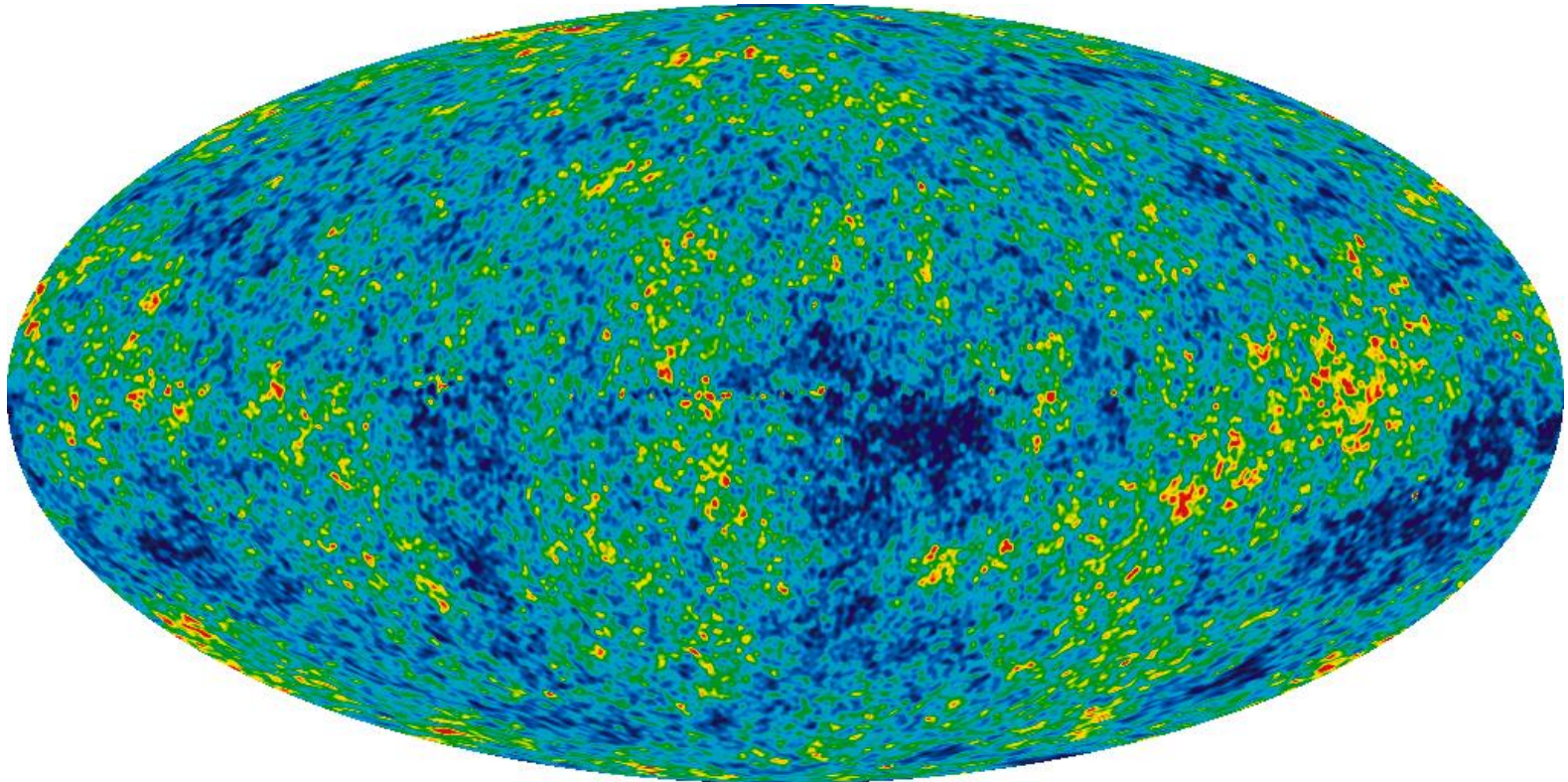


What is the **B mode**
in **CMB** polarization

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WMAP



<https://map.gsfc.nasa.gov/media/121238/index.html>

Temperature range $\pm 200 \mu\text{K}$

What do we see?

Tiny signals on the
“frosted glass”

- 10^{-4} in intensity
- 10^{-7} or 10^{-8} in polarization

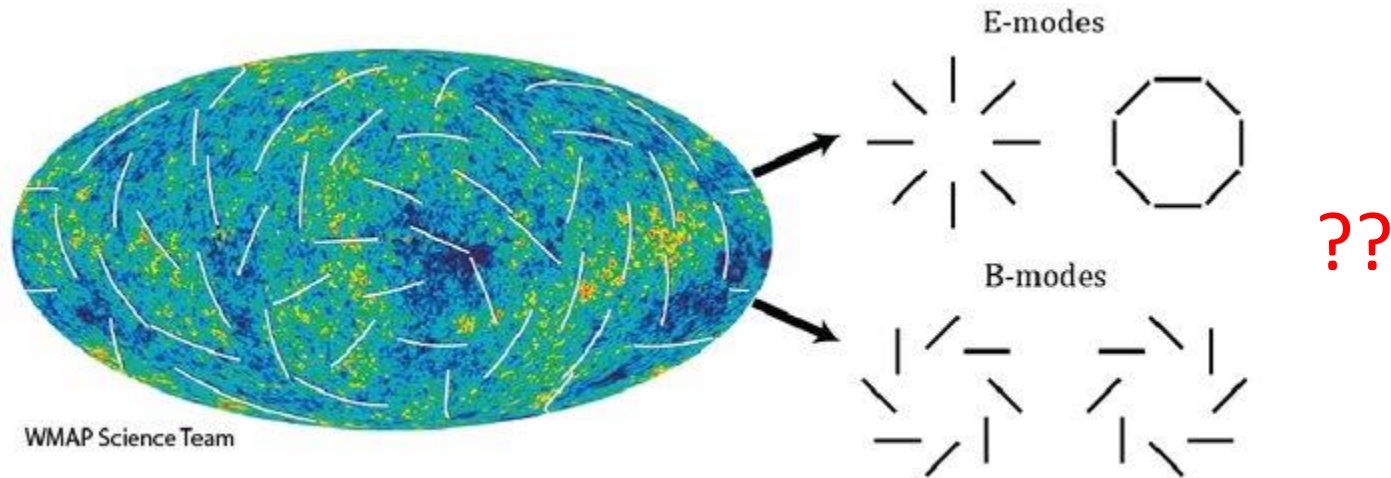


What do we see?

