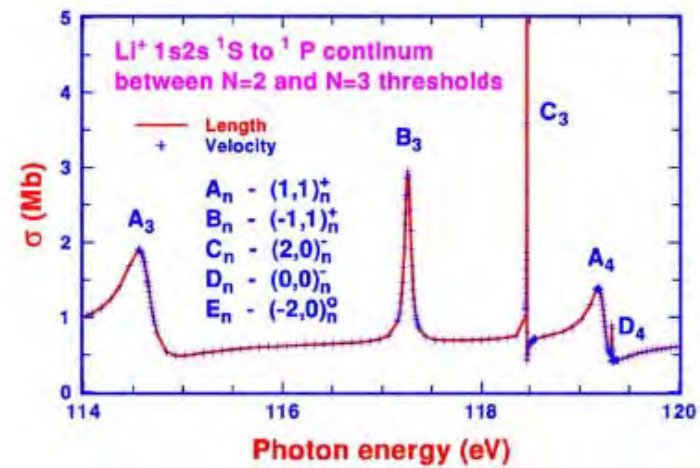


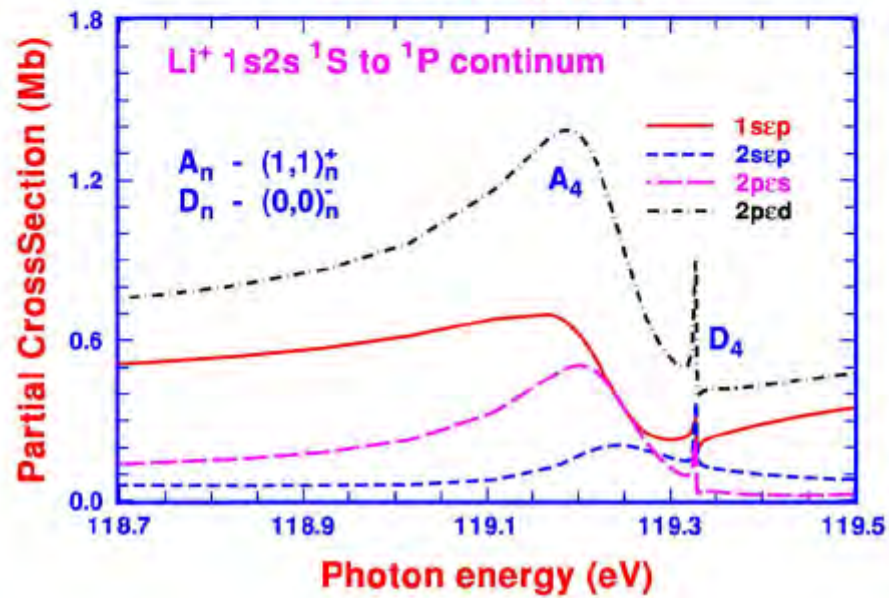
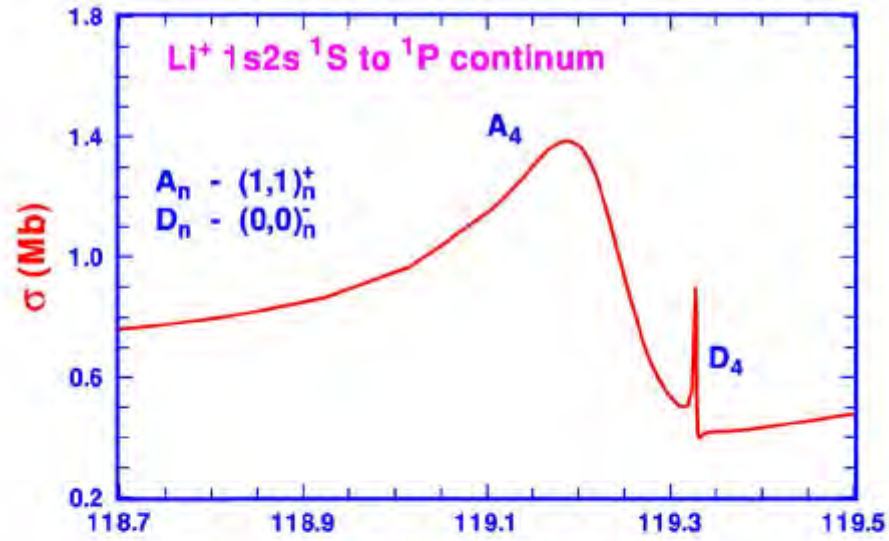


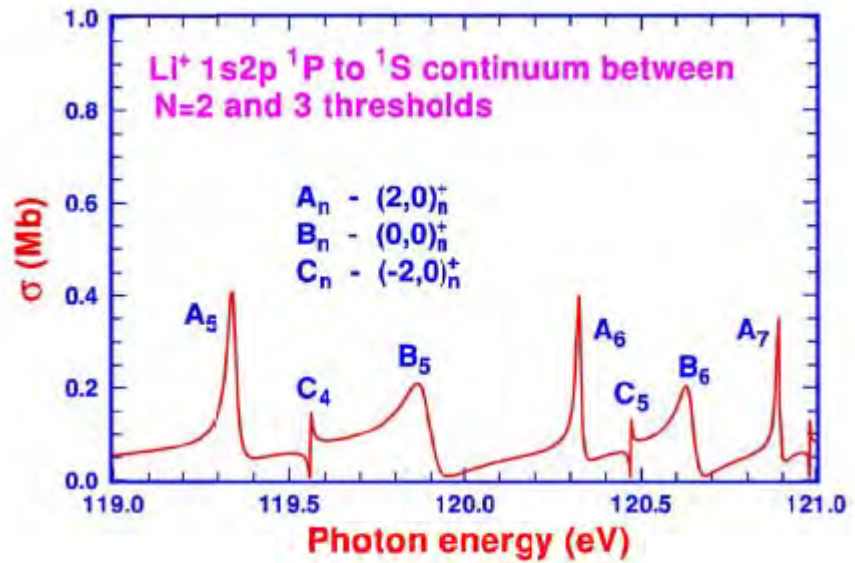
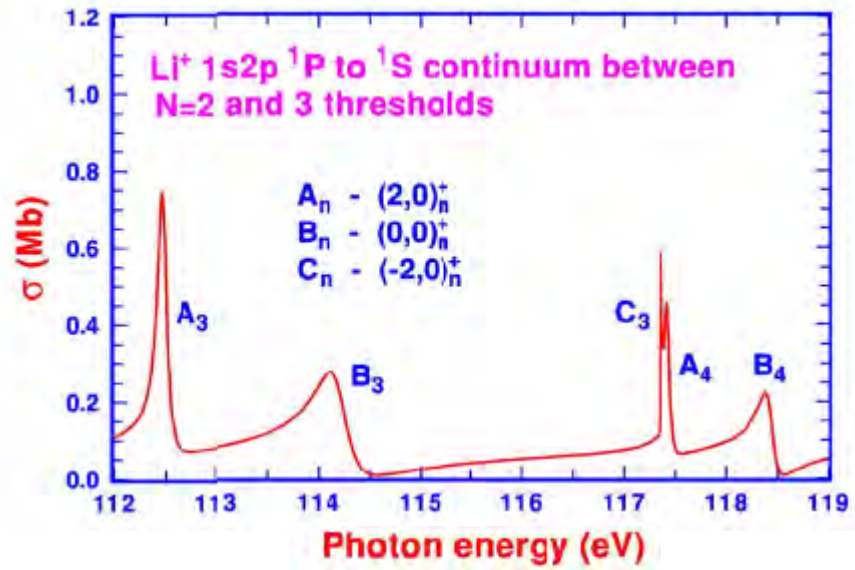


