

Reflected Entropy and Entanglement Wedge Cross Section with the First Order Correction

Mitsuhiro Nishida

(Gwangju Institute of Science and Technology)

[\[arXiv:1909.02806\]](#)

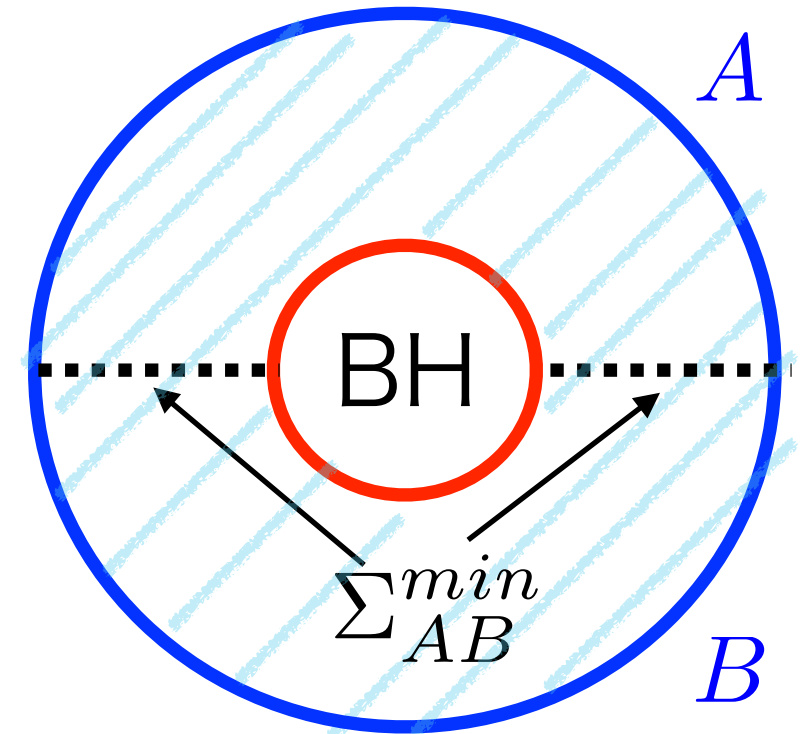
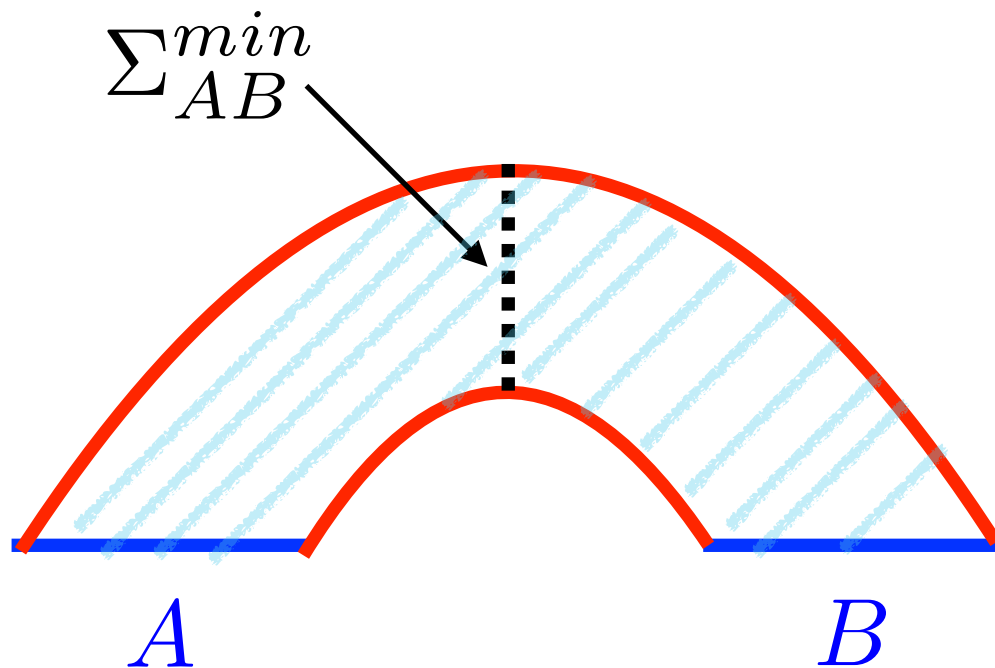
with Hyun-Sik Jeong and Keun-Young Kim

