



Correlated Bubble Crystals in Rapidly Rotating Dipole-blockaded Fermi Gases

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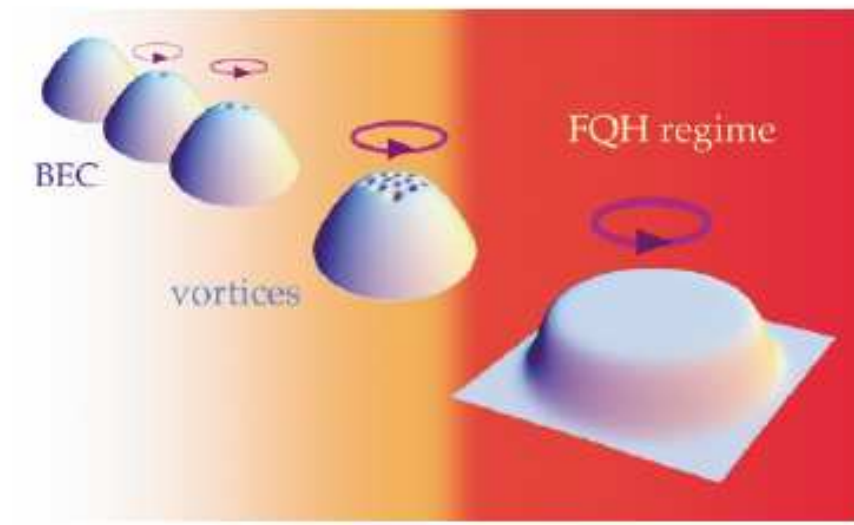
Mark-san Image photography studio

July 25, 2009, released by Mark Hsu
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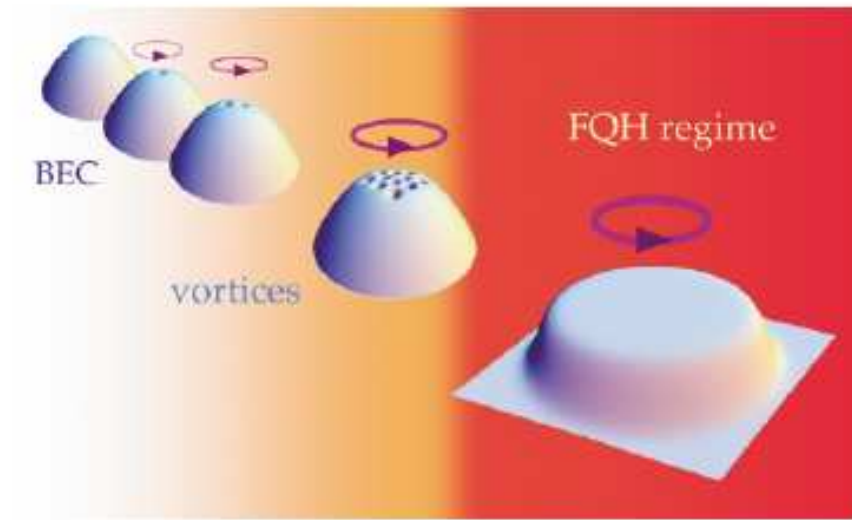
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I. Introduction:

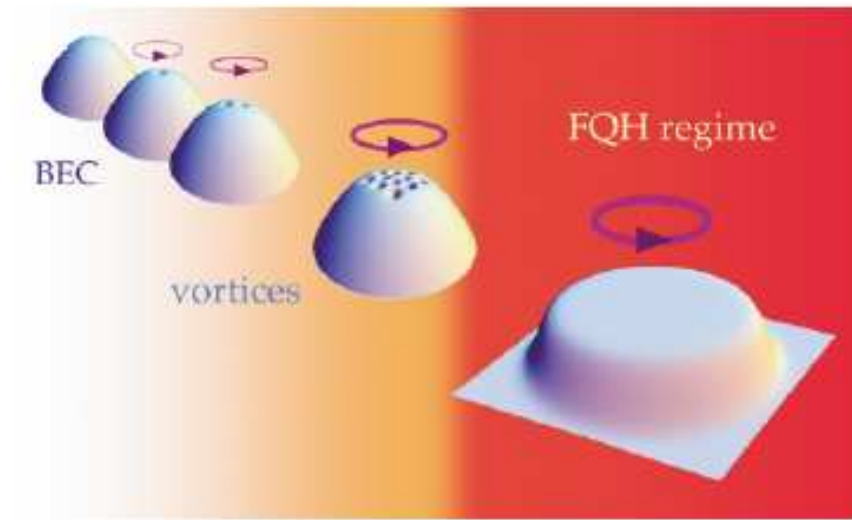


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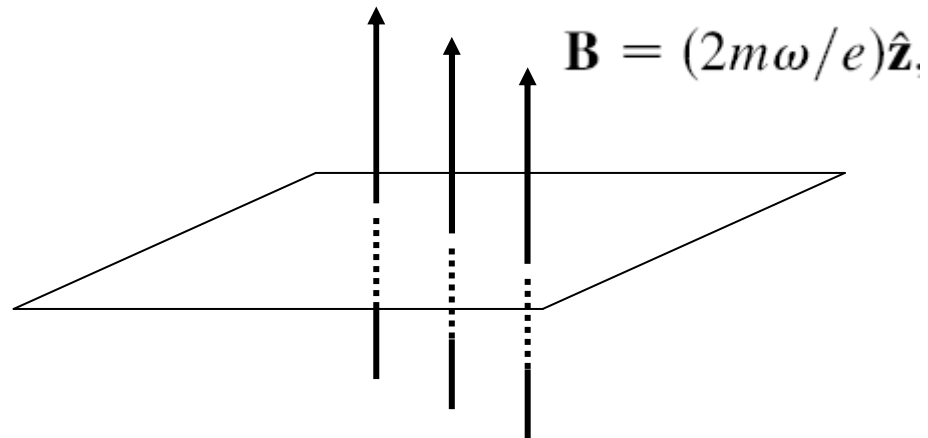
$$\mathcal{H} = \sum_{i=1}^N \frac{1}{2m} (\mathbf{p}_i - m\omega \hat{\mathbf{z}} \times \mathbf{r}_i)^2 + \frac{1}{2} m(\omega_0^2 - \omega^2)(x_i^2 + y_i^2) + \frac{1}{2} m\omega_z^2 z_i^2 + \sum_{i < j}^N V(\mathbf{r}_i - \mathbf{r}_j),$$

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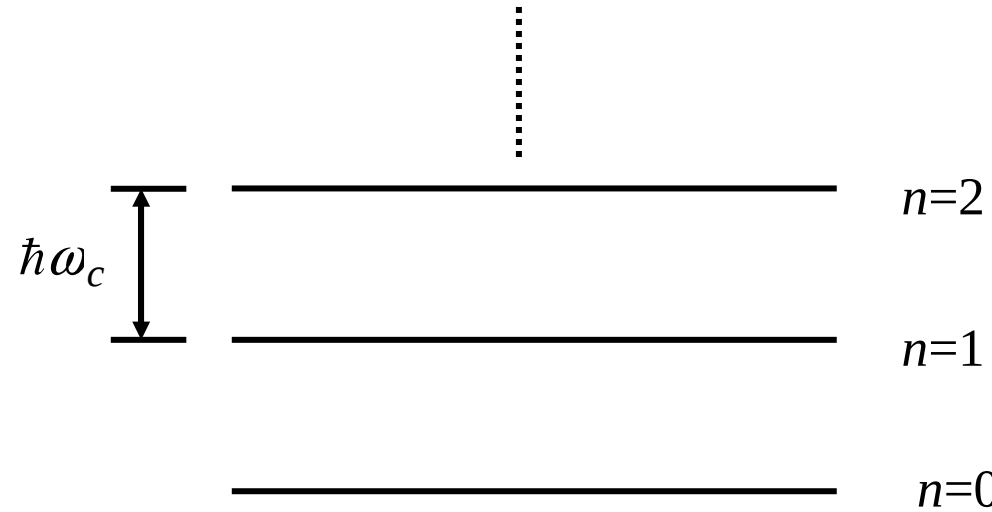


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If $\omega \approx \omega_0$



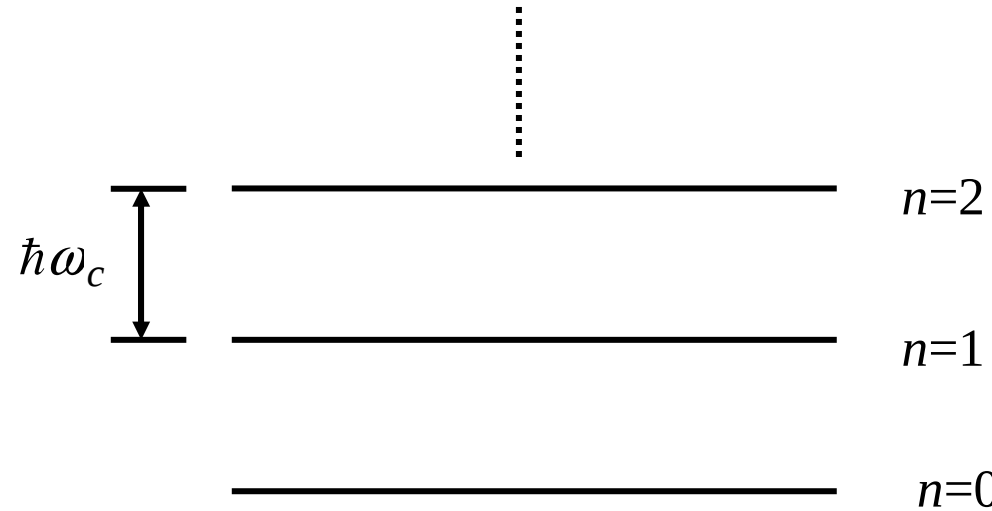
Landau levels



Cyclotron frequency: $\omega_c = 2\omega$

Magnetic length: $a = \sqrt{\hbar/2M\omega} = 1$

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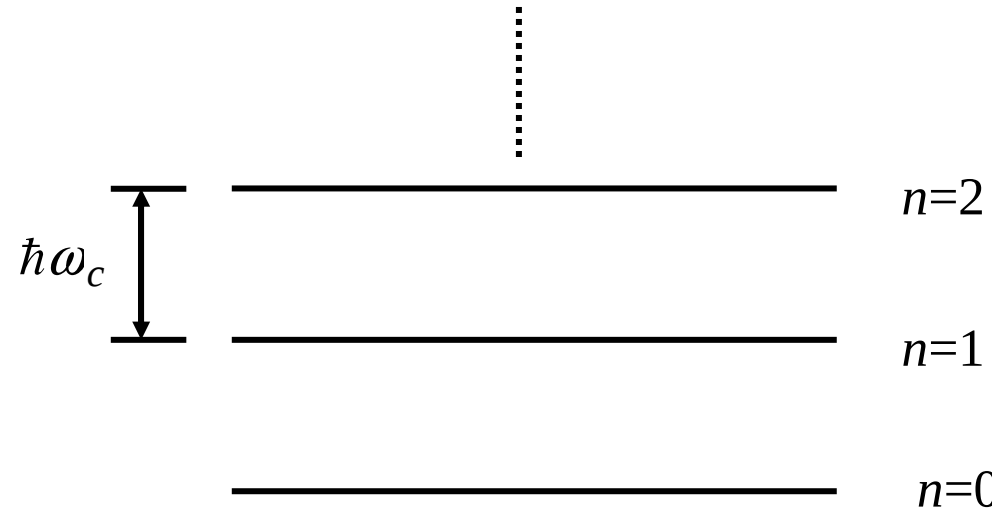