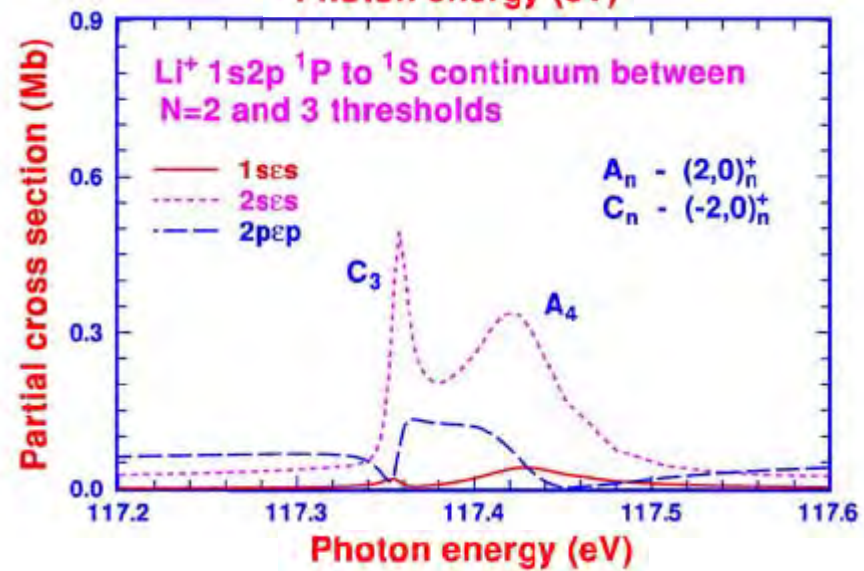
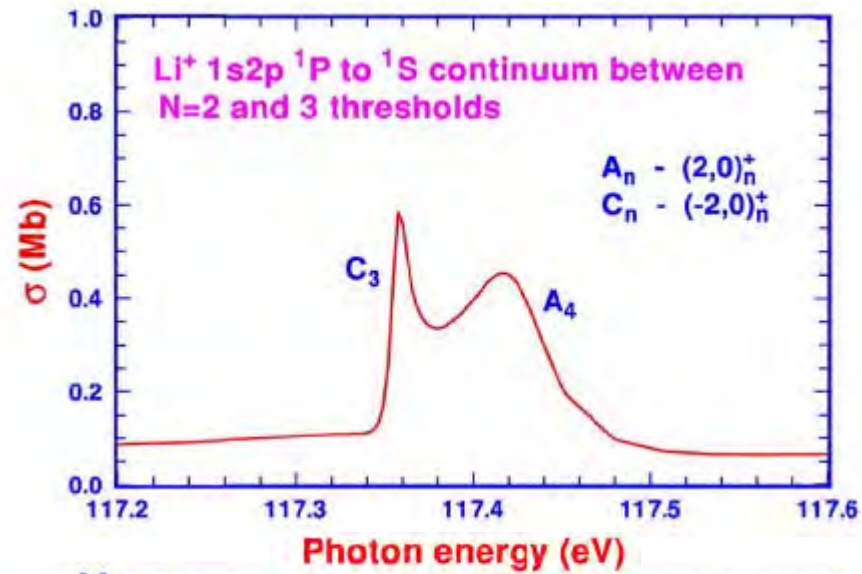


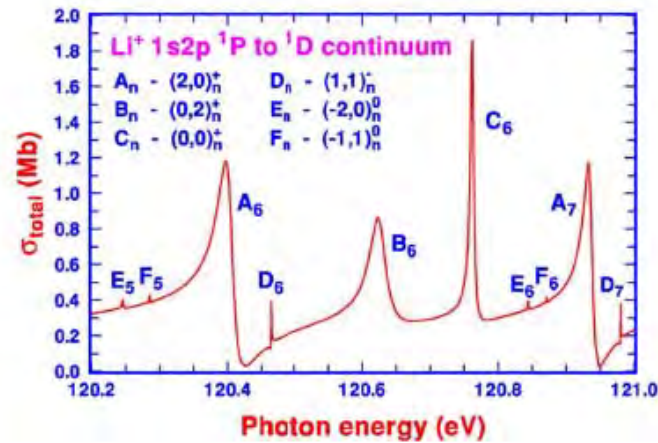
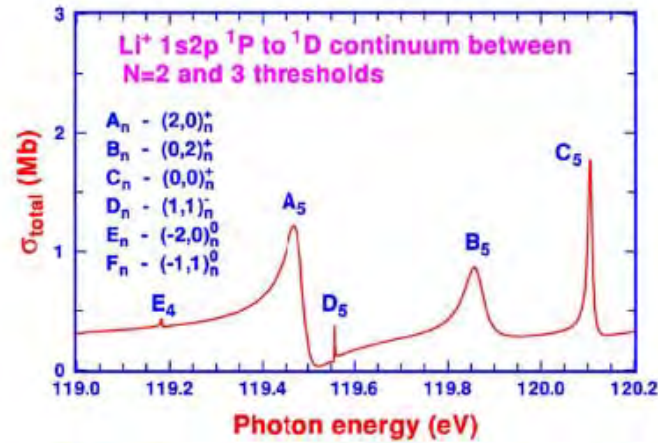
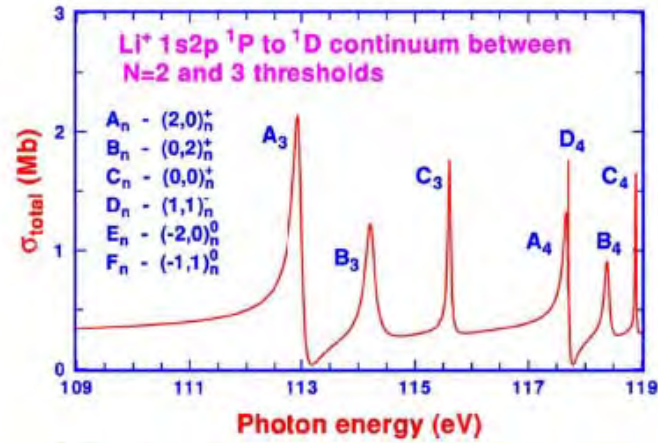


