

Efficient Quantum Memory for Photonic Polarization Qubits

Ying-Cheng Chen

Institute of Atomic and Molecular Sciences

Academia Sinica, Taiwan

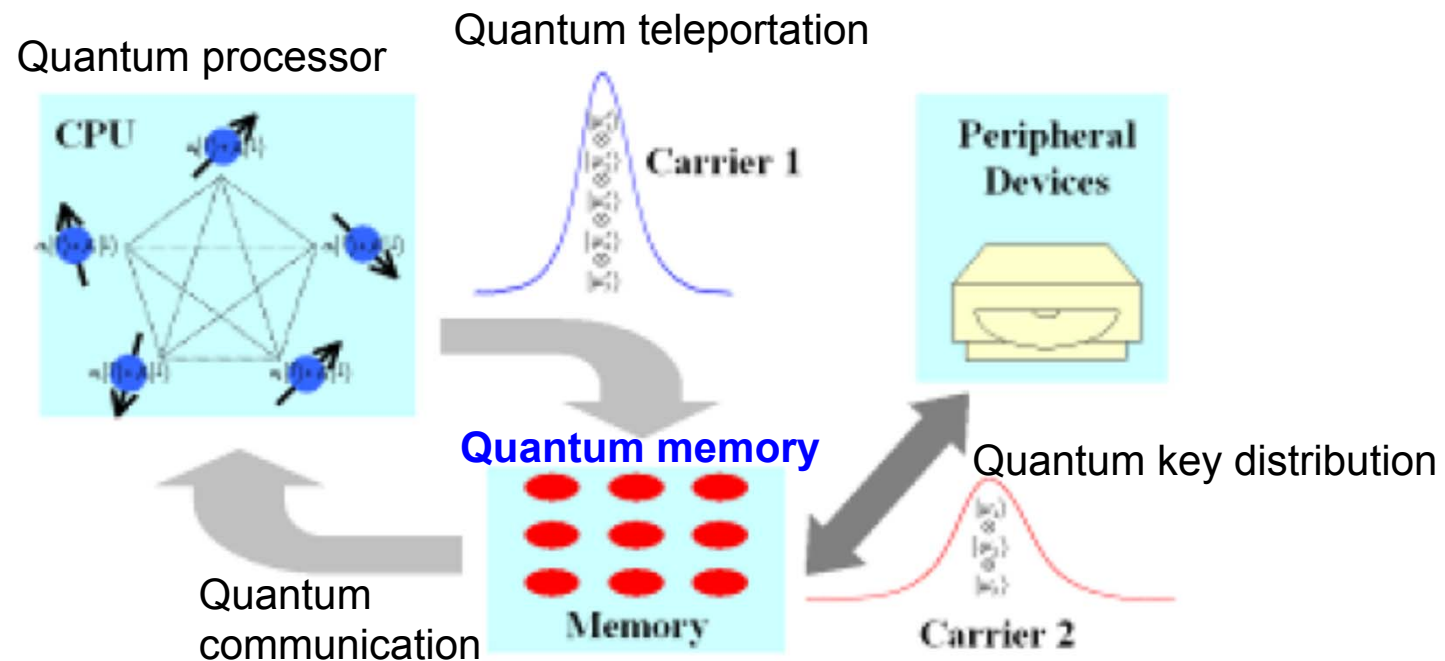
*NCTS Annual Theory Meeting, Quantum Physics,
Quantum Information & Quantum Technologies,*

Feb. 17-19, 2021



Quantum 2.0: Turning quantum weirdness into use

- Use of the fascinating properties of quantum mechanics, such as uncertainty principle, coherent superposition, quantum entanglement... as **resources** for the **information** technology.
- Spirit of quantum 2.0: Pursuit for perfection.



Outline

- Introduction to optical quantum memory ?
- Highly-efficient coherent optical memory based on electromagnetically induced transparency (EIT) [1]
- Towards broadband EIT-based memory[2]
- Development of a quantum light source based on cavity-enhanced spontaneous parametric down-conversion (SPDC)[3]
- Efficient quantum memory for heralded single photons & polarization qubitd[4,5]
- Conclusions

1, Y.-F. Hsiao,..., YCC*, Phys. Rev. Lett. 120, 183602(2018)

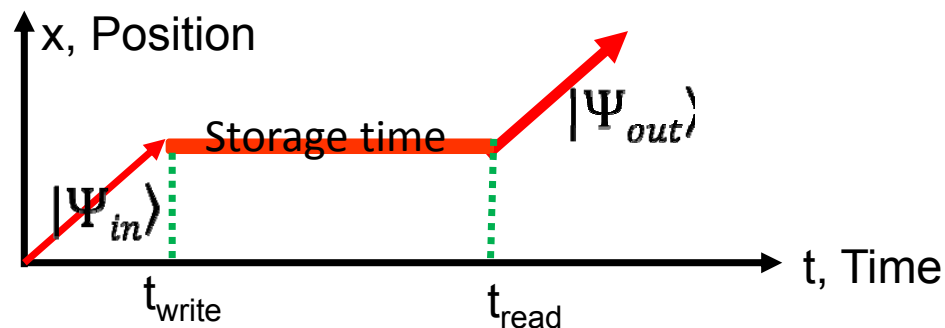
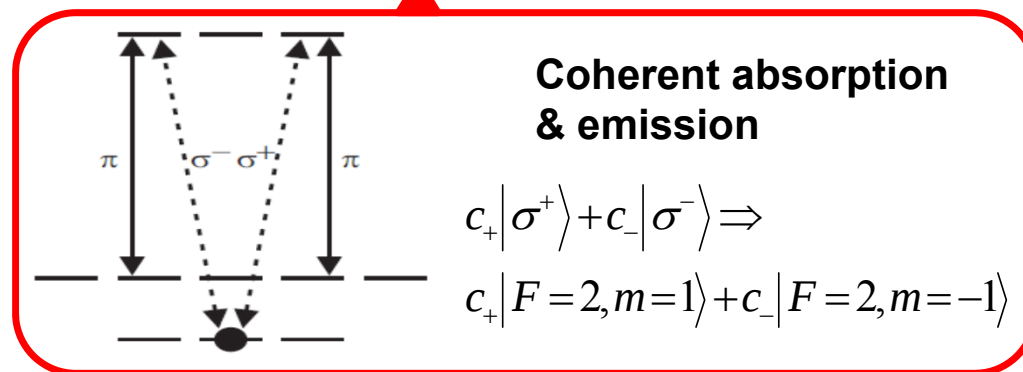
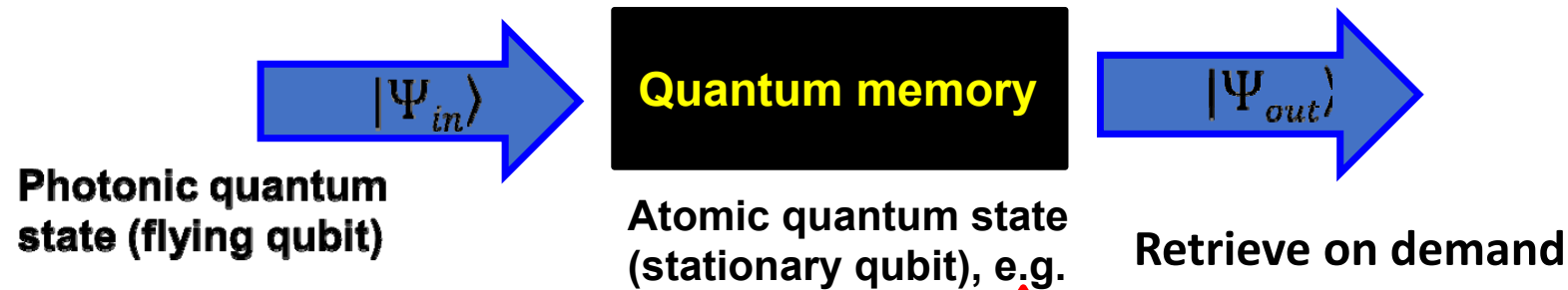
2, Y.-C. Wei,..., YCC*, Phys. Rev. A. 102, 063720(2020)

3, P.-J. Tsai, YCC*, Quantum Sci & Techno, 3, 034005(2018)

4, P.-J. Tsai,..., YCC*, Phys. Rev. Reseach, 2, 033155(2020)

5, Y.-C. Tseng, ..., YCC*, arXiv:2011.14948

What is optical quantum memory ?



QM postpone the collapse of a wavefunction.

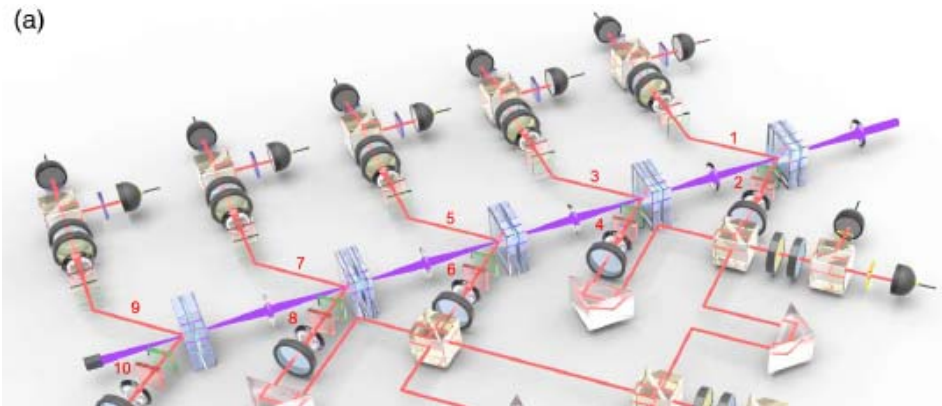
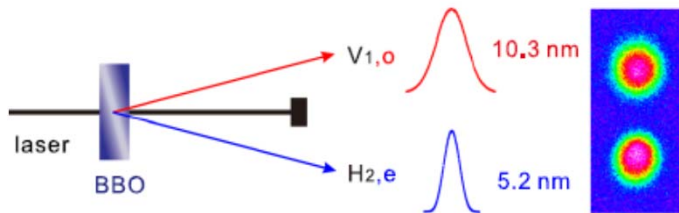
Why optical quantum memory (QM) 1/2 ?

- Synchronization of probabilistic events: QM-assisted multi-photon generation

Spontaneous parametric downconversion (SPDC)

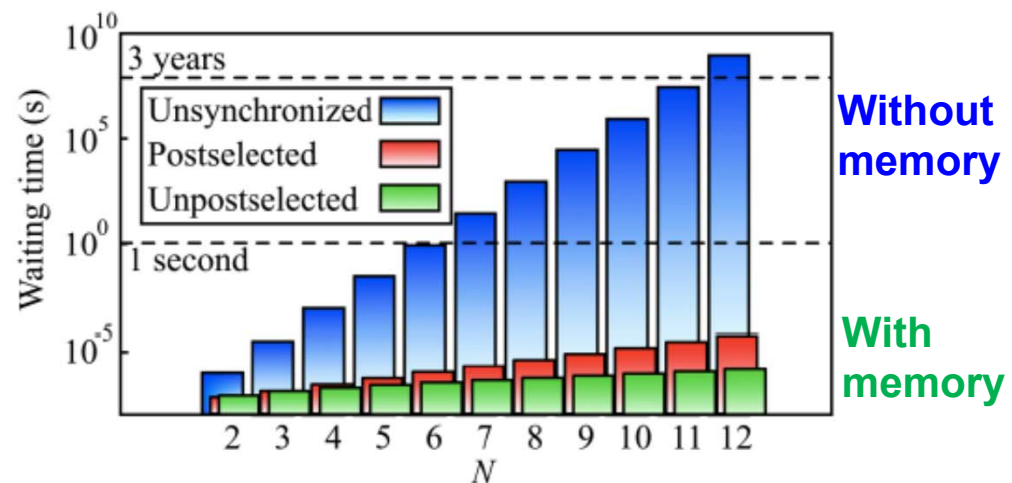
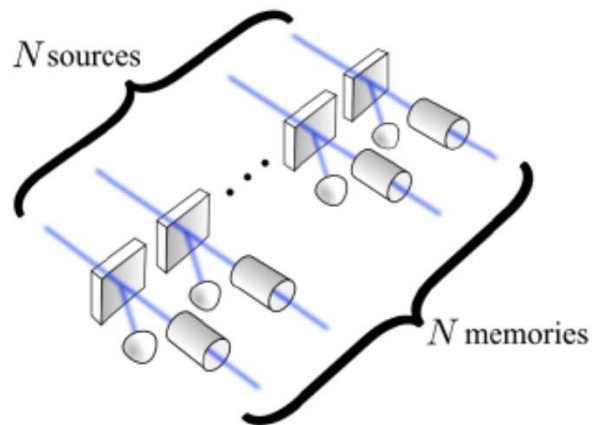
Entangled 10-photon, ~1 event/hour

X.L. Wang et al PRL 117, 210502(2016)