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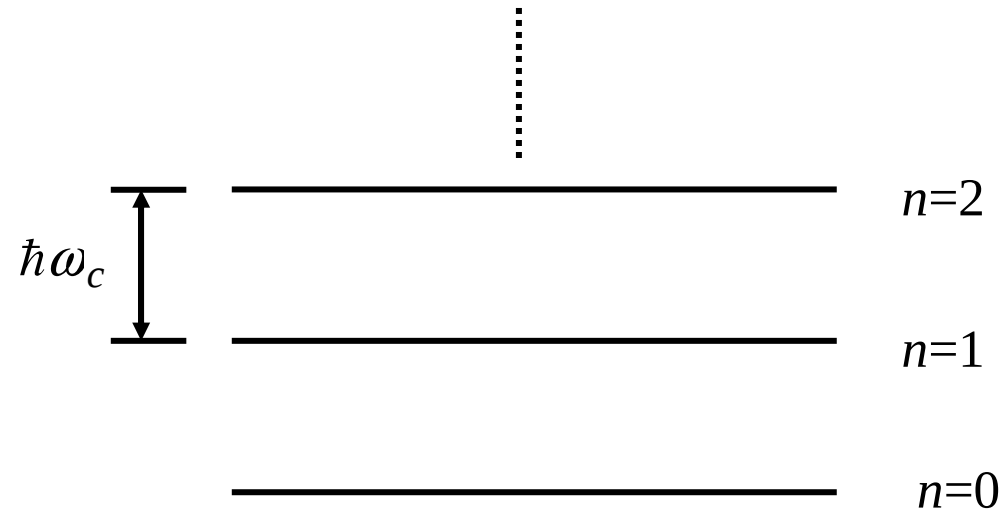
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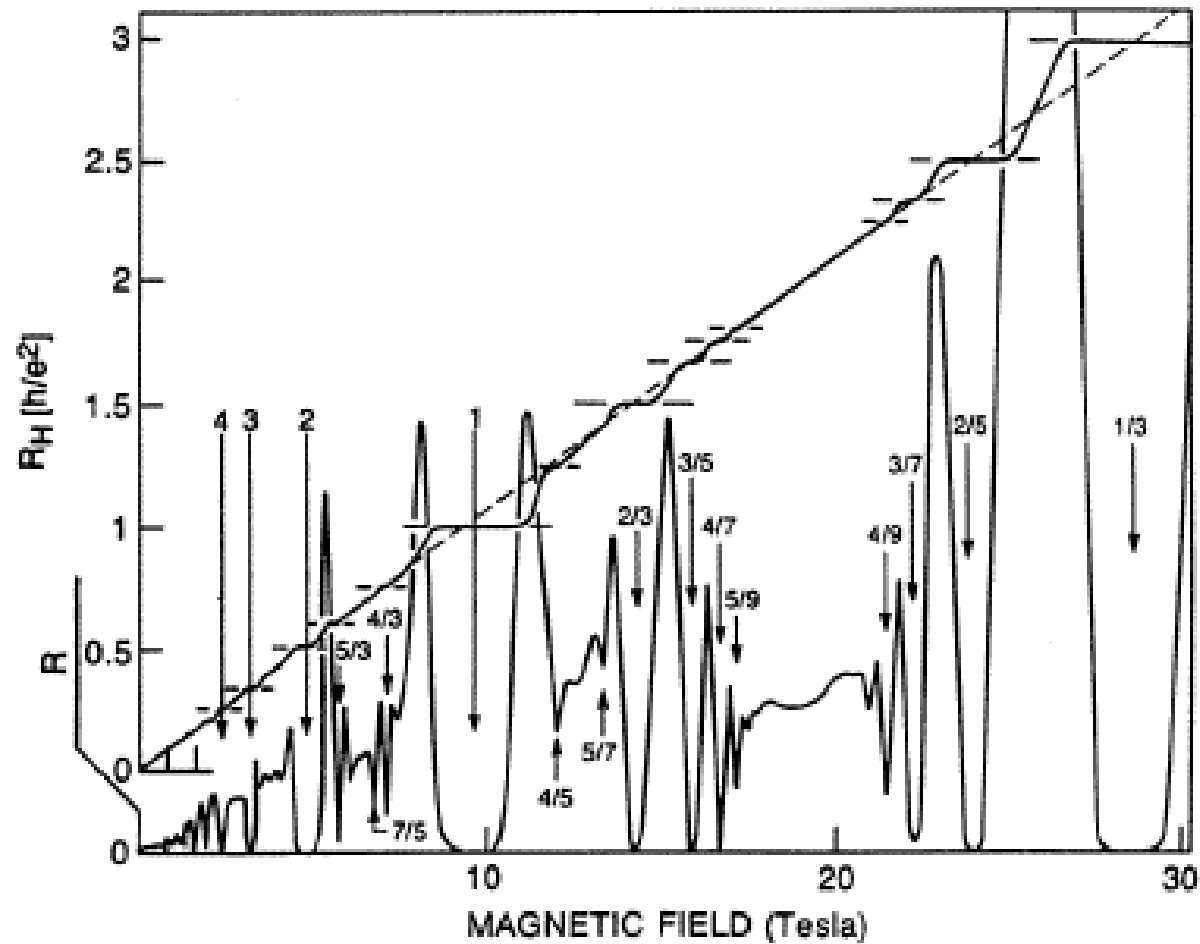
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The kinetic energy is frozen and a crystal state is favorable in the lowest Landau level.
(Wigner Crystal or WC)

But

The Laughlin liquids are favorable for $\nu=1/3, 1/5, 1/7, 1/9$.
(Fractional Quantum Hall Effect)



Eisenstein and Stormer, 1990.

(Laughlin liquid)

$$\Psi_m = \prod_{i < j}^N (z_i - z_j)^m \prod_{\ell} e^{-|z_{\ell}|^2/4}, \quad m=1, 3, 5, 7, 9 \dots \quad z=x+iy.$$

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Wigner crystals for $\nu < 1/7$ and Laughlin liquids for $\nu \geq 1/7$.

(Laughlin liquid)

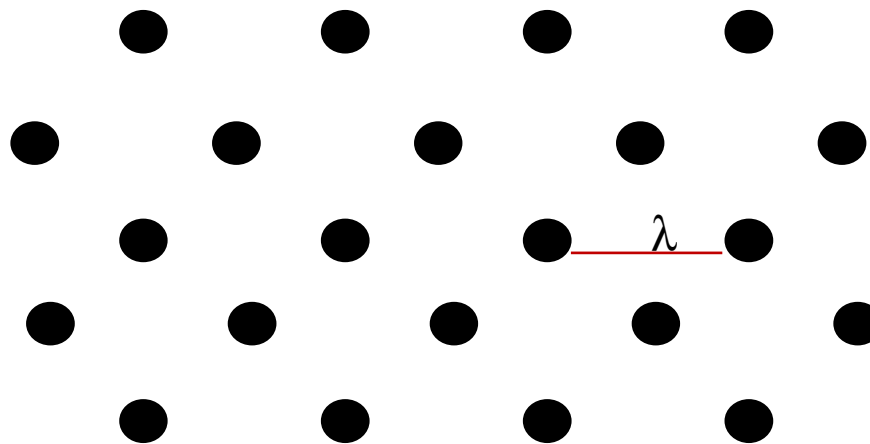
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(Wigner crystal)

$$\Psi_{WC} = A \prod_i \varphi_{\mathbf{R}_i}(\mathbf{r}_i), \quad \varphi_{\mathbf{R}}(\mathbf{r}) = \sqrt{\frac{1}{2\pi}} \exp\left[-\frac{1}{4}(\mathbf{r}-\mathbf{R})^2 + \frac{1}{2}i(\mathbf{r} \times \mathbf{R}) \cdot \mathbf{e}_z\right]$$



(Triangular structure)

$$\psi_{\vec{k}} = \bar{\rho}_{\vec{k}} \psi_m$$

Girvin, S. M., and A. H. MacDonald, and P.
M. Platzman, Phys. Rev. Lett. **54**, 581 (1985).

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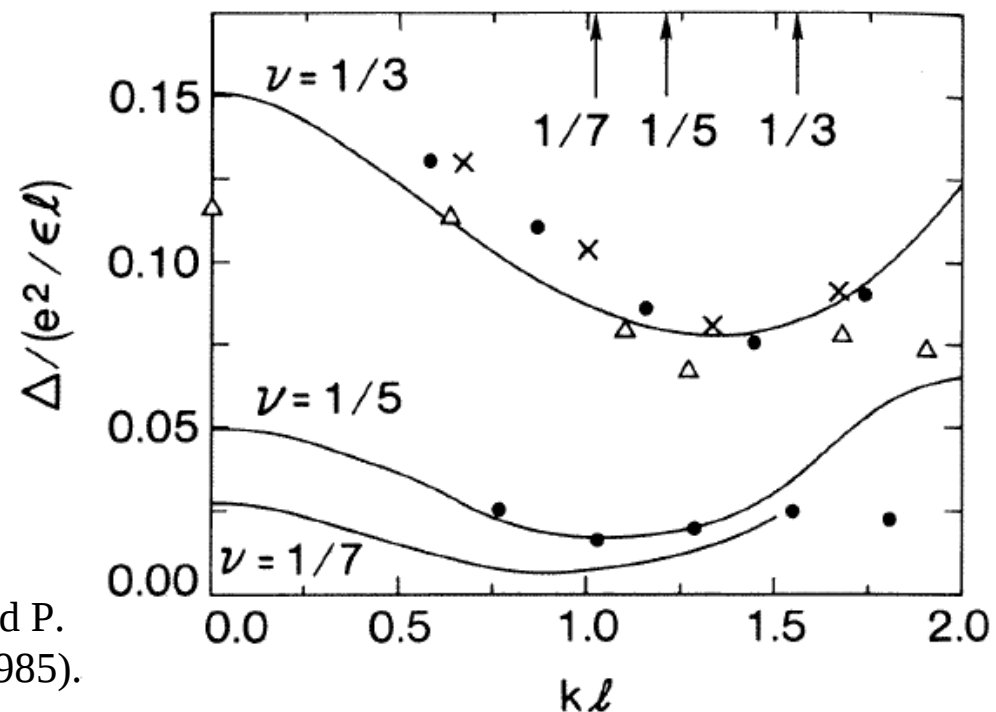
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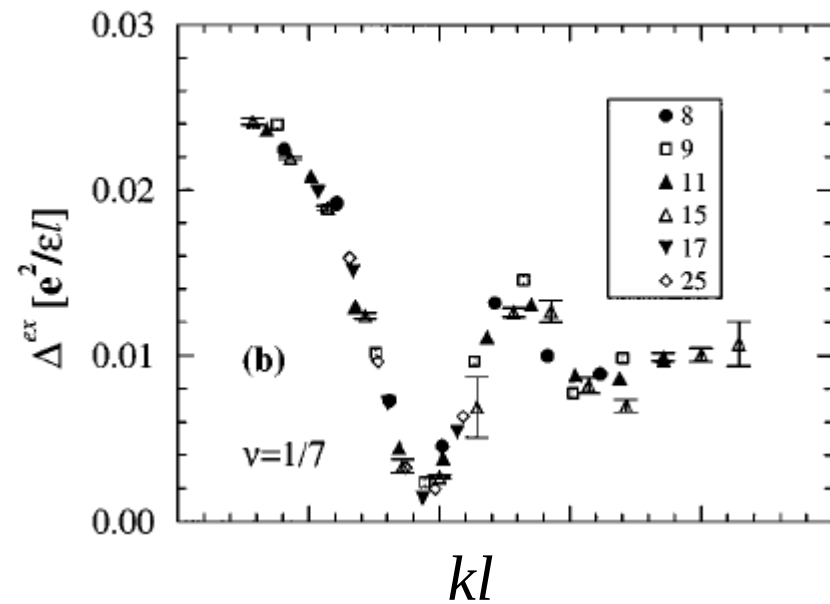
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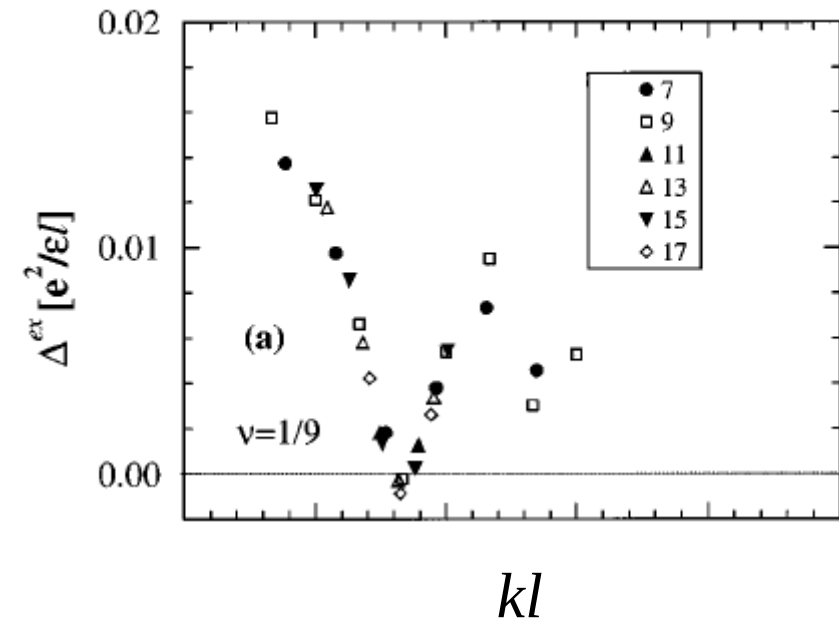
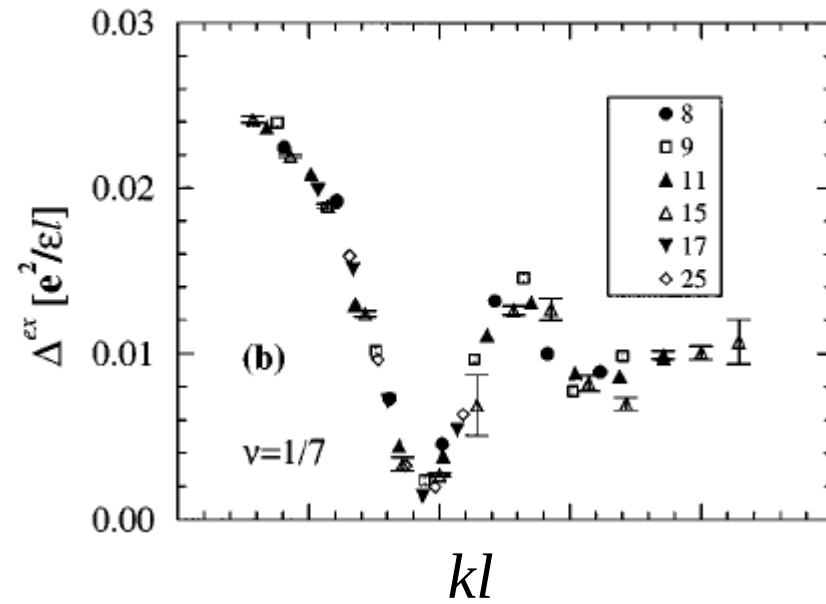
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Exciton excitation from the composite-fermion gas