(Gwangju Institute of Science and Technology)

Entanglement wedge cross section

$$E_W(A : B) = \frac{\text{Area}[\Sigma_{AB}^{min}]}{4G_N}$$

minimal cross section
in entanglement wedge



Proposals of the dual quantities

- Entanglement of purification [T. Takayanagi, K. Umemoto, 2017]
- Logarithmic negativity [P. Nguyen et al., 2017]
- Odd entropy [J. Kudler-Flam, S. Ryu, 2018]
- Odd entropy [K. Tamaoka, 2018]
- Reflected entropy [S. Dutta, T. Faulkner, 2019]

Reflected entropy $S_R(A : B)$

entanglement entropy of a canonical purification of ρ_{AB}

$$\begin{aligned} \rho_{AB} &= \frac{1}{2} |+\rangle_A |+\rangle_B \langle +|_A \langle +|_B + \frac{1}{2} |-\rangle_A |-\rangle_B \langle -|_A \langle -|_B \\ |\sqrt{\rho_{AB}}\rangle &\in \mathcal{H}_A \otimes \mathcal{H}_B \otimes \mathcal{H}_A^* \otimes \mathcal{H}_B^* \\ &:= \frac{1}{\sqrt{2}} |+\rangle_A |+\rangle_B |+\rangle_{A^*} |+\rangle_{B^*} + \frac{1}{\sqrt{2}} |-\rangle_A |-\rangle_B |-\rangle_{A^*} |-\rangle_{B^*} \end{aligned}$$

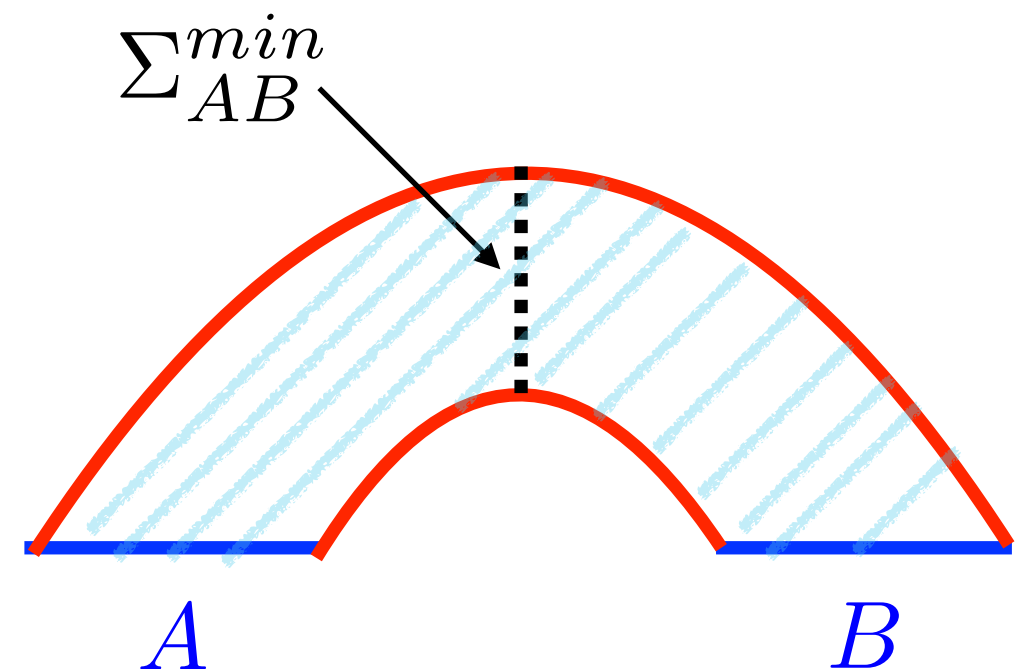
$$S_R(A : B) := -\text{Tr}_{AA^*} [\rho_{AA^*} \log \rho_{AA^*}]$$

$$\rho_{AA^*} := \text{Tr}_{BB^*} |\sqrt{\rho_{AB}}\rangle \langle \sqrt{\rho_{AB}}|$$