- Tiny signals on the “frosted glass”
- 10^{-4} in intensity
- 10^{-7} or 10^{-8} in polarization



Polarization



https://www.researchgate.net/publication/263811882_The_Primitive_Inflation_Polarization_Explorer_PIPER/figures?lo=1

“The CMB anisotropy polarization map may be decomposed into curl-free even-parity *E*-modes and divergence-free odd-parity *B*-modes. Primordial *B*-modes are only created by tensor perturbations (**inflationary gravitational waves**).”

Polarization

Polarization map may be decomposed into

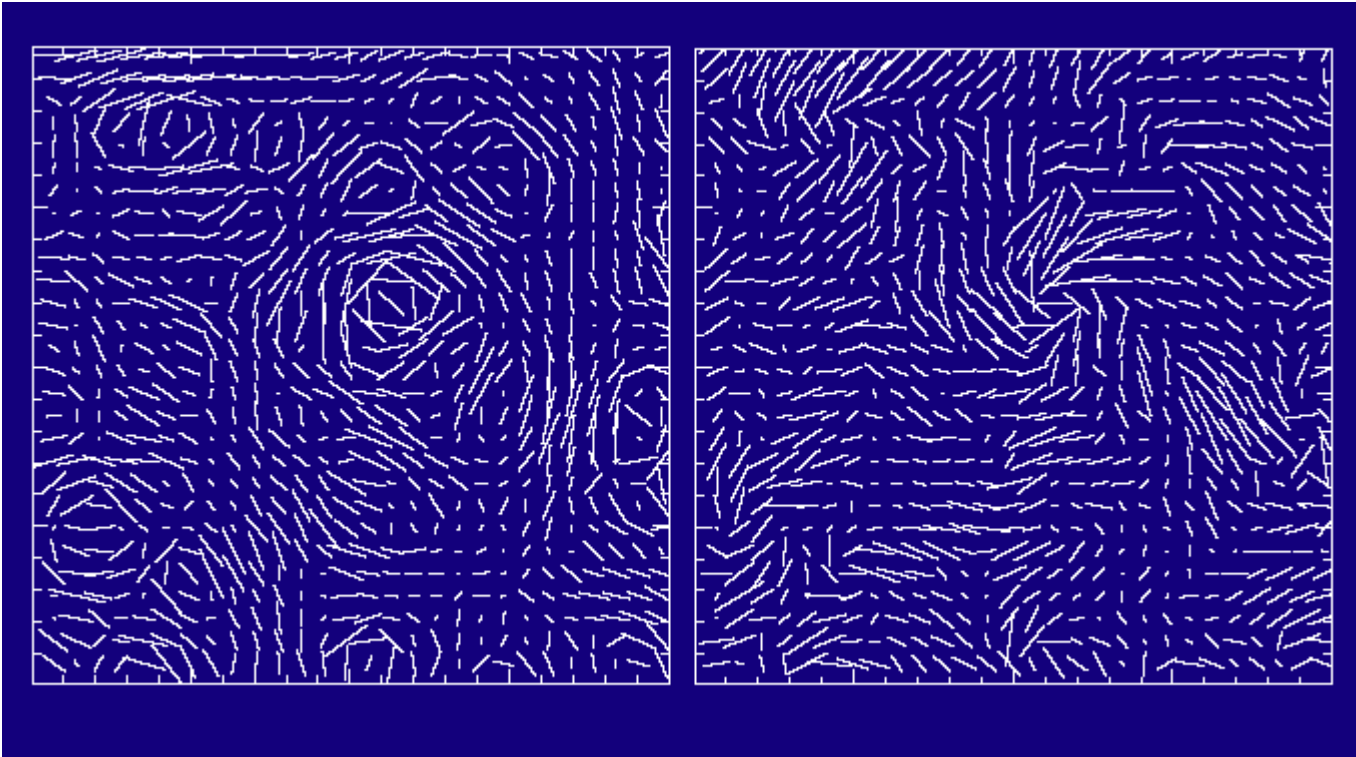
- curl-free even-parity **E modes**
- divergence-free odd-parity **B modes**

Primordial B modes are only created by tensor perturbations (inflationary gravitational waves)

Looking back

- Now = $1.4 \times 10^{10} \text{ y}$
 $\sim 4.4 \times 10^{17} \text{ s}$
 - CMB = $3.8 \times 10^5 \text{ y}$
 $\sim 1.2 \times 10^{13} \text{ s}$
 - B-mode/Inflation
 $\sim 10^{-32} \text{ s}$
- 4.5 orders
- 45 orders !!

E mode and B mode



<http://background.uchicago.edu/~whu/intermediate/Polarization/polar5.html>

Heuristic statement

E mode

- Curl free
- It is a divergence
- $\sim 10^{-7}$

B mode

- Divergence free
- It is a curl
- $\sim 10^{-8}$??

B mode

“a twisting pattern, or spin, known as a **curl** or B mode”

Breaking New, *Nature* 507, 7492 (2014)

“the **curl** component (B mode)”