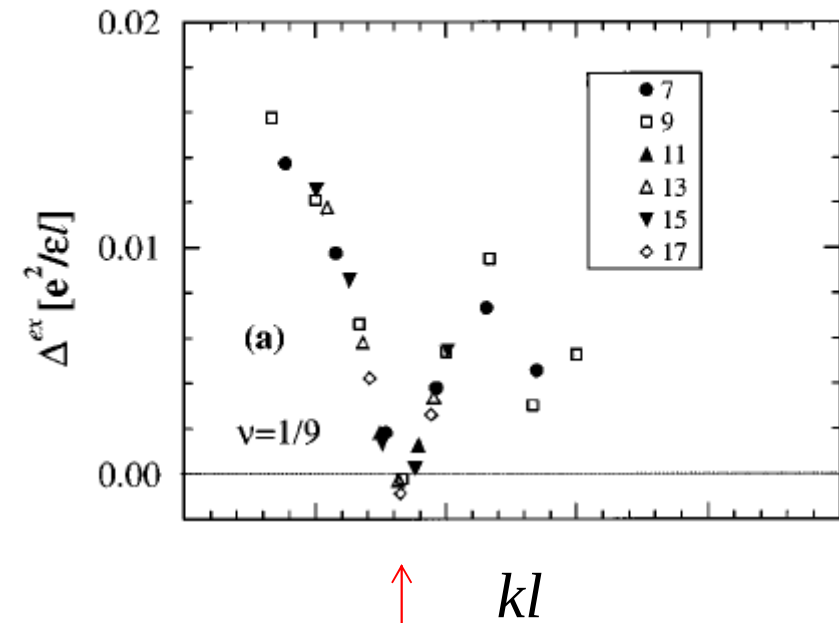
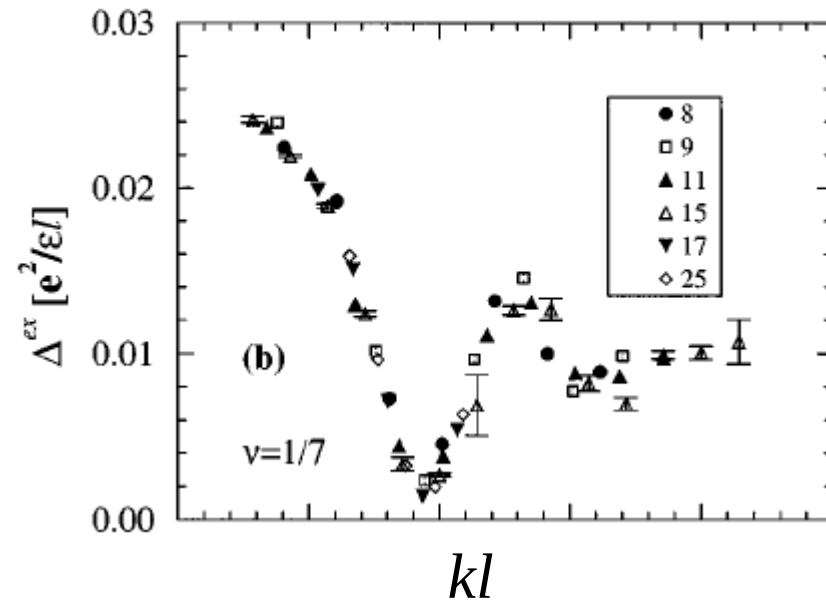
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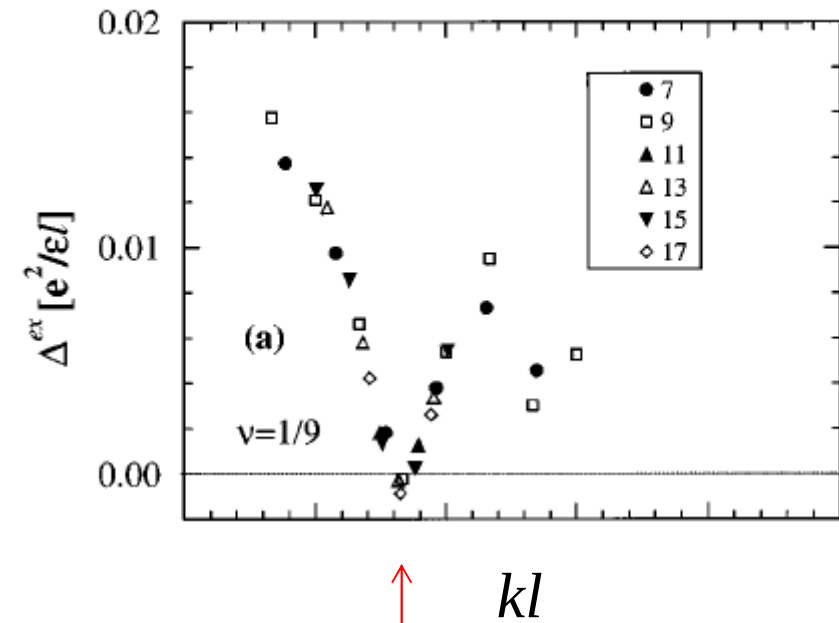
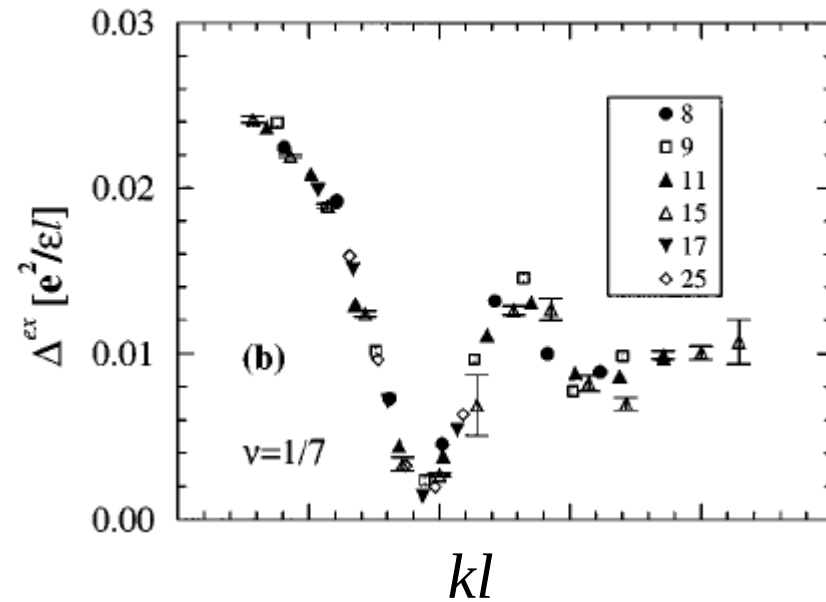
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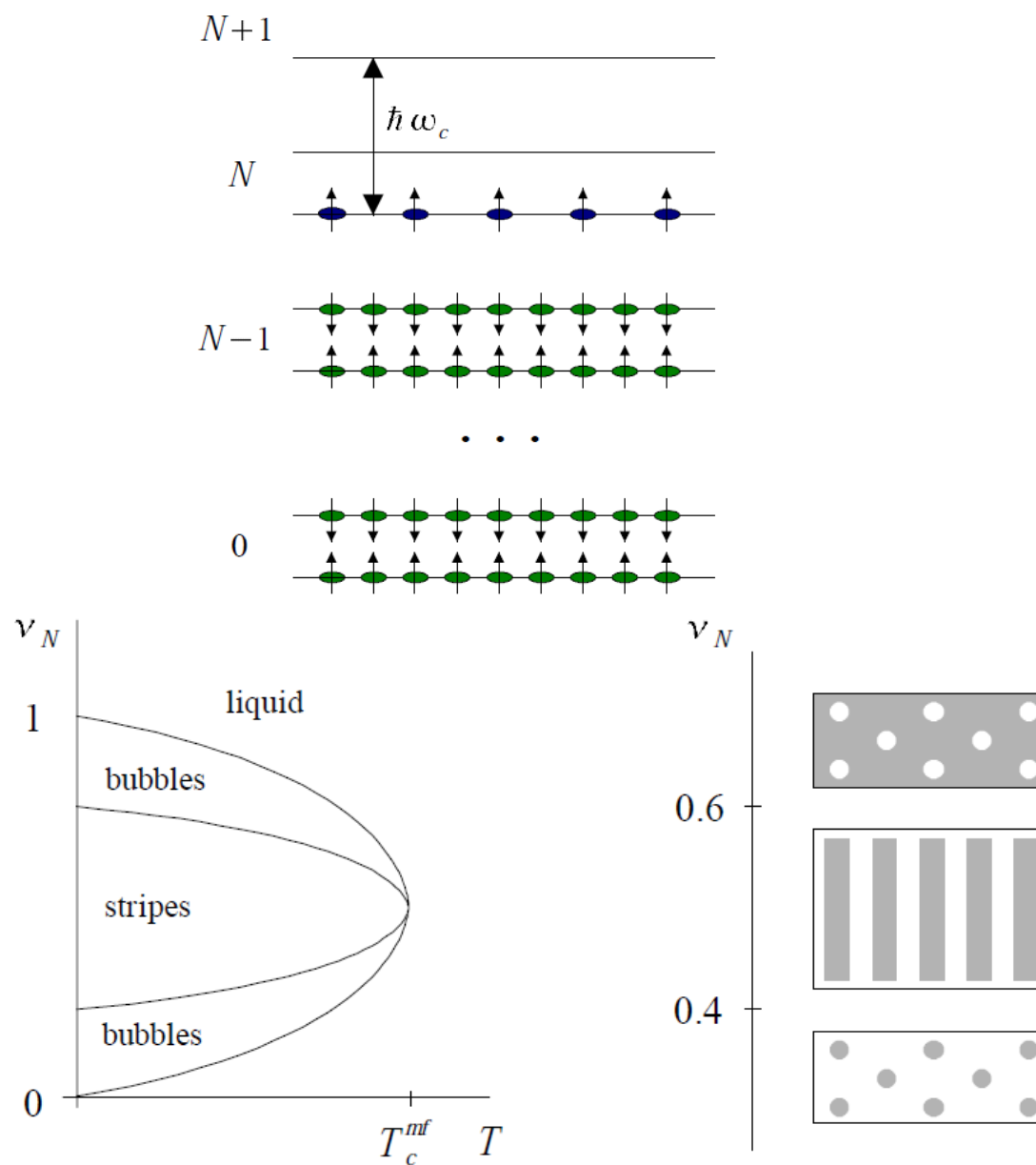