Holographic duality

[S. Dutta, T. Faulkner, 2019]

$$S_R(A : B) = 2E_W(A : B)$$



Generalization by a replica index m

[S. Dutta, T. Faulkner, 2019]

$$S_{mR}(A : B) = 2E_{mW}(A : B)$$

$S_{mR}(A : B)$: reflected entropy with
a canonical purification of ρ_{AB}^m

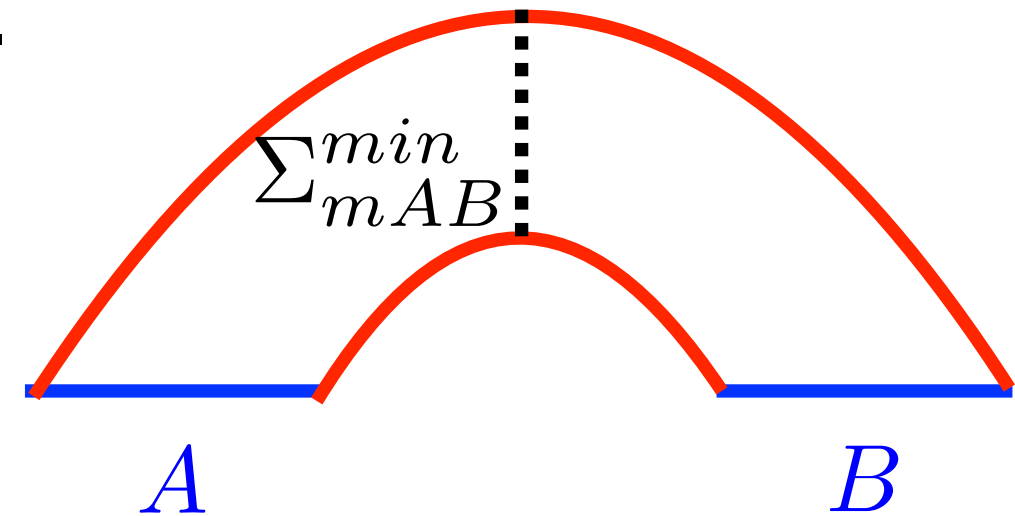
$E_{mW}(A : B)$: entanglement wedge cross section
in the gravity dual of ρ_{AB}^m

Our motivation

We want to check and understand this relation
by a simple example.

What we did

We explicitly check $S_{mR}(A : B) = 2E_{mW}(A : B)$
up to the first order in $m - 1$.



$$S_{mR}(A : B) = \frac{c}{3} \log \left[\frac{1 + \sqrt{z}}{1 - \sqrt{z}} \right] - \frac{2c(m-1)}{3} \frac{\sqrt{z} \log z}{1 - z} + \dots$$

exactly matched
with $c = \frac{3}{2G_N}$

$$2E_{mW}(A : B) = \frac{1}{2G_N} \log \left[\frac{1 + \sqrt{z}}{1 - \sqrt{z}} \right] - \frac{m-1}{G_N} \frac{\sqrt{z} \log z}{1 - z} + \dots$$

Outline

1. Review of $S_R(A : B) = 2E_W(A : B)$
with thermal state and black hole
2. Our computation of $S_{mR}(A : B)$
and $E_{mW}(A : B)$ up to the first order
in $m - 1$

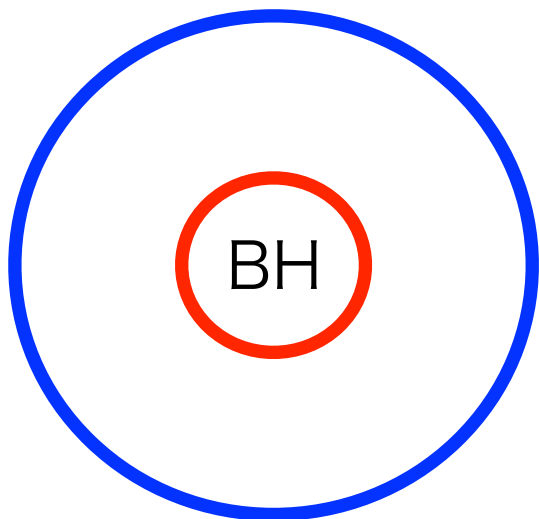
Thermofield double and two-sided AdS black hole