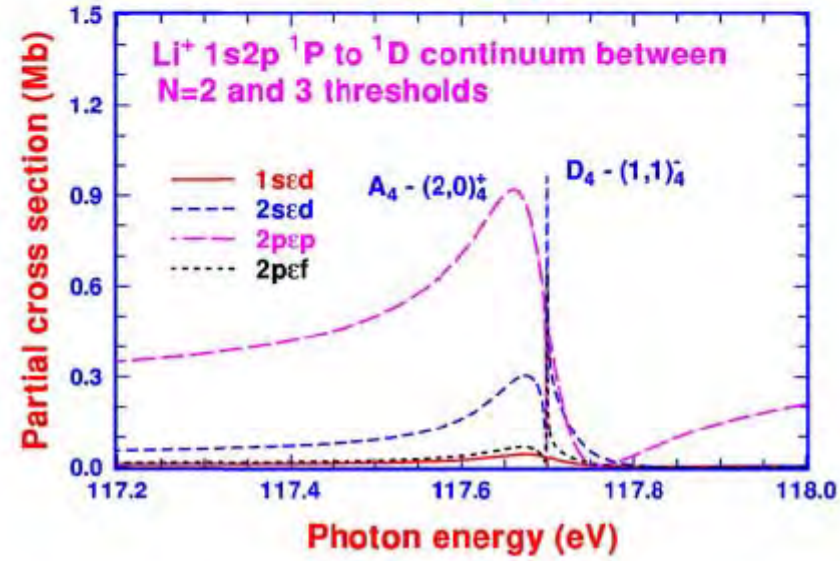
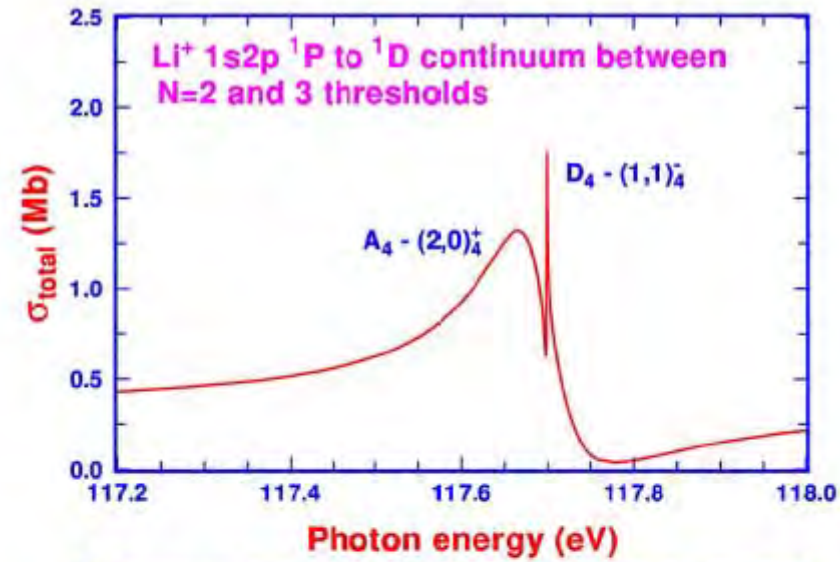


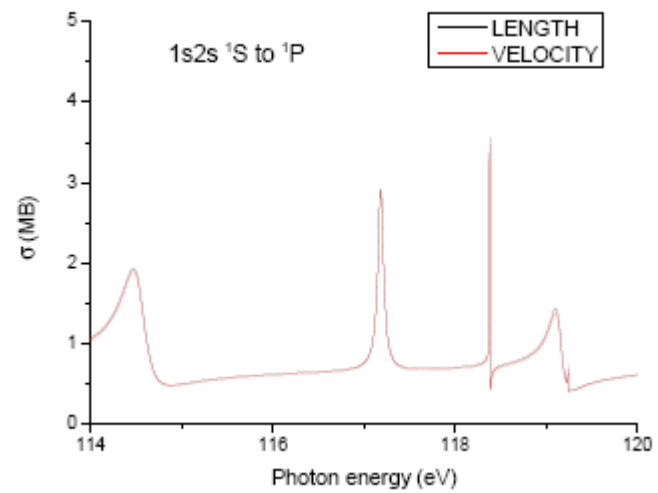
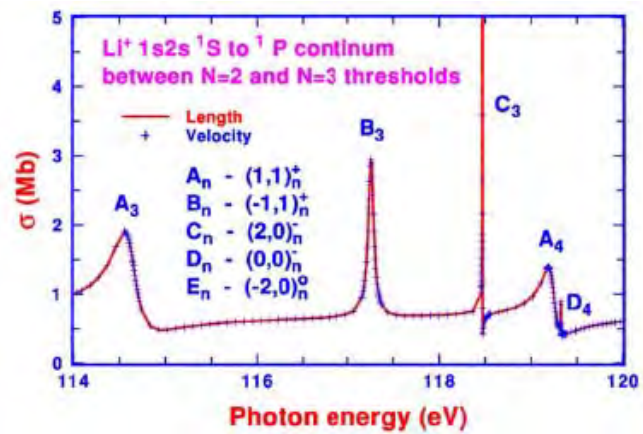


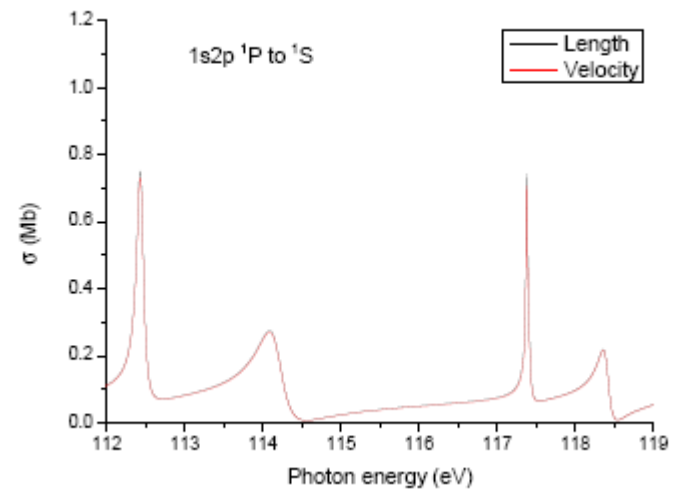
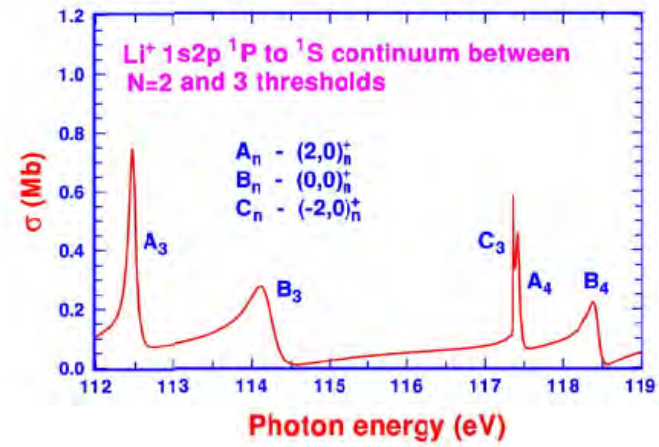


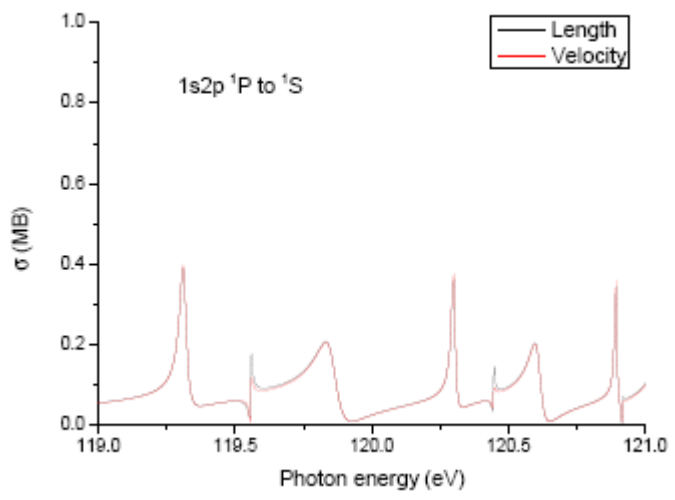
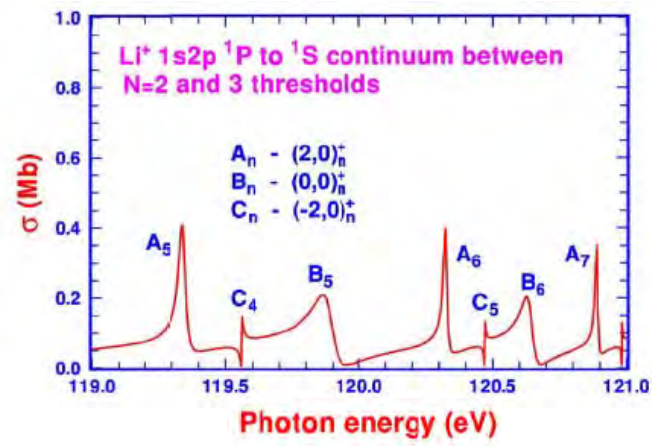