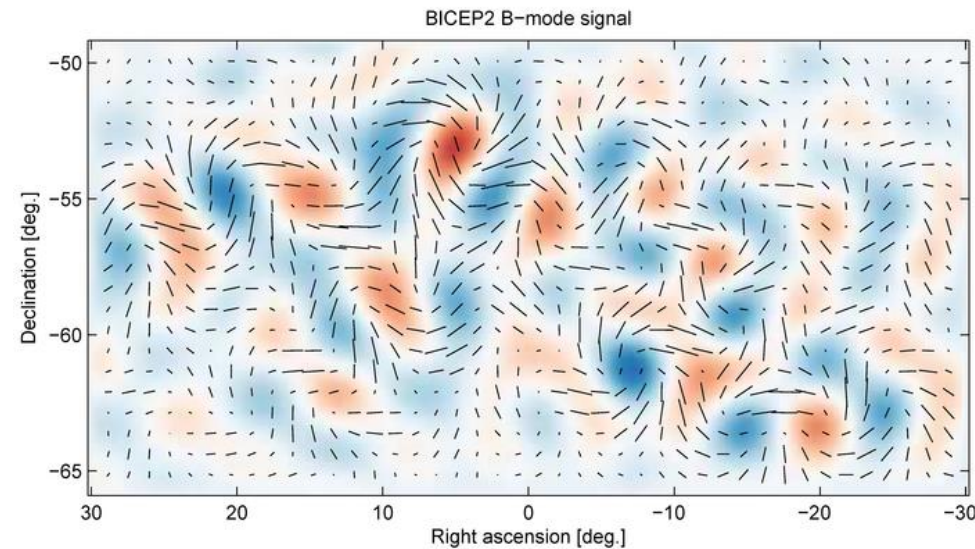
Kamionkowski & Kovetz
Ann Rev Astron & Astrophys, 54, 227 (2016)

B mode

- The “smoking gun”
- Primordial B mode can only be produced by tensor perturbations
- Would be evidence of inflation

B mode

- Seen by BICEP2
- 2014
- High profile announcement



<https://www.cfa.harvard.edu/~cbischoff/cmb/>



<https://www.nytimes.com/2014/06/20/science/space/scientist-s-debate-gravity-wave-detection-claim.html>

Announcement & Retraction

- BICEP2 (2014), Detection of B-Mode Polarization at Degree Angular Scales by BICEP2, Phys. Rev. Lett. 112, 24, 241101
- BICEP2/Keck & Planck (2015)
Joint Analysis of BICEP2/Keck Array and Planck Data, Phys. Rev. Lett. 114, 10, 101301
- Foreground dust!

Questions

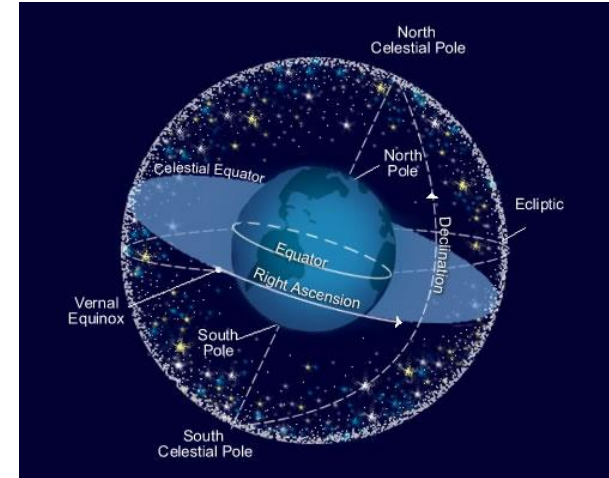
All these important physics questions rely on exquisite measurement of polarization **on a surface** and extracting the B mode

But

- How is polarization field decomposed?
- Is the E/B mode “like a grad/curl”?

What surface?

- Celestial sphere
- A plane
 - Math (toy) model
 - Small patch of sky
 - Any tangent plane of celestial sphere



<http://planetary-science.org/astronomy/the-celestial-globe/>

Decompose?

Literature seems hard, not accessible ..., but

- No need for inflation, cosmology, GR, ...
- Simple concept of plane surface (not sphere)
- Forget pixels
- It is a problem in classical electrodynamics in vacuum!

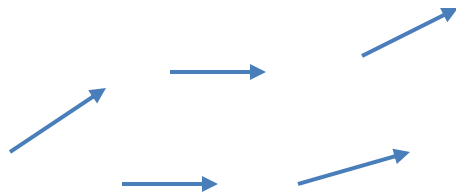
Like grad/ curl?

- No!
- Polarization is not a vector field, but a tensor field
- What does grad/ curl mean?
- Not like E and B field
- (E, B) modes $\rightarrow$ (A, B) modes

Vector / tensor

Vector field

- Arrows
- Returns after 360°



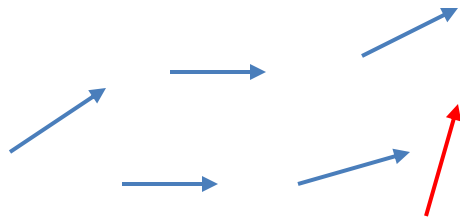
Polarization field

- Unsigned lines
- Returns after 180°



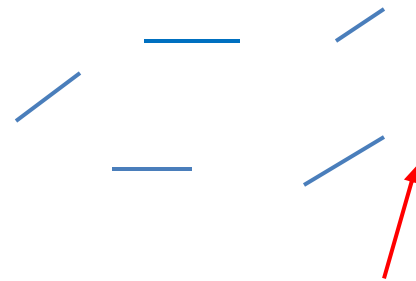
Curl?