[J. M. Maldacena, 2001]

Thermal state on $\mathcal{H}_{AB}$

$$\rho_{AB} = \sum_n \frac{e^{-\beta E_n}}{Z} |n\rangle_{AB} \langle n|_{AB}$$

one-sided AdS BH

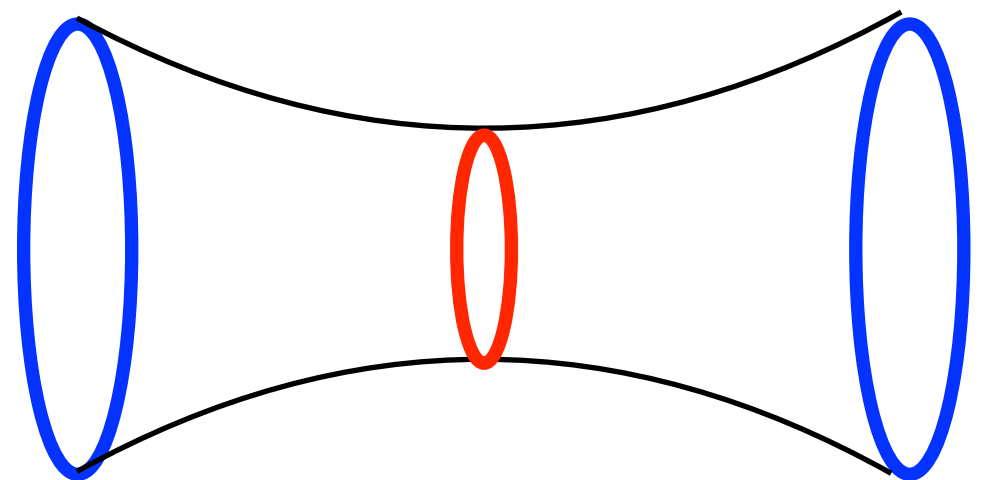


Thermofield double

on $\mathcal{H}_{AB} \otimes \mathcal{H}_{AB}^*$

$$\begin{aligned} & |\sqrt{\rho_{AB}}\rangle \\ &= \sum_n \frac{e^{-\beta E_n/2}}{\sqrt{Z}} |n\rangle_{AB} |n\rangle_{A^* B^*} \end{aligned}$$