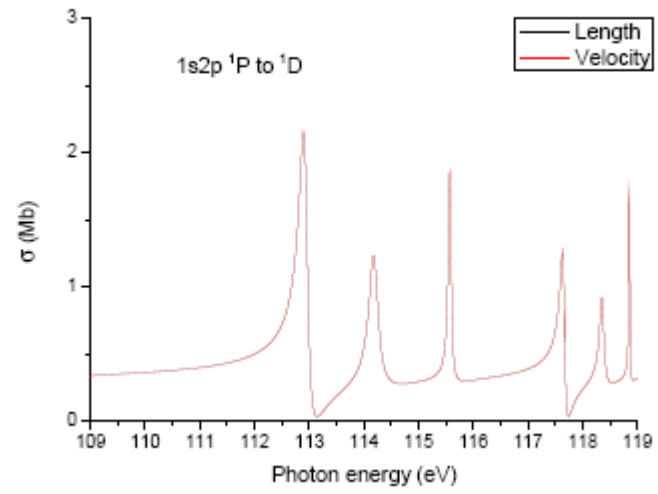
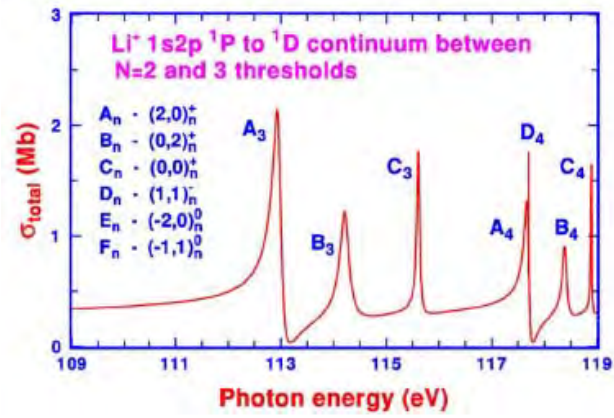


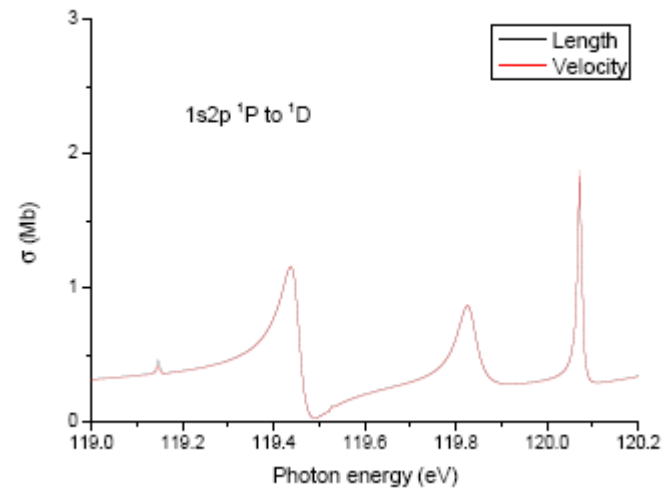
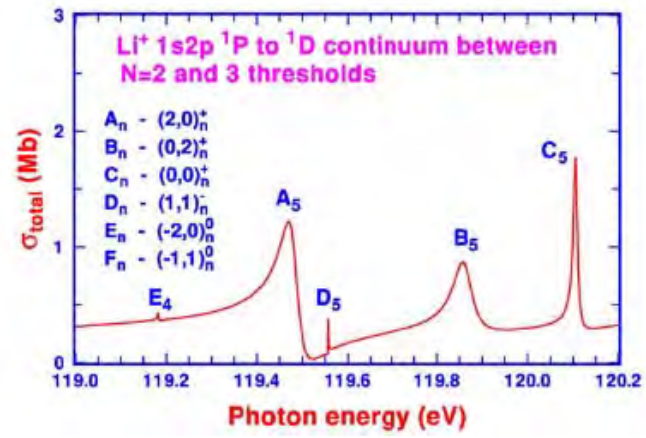












Absolute Cross Sections for the Photoionization of the $6s6p\ ^1P$ Excited State of Barium

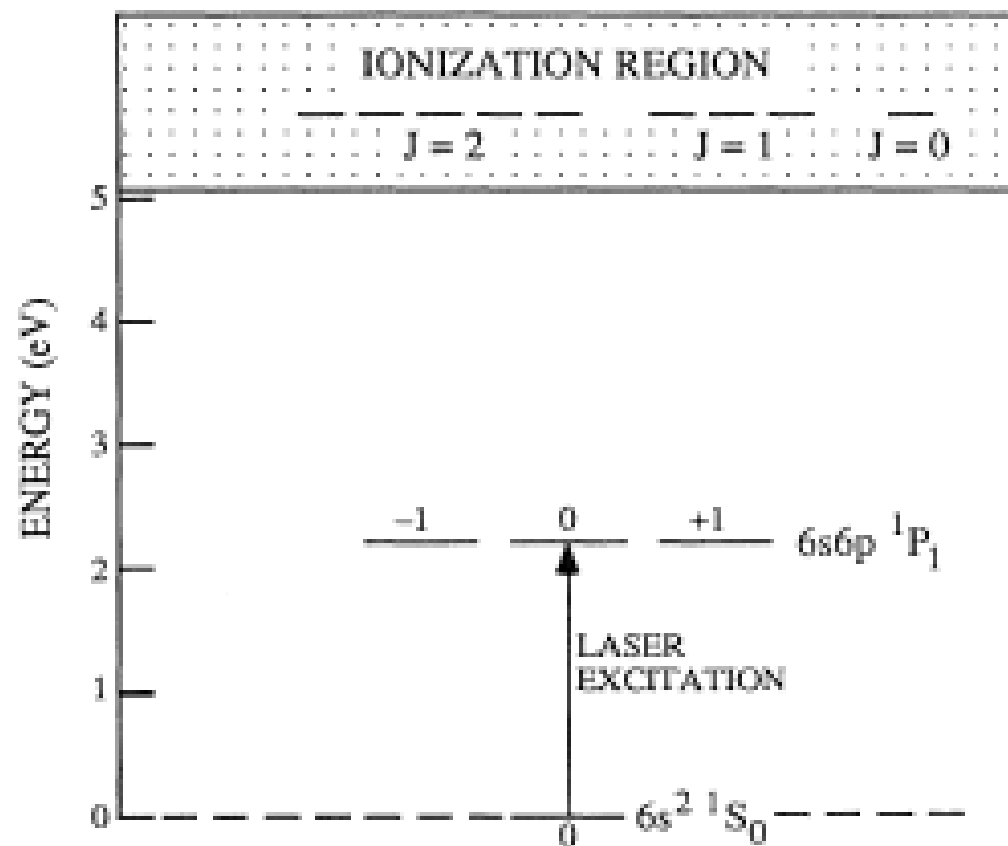
L.-W. He, C. E. Burkhardt, M. Ciocca, and J. J. Leventhal