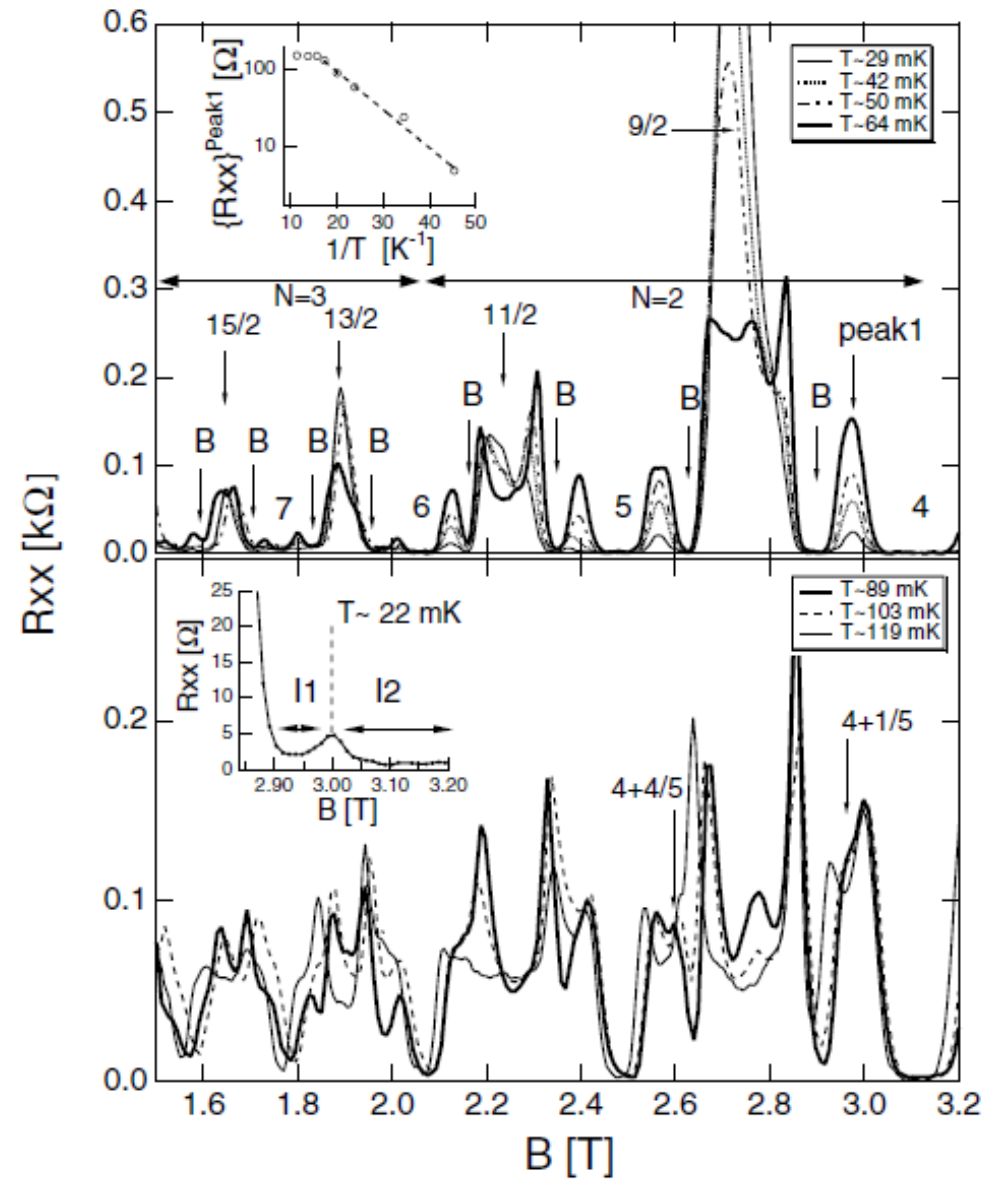
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Wigner Crystallization

R. K. Kamilla and J. K. Jain,
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Bubble crystals for a higher Landau level





G. Gevais *et al.*, Phys. Rev. Lett. **93**,
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Any quantum Hall effect in the atomic Fermi gas?

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Atoms with long-range interactions.

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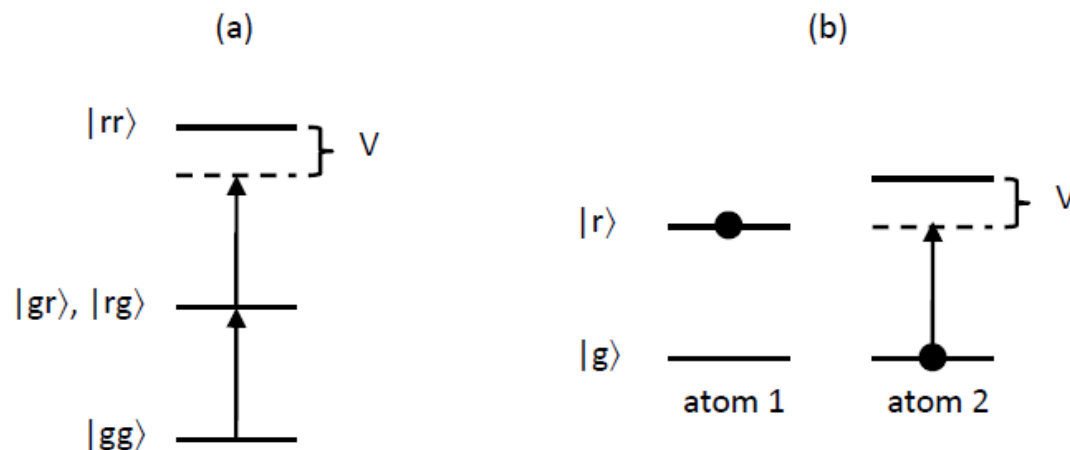
Rydberg atoms have long-range interactions.

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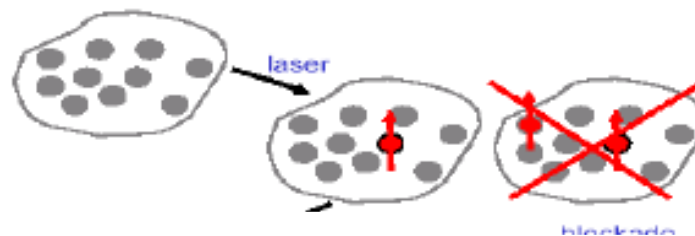
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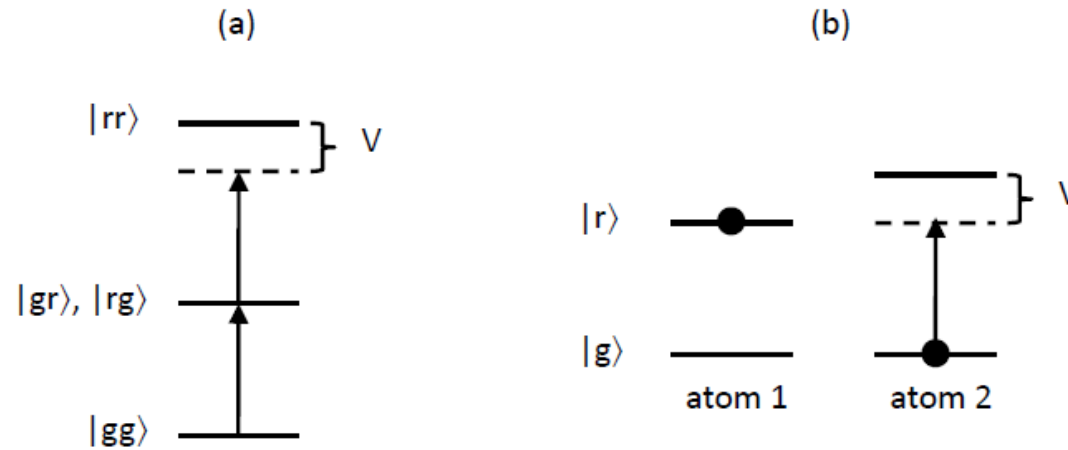
Dipole-blockaded Effect for Rydberg atoms



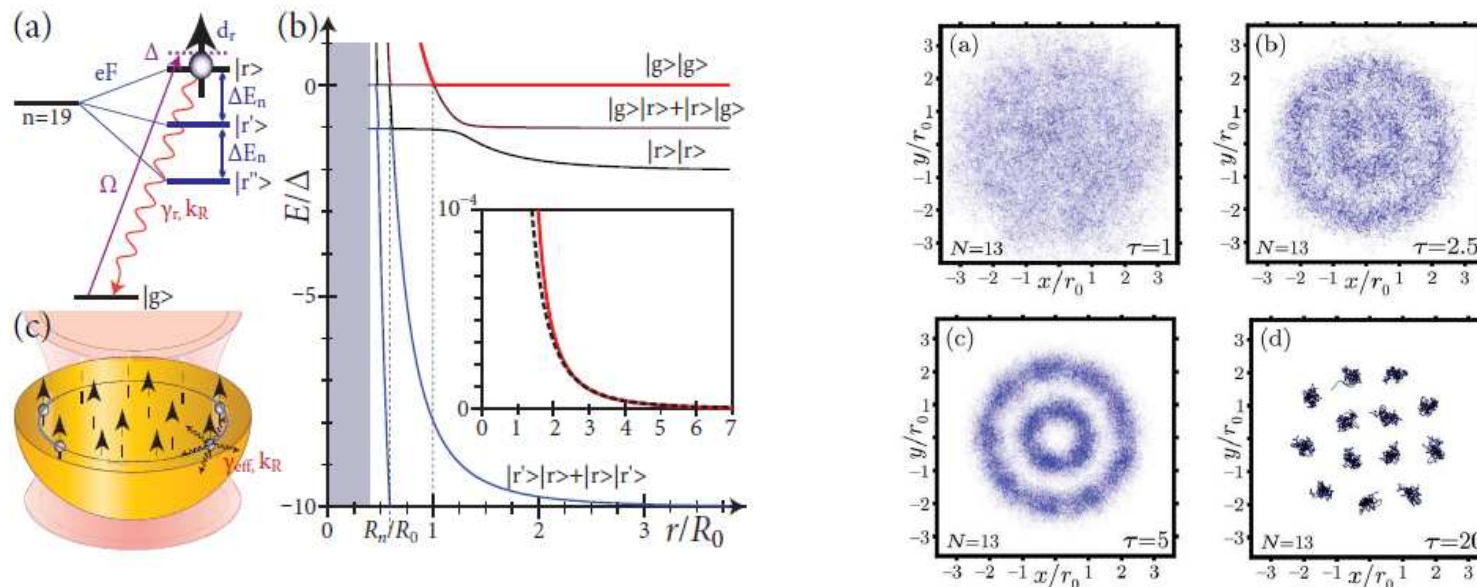
(Taken from the thesis of Tony E. Lee)