Vector field



$$\sum \vec{V} \cdot \Delta \vec{r}$$

Polarization field



$$\sum \vec{P} \cdot \Delta \vec{r} \quad ?? \text{ No such thing}$$

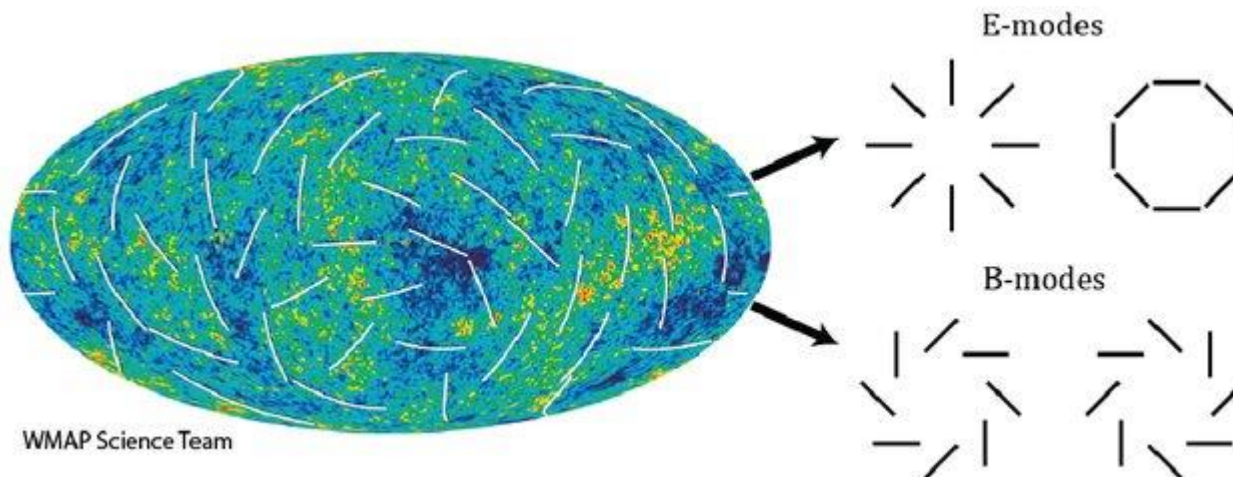
Factor of 2

Vector field

- V_i spin 1
- Returns after 360°
- Makes $0/90^\circ$ with $\mathbf{k}$

Polarization field

- $\text{Av } E_i * E_j$ spin 2
- Returns after 180°
- Makes $\pm 45^\circ$ with $\mathbf{k}$



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Vector field on a plane

$\mathbf{V}(\mathbf{r})$ on a "patch" Ω on the plane

$$\mathbf{r} = (x, y)^T = (x_1, x_2)^T$$

$$\mathbf{d} = (d_1, d_2)^T = (\partial_x, \partial_y)^T = \nabla_{\perp}$$

$$\tilde{\mathbf{d}} = (\tilde{d}_1, \tilde{d}_2)^T = (\partial_y, -\partial_x)^T$$

$$\tilde{\mathbf{d}} = -\hat{\mathbf{n}} \times \mathbf{d} = [\mathbf{E}] \mathbf{d}$$

$$[\mathbf{E}] = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

Vector field on a plane

$$\mathbf{d} = (d_1, d_2)^T = (\partial_x, \partial_y)^T$$

$$\tilde{\mathbf{d}} = (\tilde{d}_1, \tilde{d}_2)^T = (\partial_y, -\partial_x)^T$$

$$\mathbf{d}^T \mathbf{d} = \tilde{\mathbf{d}}^T \tilde{\mathbf{d}} = \nabla_{\perp}^2$$

$$\tilde{\mathbf{d}}^T \mathbf{d} = \mathbf{d}^T \tilde{\mathbf{d}} = 0$$

Vector field on a plane

Given $\mathbf{V}(\mathbf{r})$, seek representation

$$\mathbf{V} = \mathbf{d}\Phi^a - \tilde{\mathbf{d}}\Phi^b = \nabla_{\perp}\Phi^a + \nabla_{\perp} \times (\hat{\mathbf{n}}\Phi^b)$$

$$\mathbf{d}^T : \quad \nabla_{\perp}^2 \Phi^a = \mathbf{d}^T \mathbf{V} \equiv \varphi^a$$

$$\tilde{\mathbf{d}}^T : \quad \nabla_{\perp}^2 \Phi^b = -\tilde{\mathbf{d}}^T \mathbf{V} \equiv \varphi^b$$

Can laplacian be inverted?

Yes for global data

Some ambiguity for a patch

Vector field on a plane

$$\mathbf{V} = \mathbf{d}\Phi^a - \tilde{\mathbf{d}}\Phi^b$$

$$V_1 = \partial_x \Phi^a - \partial_y \Phi^b$$

$$V_2 = \partial_y \Phi^a + \partial_x \Phi^b$$

$$V_1 \pm iV_2 = \partial_{\pm}(\Phi^a \pm i\Phi^b)$$

$$\mathbf{V} \pm i\tilde{\mathbf{V}} = (\mathbf{d} \pm i\tilde{\mathbf{d}})(\Phi^a \pm i\Phi^b)$$

Vector field on a plane

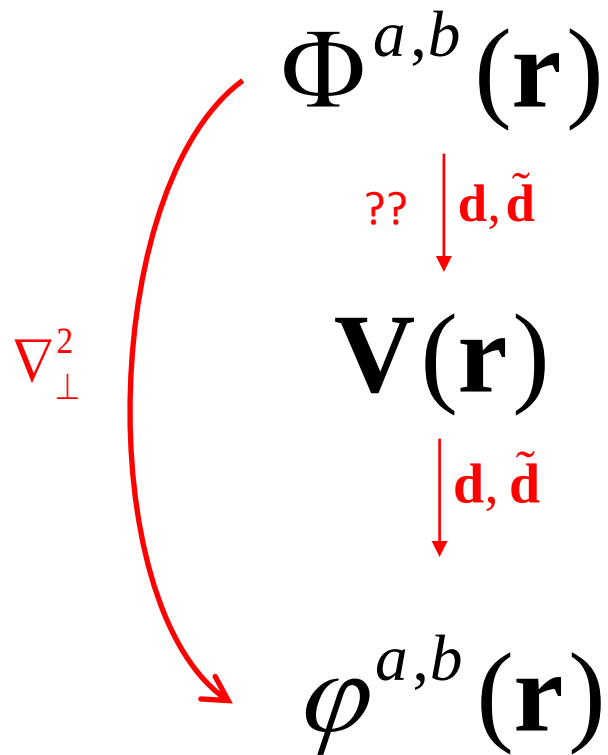
$$V_1 \pm iV_2 = \partial_{\pm}(\Phi^a \pm i\Phi^b)$$

$$\partial_{\mp}(V_1 \pm iV_2) = \nabla_{\perp}^2(\Phi^a \pm i\Phi^b) = \varphi^a \pm i\varphi^b$$

$$\mathbf{V} \pm i\tilde{\mathbf{V}} = (\mathbf{d} \pm i\tilde{\mathbf{d}})(\Phi^a \pm i\Phi^b)$$

$$(\mathbf{d} \mp i\tilde{\mathbf{d}})(\mathbf{V} \pm i\tilde{\mathbf{V}}) = \nabla_{\perp}^2(\Phi^a \pm i\Phi^b) = \varphi^a \pm i\varphi^b$$

Decomposition



Using two scalar fields

- Independent of axes
- Leads to expansion in terms of harmonic functions
- This method is unique to 2D