Department of Physics, University of Missouri-St. Louis, St. Louis, Missouri 63121

S. T. Manson

Department of Physics and Astronomy, Georgia State University, Atlanta, Georgia 30303

(Received 8 July 1991)



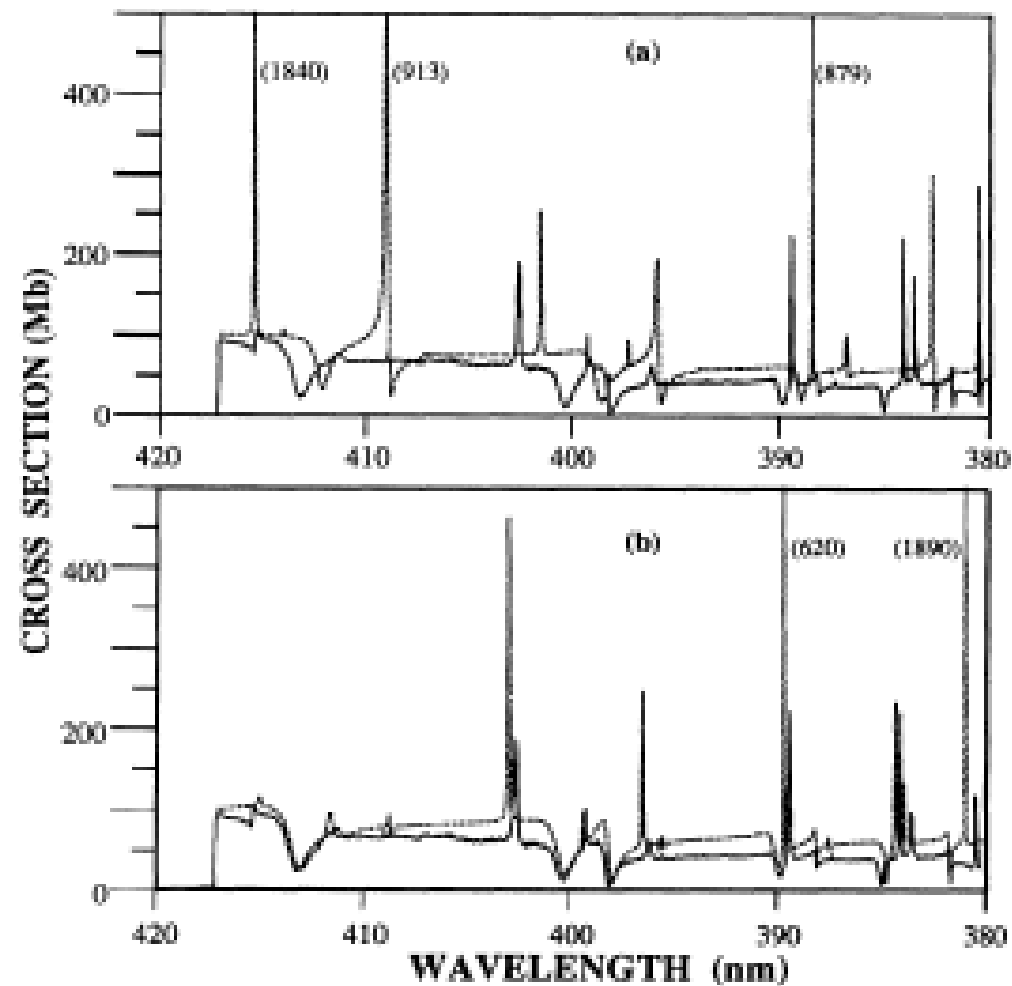
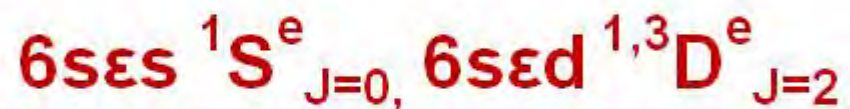


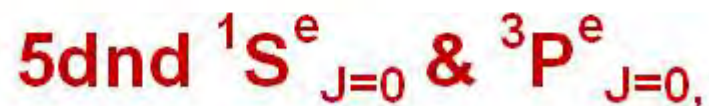
FIG. 2. Photoionization cross section of Ba $6s6p\ ^1P$ state near threshold. The solid lines are the data reported in this work. The dashed lines represent the calculations of (a) Ref. [10] and (b) Ref. [11].



Open chs:



Closed chs:



Atomic photoionization in changing plasma environment

$$H(r_1, r_2, \dots; D) = \sum_i h_0(r_i; D) + \sum_{i>j} e^2 / r_{ij},$$

where $h_0(r; D) = (p^2/2m) - (Z e^2 / r) e^{-r/D}$,

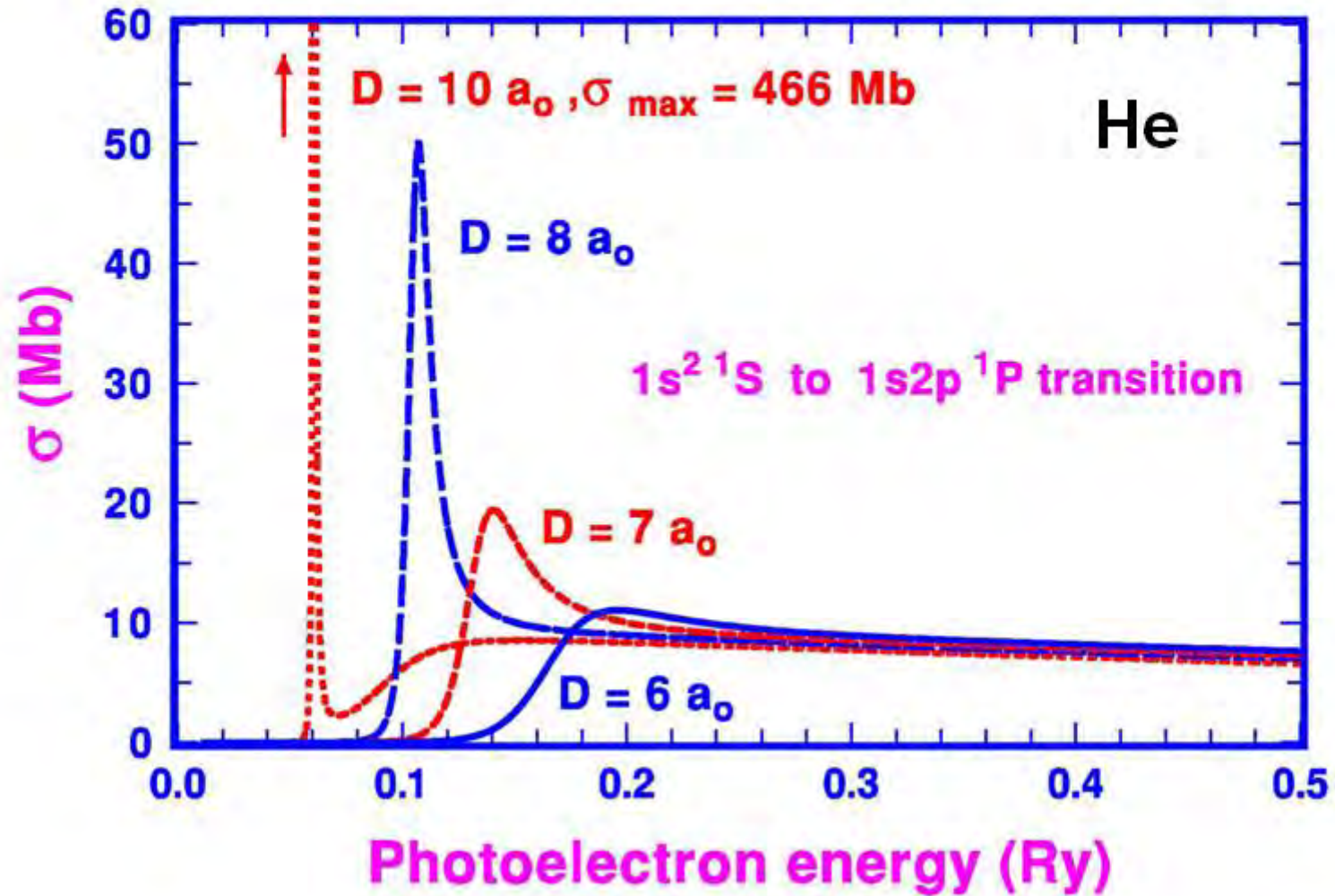
p : electron momentum, D : Debye length given by