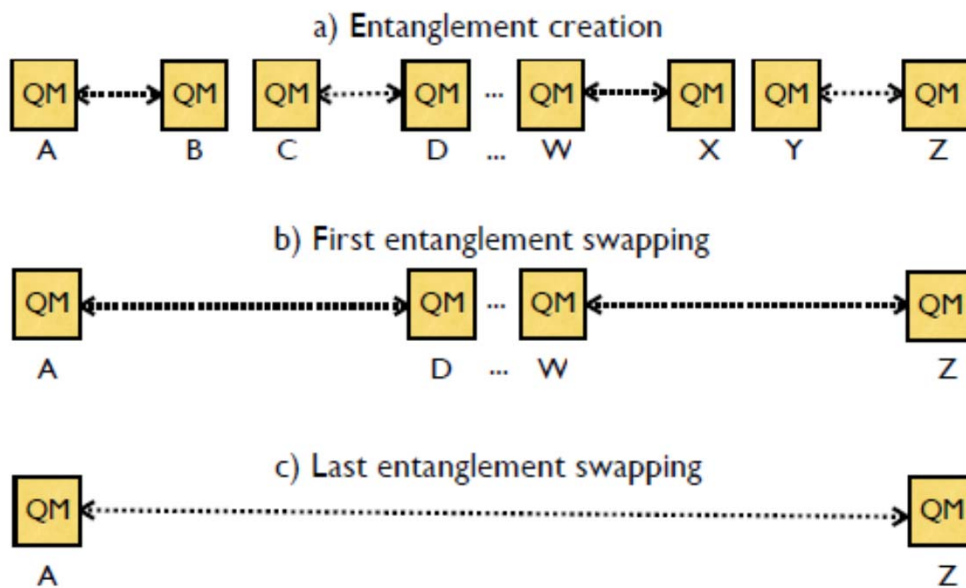
Enhancing multiphoton rate by QM

J. Nunn et. al. PRL 110, 133601(2013)

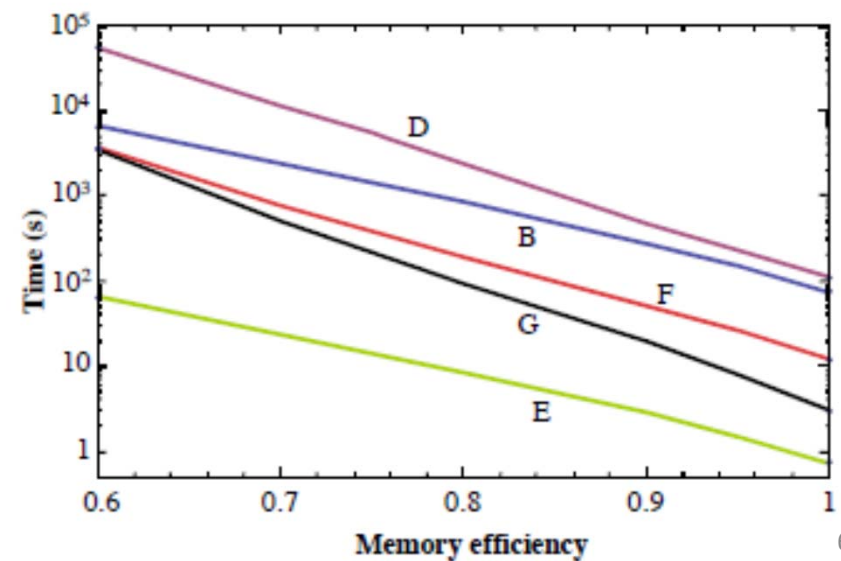


Why optical quantum memory (2/2) ?

- **Photon loss** is a problem in quantum key distribution (**no-cloning theorem**): 0.2 dB/km fiber loss and 10GHz photon rate $\rightarrow$ 1 Hz detection rate for 500 km and 10^{-10} Hz for 1000 km (**1 photon/300 years**)!
- High-performance quantum memory a key component for **quantum repeater** protocol which allows the **wait-until-success** strategy !



- An increase of 1% in SE decrease the entanglement distribution time by 7-18%, depending on the protocol.



N. Sangouard et al, Rev. Mod. Phys. 83, 33(2011)

Performance Parameters for a Quantum Memory

Fidelity: How well the quantum state remains?

$$F_{|\Psi_{\text{in}}\rangle} = \inf_{|\Psi_{\text{in}}\rangle} \sqrt{\langle \Psi_{\text{in}} | \rho_{\text{out}} | \Psi_{\text{in}} \rangle}.$$

Noise level : noise reduces the non-classical feature

Storage time : How long the quantum state can be stored ?

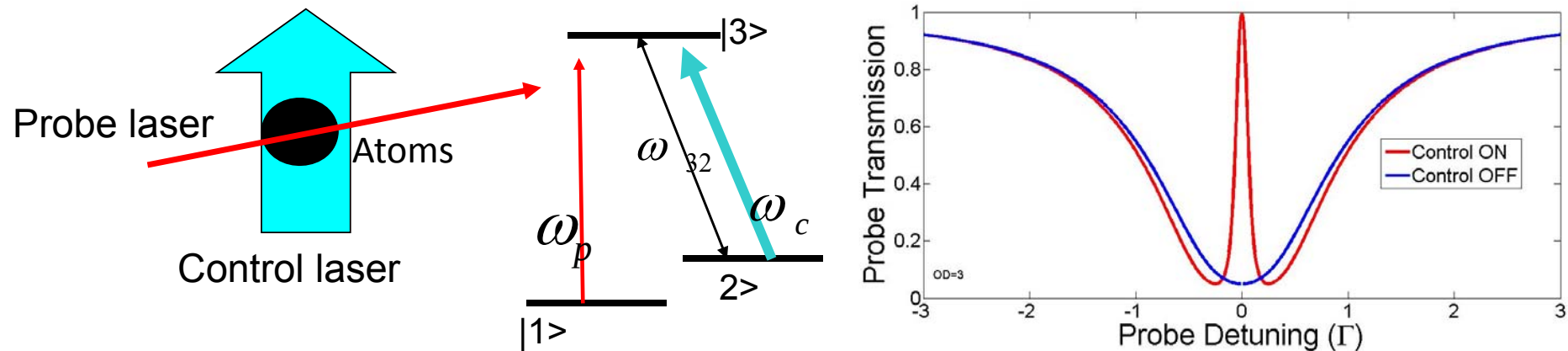
Bandwidth (capacity) : How many states can one store?

Our focus

Efficiency : What percentage of photon left after the storage?

$$SE = \frac{N_{out}}{N_{in}}$$

Electromagnetically induced transparency (EIT)

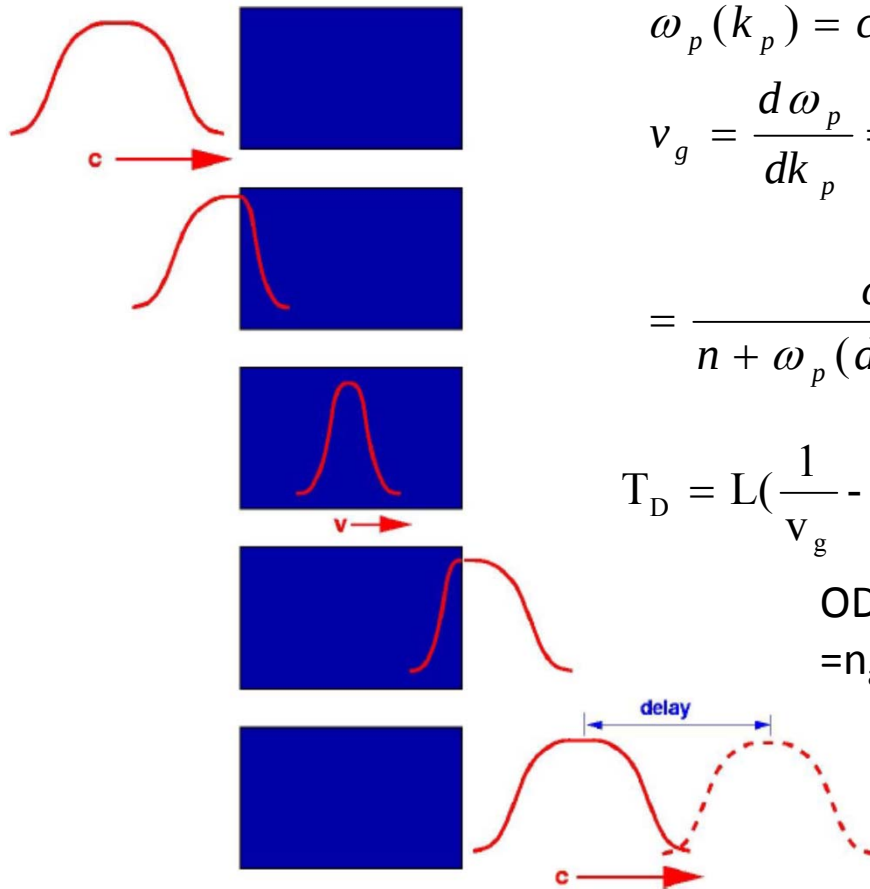


- A **dark state** of the Hamiltonian containing atom-photon interactions and decouples from the excited state.

$$H = \begin{pmatrix} 0 & 0 & -\frac{\hbar\Omega_p^*}{2} \\ 0 & -\hbar\delta_2 & -\frac{\hbar\Omega_c^*}{2} \\ -\frac{\hbar\Omega_p}{2} & -\frac{\hbar\Omega_c}{2} & -\hbar\delta_p \end{pmatrix}, |D\rangle = \frac{\Omega_c}{\sqrt{\Omega_c^2 + \Omega_p^2}}|1\rangle - \frac{\Omega_p}{\sqrt{\Omega_c^2 + \Omega_p^2}}|2\rangle$$

M. Fleischhauer et al. Rev. Mod. Phys. 77, 633, 2005 8

EIT and Slow Light



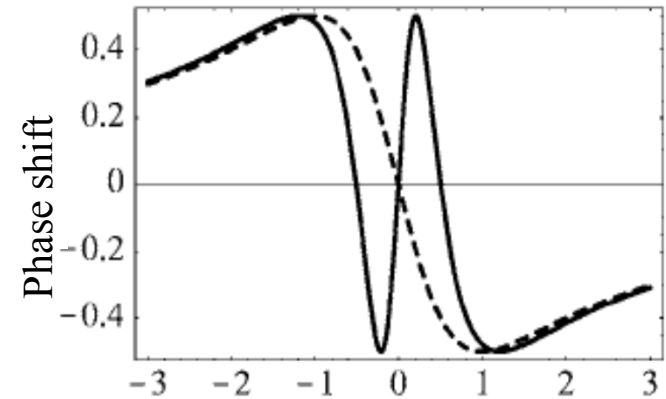
$$\omega_p(k_p) = ck_p / n(\omega_p)$$

$$v_g = \frac{d\omega_p}{dk_p} = \frac{1}{dk_p / d\omega_p}$$

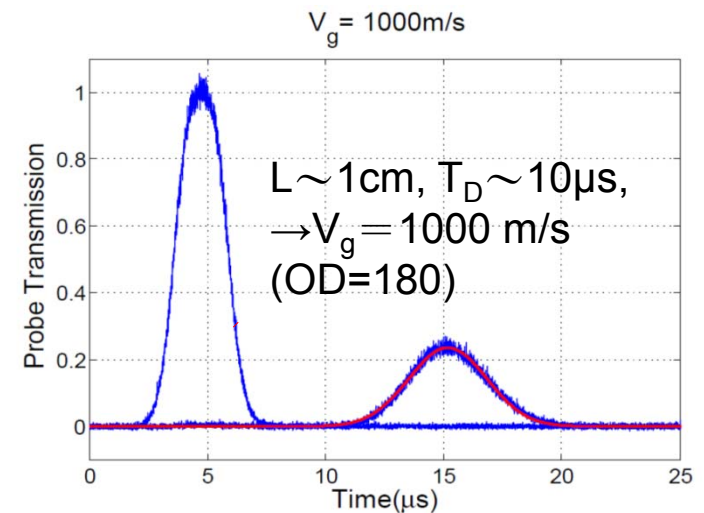
$$= \frac{c}{n + \omega_p (dn / d\omega_p)}$$

$$T_D = L \left(\frac{1}{v_g} - \frac{1}{c} \right) \cong \frac{(OD) \Gamma}{\Omega_c^2}$$

OD= optical depth
 $= n_a \sigma_{abs} L$



Detuning (Γ)

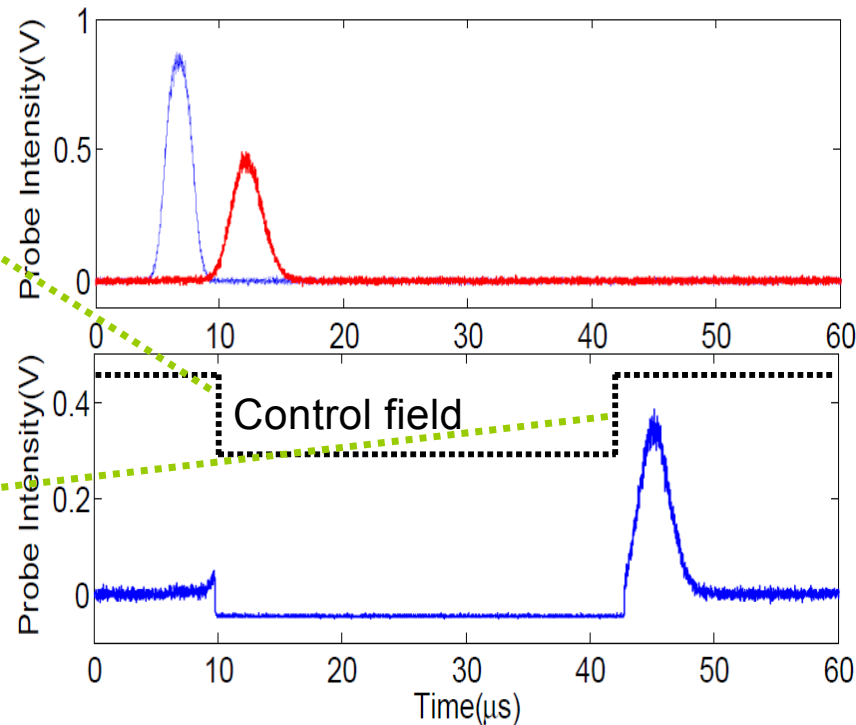
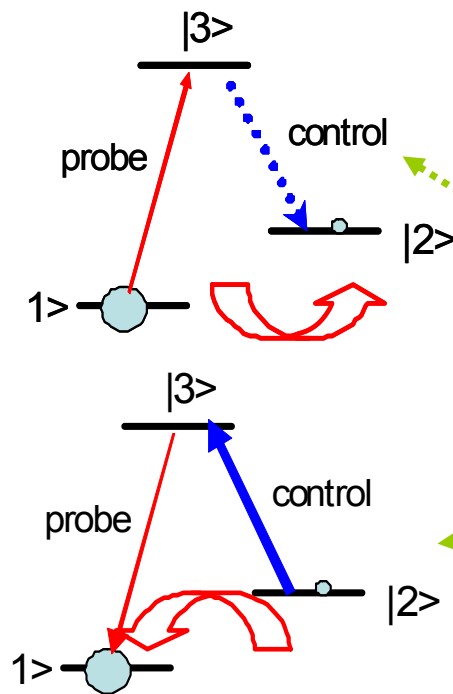