Dipole Blocked Effect



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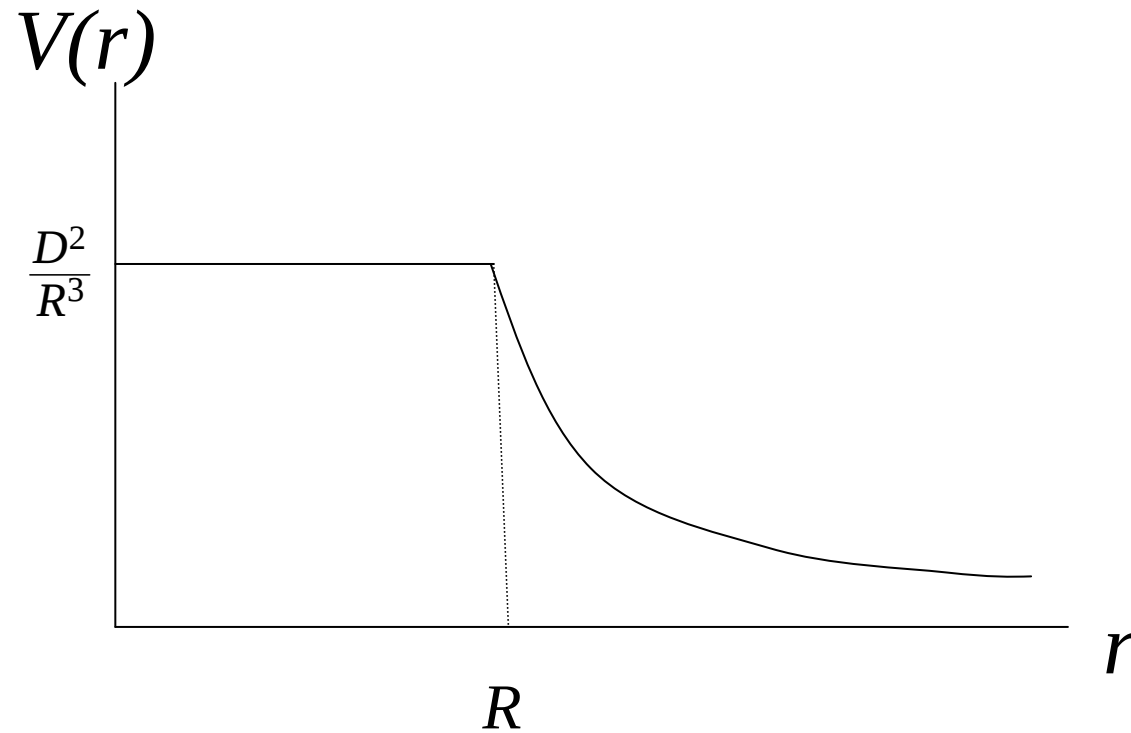


G. Pupillo *et al.*, PRL 104, 223002 (2010)

Dipole-blockaded Potential:

$$V(r) = \begin{cases} \frac{D^2}{R^3}, & r \leq R \\ \frac{D^2}{r^3}, & r > R \end{cases}$$

R : blockade radius



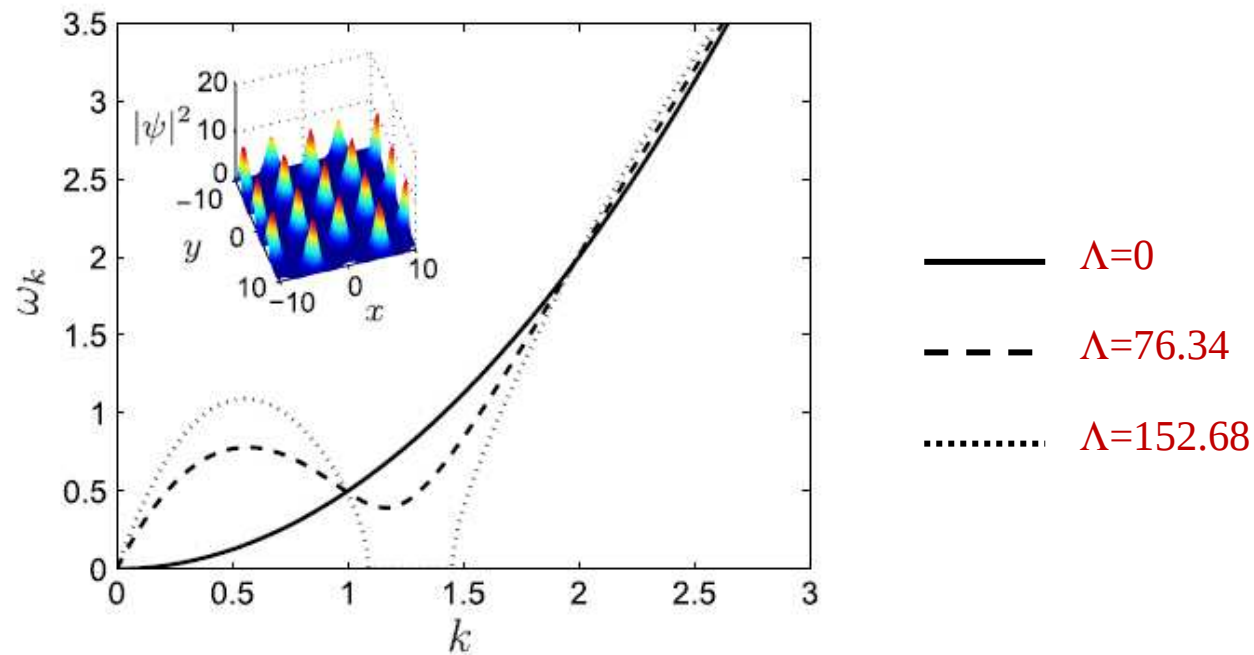
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Excitation spectra of dipole-blockaded Bose gas:

Peter Mason *et al.*, PRL 109, 045301 (2012)



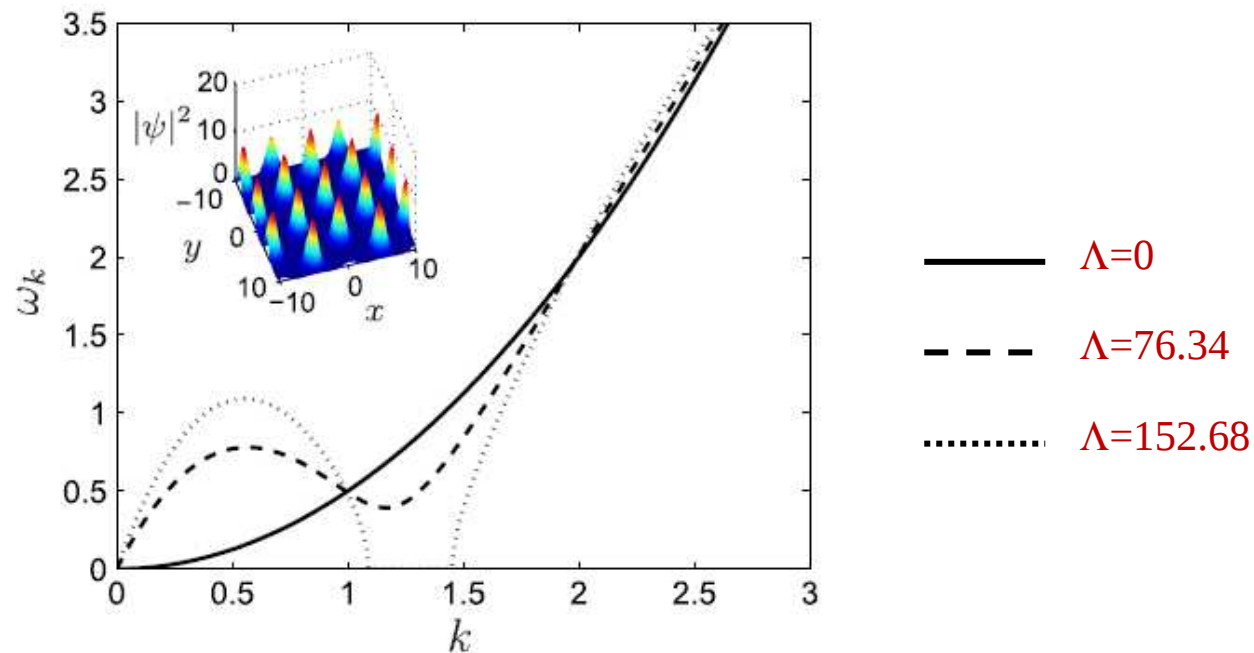
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$\omega_k < 0$ at finite k $\longrightarrow$ A crystal state

Monte Carlo simulations: F. Cinti *et al.*, PRL 105, 135301 (2010)

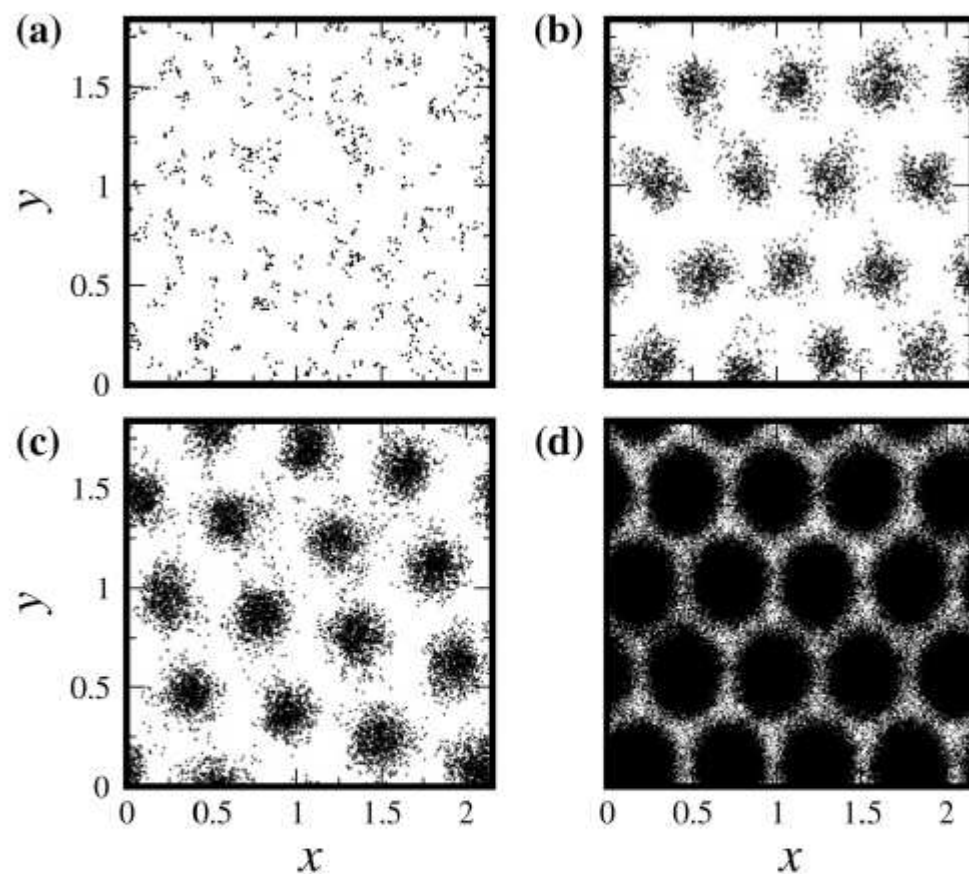
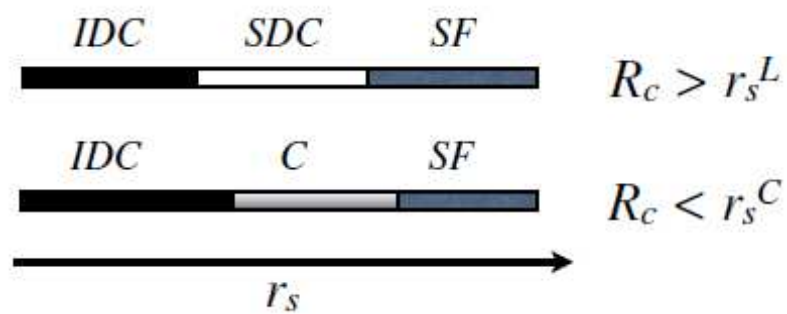


FIG. 1. Snapshots of a system of bosons interacting via potential (1), at the four different temperatures 200 (a), 20 (b), 1.0 (c) and 0.1 (d), expressed in units of ϵ_0 . Points shown are taken along individual particle world lines. The nominal value of r_s in this case is 0.14, whereas the cutoff of the potential (1) is $R_c = 0.3$.