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3. Polarization field

3.1 Review of polarization

3.2 Decomposition of polarization field

Review of polarization

Propagate along z , measure on xy plane

$$E_j = a_j e^{i\delta_j}, \quad j = 1, 2$$

$$S_{ij} = \text{Av } E_i^* E_j$$

$$|\Psi\rangle = \begin{pmatrix} E_1 \\ E_2 \end{pmatrix}$$

$$[\mathbf{S}] = \text{Av } |\Psi\rangle\langle\Psi| = \begin{pmatrix} I + Q & U + iV \\ U - iV & I - Q \end{pmatrix}$$

Stokes' parameters

Review of polarization

$$[\mathbf{S}] = A \mathbf{v} |\Psi\rangle \langle \Psi| = \begin{pmatrix} I + Q & U + iV \\ U - iV & I - Q \end{pmatrix}$$

$I \sim$ total intensity

$Q, U \sim$ linear polarization

$V \sim$ circular polarization

Review of polarization

Ignore I

Ignore V because

(a) need CP violation,

(b) instruments cannot measure

$$[\mathbf{S}] \rightarrow [\mathbf{P}] = \begin{pmatrix} Q & U \\ U & -Q \end{pmatrix} = (\mathbf{P}, \tilde{\mathbf{P}})$$

$[\mathbf{P}] = 2 \times 2$ sym traceless tensor

$\mathbf{P}$ = doublet $\neq$ vector

Review of polarization

$$[\mathbf{P}] = \begin{pmatrix} Q & U \\ U & -Q \end{pmatrix} = (\mathbf{P}, \tilde{\mathbf{P}})$$

$[\mathbf{P}] = 2 \times 2$ sym traceless tensor

Under rotations by angle α

$$[\mathbf{P}] \rightarrow [\mathbf{P}'] = [\mathbf{R}]^T [\mathbf{P}] [\mathbf{R}] \quad [\mathbf{R}] = \begin{pmatrix} \cos \alpha & -\sin \alpha \\ \sin \alpha & \cos \alpha \end{pmatrix}$$

$$Q \pm iU \rightarrow Q' \pm iU' = e^{\pm 2i\alpha} (Q \pm iU)$$

Spin 2 , returns to itself under 180° rotation

Review of polarization

$$Q \sim E_x^2 - E_y^2 = E^2 (\cos^2 \chi - \sin^2 \chi) \sim \cos 2\chi$$

$$U \sim 2E_x E_y = 2E^2 \cos \chi \sin \chi \sim \sin 2\chi$$

$$\tan 2\chi = U / Q$$

2χ is determined mod 2π

3. Polarization field

3.1 Review of polarization

3.2 Decomposition of polarization field

Decompose polarization

$$V_1 \pm iV_2 = \partial_{\pm}(\Phi^a \pm i\Phi^b) \quad \text{spin 1}$$

$$Q \pm iU = \partial_{\pm}^2(\Phi^a \pm i\Phi^b) \quad \text{spin 2}$$

- Conceptually, this is all we need
- Scalars are independent of axes on tangent planes
- But this representation is not convenient for mapping to other (tangent) planes

Decompose polarization

$$V_1 \pm iV_2 = \partial_{\pm}(\Phi^a \pm i\Phi^b) \quad \text{spin 1}$$

$$Q \pm iU = \partial_{\pm}^2(\Phi^a \pm i\Phi^b) \quad \text{spin 2}$$

$$\partial_{\pm}^2 = D_1 \pm iD_2$$

$$D_1 = \partial_x^2 - \partial_y^2 \quad D_2 = 2\partial_x\partial_y$$

$$Q = D_1\Phi^a - D_2\Phi^b$$

$$U = D_2\Phi^a + D_1\Phi^b$$

Decompose polarization

$$Q = D_1\Phi^a - D_2\Phi^b$$

$$U = D_2\Phi^a + D_1\Phi^b$$

$$\mathbf{D} = (D_1, D_2)^T \quad \tilde{\mathbf{D}} = (D_2, -D_1)^T = [\mathbf{E}] \mathbf{D}$$

$$\mathbf{P} = \mathbf{D}\Phi^a - \tilde{\mathbf{D}}\Phi^b$$

Compare

$$\mathbf{V} = \mathbf{d}\Phi^a - \tilde{\mathbf{d}}\Phi^b$$

Decompose polarization

$$\mathbf{D} = (D_1, D_2)^T \quad \tilde{\mathbf{D}} = (D_2, -D_1)^T$$

$$\mathbf{D}^T \mathbf{D} = \tilde{\mathbf{D}}^T \tilde{\mathbf{D}} = \nabla_{\perp}^4$$

$$\tilde{\mathbf{D}}^T \mathbf{D} = \mathbf{D}^T \tilde{\mathbf{D}} = 0$$

Compare

$$\mathbf{d}^T \mathbf{d} = \tilde{\mathbf{d}}^T \tilde{\mathbf{d}} = \nabla_{\perp}^2$$

$$\tilde{\mathbf{d}}^T \mathbf{d} = \mathbf{d}^T \tilde{\mathbf{d}} = 0$$

Decompose polarization

Given $\mathbf{P}(\mathbf{r})$, seek representation

$$\mathbf{P} = \mathbf{D}\Phi^a - \tilde{\mathbf{D}}\Phi^b$$

$$\mathbf{D}^T : \quad \nabla_{\perp}^4 \Phi^a = \mathbf{D}^T \mathbf{P} \equiv \varphi^a$$