two-sided AdS BH



Reflected entropy of thermal state

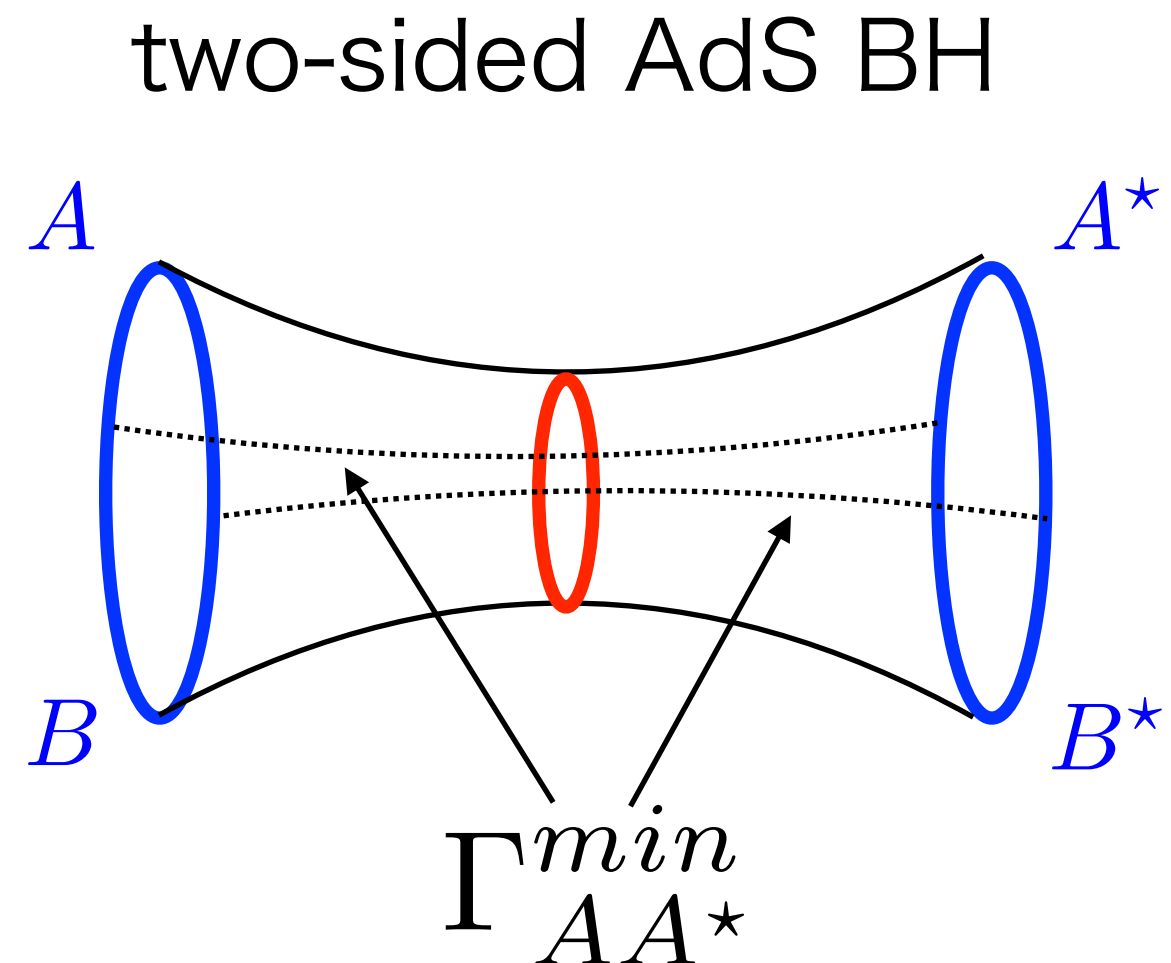
= Entanglement entropy of thermofield double