F. Cinti *et al.*, PRL 105, 135301 (2010)

$$r_0 = mD/\hbar^2, \quad r_s = 1/\sqrt{nr_0^2}.$$

$$r_s^L = 0.08 \quad r_s^C = 0.06$$



IDC: insulting droplet crystal

SDC: superfluid droplet crystal

SF: superfluid

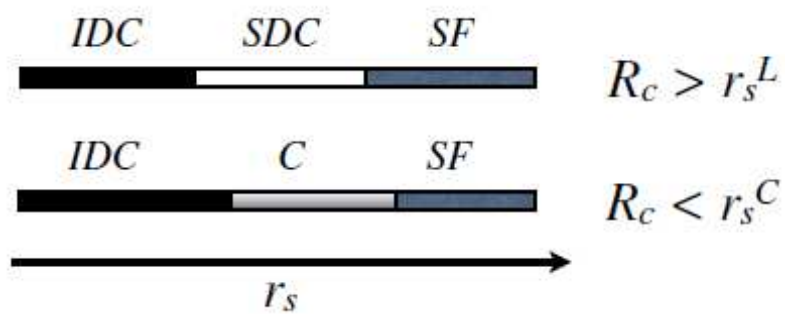
C: a single-particle crystal

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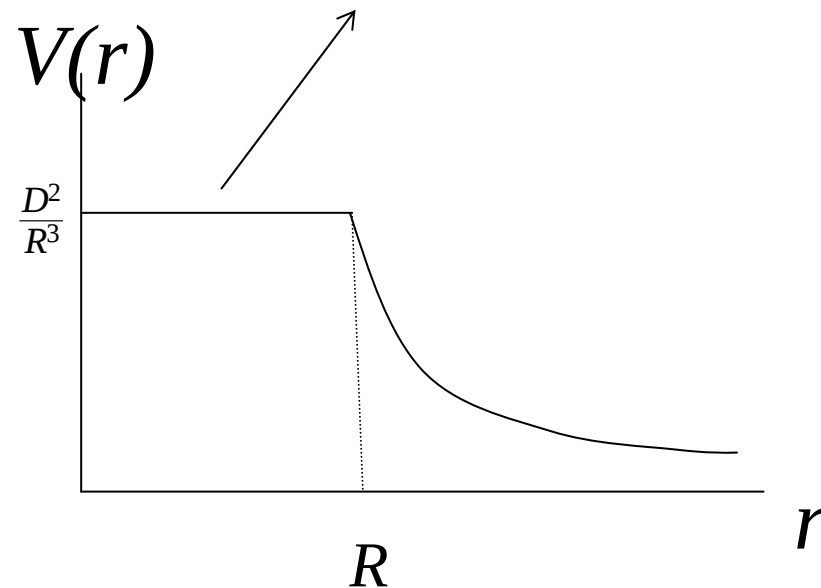
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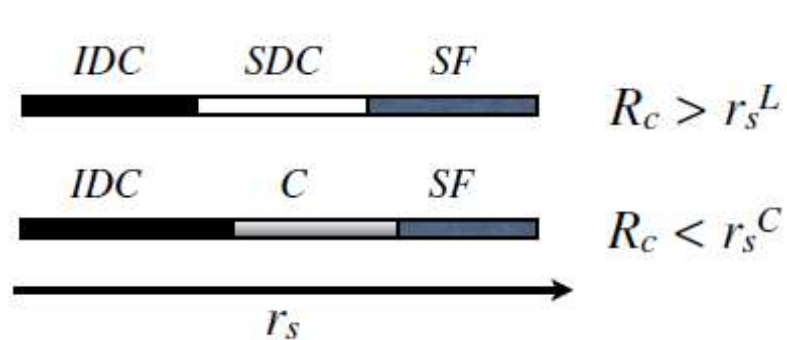
Rydberg atoms can easily form a bubble crystal.



F. Cinti *et al.*, PRL 105, 135301 (2010)

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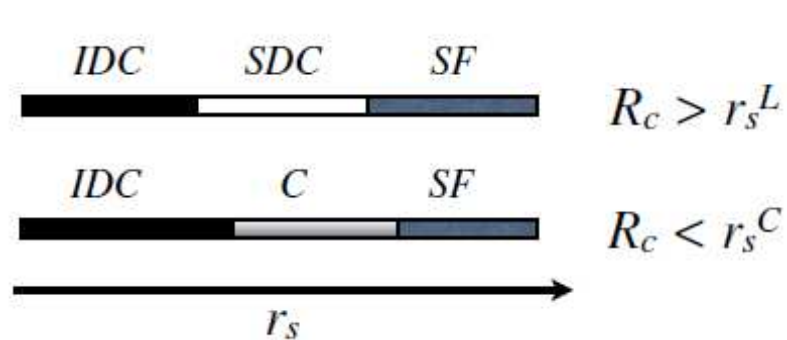
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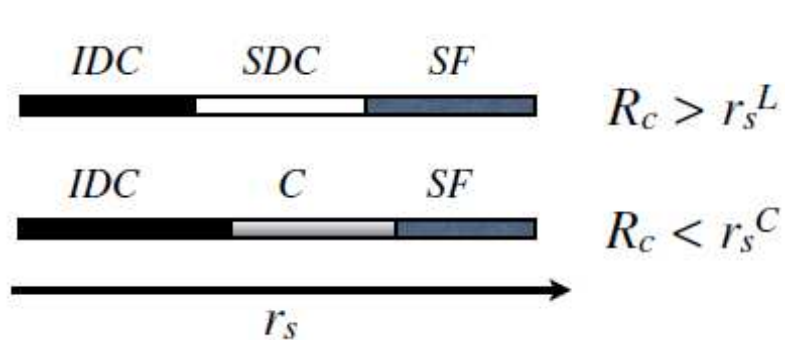
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How rotation affect crystallization?

II. Correlated Wigner Crystals for Rotating Dipolar Fermions:

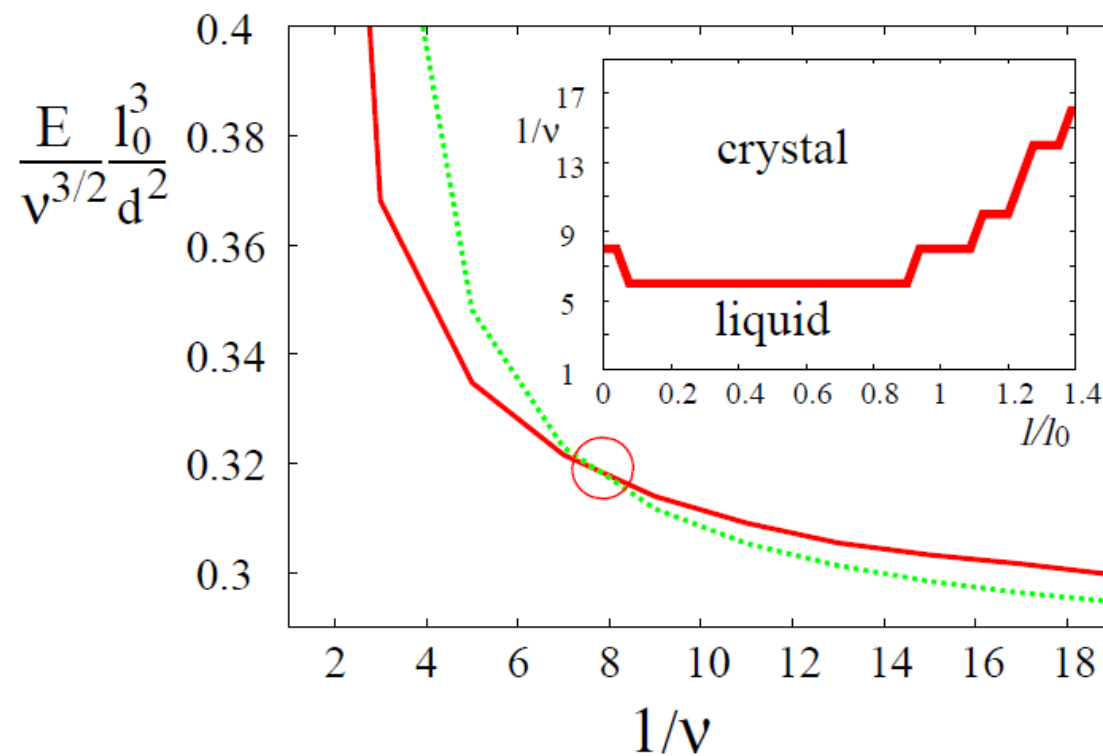
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M. A. Baranov, *et. al.*
PRL **100**, 200402 (2008).

The Wigner-crystal states exist in the regime $v < 1/7$ for zero thickness.

Collective Excitation Spectra of the Laughlin Liquid:

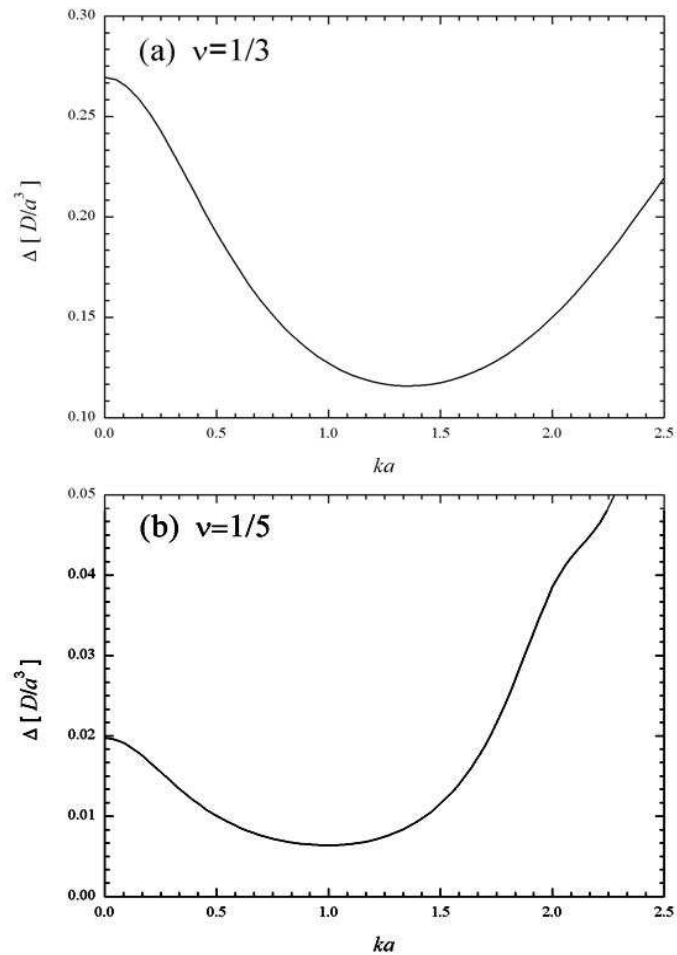
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(arXiv:cond-mat/1305.3494)