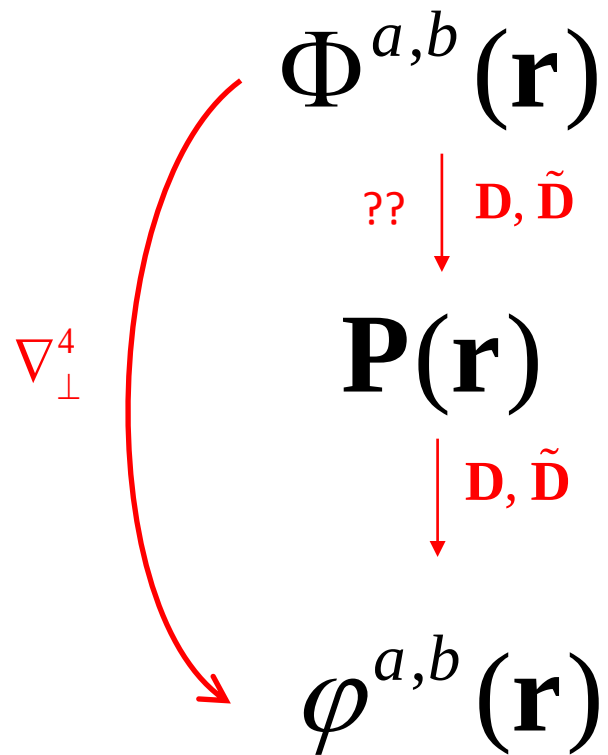
$$\tilde{\mathbf{D}}^T : \quad \nabla_{\perp}^4 \Phi^b = -\mathbf{D}^T \mathbf{P} \equiv \varphi^b$$

Can **bi**laplacian be inverted?

Yes for global data

Some ambiguity for a patch

Decomposition



Decompose polarization

Given $\mathbf{D}$ (doublet), introduce $[\mathbf{D}]$ (2×2 matrix)

$$[\mathbf{D}] = (\mathbf{D}, \tilde{\mathbf{D}})$$

In Cartesian representation

$$\begin{aligned} [\mathbf{D}] \pm i[\tilde{\mathbf{D}}] &= (\mathbf{d} \pm i\tilde{\mathbf{d}}) \otimes (\mathbf{d} \pm i\tilde{\mathbf{d}}) \\ &= (\mathbf{d} \otimes \mathbf{d} - \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \pm i(\mathbf{d} \otimes \tilde{\mathbf{d}} + \tilde{\mathbf{d}} \otimes \mathbf{d}) \end{aligned}$$

in terms of which

$$[\mathbf{P}] \pm i[\tilde{\mathbf{P}}] = ([\mathbf{D}] \pm i[\tilde{\mathbf{D}}])(\Phi^a \pm i\Phi^b)$$

Decompose polarization

$$[\mathbf{P}] = \begin{pmatrix} Q & U \\ U & -Q \end{pmatrix} = (\mathbf{P}, \tilde{\mathbf{P}})$$

Tensor representation says the same thing twice. Advantage?

- Indices are Cartesian
- Can regard all quantities as in 3D space
- Specified by normal vector
- Normal components zero

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Global analysis (1)

$$\mathbf{V} \sim (\mathbf{d}, \tilde{\mathbf{d}}) \Phi^{a,b}$$

local

$$(\mathbf{d}, \tilde{\mathbf{d}}) \mathbf{V} \sim \varphi^{a,b}$$

$$\nabla_{\perp}^2 \Phi^{a,b} = \varphi^{a,b}$$

global

$$\varphi^{a,b}(\mathbf{r}) = \sum_{\alpha} c_{\alpha}^{a,b} F_{\alpha}(\mathbf{r}) \quad , \text{ e.g., plane waves } \mathbf{k}$$

$$\Phi^{a,b}(\mathbf{r}) = \sum_{\alpha} C_{\alpha}^{a,b} F_{\alpha}(\mathbf{r})$$

$$C_{\alpha}^{a,b} = \lambda_{\alpha}^{-1} c_{\alpha}^{a,b}$$

$$\lambda_{\alpha} = \text{eigenvalue of } -\nabla_{\perp}^2, \text{ e.g., } \mathbf{k}^2$$

Global analysis (2)

$$\mathbf{V} = \mathbf{d}\Phi^a - \tilde{\mathbf{d}}\Phi^b$$

$$\Phi^{a,b}(\mathbf{r}) = \sum_{\alpha} C_{\alpha}^{a,b} F_{\alpha}(\mathbf{r})$$

$$\mathbf{V}(\mathbf{r}) = \sum_{\alpha} C_{\alpha}^a \mathbf{d}F_{\alpha}(\mathbf{r}) - \sum_{\alpha} C_{\alpha}^b \tilde{\mathbf{d}}F_{\alpha}(\mathbf{r})$$

vector harmonics

Global analysis (3)

$$\mathbf{P} \sim (\mathbf{D}, \tilde{\mathbf{D}}) \Phi^{a,b} \quad \text{local}$$

$$(\mathbf{D}, \tilde{\mathbf{D}}) \mathbf{P} \sim \varphi^{a,b}$$

$$\nabla_{\perp}^4 \Phi^{a,b} = \varphi^{a,b} \quad \text{global}$$

$$\varphi^{a,b}(\mathbf{r}) = \sum_{\alpha} c_{\alpha}^{a,b} F_{\alpha}(\mathbf{r}) \quad , \text{ e.g., plane waves } \mathbf{k}$$

$$\Phi^{a,b}(\mathbf{r}) = \sum_{\alpha} C_{\alpha}^{a,b} F_{\alpha}(\mathbf{r})$$

$$C_{\alpha}^{a,b} = \lambda_{\alpha}^{-2} c_{\alpha}^{a,b}$$

$$\lambda_{\alpha} = \text{eigenvalue of } -\nabla_{\perp}^2, \text{ e.g., } \mathbf{k}^2$$

Global analysis (4)

$$\mathbf{P} = \mathbf{D}\Phi^a - \tilde{\mathbf{D}}\Phi^b$$

$$\Phi^{a,b}(\mathbf{r}) = \sum_{\alpha} C_{\alpha}^{a,b} F_{\alpha}(\mathbf{r})$$

$$\mathbf{P}(\mathbf{r}) = \sum_{\alpha} C_{\alpha}^a \boxed{\mathbf{D}F_{\alpha}(\mathbf{r})} - \sum_{\alpha} C_{\alpha}^b \boxed{\tilde{\mathbf{D}}F_{\alpha}(\mathbf{r})}$$

tensor harmonics

- Ordering of operators?