Dark-state polariton and Optical quantum memory

- Slow light can be viewed as a dark-state polariton, which is a coherent superposition of light field and collective atomic coherence and tuned by the control field.

$$\Psi(z, t) = \cos \theta E_p(z, t) - \sin \theta \sqrt{N} \rho_{21}(z, t) e^{i\Delta k z}$$

$$\tan \theta = N \frac{3\lambda^2}{2\pi} c \frac{\Gamma_{31}}{\Omega_c^2}; \left[\frac{\partial}{\partial t} + c \cos^2 \theta \frac{\partial}{\partial z} \right] \Psi(z, t) = 0; v_g = c \cos^2 \theta$$



M Fleischhauer & M. Lukin, PRL84,5092(2000)

EIT quantum memory for high efficiency (1/2)

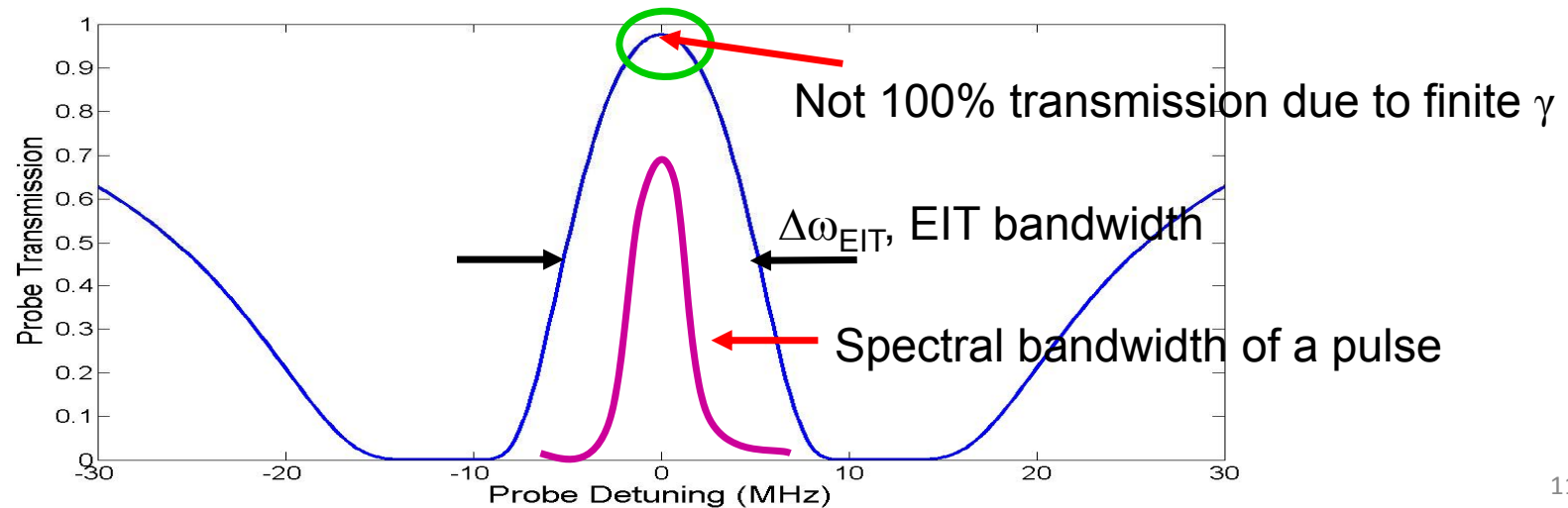
- By solving the Maxwell-Bloch equation, one gets

$$\eta_{\text{trans}} = \frac{\exp(-2\gamma T_D)}{\sqrt{1 + \left(\frac{4\ln 2}{T_p \Delta\omega_{\text{EIT}}}\right)^2}} = \frac{\exp(-2\gamma T_D)}{\sqrt{1 + \frac{16\ln 2}{OD} \left(\frac{T_D}{T_p}\right)^2}} \quad T_D = \frac{(OD)\Gamma}{\Omega_c^2}$$

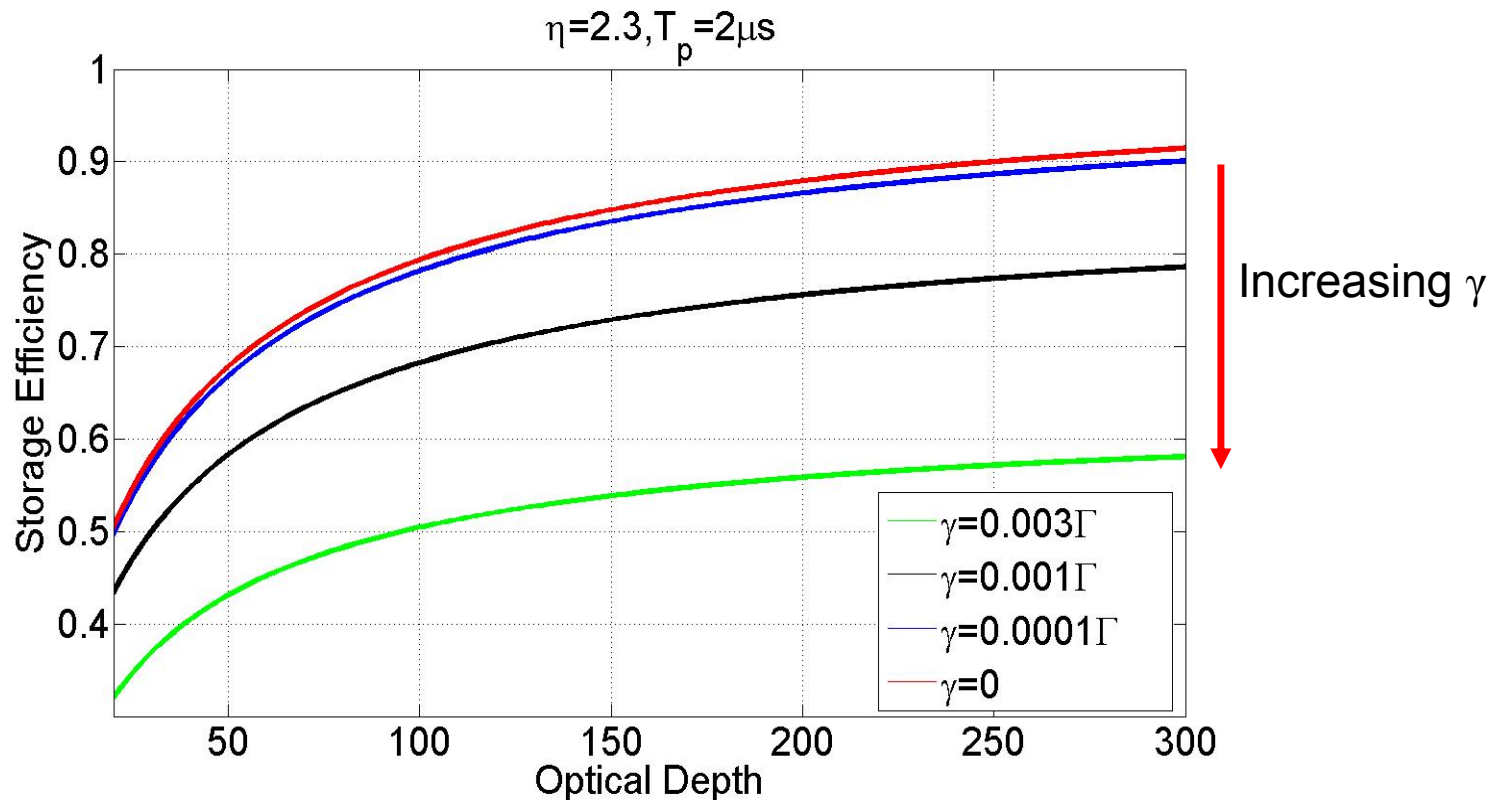
Loss due to finite ground-state decoherence rate γ

Loss due to finite EIT bandwidth

Keep $T_D/T_p \sim 2.5$ in order to compress > 99% of pulse into media (minimize leakage loss)

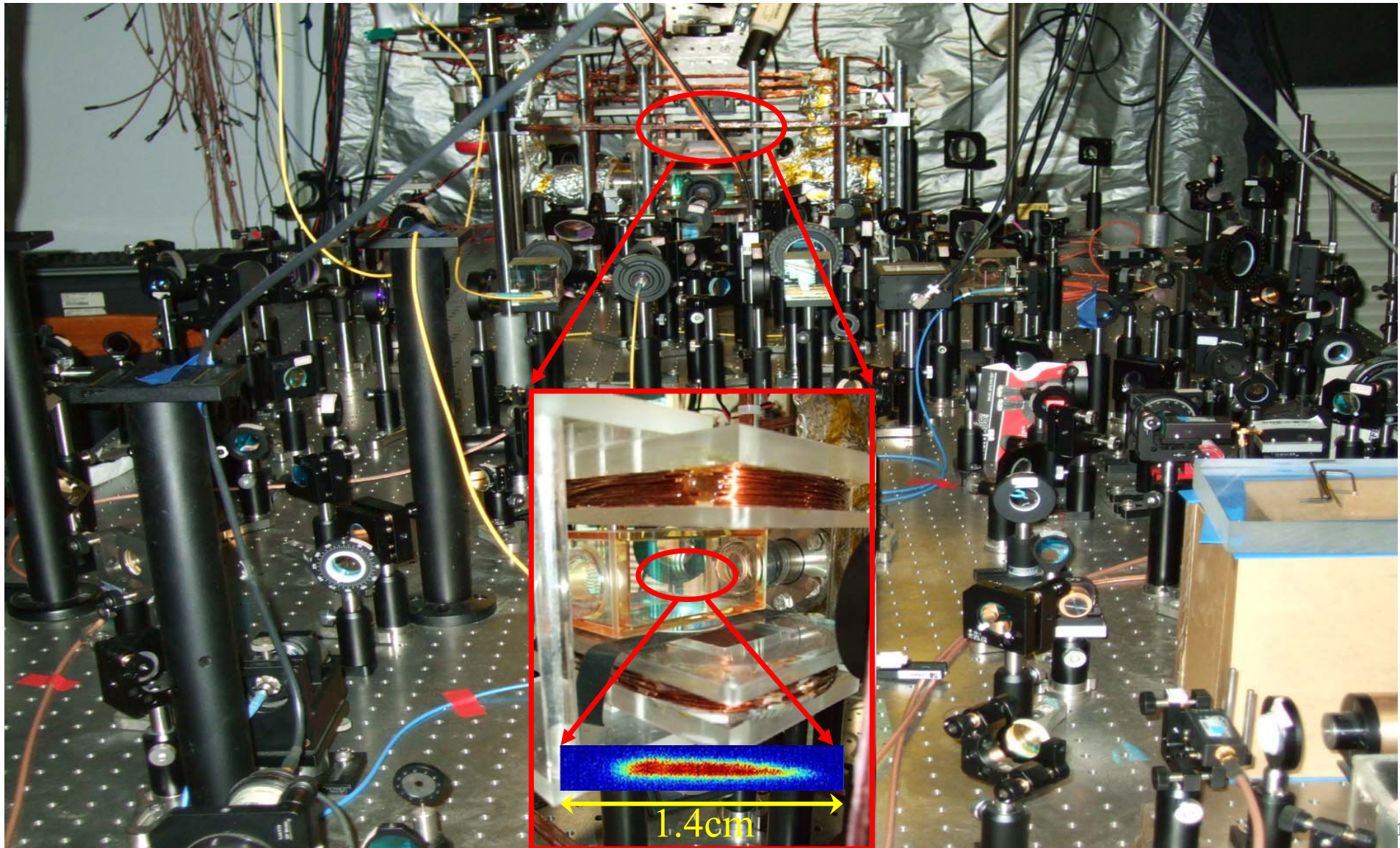


EIT quantum memory for high efficiency (2/2)



- **High optical depth** and **low ground-state decoherence rate** are two keys to achieve a high storage efficiency !
- Technically, one also needs high control intensity when increasing the optical depth.