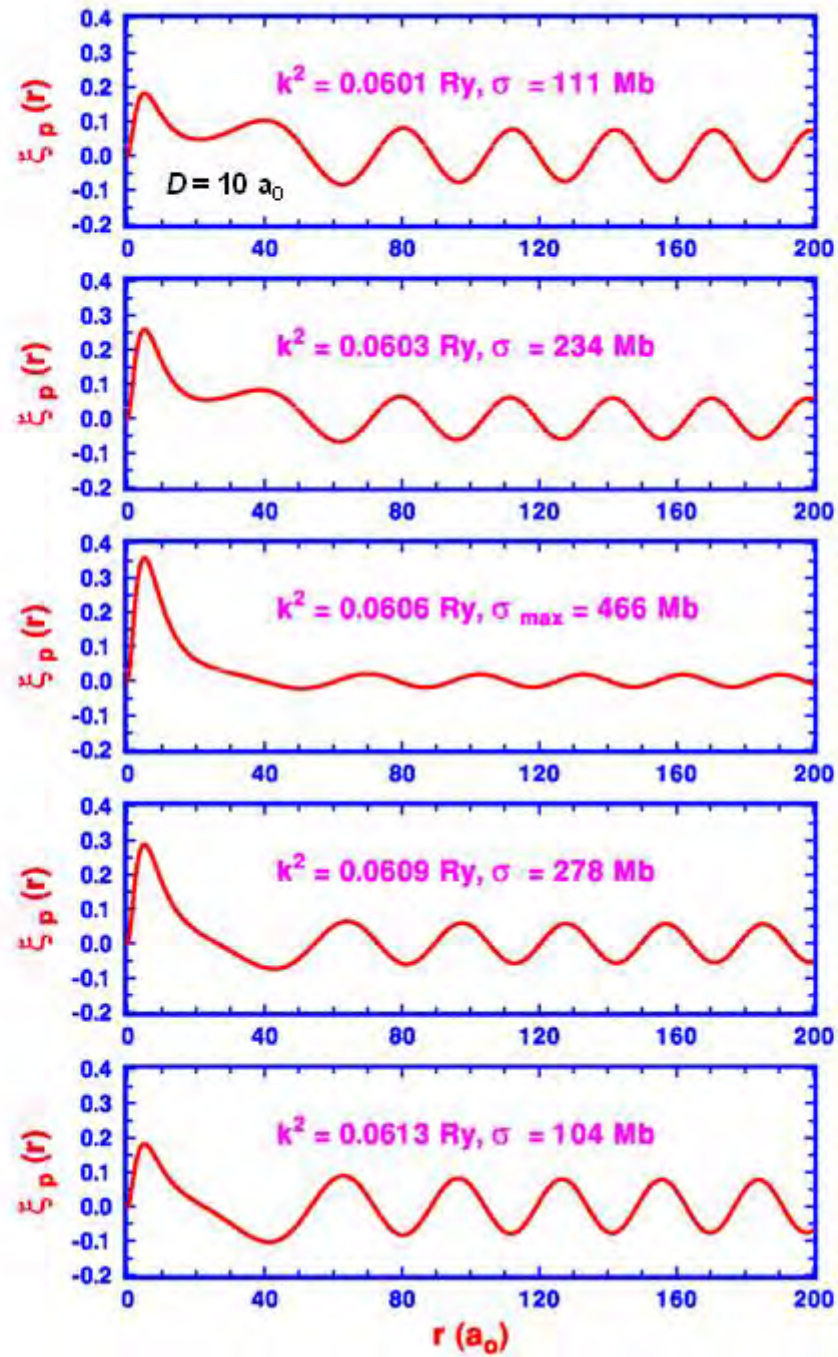
$$D = 1.304 \times 10^9 (T/n)^{1/2}, \text{ (in unit of } a_0 \text{)}$$