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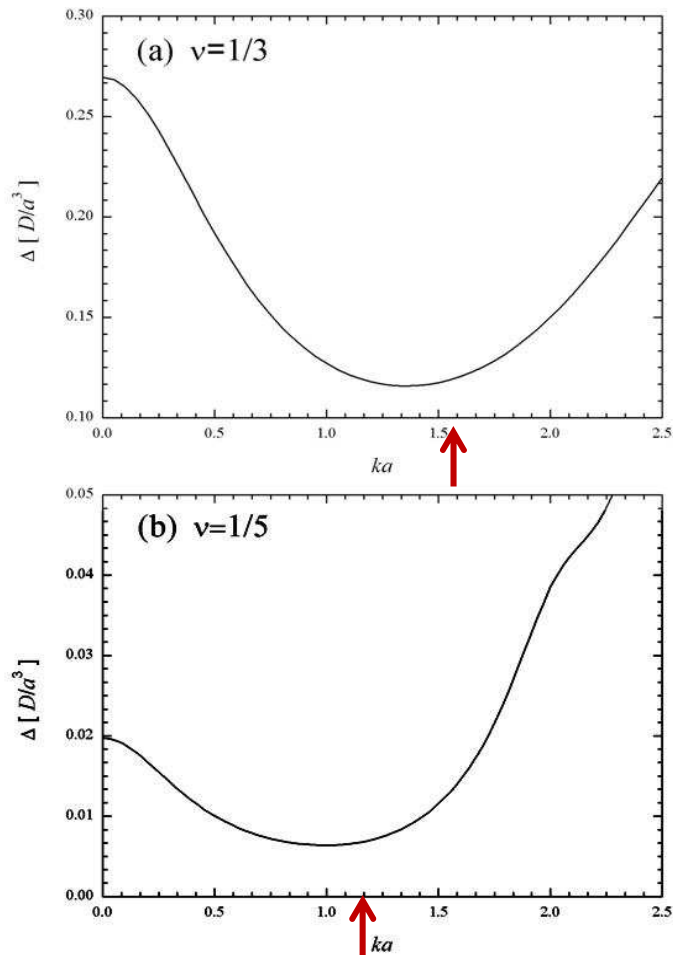
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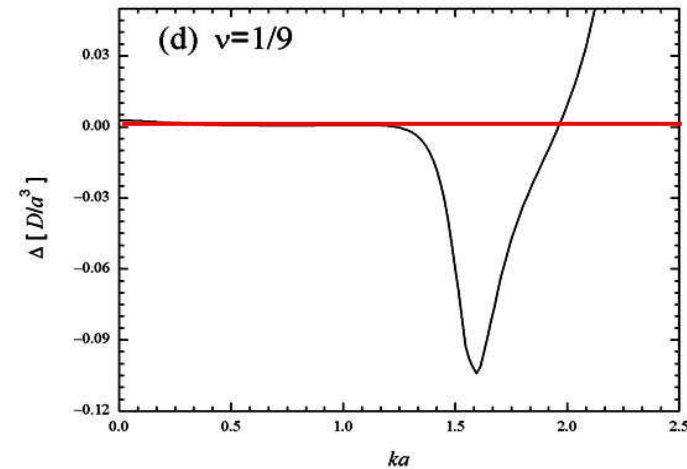
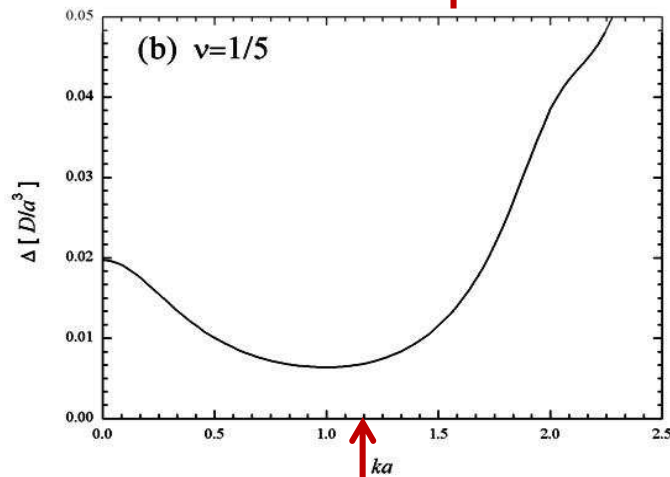
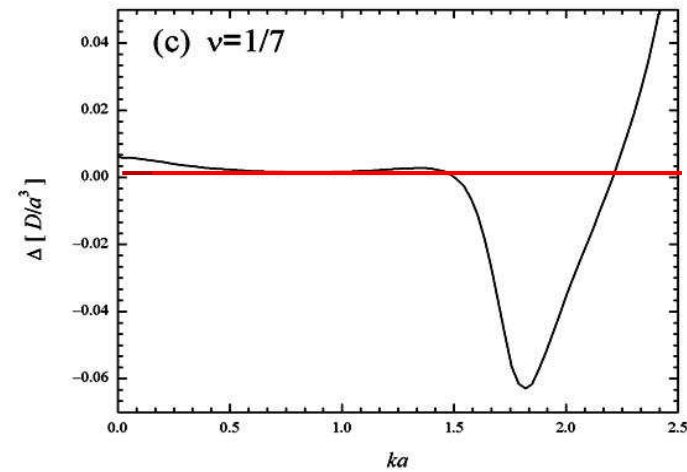
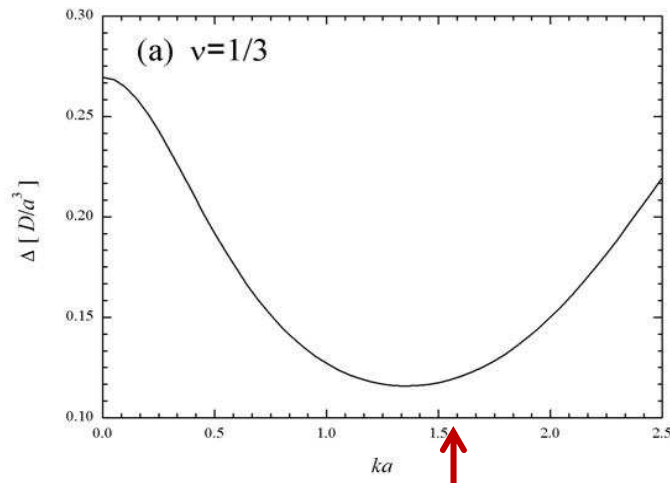


Reciprocal lattice vector of the Wigner crystal: $k_{WC} a = \sqrt{\frac{4\pi\nu}{\sqrt{3}}}$

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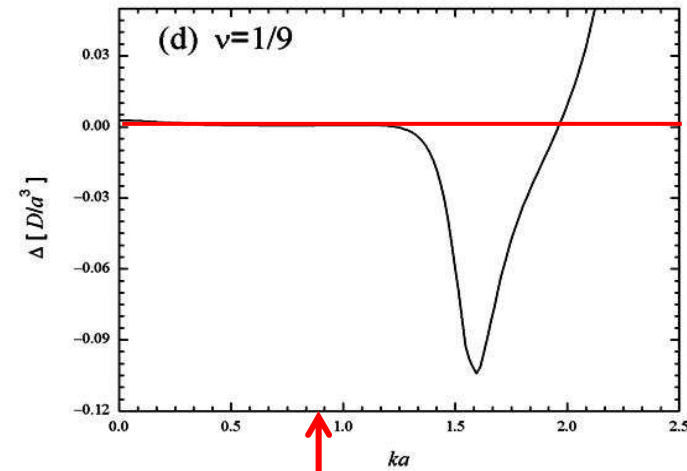
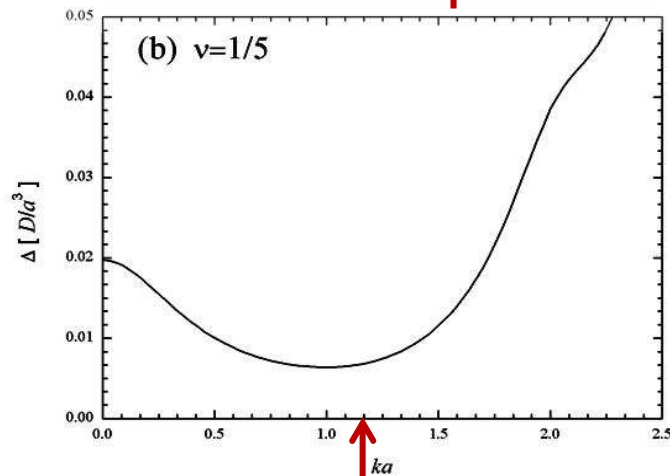
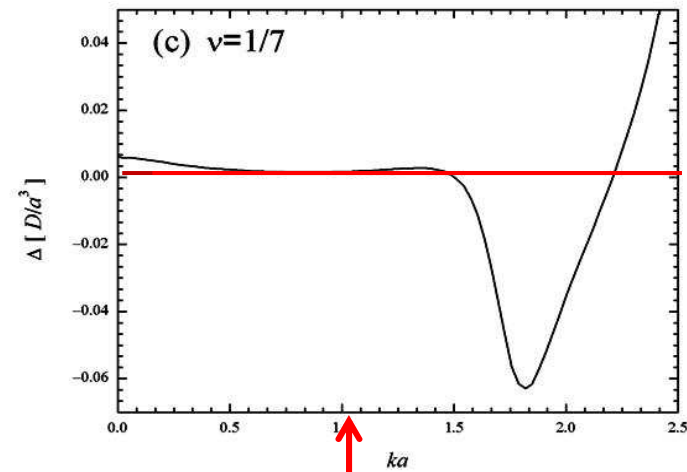
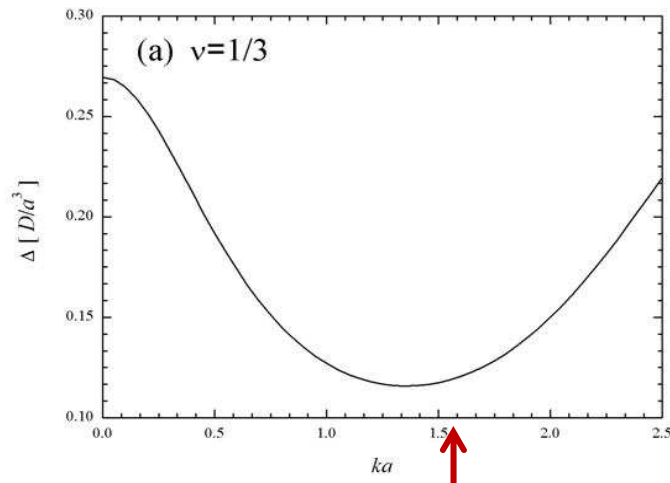


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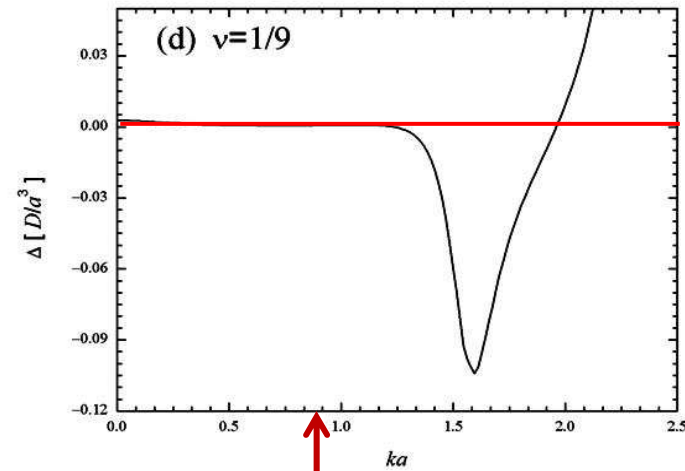
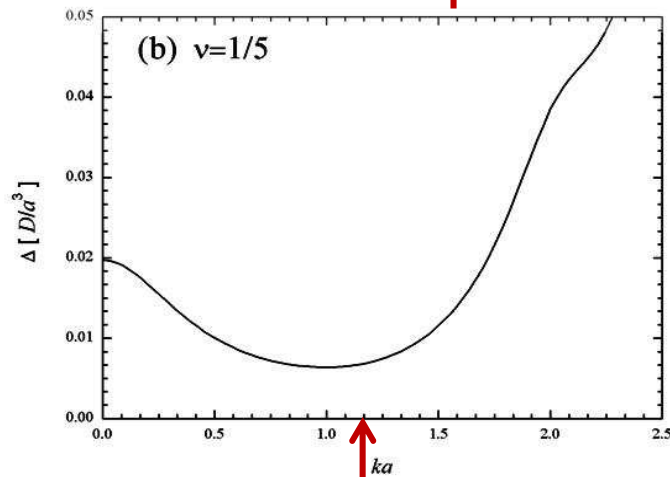
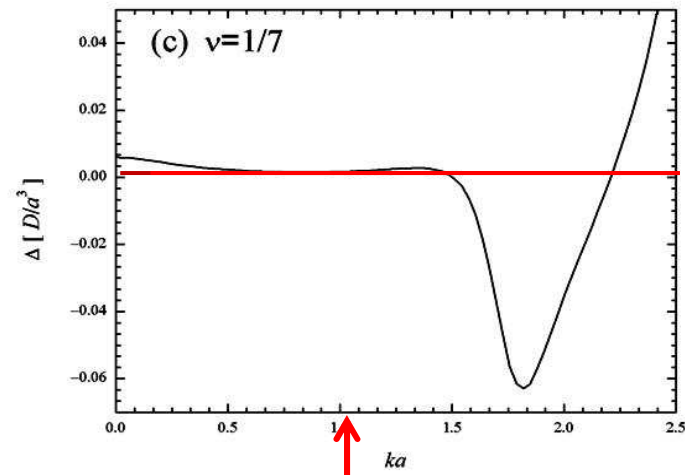
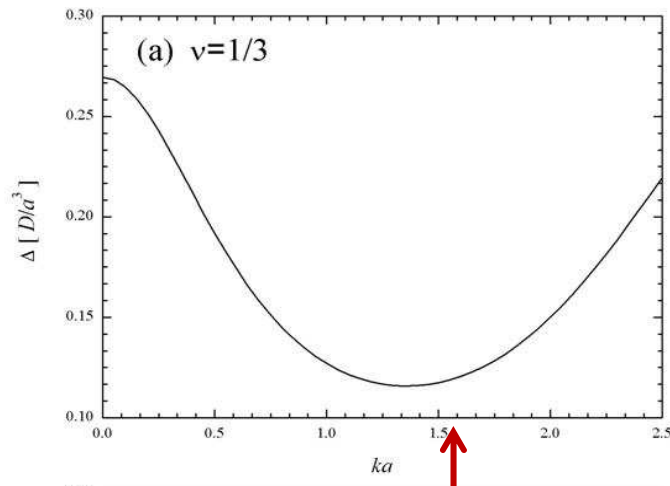
(arXiv:cond-mat/1305.3494)

$$\bar{\rho}_{\mathbf{k}} = \sum_{j=1}^N \exp(-ik\partial/\partial z_j) \exp(-ik^* z_j/2)$$