Setup: Cesium elongated magneto-optical trap

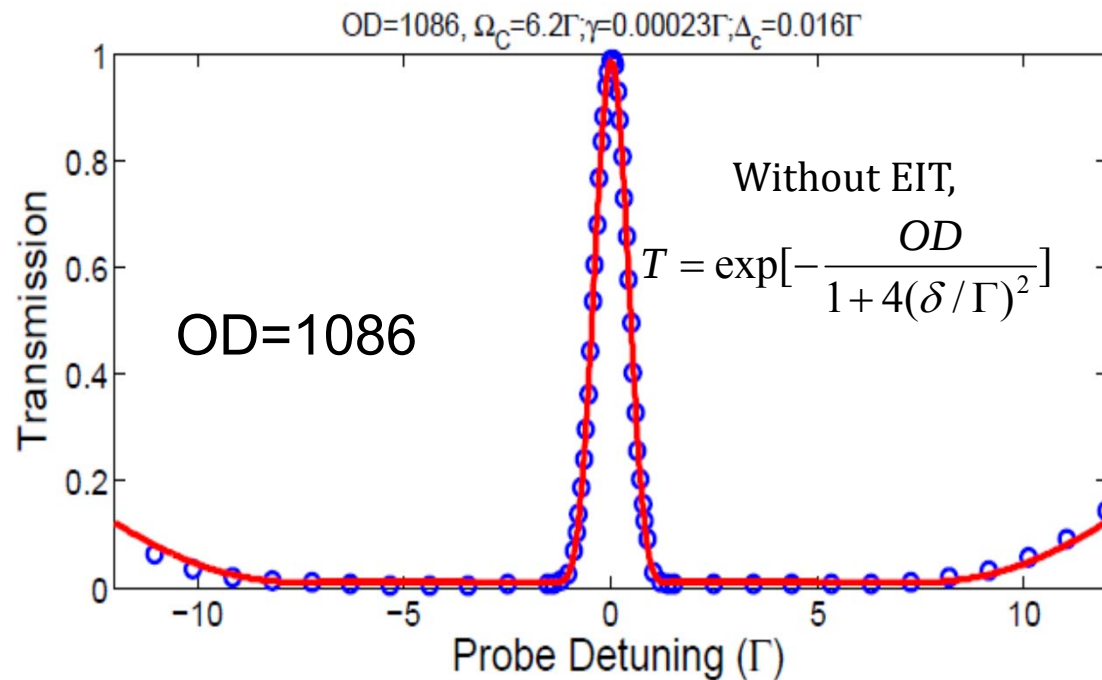


Achieve high-quality EIT atomic media

- Optical depth > 1000 and decoherence rate $\sim 10^{-4} \Gamma$ in a magneto-optical trap.

$$T = \exp(-1086) \approx 1.66 \times 10^{-477} \Rightarrow T \approx 1, \text{ due to EIT}$$

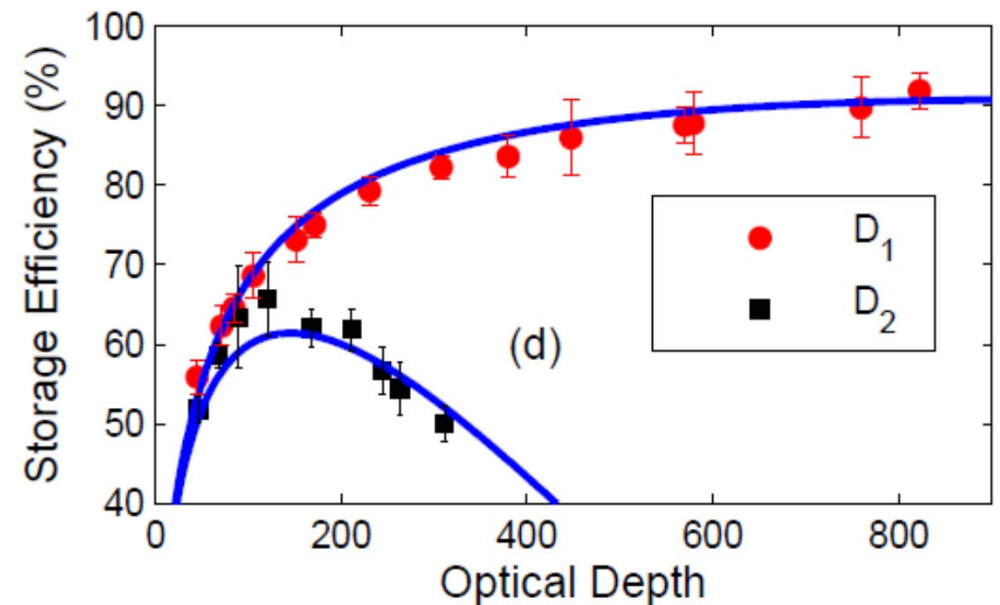
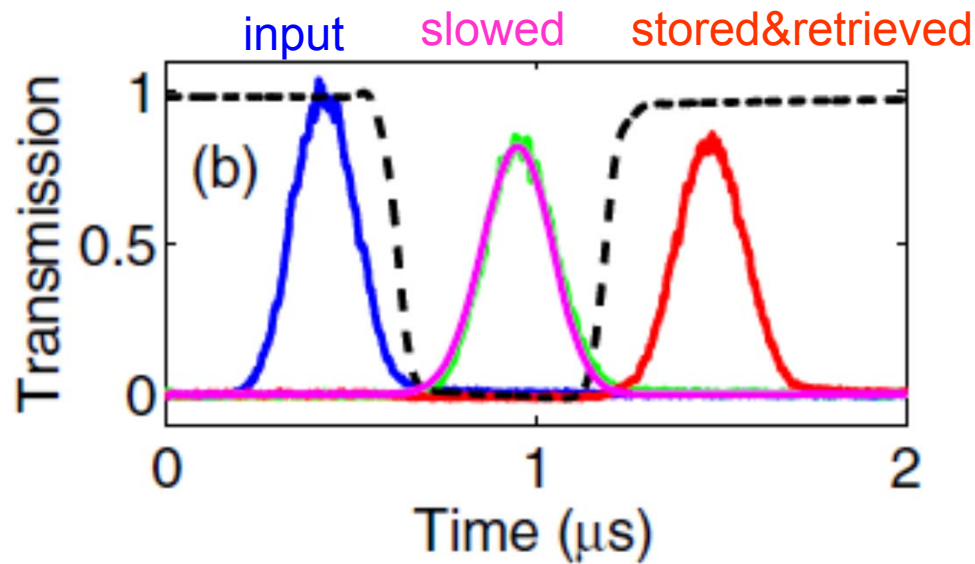
$$\gamma \approx 1 \times 10^{-4} \Gamma$$



Y.-F. Hsiao..., YCC*, Phys. Rev. A, 90, 055401(2014)

Record-high storage efficiency for EIT memory

- The highest record for storage efficiency (92%) for all kinds of optical memory.



Y.-F. Hsiao...YCC*, Phys Rev. Lett. 120, 183602(2018), cited 106 times @Feb. 17, 2021

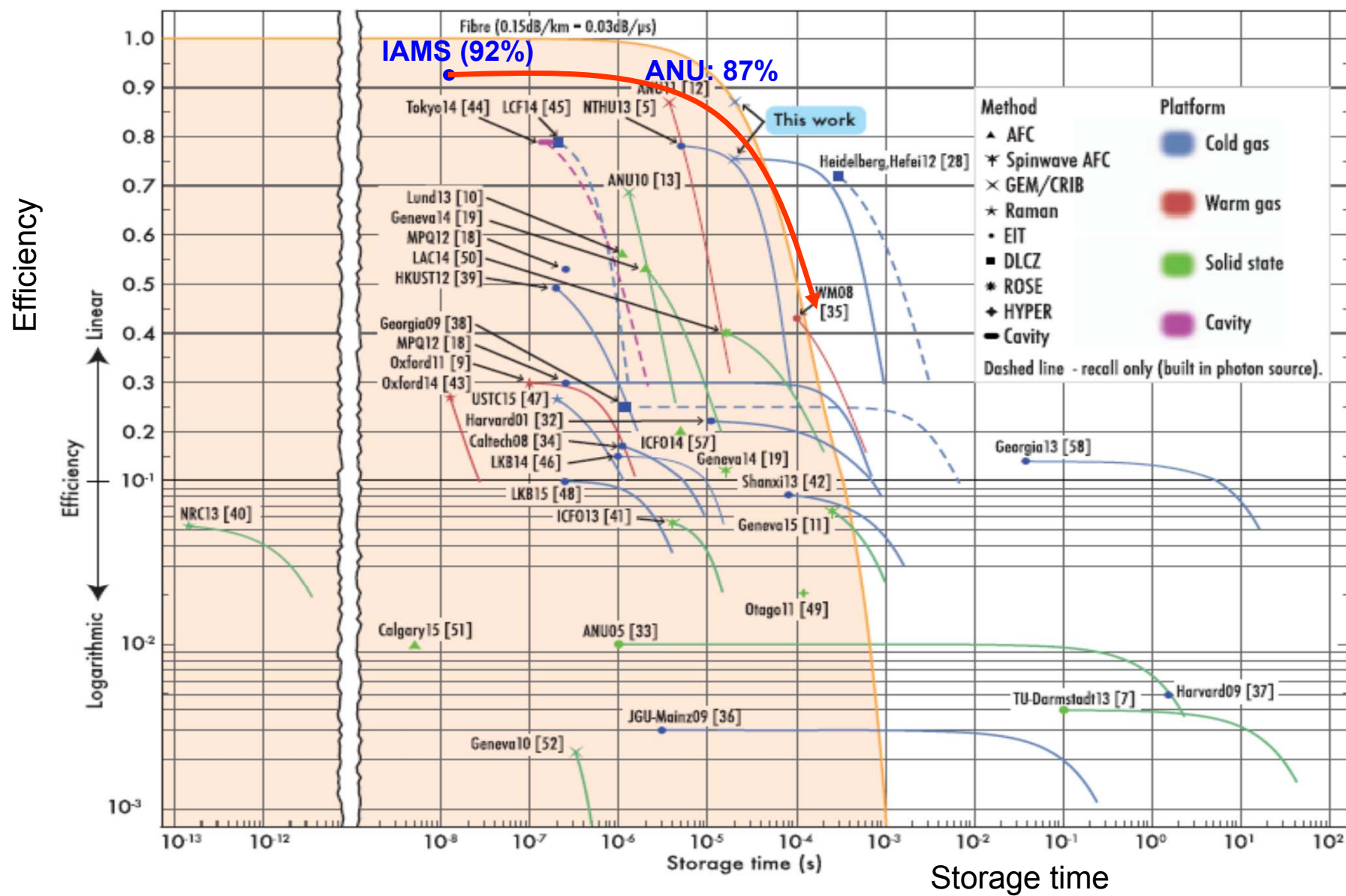
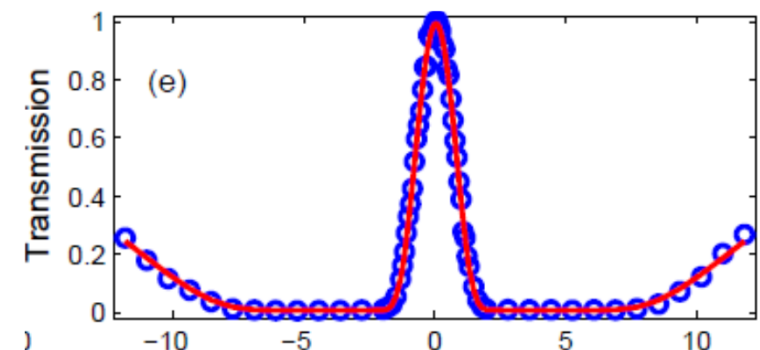
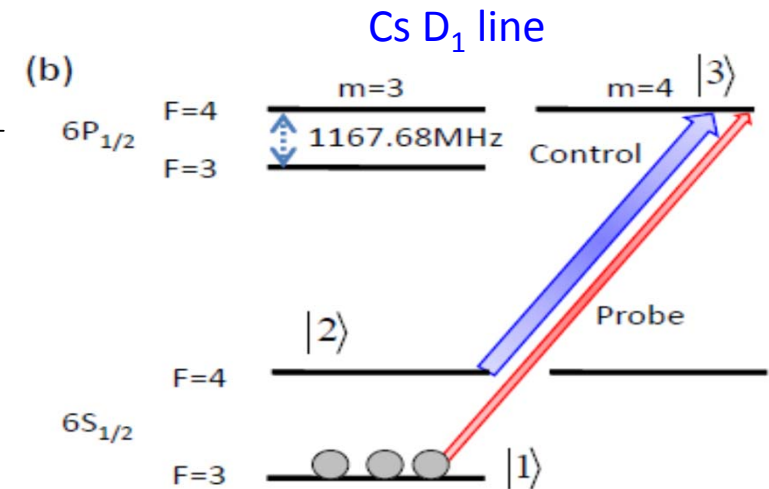
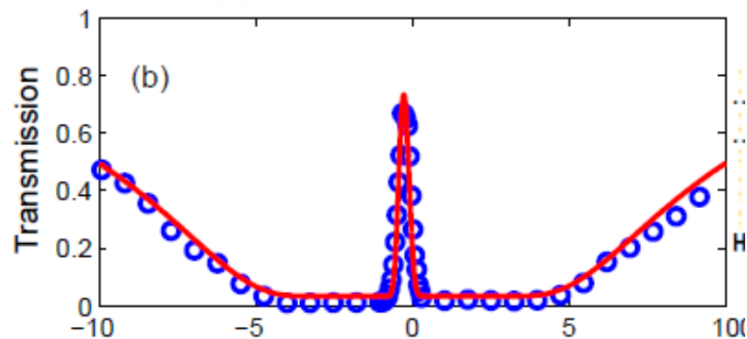
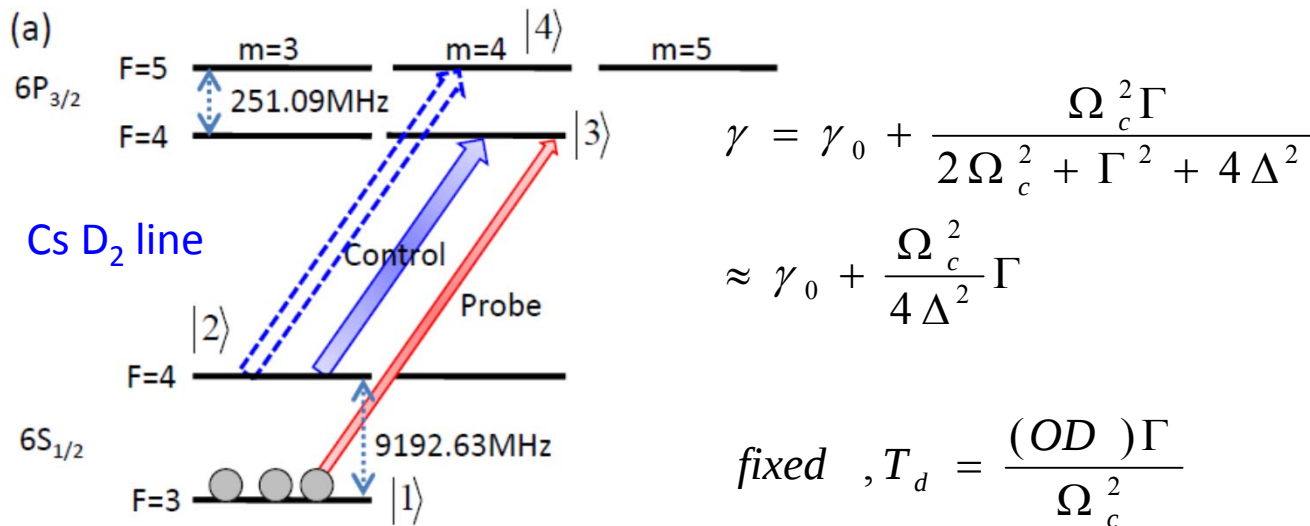