= Holographic entanglement entropy
in two-sided AdS black hole

$$S_R(A : B)$$

$$= S_{AA^*}$$

$$= \frac{\text{Area}[\Gamma_{AA^*}^{min}]}{4G_N}$$

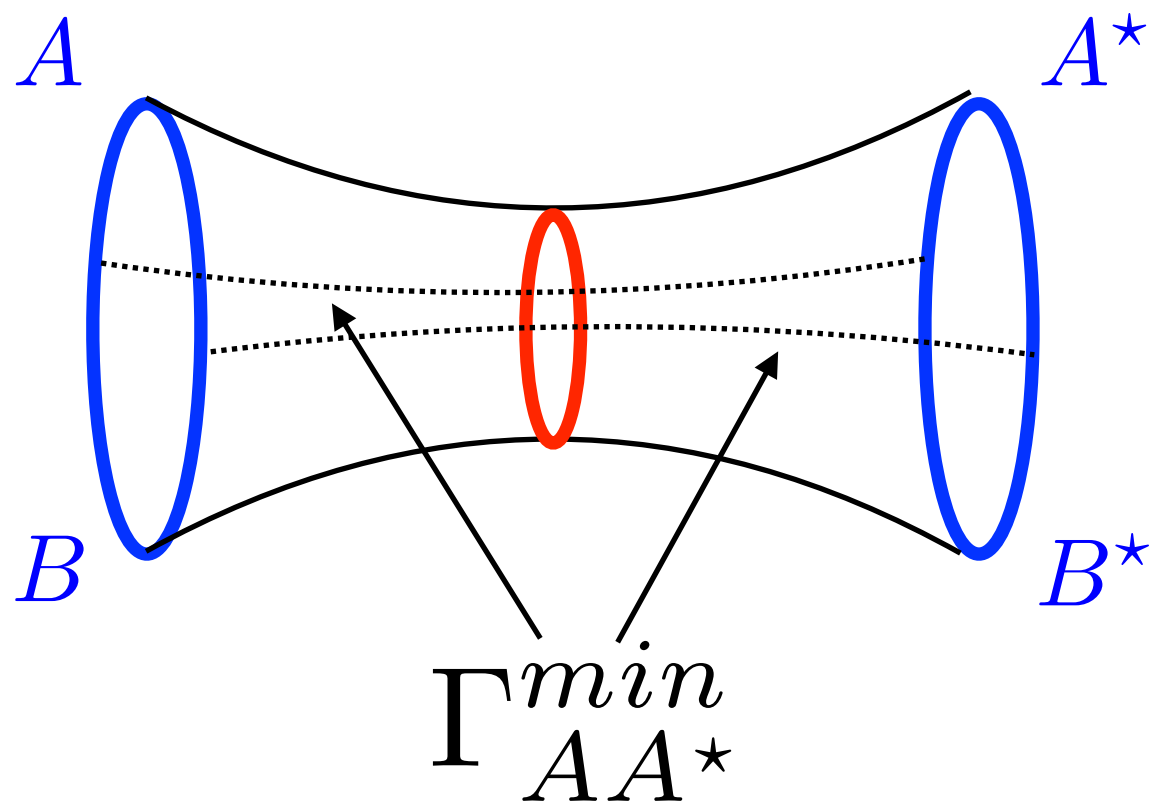


Holographic duality

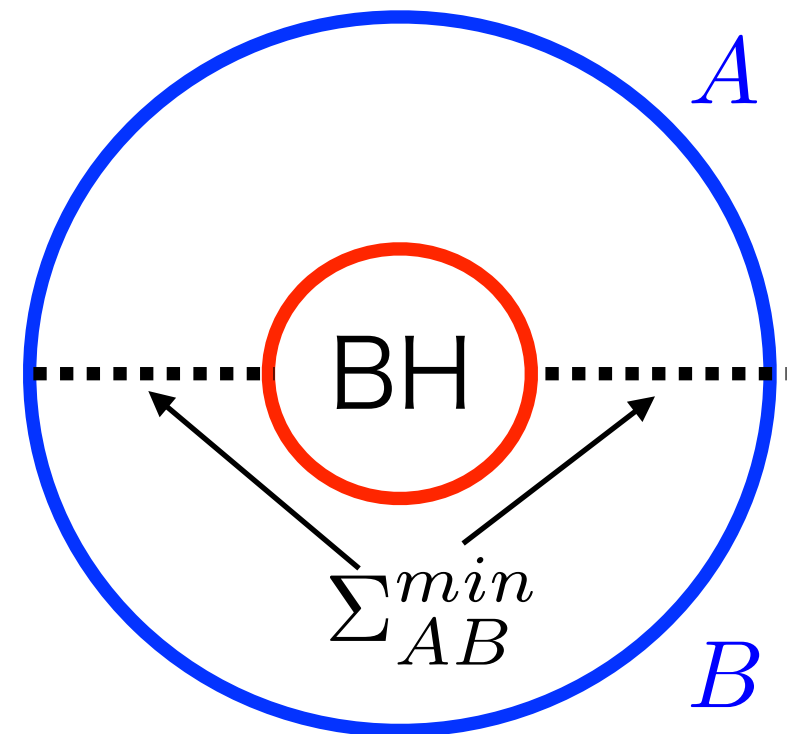
[S. Dutta, T. Faulkner, 2019]

$$S_R(A : B) = \frac{\text{Area}[\Gamma_{AA^*}^{min}]}{4G_N} = 2 \frac{\text{Area}[\Sigma_{AB}^{min}]}{4G_N} = 2E_W(A : B)$$

two-sided AdS BH



one-sided AdS BH



Outline

1. Review of $S_R(A : B) = 2E_W(A : B)$
with thermal state and black hole
2. Our computation of $S_{mR}(A : B)$
and $E_{mW}(A : B)$ up to the first order
in $m - 1$

Generalization by a replica index m

$$S_{mR}(A : B) = 2E_{mW}(A : B)$$

[S. Dutta, T. Faulkner, 2019]

$S_{mR}(A : B)$: reflected entropy with
a canonical purification of ρ_{AB}^m

$E_{mW}(A : B)$: entanglement wedge cross section
in the gravity dual of ρ_{AB}^m

Our motivation

We want to check and understand this relation
by a simple example
(the first order correction in $m - 1$).