Reciprocal lattice vector of the Wigner crystal: $k_{WC} a = \sqrt{\frac{4\pi\nu}{\sqrt{3}}}$

The Laughlin liquid is unstable for $\nu < 1/5$!



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Correlated Wigner Crystal

Proposed by H. Yi and H. A. Fertig, Phys. Rev. B **58**, 4019 (1998)

$$\Psi = A \prod_{\ell < j} (Z_{\ell} - Z_j)^m \prod_k \exp[-\frac{1}{4}(\vec{r}_k - \vec{R}_k)^2 + \frac{i}{2}\vec{r}_k \times \vec{R}_k \bullet \vec{z}]$$

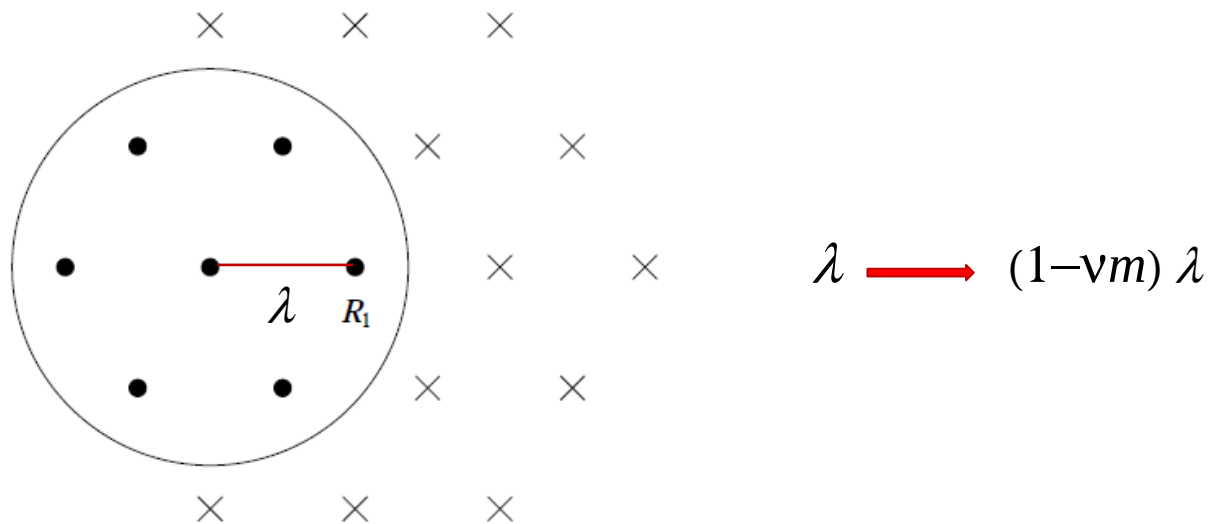
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To obtain the energy: Monte Carlo Calculations.

$N=20, 40, 60, 80, 100, 120, 140, 160$ and 180 particles.



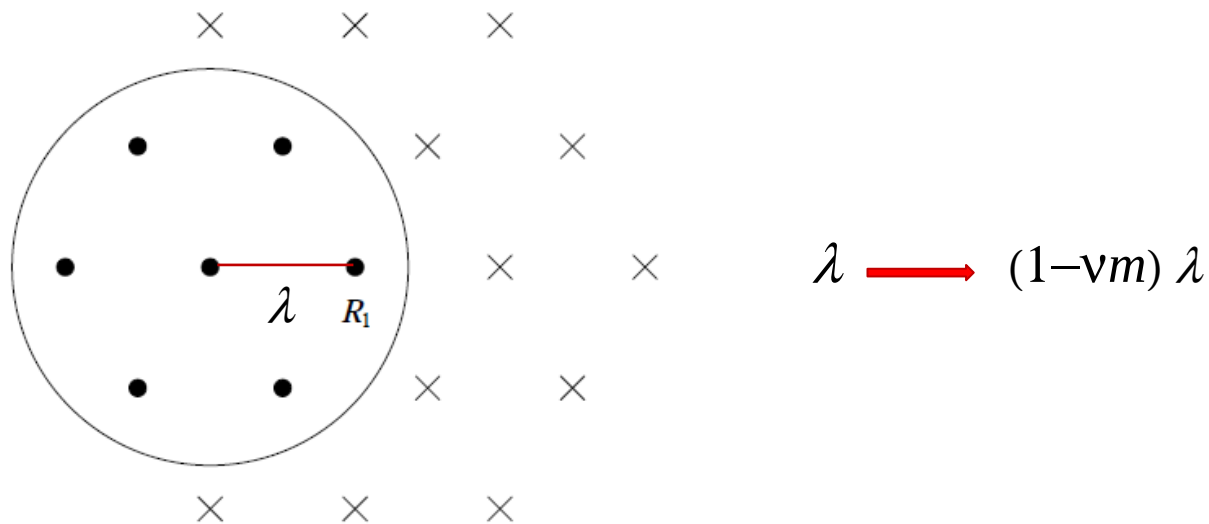
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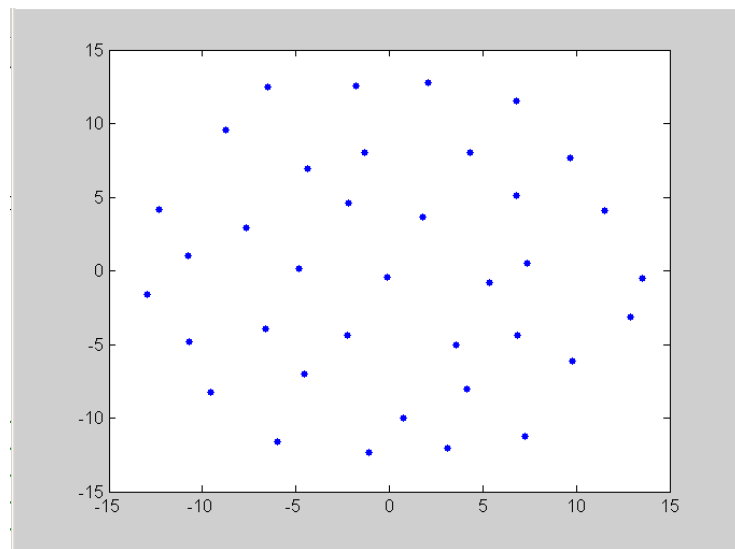
$$\Psi = A \prod_{\ell < j} (Z_{\ell} - Z_j)^m \prod_k \exp[-\frac{1}{4}(\vec{r}_k - \vec{R}_k)^2 + \frac{i}{2}\vec{r}_k \times \vec{R}_k \cdot \vec{z}]$$

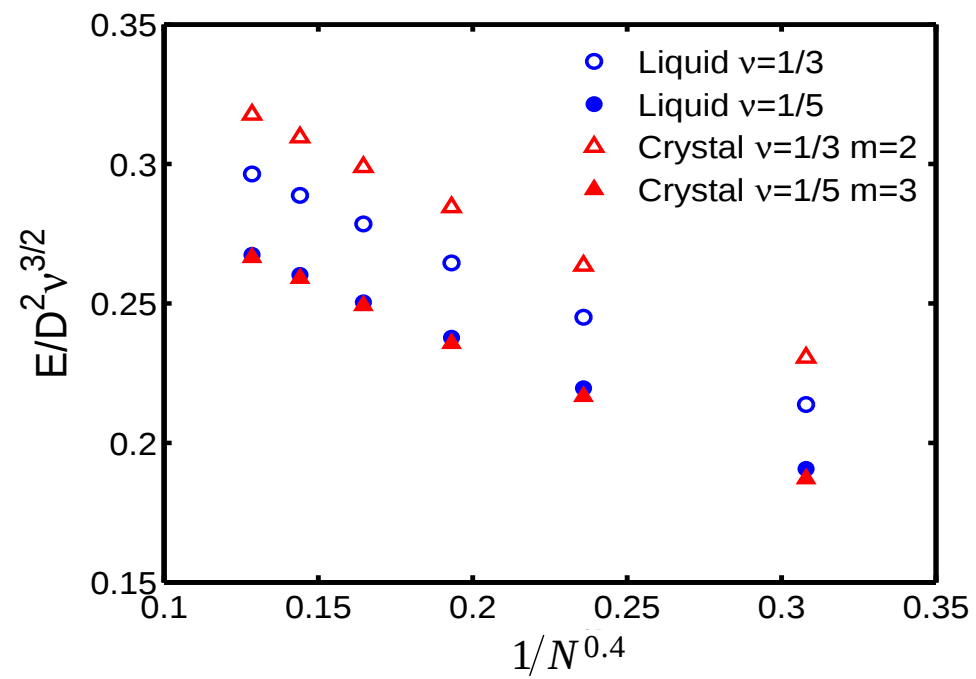
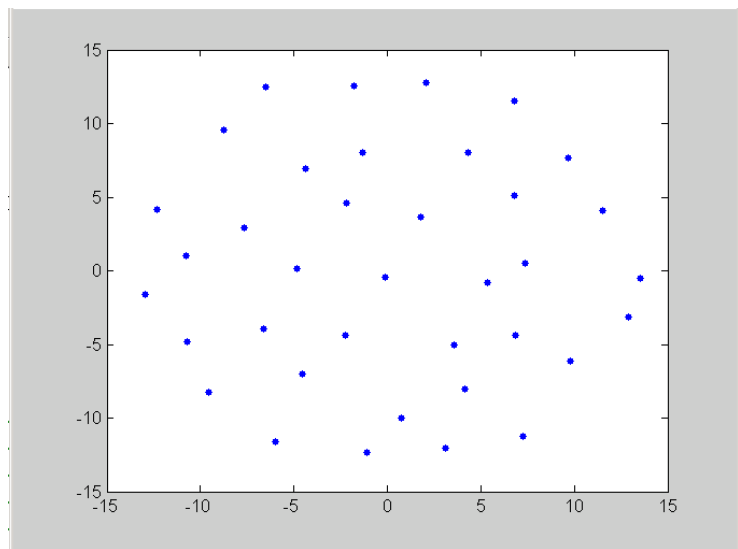
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$$\Psi = A \prod_{\ell < j} (Z_{\ell} - Z_j)^m \prod_k \text{Exp}[-\frac{1}{4}(\vec{r}_k - (1-\nu m)\vec{R}_k)^2 + \frac{i}{2}\vec{r}_k \times (1-\nu m)\vec{R}_k \cdot \vec{z} - \nu m(1-\nu m)|\vec{R}_k|^2/4]$$





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