Figure from PK Lam et al. Optica, 3, 1(2016)

Complications 1/2; photon switching effect

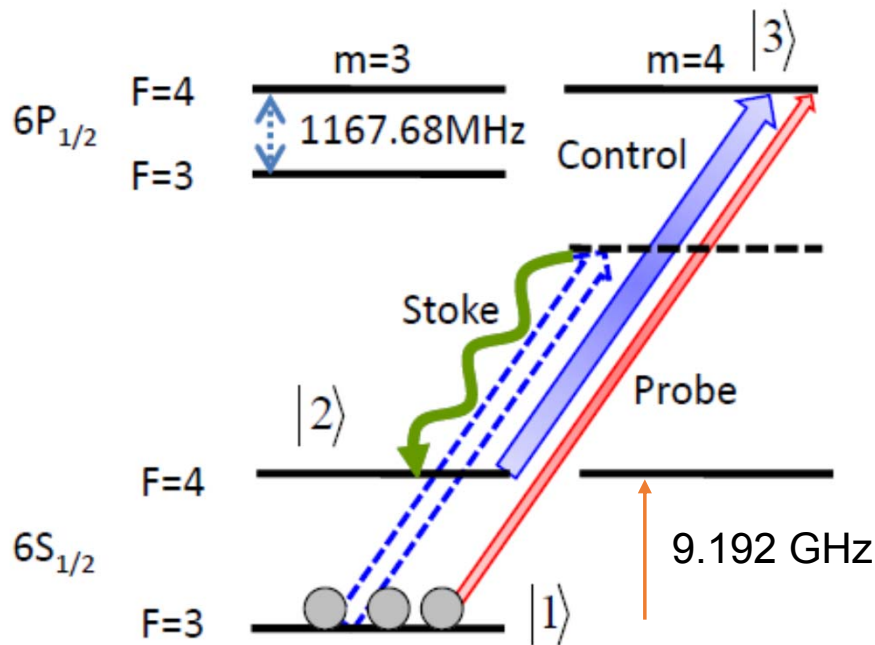
- Off-resonant excitation of the control field to nearby transitions induces a control-intensity-dependent decoherence rate.
- Resolve this problem by implementing EIT in cesium D₁ line and conducting Zeeman-state optical pumping.



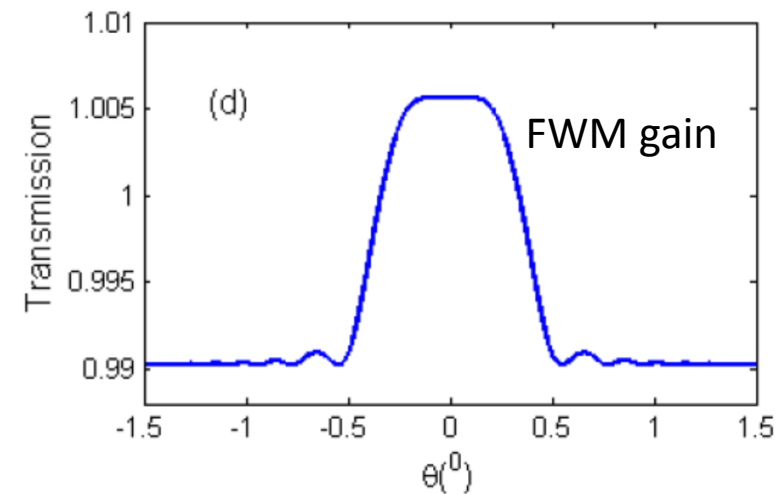
fixed, $T_d = \frac{(OD)\Gamma}{\Omega_c^2}$

Complications 2/2 : Four-wave mixing

- Off-resonant excitation of the control field on the probe transition induces the FWM, which introduce **quantum noise** to reduce the memory fidelity.
- Reduce the FWM by breaking the phase matching condition through the geometry arrangement between the control and probe beams.



$$\Delta k_z = (\vec{k}_{probe} - 2\vec{k}_{control} + \vec{k}_{Stoke}) \cdot \hat{z}$$



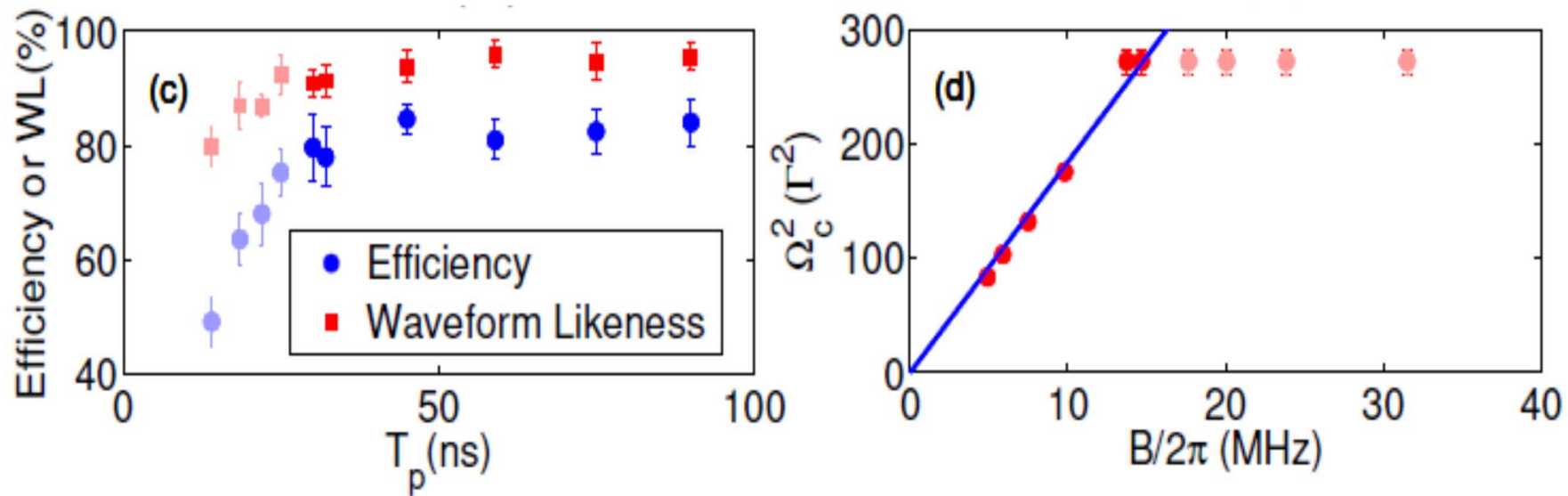
Y.-F. Hsiao..., YCC* Opt. Lett. 39, 3394(2014)

Y.-F. Hsiao...YCC*, PRL. 120, 183602(2018)

Fleischhauer et al., PRA 88, 013823(2013);

Towards broadband EIT-based optical memory

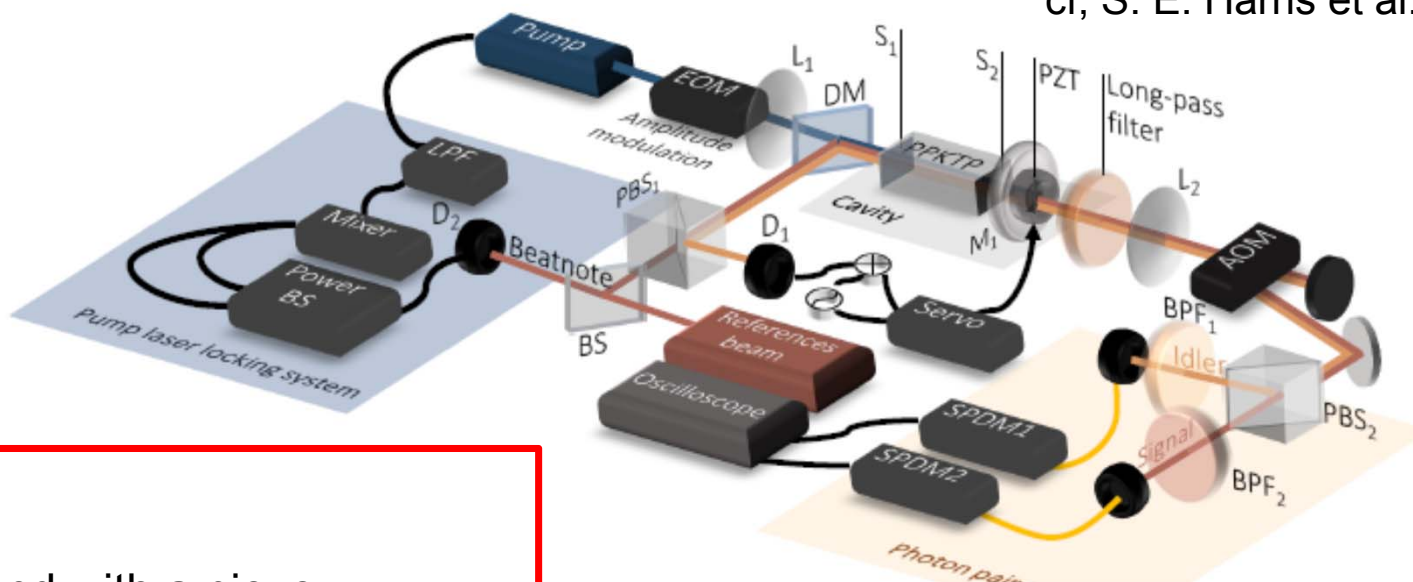
- High optical depth & strong control field are requirements to achieve a high bandwidth.
- Achieved > 80% efficiency for 30-ns pulses ($\sim 15\text{MHz}$) and > 50% efficiency for 14 ns pulses (31MHz).
- Currently limited by available control power towards even higher bandwidth.



Y.-C. Wei, ..., YCC*, Phys. Rev. A. 102, 063720(2020)

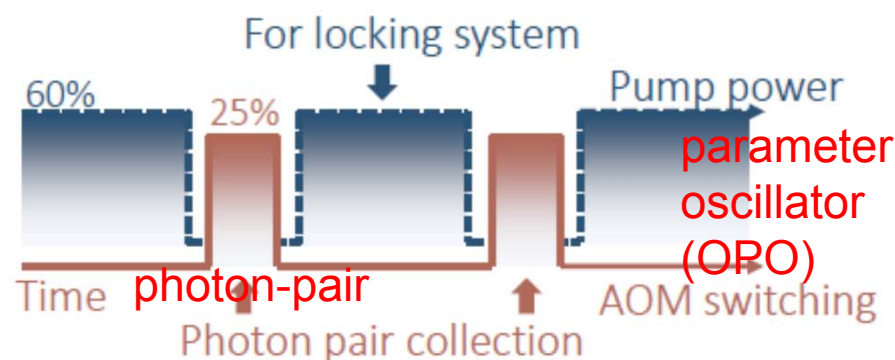
Photon-pair source by cavity-enhanced spontaneous parametric down-conversion (SPDC)

cf, S. E. Harris et al. APL 101,051108

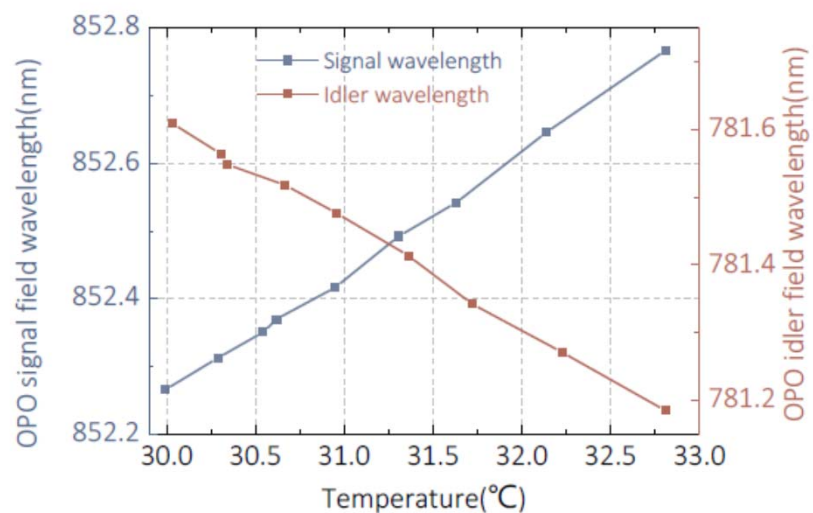


Features

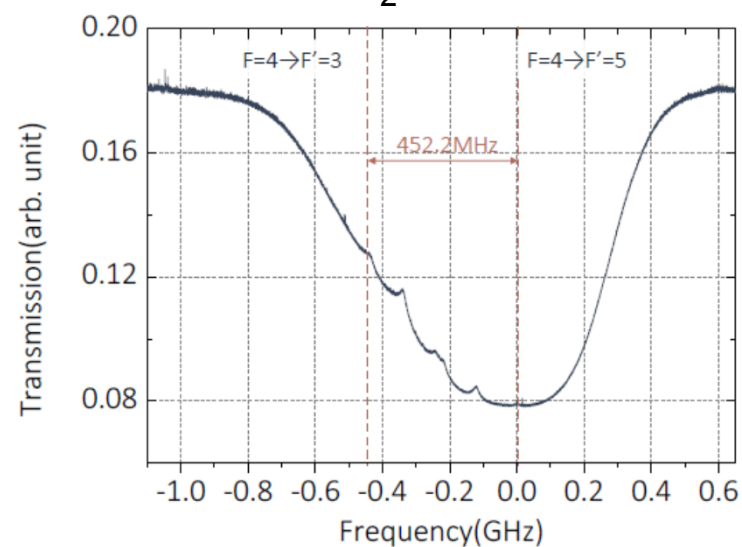
- 1, One cavity end with a piezo, easy tuning & frequency locking
- 2, Switching OPO/photon-pair modes, OPO allow easy wavelength monitoring & alignment