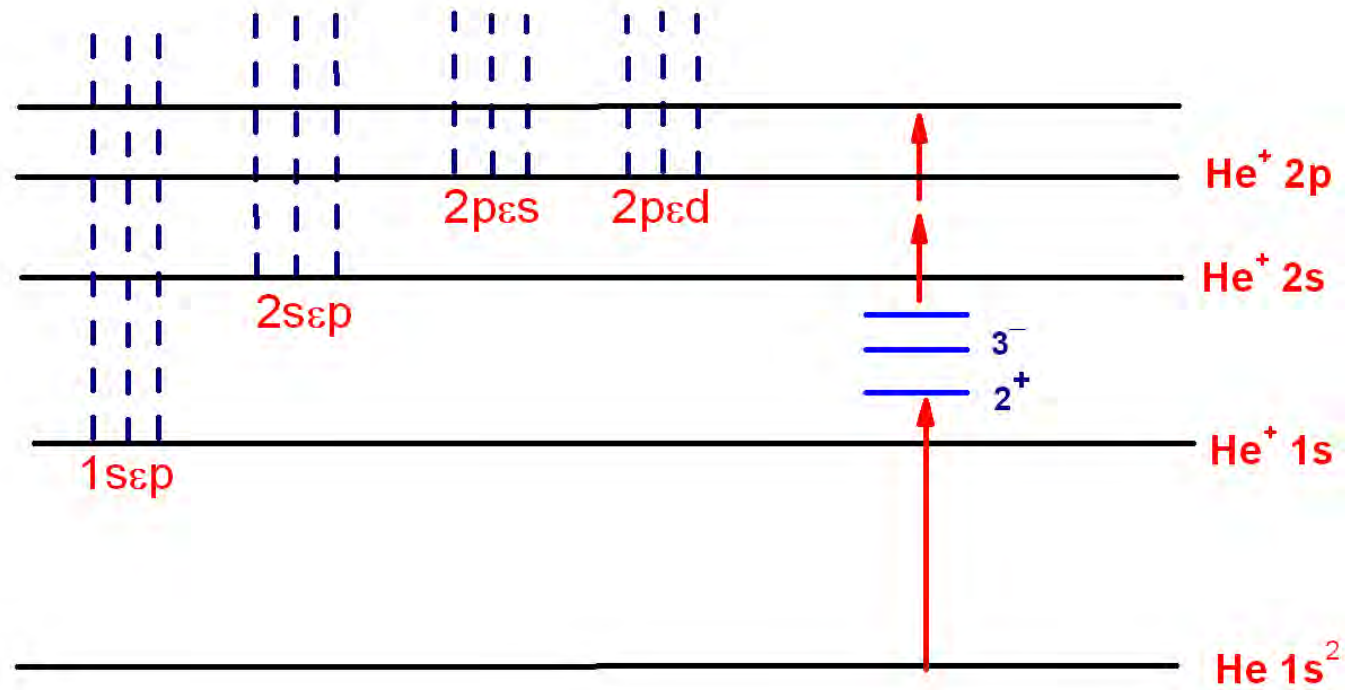
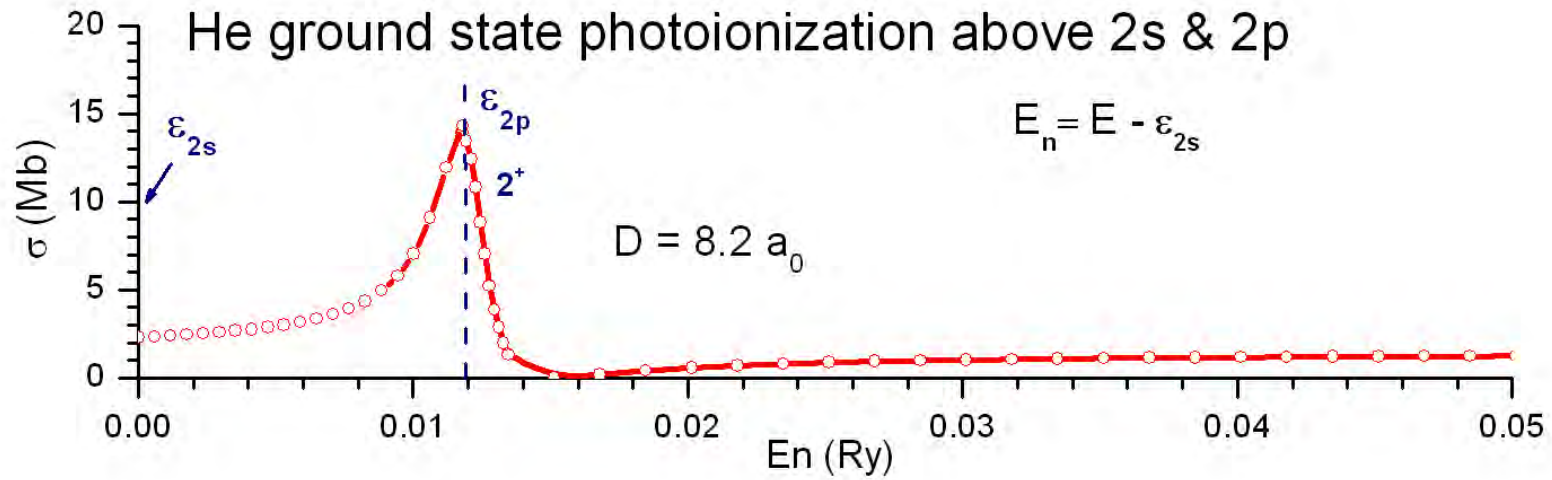
T : plasma temperature (in degree Kelvin)

n : density (in cm^{-3})

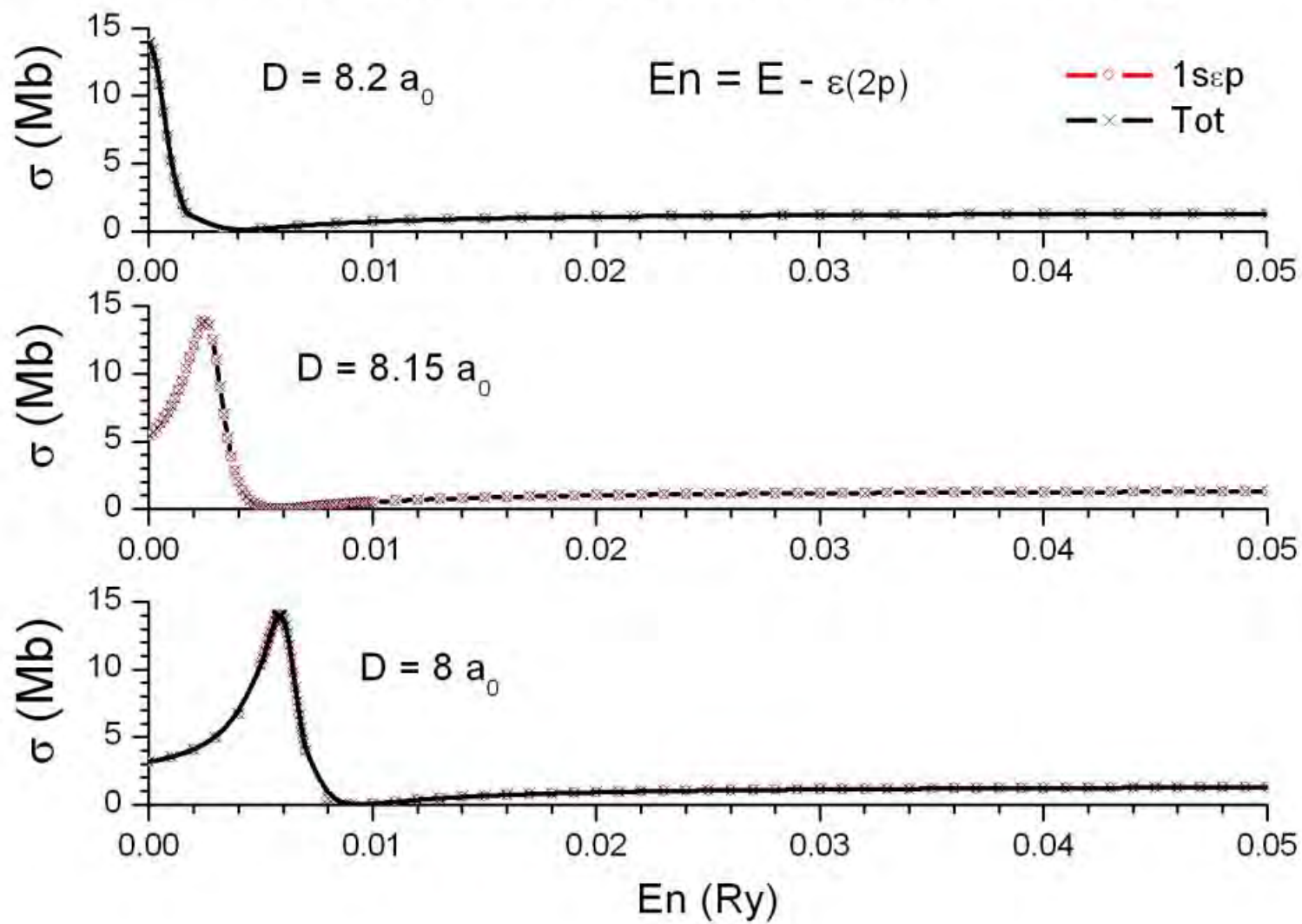
Phys. Rev. A88,023406 (2013)