$$\mathbf{P}(\mathbf{r}) = \sum_{\alpha} C_{\alpha}^a \mathbf{D} F_{\alpha}(\mathbf{r}) - \sum_{\alpha} C_{\alpha}^b \tilde{\mathbf{D}} F_{\alpha}(\mathbf{r})$$

$$\mathbf{D} \sim \begin{pmatrix} \partial_x^2 - \partial_y^2 \\ 2\partial_x \partial_y \end{pmatrix} \sim - \begin{pmatrix} k_x^2 - k_y^2 \\ 2k_x k_y \end{pmatrix} \sim -k^2 \begin{pmatrix} \cos^2 \phi - \sin^2 \phi \\ 2 \cos \phi \sin \phi \end{pmatrix}$$

$$\begin{pmatrix} Q \\ U \end{pmatrix} = - \sum_{\mathbf{k}} k^2 \times$$

$$\left[C_{\mathbf{k}}^a \begin{pmatrix} \cos 2\phi \\ \sin 2\phi \end{pmatrix} + C_{\mathbf{k}}^b \begin{pmatrix} -\sin 2\phi \\ \cos 2\phi \end{pmatrix} \right] F_{\mathbf{k}}$$

$\tan 2\chi = U / Q$ direction of polarization

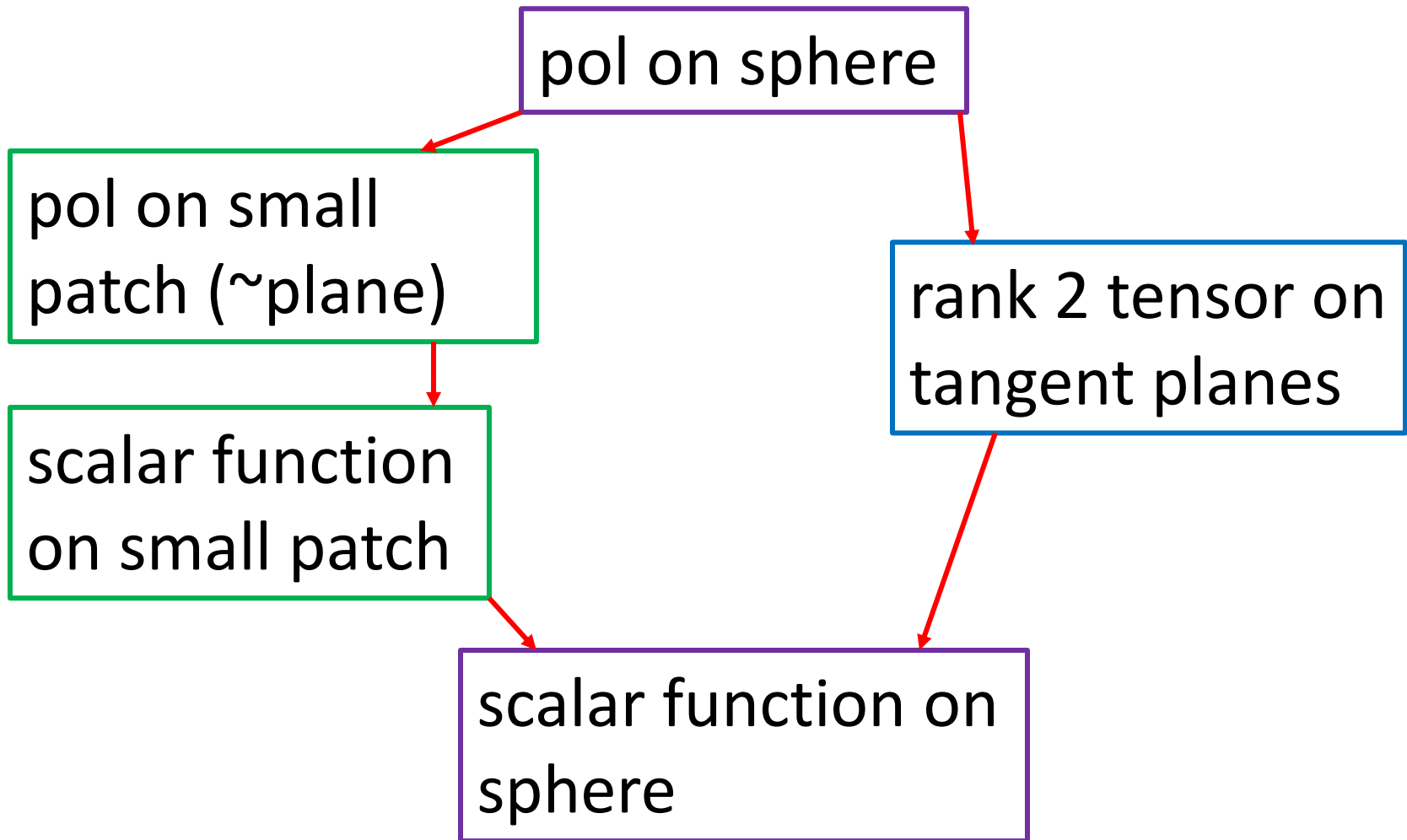
A mode: $\tan 2\chi = \tan 2\phi$, $\chi - \phi = n\pi / 2$

B mode: $\tan 2\chi = -\cot 2\phi$, $\chi - \phi = \pi / 4 + n\pi / 2$

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Two paths



Sphere

Recall

$$[\mathbf{P}] \pm i[\mathbf{P}] = ([\mathbf{D}] \pm i[\tilde{\mathbf{D}}])(\Phi^a \pm i\Phi^b)$$

$$\begin{aligned} [\mathbf{D}] \pm i[\tilde{\mathbf{D}}] &= (\mathbf{d} \pm i\tilde{\mathbf{d}}) \otimes (\mathbf{d} \pm i\tilde{\mathbf{d}}) \\ &= (\mathbf{d} \otimes \mathbf{d} - \tilde{\mathbf{d}} \otimes \tilde{\mathbf{d}}) \pm i(\mathbf{d} \otimes \tilde{\mathbf{d}} + \tilde{\mathbf{d}} \otimes \mathbf{d}) \end{aligned}$$

All of these can be regarded as objects in 3D

But having normal component zero

But what are $\mathbf{d}$, $\tilde{\mathbf{d}}$?