Computation of $S_{mR}(A : B)$

$$S_{mR}(A : B) = \lim_{n \rightarrow 1} \frac{1}{1-n} \log \frac{\left\langle \sigma_{g_A}(x_1) \sigma_{g_A^{-1}}(x_2) \sigma_{g_B}(x_3) \sigma_{g_B^{-1}}(x_4) \right\rangle_{CFT^{\otimes mn}}}{\left(\left\langle \sigma_{g_m}(x_1) \sigma_{g_m^{-1}}(x_2) \sigma_{g_m}(x_3) \sigma_{g_m^{-1}}(x_4) \right\rangle_{CFT^{\otimes m}} \right)^n}$$

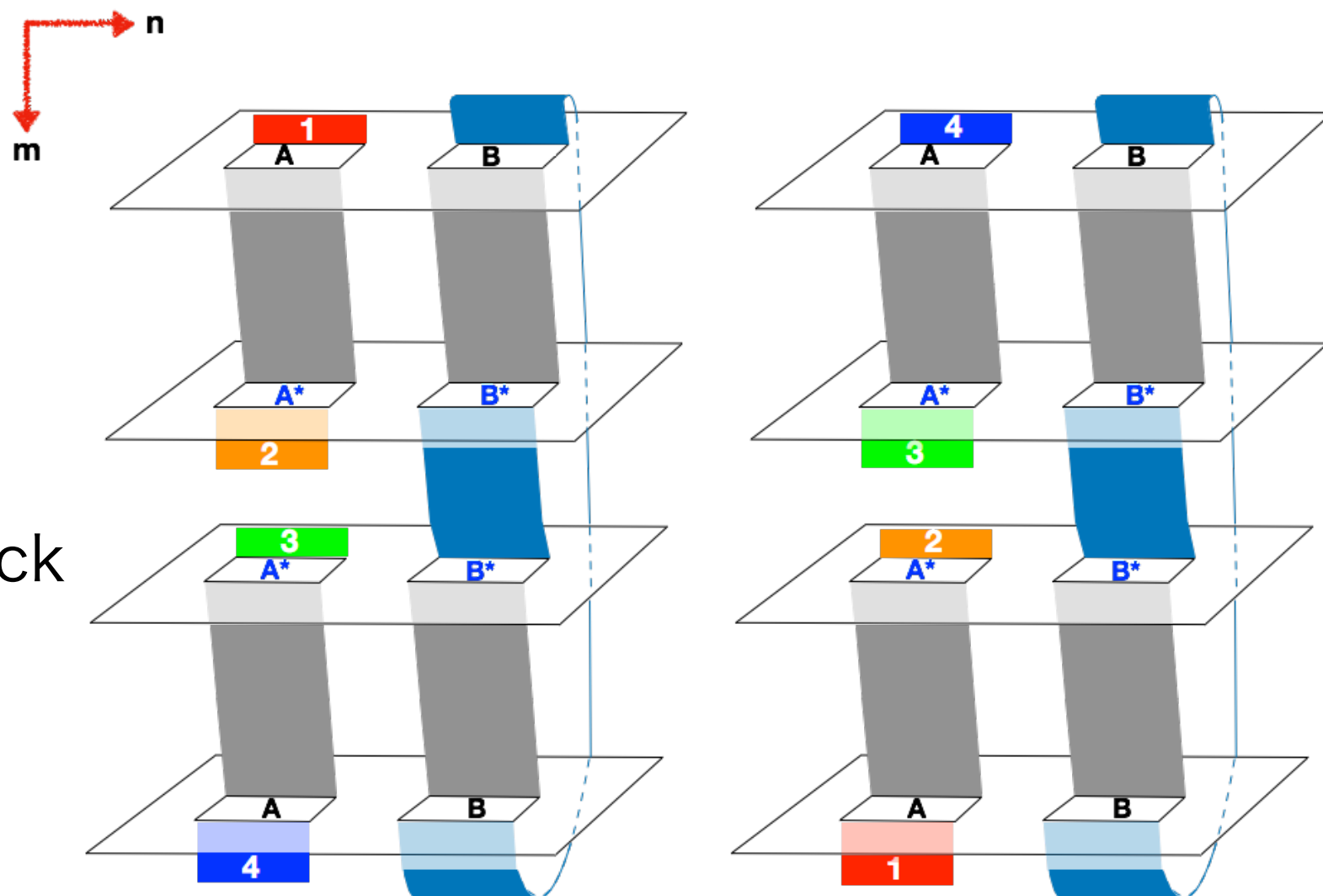
• In the 2d holographic CFT, the four point function can be approximated by a single conformal block.

