

NCTS Workshop on Quantum Condensation (QC13) 2013/08/29

Symmetry-protected topological phases and gapped vector-spin-chirality phases in a dimerized spin-1/2 XXZ zigzag chain

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H.U. and S.O., arXiv:1308.5747

Introduction: Spiral magnet and ferroelectricity



Vector-chiral (VC) order in quasi 1D system



Vector-chiral (VC) order in quasi 1D system



Vector-chiral (VC) order in quasi 1D system



Field-induced ferroelectric system Rb₂Cu₂Mo₃O₁₂



By courtesy of Yanagisawa et al.

Buckling of CuO₂ plane
 bond alternation

Solodovnikov et al. (1997)

•
$$J_1/J_2 \sim -2.7 \ (J_1 = -138 \text{K}, J_2 = 51 \text{K})$$

 $\mathcal{H} = \sum_{\ell=1,2} J_\ell \sum_j \hat{S}_j \cdot \hat{S}_{j+\ell}$

No magnetic order (*T*>2K) Hase et al. (2004)

 Field-induced ferroelectricity (T<8K)
 → possible VC LRO w/o spin spiral Yasui et al. (2013)

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$$CuO_2$$
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Y. Yasui et al., J. Appl. Phys. 113, 17D910 (2013).

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Main goals

✓Theoretical

- ✓ Propose a model which has gapped VC phases w/o any spiral spin LRO over a wide region of its parameter space.
 <u>Key: Bond alternation</u>
 ✓ Understand the model under low-magnetic field
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 Key: Staggered scalar chirality
- ✓ Understand symmetry-protected topological (SPT) phases.
 <u>Key: Neel LRO</u>

✓Experimental

- ✓ Find a new type of ferroelectric transition
 - ✓ $Rb_2Cu_2Mo_3O_{12}$: Weak bond alternation
 - ✓ spin gap: ~ 0.2 meV ⇔ ~2 T
 - ✓ Ferroelectric transition at $H \sim 0.1$ T



Zero Field Case

$$\hat{\mathcal{H}}_{\delta_{\text{XXZ}}} = J_1 \sum_{j} (1 - (-1)^j \delta) \left[\sum_{\mu = x, y} \hat{S}_j^{\mu} \hat{S}_{j+1}^{\mu} + \Delta \hat{S}_j^z \hat{S}_{j+1}^z \right] + J_2 \sum_{j} \left[\sum_{\mu = x, y} \hat{S}_j^{\mu} \hat{S}_{j+2}^{\mu} + \Delta \hat{S}_j^z \hat{S}_{j+2}^z \right]$$

$$(J_1 < 0, \quad \delta \ge 0, \quad 0 \le \Delta \le 1)$$

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$$J_1 (1 - \delta)$$

$$J_2$$



S. Furukawa, M. Sato, SO, and A. Furusaki, PRB **86**, 094417 (2012).

Phase diagrams (δ =0)

✓ Even-parity dimer (D_) phase



Dimer LRO: Gapped phase $\hat{D}_{j}^{\mu} = (-1)^{j} (\hat{S}_{j-1}^{\mu} \hat{S}_{j}^{\mu} - \hat{S}_{j}^{\mu} \hat{S}_{j+1}^{\mu}),$ $(\underline{\operatorname{Ave}}_{j} \langle \hat{D}_{j}^{x} \rangle) (\operatorname{Ave}_{j} \langle \hat{D}_{j}^{z} \rangle) < 0,$ a spatial average only over *N*/2 spins located at

for excluding the boundary effects^[1].

[1] T. Hikihara *et al.*, PRB **78**, 144404 (2008).

Phase diagrams (δ =0)



Phase diagrams ($\delta = 0$)



✓Haldane-dimer (D+) phase

Dimer LRO: Gapped phase

$$\left(\operatorname{Ave}_{j}\langle \hat{D}_{j}^{x}\rangle\right)\left(\operatorname{Ave}_{j}\langle \hat{D}_{j}^{z}\rangle\right) > 0,$$

Phase diagrams (δ =0)



[*] S. Furukawa, M. Sato, SO, and A. Furusaki, PRB 86, 094417 (2012).

zero Field Case (7/19)

Phase diagrams (δ >0)



Zero Field Case (8/19)

Phase diagrams (δ >0)



Abelian bosonization analysis

Start from decoupled XXZ chains:

1.
$$J_1$$
 coupled two chain: $\mathcal{H}_+ + \mathcal{H}_- + \mathcal{H}_{int}$
 $\mathcal{H}_{\pm} = \sum_{\nu=\pm} \frac{v_{\nu}}{2} \left[K_{\nu} (\partial_x \theta_{\nu})^2 + K_{\nu}^{-1} (\partial_x \phi_{\nu})^2 \right]$

- 2. Mean field decoupling of the \pm sector (valid in VC, Neel, and/or dimer phases)
- 3. Self-consistent solution to two sine-Gordon problems

for
$$\left\{ \begin{array}{l} \langle \sin\sqrt{4\pi}\phi_+ \rangle \\ \langle \cos\sqrt{4\pi}\phi_+ \rangle \end{array} \right\}$$
 and $\left\{ \begin{array}{l} \langle \sin\sqrt{4\pi}\theta_- \rangle \\ \langle \cos\sqrt{4\pi}\theta_- \rangle \end{array} \right\}$

This analysis for δ =0: Nersesyan. *et al.*, (1998)

S. Furukawa, M. Sato, SO, and A. Furusaki, PRB 86, 094417 (2012).

 $\text{VCD}_{\pm} \text{ Phase} \begin{bmatrix} \text{Absence of Neel LRO: } \langle \sin \sqrt{4\pi}\phi_{+} \rangle = 0 \\ \text{VC LRO: } \langle \sin \sqrt{4\pi}\theta_{-} \rangle \neq 0 \\ \text{z-component of dimer LRO: } \langle \cos \sqrt{4\pi}\phi_{+} \rangle \neq 0 \\ \text{x, y-component of dimer LRO: } \langle \cos \sqrt{4\pi}\theta_{-} \rangle \neq 0 \end{bmatrix}$

✓Transition from VCD₊ to VCD₋

$$\begin{array}{c} & \bigoplus \quad \text{Sign changing of } \langle \cos \sqrt{4\pi}\phi_{+} \rangle \\ \propto J_{1}\Delta\delta & \propto J_{1} \\ \langle \cos \sqrt{4\pi}\phi_{+} \rangle_{0} = C(K_{+}) \begin{pmatrix} \gamma_{\delta}^{+} + \gamma_{1}(\cos \sqrt{4\pi}\theta_{-})_{0} \\ v_{+}/a^{2} \end{pmatrix}^{\frac{1}{2K_{+}^{-1}-1}} \\ \left| \frac{\gamma_{\delta}^{z+} + \gamma_{1}(\cos \sqrt{4\pi}\theta_{-})_{0}}{v_{+}/a^{2}} \right| \geq \left[C(K_{+})K_{-}\frac{(\gamma_{\text{tw}}^{-+})^{2}}{v_{+}v_{-}/a^{2}} \right]^{\frac{2K_{+}^{-1}-1}{2(K_{+}^{-1}-1)}} \\ \text{for absence of Neel LRO} \end{aligned}$$

VCD_± Phase $\begin{cases}
\text{Absence of Neel LRO: } \langle \sin \sqrt{4\pi}\phi_+ \rangle = 0 \\
\text{VC LRO: } \langle \sin \sqrt{4\pi}\theta_- \rangle \neq 0 \\
\text{z-component of dimer LRO: } \langle \cos \sqrt{4\pi}\phi_+ \rangle \neq 0 \\
\text{x, y-component of dimer LRO: } \langle \cos \sqrt{4\pi}\theta_- \rangle \neq 0
\end{cases}$

✓Possibilities:

(1) An accidental direct continuous transition

 $\Rightarrow \gamma_{\delta}^{z+} + \gamma_1 \langle \cos \sqrt{4\pi} \theta_- \rangle = 0, K_+ = 1$ are satisfied at a single point.

(2) Gapless VCD₀ phase (ϕ_+ is unlocked.)

(3) Neel LRO at a transition point

(4) VCND phase sandwiched by VCD_{\pm} phases



String correlation functions

$$C_N^{(O_\ell^z)}(2r) = \operatorname{Ave}_j \langle \hat{O}_{2j+\ell,2j+\ell+2r}^z \rangle \quad (\ell = 1, 2),$$

where $\hat{O}_{j,j+2r}^z = -(\hat{S}_j^z + \hat{S}_{j+1}^z) \exp(i\pi \sum_{k=2}^{2r-1} \hat{S}_{j+k}^z) (\hat{S}_{j+2r}^z + \hat{S}_{j+2r+1}^z).$

E. H. Kim et.al, PRB 62, 14965 (2000), and references therein.

Distinction of two dimer phases:

Phase	LRO	A particular spin configuration
D ₊ VCD ₊	$\lim_{r \to \infty} C_N^{(O_1^z)}(2r) > 0$ $\lim_{r \to \infty} C_N^{(O_2^z)}(2r) = 0$	$S_{2j+1}^{z} + S_{2j+2}^{z} : 0 1 -1 0 0 1 0 -1$
D_ VCD_	$\lim_{r \to \infty} C_N^{(O_1^z)}(2r) = 0$ $\lim_{r \to \infty} C_N^{(O_2^z)}(2r) > 0$	$S_{2j+2}^{z} + S_{2j+3}^{z} : 1 0 -1 0 1 0 -1$



✓
$$J_1/J_2 = -2.5, \delta = 0.02$$

Entanglement Entropy: $S^{vN} = -\text{Tr}_A \left[\rho_A \log \rho_A \right]$

around critical:

$$S^{\rm vN} = \frac{c}{6} \log \frac{\xi}{\pi} + S_0$$

Critical:

$$S^{\mathbf{vN}} = \frac{c}{6} \log \frac{N}{\pi} + S_0$$

c: central charge







Criticality at the D₋-VCD₋ boundary

 $S_A(N,\ell) = \frac{\tilde{c}}{6} \log \left[\frac{N}{\pi} \sin \frac{\pi \ell}{N} \right] + S_0$ P. Calabrese and J. Cardy, J. Stat. Mech. (2004) P06002.



Entanglement Spectrum

$$S^{\mathrm{vN}} = -\mathrm{Tr}_A[\rho_A \log \rho_A] = \sum_i \zeta_i \exp(-\zeta_i)$$

 ζ_i : Entanglement Spectrum



 D_+ and D_- are SPT phases distinct in the absence of the Neel LRO.

$$\checkmark : \text{like } (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)/\sqrt{2}$$

(odd-parity dimer)

$$\begin{array}{c} \checkmark \vdots \text{ like } (|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle)/\sqrt{2} \\ \text{(even-parity dimer)} \end{array}$$

Entanglement Spectrum



VCD+ and VCD- have different string orders.➔ Distinct SPT?

Low-field case without Magnetization

$$\hat{\mathcal{H}}_{\rm ssc} = gh \sum_{j} (\hat{\boldsymbol{S}}_{2j-1} - \hat{\boldsymbol{S}}_{2j+2}) \cdot \hat{\boldsymbol{S}}_{2j} \times \hat{\boldsymbol{S}}_{2j+1}$$





Bosonization analysis

	Neel LRO: $\langle \sin \sqrt{4\pi}\phi_+ \rangle \neq 0$
	VC LRO: $\langle \sin \sqrt{4\pi}\theta_{-} \rangle \neq 0$
VCND Phase -	z-component of Dimer LRO: $\langle \cos \sqrt{4\pi} \phi_+ \rangle \neq 0$
	xy-component of Dimer LRO: $\langle \cos \sqrt{4\pi}\theta_{-} \rangle \neq 0$







 $J_1/J_2 = -2.5$

Case at moderately large magnetic field: Zeeman interaction

$$\hat{\mathcal{H}}_Z = -h \sum_j \hat{S}_j^z$$

Field induced VCND Phase

Phase diagram without $\hat{\mathcal{H}}_Z$:

 VCD_{-} 0.716 0.714 0.712 VCND -0.7100.708-0.7060.704 D_{-} 0.702 0.700 0.1 0.20.3 0.0 0.4 gh

- \checkmark If the system starts from D_+
 - ➔ The field Induced VCND does not appear.





Summary

$$\checkmark$$
Investigate $\hat{\mathcal{H}}_{\delta_{\mathrm{XXZ}}} + \hat{\mathcal{H}}_{\mathrm{ssc}} + \hat{\mathcal{H}}_Z$



✓ Gapped phases:

✓ Dimer phases at zero-filed

Gapped vector chiral dimer phase at weak field

✓ A possible relevance to experiments of $Rb_2Cu_2Mo_3O_{12}$ ✓ Our model: Bond alternation of J_1 : 2 %, Spin gap: ~ 0.02 J_2 ✓ $Rb_2Cu_2Mo_3O_{12}$: $J_1 = -138K$, $J_2 = 51K$

Alternation of the Cu-O-Cu angles: ~1 % @ 300K

(Intensity of the alternation may increase at a low-T.)

Spin gap: 0.2meV $\Leftrightarrow \sim 2K \sim 0.04 J_2$

Thank you for your attention! 謝謝!