Sphere

No need to say perp

Now write as

$$\tilde{\mathbf{d}} = -\hat{\mathbf{n}} \times \nabla \quad \mathbf{d} = \mathbf{n} \times \hat{\mathbf{d}}$$

Can be ported to any tangent plane

On a sphere

$$\tilde{\mathbf{d}} = -i\mathbf{L} \quad \mathbf{d} = i\tilde{\mathbf{L}}$$

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Summary (1)

$$V_1 \pm iV_2 = \partial_{\pm} (\Phi^a \pm i\Phi^b)$$

is the curl

$$Q \pm iU = \partial_{\pm}^2 (\Phi^a \pm i\Phi^b)$$

“like the curl”
But second-order
derivative

Like the curl only in this formal sense

Summary (2)

schematic



$\mathbf{V} \rightarrow \mathbf{dV} = \varphi^{a,b} \rightarrow \text{invert laplacian gives } \Phi^{a,b}$

$\mathbf{P} \rightarrow \mathbf{DV} = \varphi^{a,b} \rightarrow \text{invert bilaplacian gives } \Phi^{a,b}$

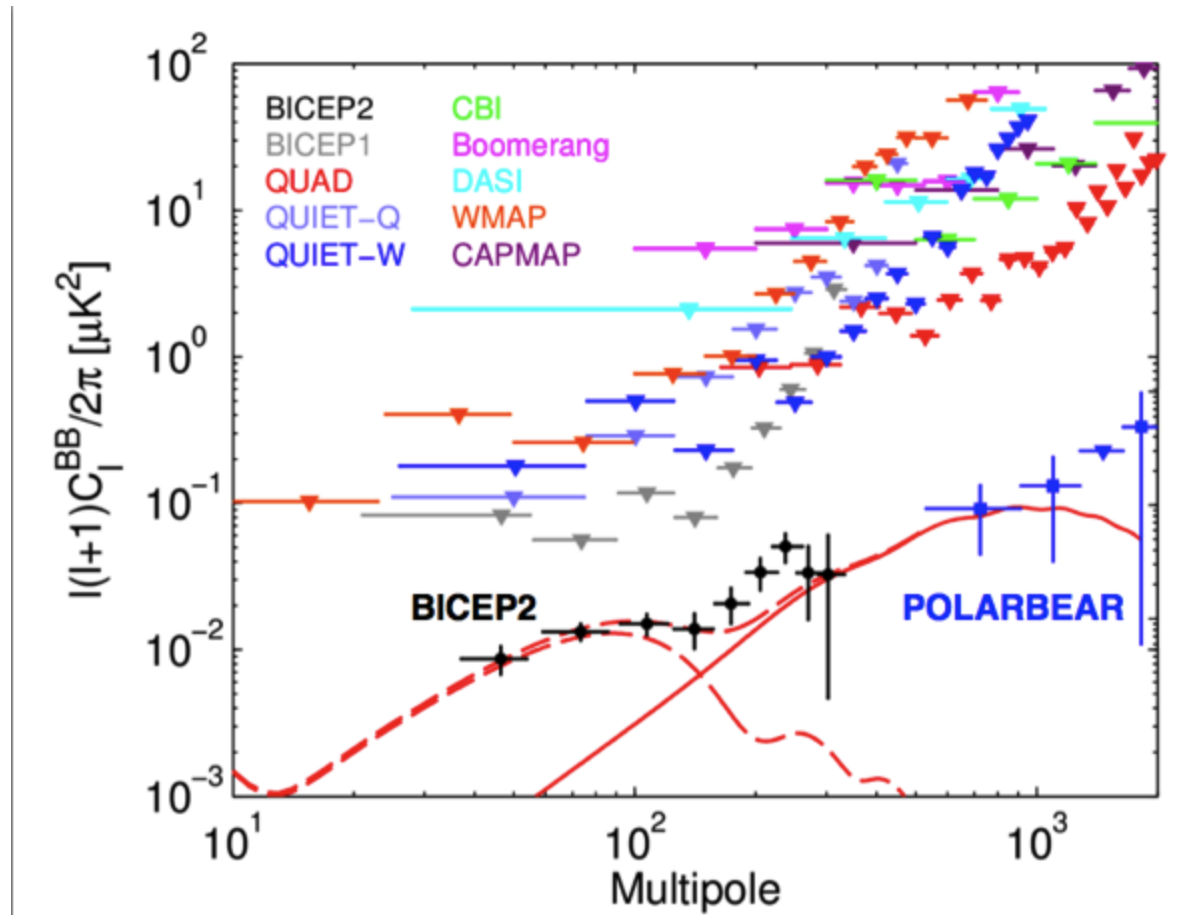
"Amount of B mode" given by

$|\varphi^b|$ locally or $|\Phi^b|$ globally

Or globally in terms of the mode

amplitudes of Φ^b

Observational limits (2014)

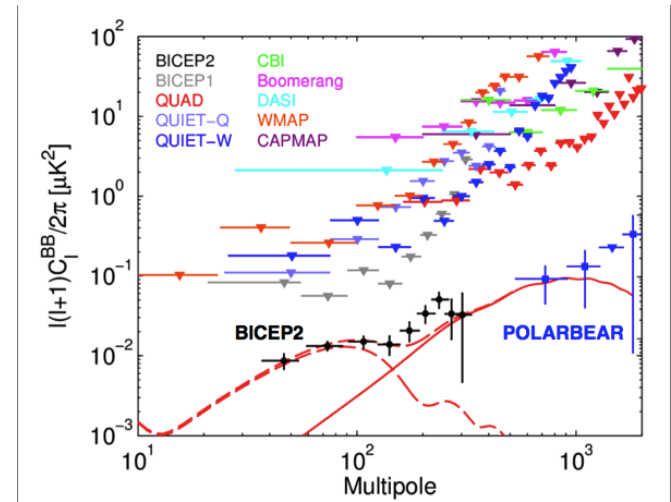


<https://www.andrewjaffe.net/blog/2014/03/gravitational-w.html>

Observational limits (2014)

Vertical axis is essentially
the squared mode coefficients in

$$\Phi^b = \sum_{\ell m} C_{\ell m}^b Y_{\ell m}$$



<https://www.andrewjaffe.net/blog/2014/03/gravitational-w.html>

Conclusion

- Polarization
- Measurement very difficult
- Separation into two modes conceptually simple
- Removing other effects (dust, lensing, ...) tricky
- Potential implication is huge