[T. Hartman, 2013]

• We used a perturbative expansion of conformal block in $m-1$ and $n-1$.

[A. L. Fitzpatrick, J. Kaplan, M. T. Walters, 2014]

Replica manifold ($m = 4, n = 2$)

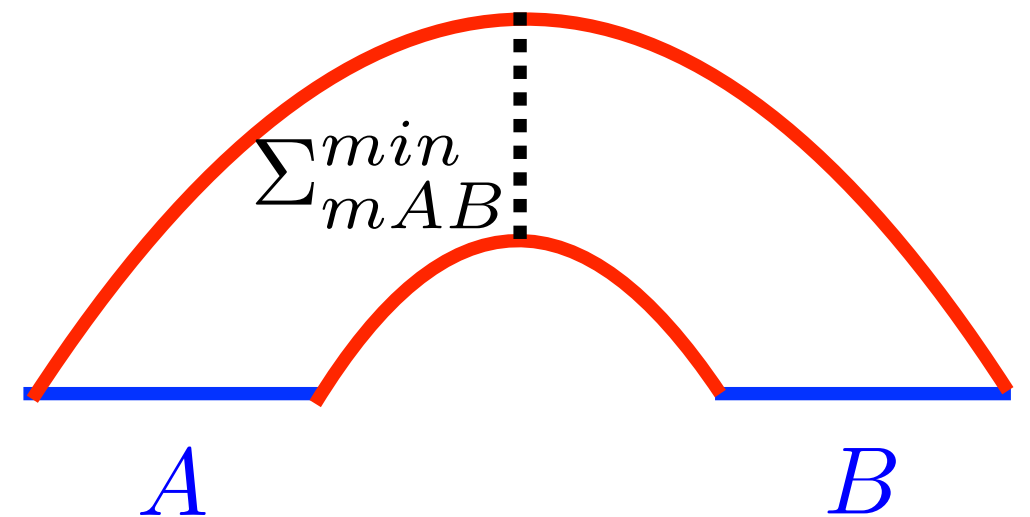


Computation of $E_{mW}(A : B)$

$$E_{mW}(A : B) = \frac{\text{Area}[\Sigma_{mAB}^{min}]}{4G_N}$$

entanglement wedge cross section
in the gravity dual of ρ_{AB}^m

- The gravity dual of ρ_{AB}^m includes the backreaction from cosmic branes.
[X. Dong, 2016]
- At the first order in $m - 1$, we can apply the computation method in [X. Dong, 2016] .



Red surface are cosmic branes
with tension T_m .

Check of $S_{mR}(A : B) = 2E_{mW}(A : B)$

reflected entropy with a canonical purification of ρ_{AB}^m

$$S_{mR}(A : B) = \frac{c}{3} \log \left[\frac{1 + \sqrt{z}}{1 - \sqrt{z}} \right] - \frac{2c(m-1)}{3} \frac{\sqrt{z} \log z}{1-z} + \dots$$

exactly matched
with $c = \frac{3}{2G_N}$

$$2E_{mW}(A : B) = \frac{1}{2G_N} \log \left[\frac{1 + \sqrt{z}}{1 - \sqrt{z}} \right] - \frac{m-1}{G_N} \frac{\sqrt{z} \log z}{1-z} + \dots$$

entanglement wedge cross section

in the gravity dual of ρ_{AB}^m

Summary

$$S_{mR}(A : B) = 2E_{mW}(A : B)$$

- We study the holographic duality between reflected entropy and entanglement wedge cross section.
- We explicitly compute and check this duality up to the first order in $m - 1$.