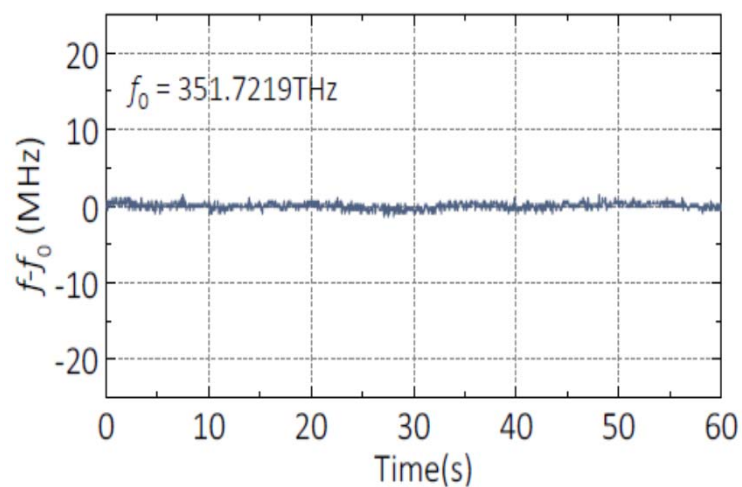
Wavelength tuning



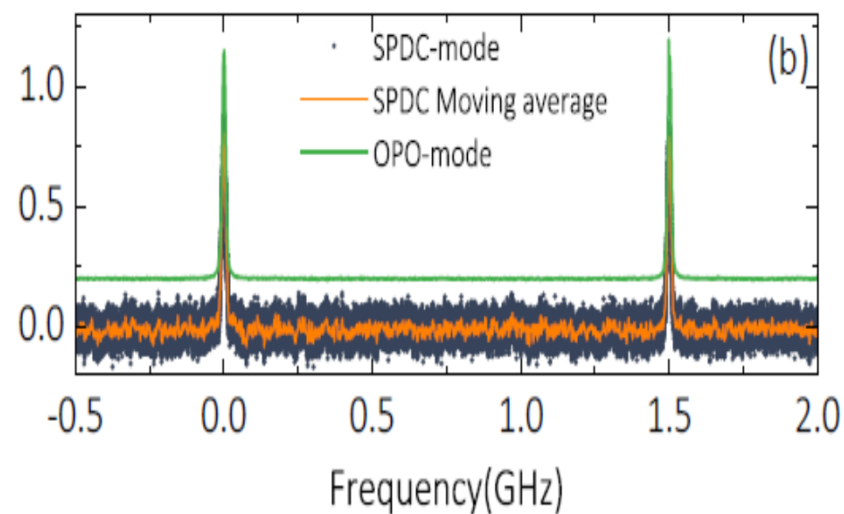
Cesium D₂ line



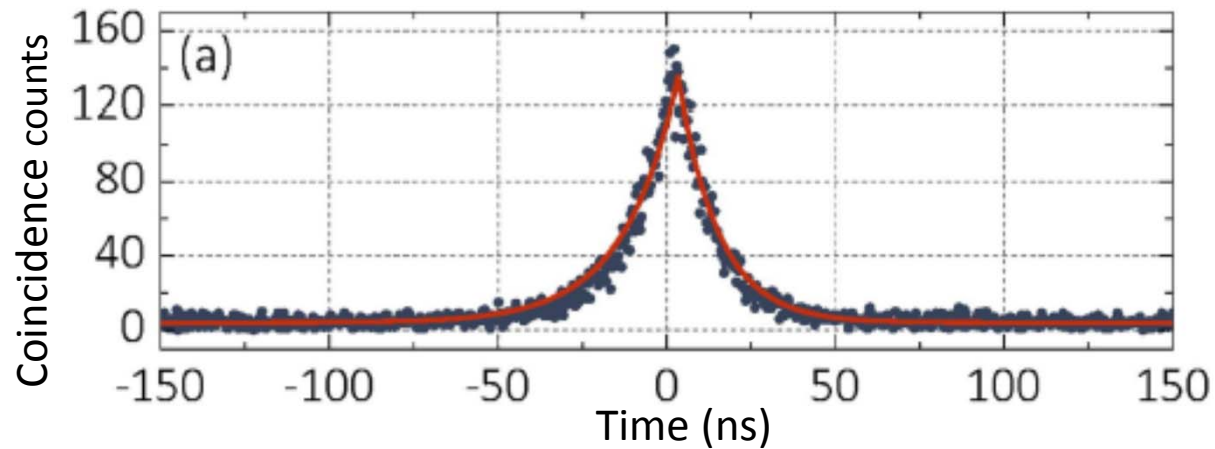
Frequency stabilization



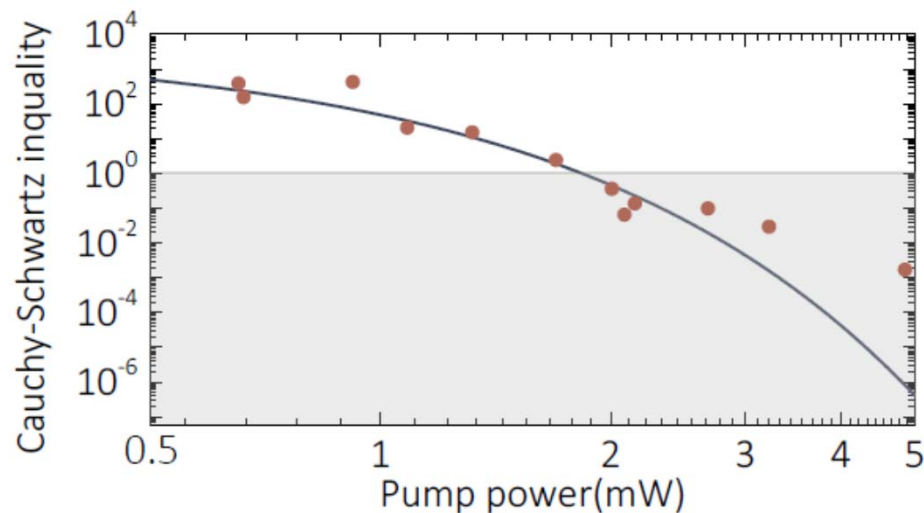
Single longitudinal mode



Glauber correlation function $G_{s,i}^{(2)}(\tau) = \langle \hat{a}_i^\dagger(t+\tau) \hat{a}_s^\dagger(t) \hat{a}_s(t) \hat{a}_i(t+\tau) \rangle$



- Bandwidth~ 6.6 MHz, compared to Cs D₂ natural linewidth of 5.2 MHz.
- Nonclassicality: violation of Cauchy-Schwarz inequality by a factor of up to 425.

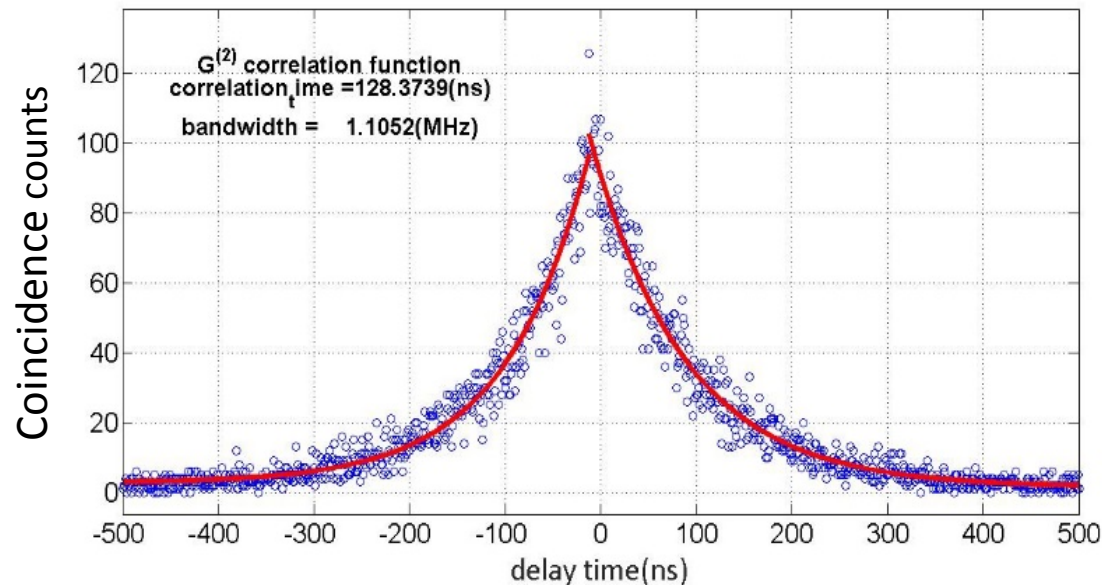


$$[g_{s,i}^{(2)}(0)]^2 / [g_{s,s}^{(2)}(0) g_{i,i}^{(2)}(0)] \leq 1$$

P.-J. Tsai, YCC*, Quantum Sci & Techno, 3, 034005(2018)

An improved photon-pair source at Cs D₁ line

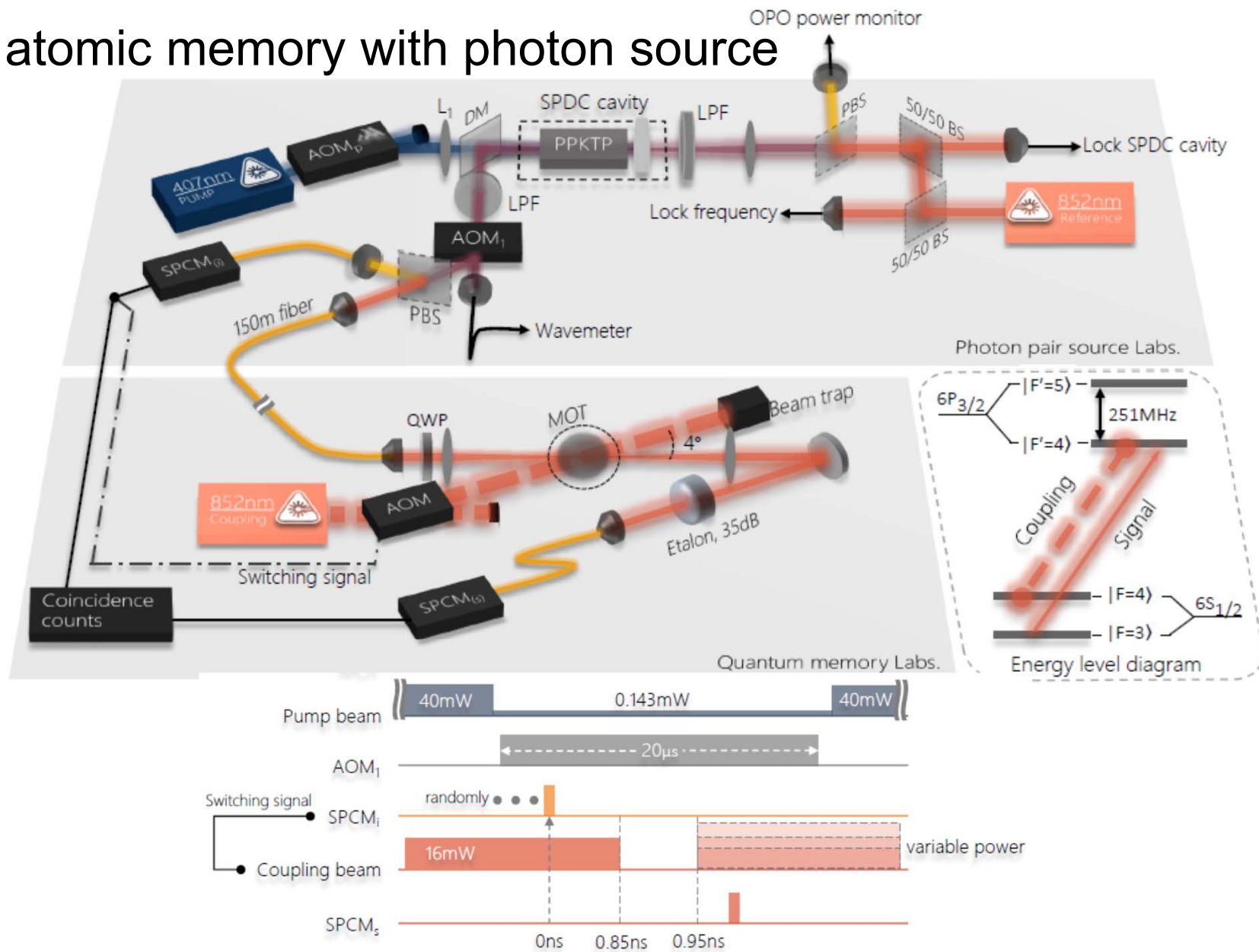
- Quantum memory operating at Cs D₁ line (~894.6 nm) is better due to a smaller photon-switching effect.
- We recently built a new photon-pair source at 894.6 nm with a bandwidth of ~2 MHz.