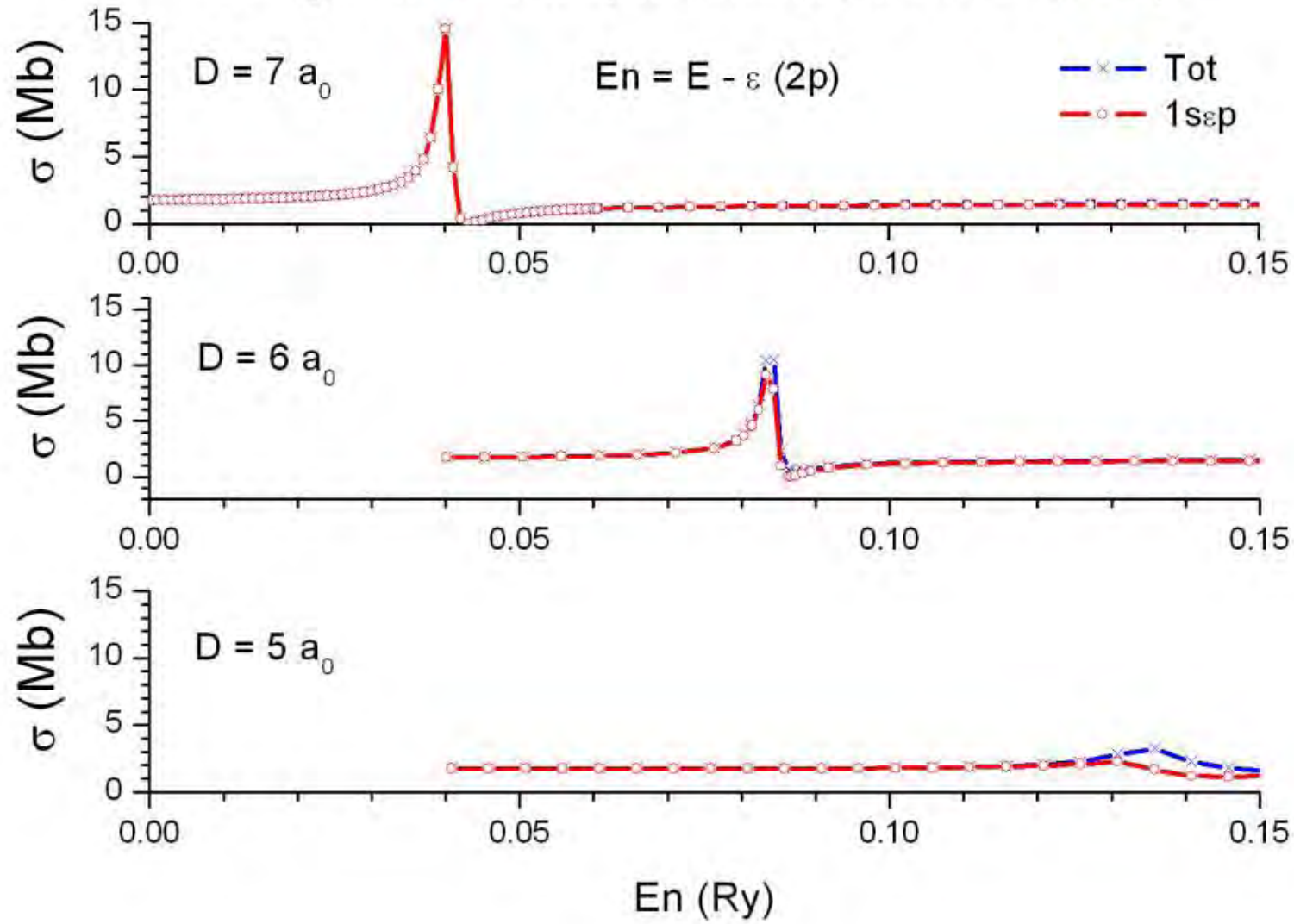




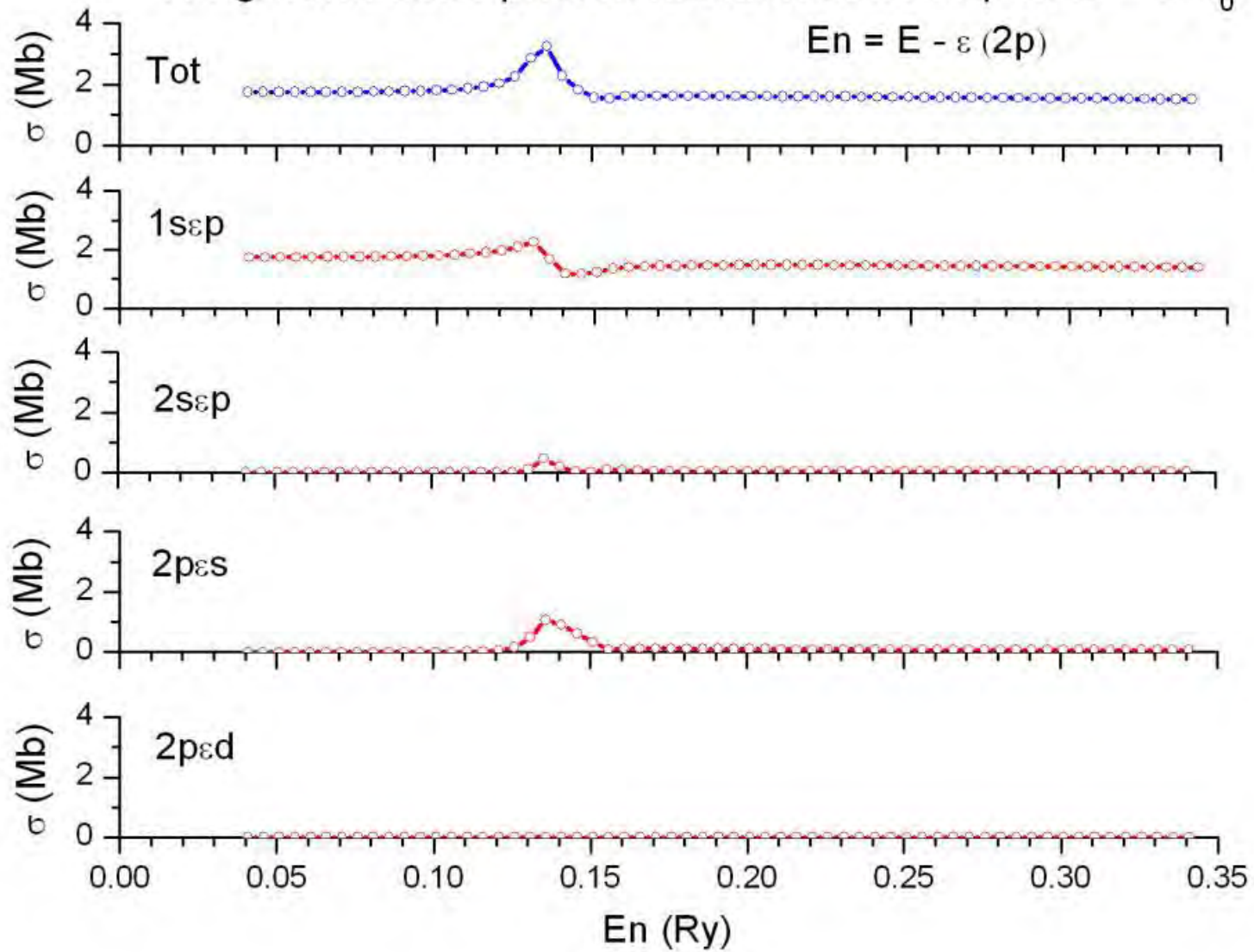
He ground state photoionization above 2p



He ground state photoionization above 2p



He ground state photoionization above 2p at $D = 5 a_0$



Triply excited lithium: Phys. Rev. A 63, 020702 (2001).