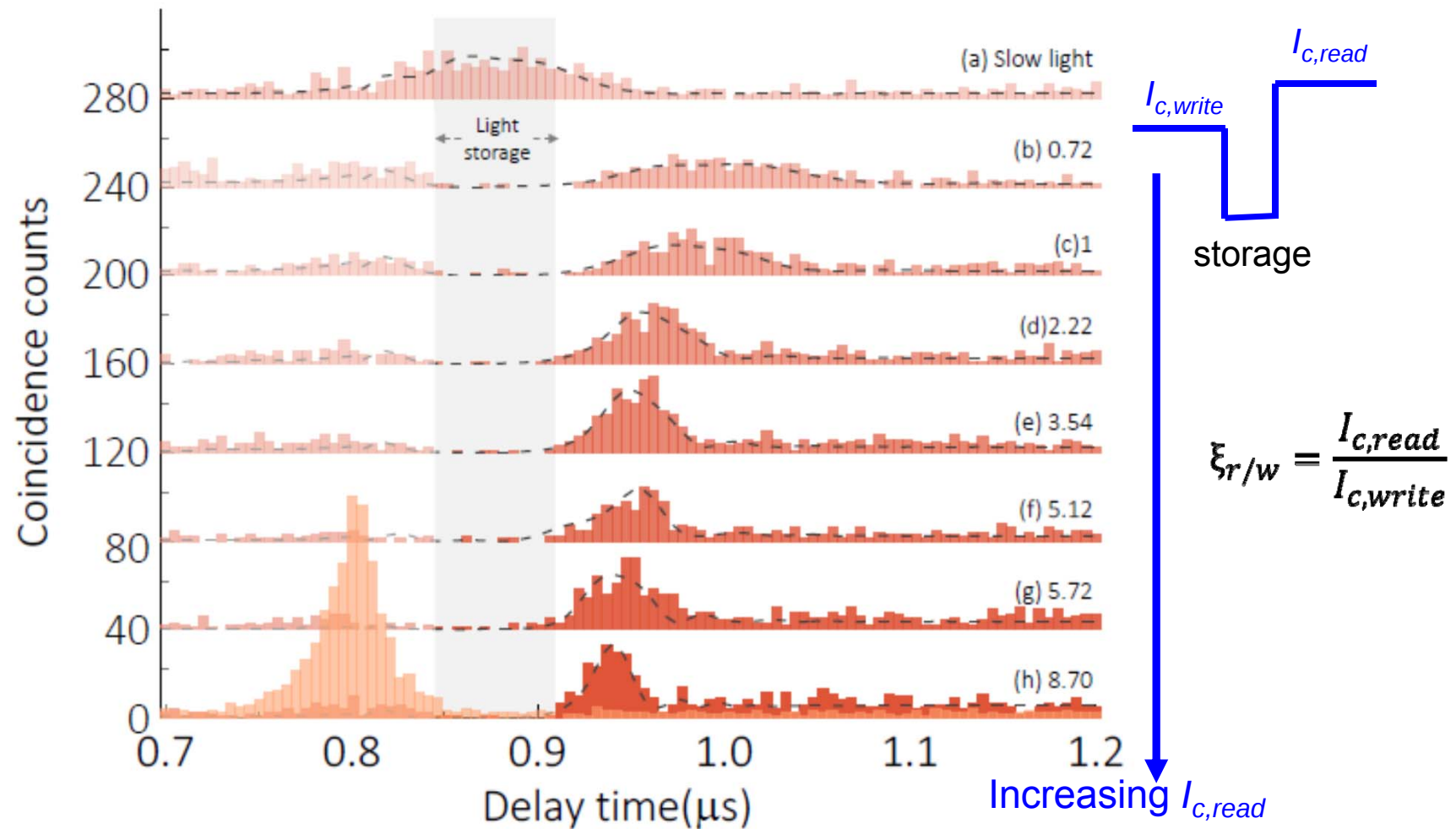
Bandwidth : 2.2 MHz

Efficient quantum memory for heralded single photons

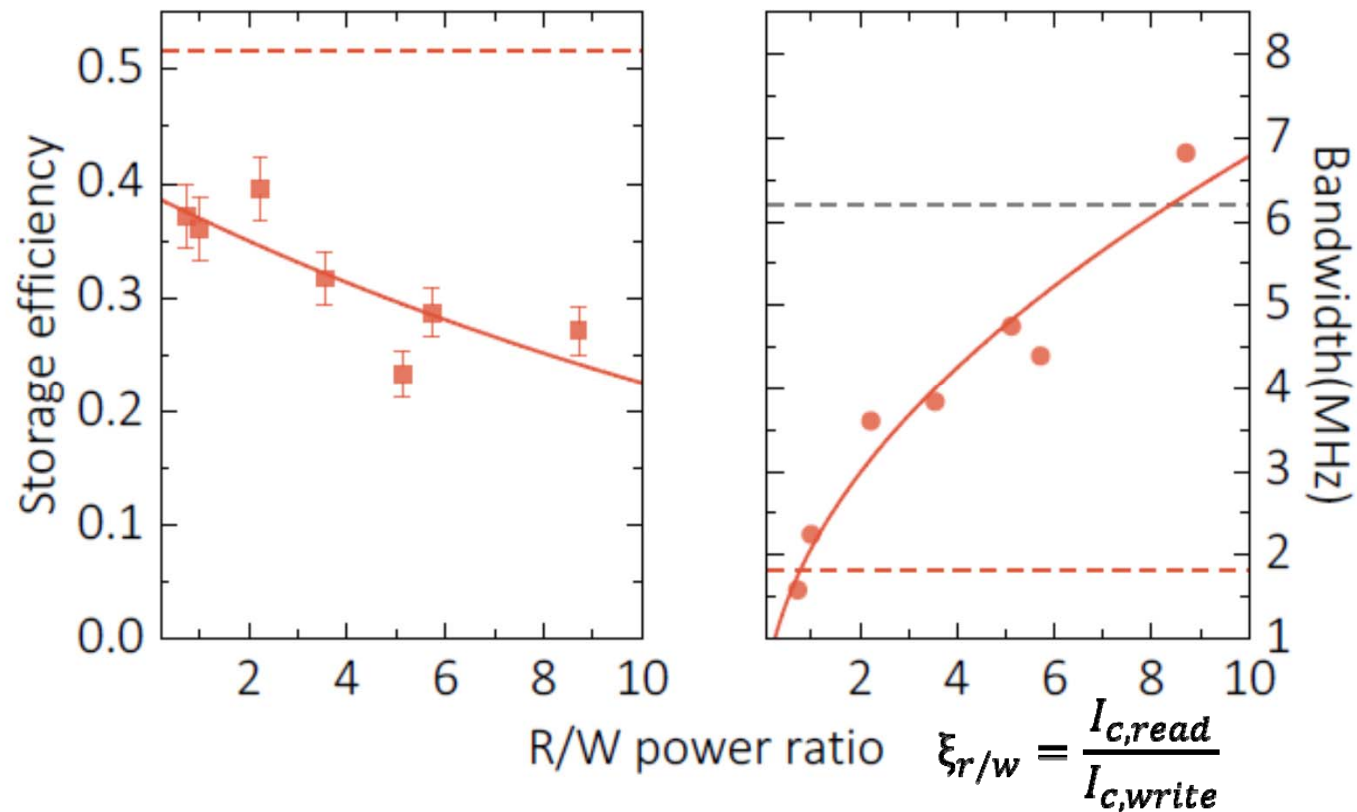
Combining atomic memory with photon source



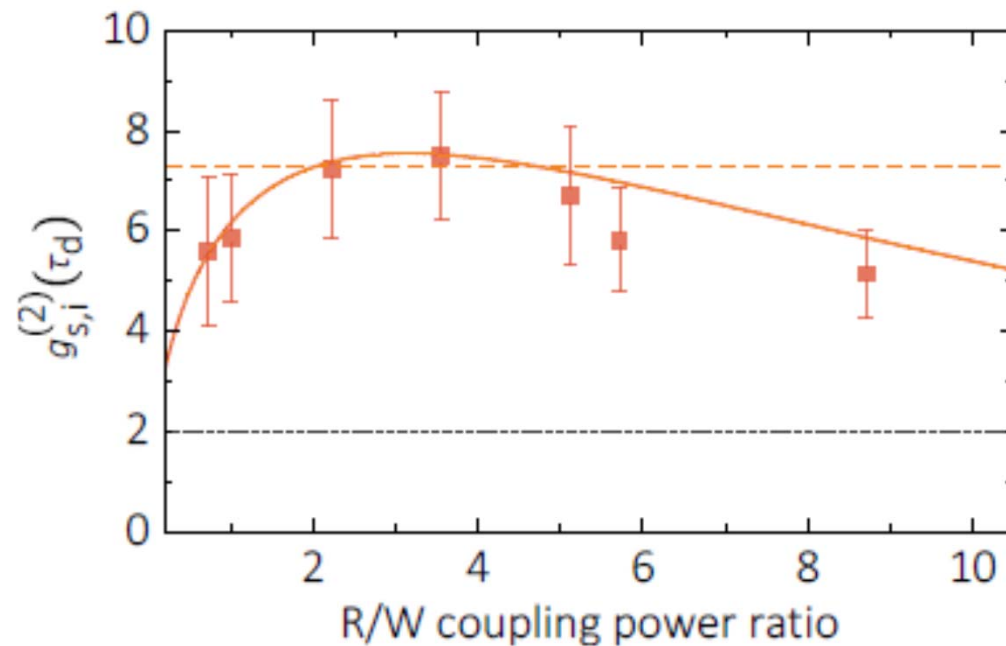
Storage and manipulation of heralded single photons



- Relatively low efficiency ($\sim 36\%$) @ Cs D₂ line:
 - 1) bandwidth of single photons is too large
 - 2) D₂ line photon switching effect
 - 3) Optical depth not enough in the new MOT system (up to 80).
- Waveform (bandwidth) manipulation of single photons by memory



Manipulate the non-classical cross correlation



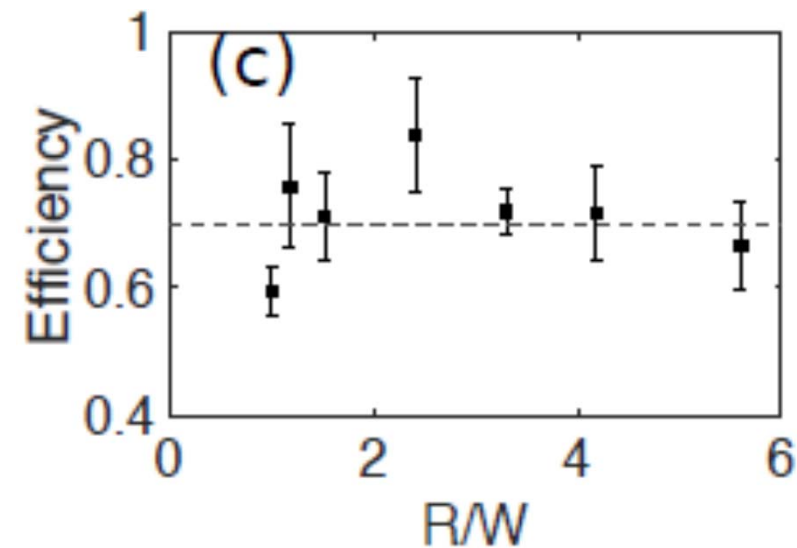
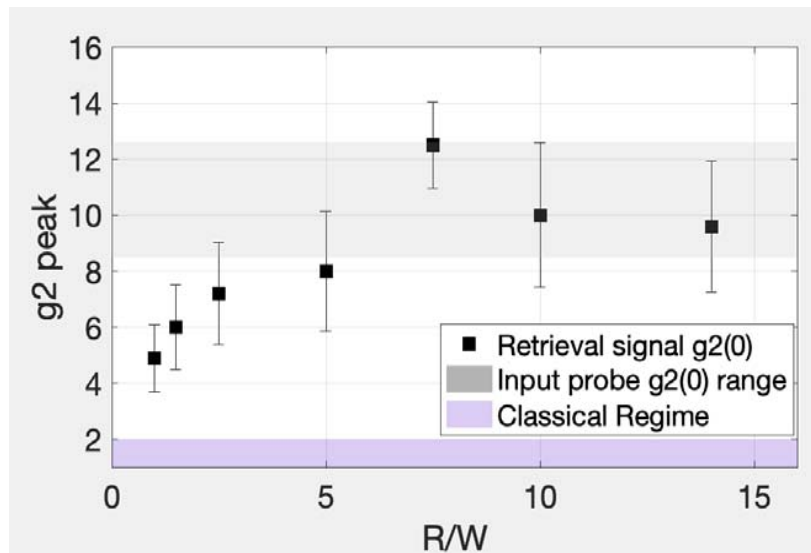
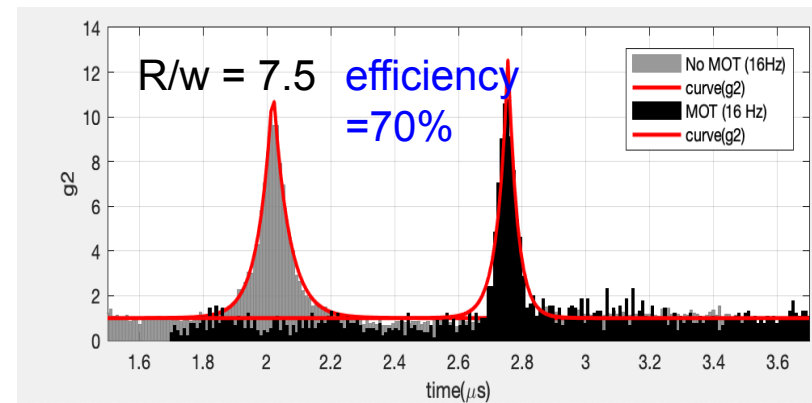
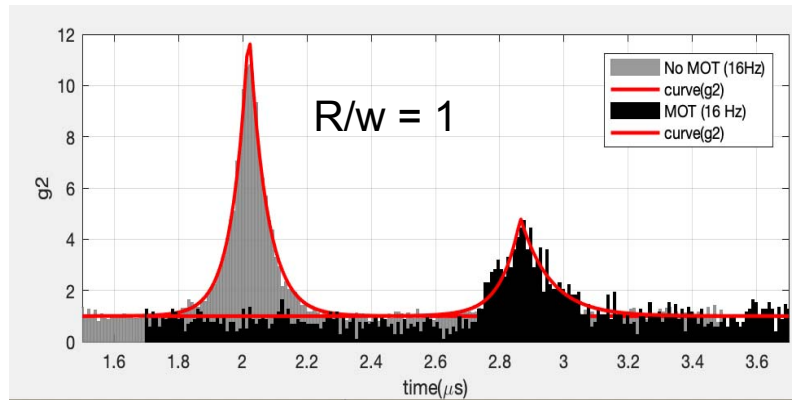
$$\xi_{r/w} = \frac{I_{c,read}}{I_{c,write}}$$

$$g_{s,i}^{(2)}(\tau_d) = \frac{N_{s,i} \sqrt{\xi_{r/w}} e^{-\gamma_s \xi_{r/w}}}{\alpha \xi_{r/w} + N_b},$$

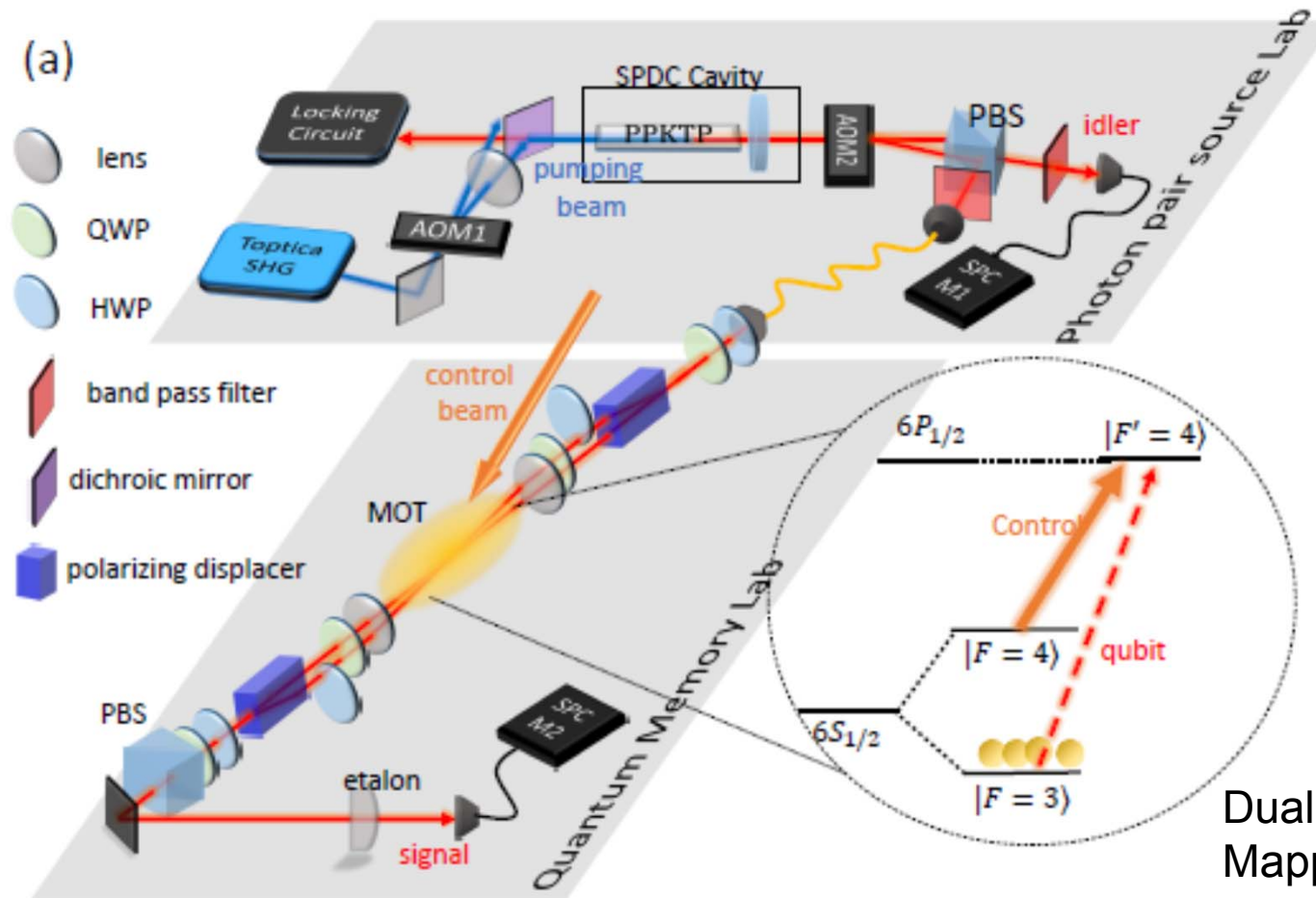
enhancement $\rightarrow$ $\sqrt{\xi_{r/w}}$ $\leftarrow$ Photon switching effect $e^{-\gamma_s \xi_{r/w}}$

$\alpha \xi_{r/w}$ $\leftarrow$ Control leakage N_b $\leftarrow$ Stray light/detector dark counts

Recent results for quantum storage @ D_1 EIT System



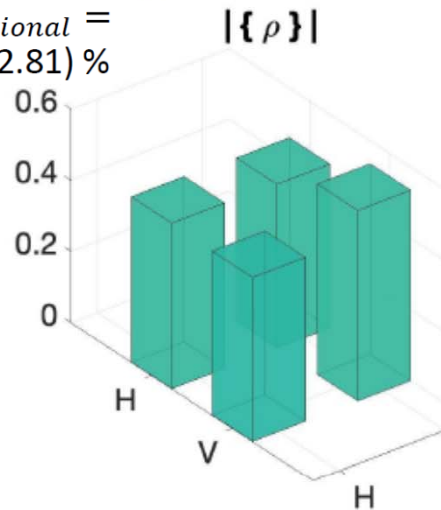
Quantum memory for polarization qubit



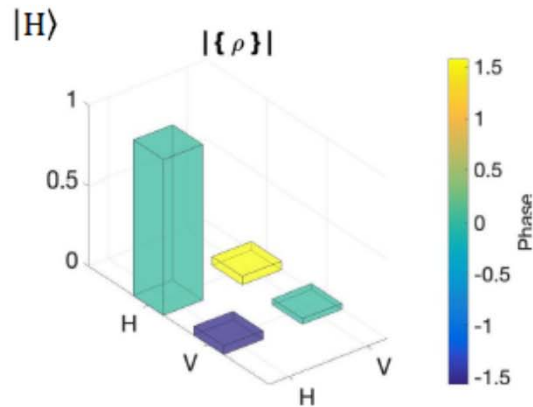
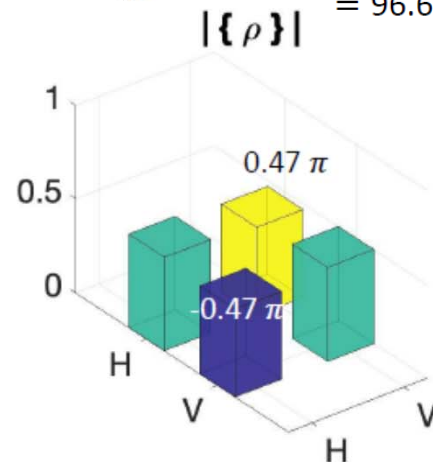
Dual rail scheme:
Mapping the polarization modes
into spatial modes

Fidelity for polarization qubits after storage

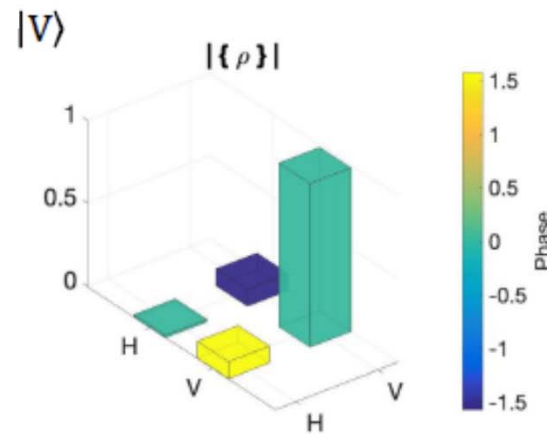
Consider $|D\rangle = \frac{1}{\sqrt{2}}(|H\rangle + |V\rangle)$
 $\mathcal{F}_{conditional} = 95.96 (2.81) \%$



Consider $|R\rangle = \frac{1}{\sqrt{2}}(|H\rangle + i|V\rangle)$ $\mathcal{F}_{conditional} = 96.64 (3.12) \%$

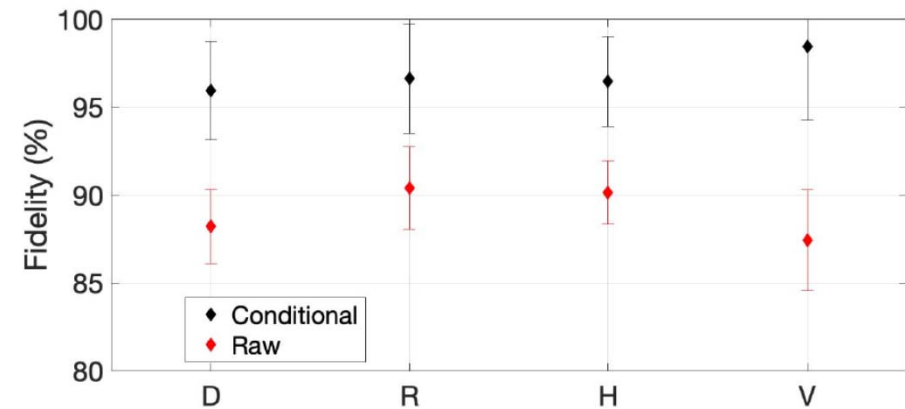


$\mathcal{F}_{conditional} = 96.47 (2.55) \%$



$\mathcal{F}_{conditional} = 98.45 (4.17) \%$

Fidelity > ~96% for four polarization states



Y.-C. Tseng, ..., YCC*, [arXiv:2011.14948](https://arxiv.org/abs/2011.14948)

Conclusions

- Obtain a high optical depth (>1000) and a low decoherence rate ($10^{-4} \Gamma$) with a MOT setup.
- Obtain a record-high storage efficiency of 92% for EIT-based coherent optical memory!
- Push the EIT-memory bandwidth to ~ 30 MHz .
- Developed a bright, narrowband, single-mode, and convenient photon-pair source lockable to atomic transition.
- Demonstrated quantum storage and manipulation of single photons with EIT atomic memories (36% efficiency @ Cs D_2 line and $>70\%$ @ D_1 line)
- Demonstrated quantum memory for polarization qubits with a fidelity $> 96\%$.

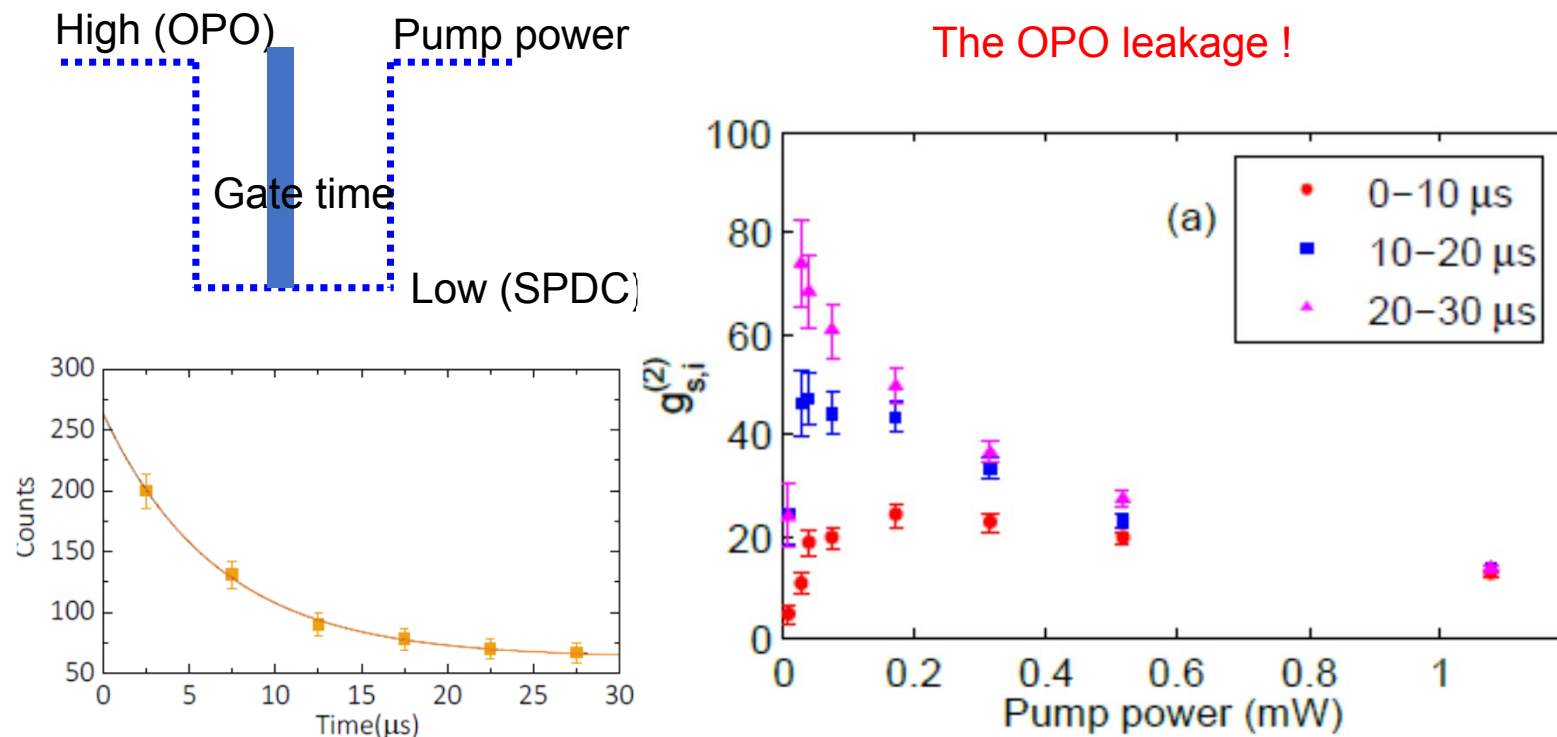
Acknowledgements

- IAMS: Ya-Fen Hsiao, Pin-Ju Tsai, Hung-Shiue Chen, Sheng-Xiang Lin, Bo-Han Wu, Yan-Cheng Wei, Yu-Chih Tseng, ...
- NTHU: Ite A. Yu, NSYSU: Yi-Hsin Chen
- NCKU: Yong-Fan Chen
- MOST and IAMS for budget support, help from NCTS.



Noise Sources

- Noise is a serious issue in quantum storage which degrades the quantum feature.
- Noises from SPDC photon source, stray light, control leakage, Raman-generated photons due to cold atom and hot vapor...



Noise from Raman-generated Photons

- These highlight the absolute requirement of quantum technology!

