Neutrinos Meet SIMP

Koji Tsumura (Kyoto U.) Energy Frontier in Particle Physics: LHC and Future Colliders National Taiwan University, Sep 29-30, 2017

> Based on JHEP 1708, 101 (2017) [arXiv:1705.00592] "A Radiative Neutrino Mass Model with SIMP Dark Matter" In collaboration with Shu-Yu Ho (Caltech), Takashi Toma (TUM)

<u>Contents</u>



Dark Matter

Dark Matter

Many Evidences for DM

- Galaxy Rotation Curve
- Velocity Dispersion of Galaxies
- Galaxy Clusters and Gravitational Lensing
- Sky surveys and baryon acoustic oscillations
- Cosmic Microwave Background (CMB) ____
- Type la supernovae distance measurements
- Lyman-Alpha Forest
- Structure Formation





Need DM

Dark Matter

Many Candidates for DM

- Primordial Black Hole
- Neutrino (Hot)
- WIMP [Weakly Interacting Massive Particle]
- SIMP [Strongly Interacting Massive Particle]
- Axion, Axion cluster
- Soliton (Q-ball, B-ball, ...)
- Super Massive Relic (WIMPzilla, ...)



Some Dark Matter Candidate Particles

...

WIMP Paradigm

A promising candidate for Thermal DM



WIMP Miracle

A promising candidate for Thermal DM

Annihilation



SIMP Paradigm

New promising candidate for Thermal DM $\dot{n} + 3Hn = -(n^3 - n^2 n_{\rm eq}) \langle \sigma_{3 \to 2} v_{\rm rel} \rangle$



<u>SIMP Miracle</u>

New promising candidate for Thermal DM

Annihilation



3→2 annihilation in DM sector

$$\langle \sigma_{3 \to 2} v_{\rm rel}^2 \rangle \equiv \frac{\alpha_{3 \to 2}^3}{M_{\rm DM}^5}$$

 $M_{\rm DM} \simeq \alpha_{3 \to 2} \times 100 {\rm MeV}$

Strong int. + Λ_{QCD} (SIMP Miracle)

SIMP Stability? Strong Cubic int. $|\mathsf{ex. } \mathsf{Z}_3 \colon \chi \to e^{(2\pi/3)i}\chi \quad \mathcal{L}_{Z_3} = m^2 |\chi|^2 + \kappa(\chi^3 + \mathrm{H.c.}) + \lambda |\chi|^4$ $\sim \frac{\kappa^3}{M^4}$ DM DM DM

$3 \rightarrow 2$ annihilation in DM sector

$$\langle \sigma_{3\to 2} v_{\rm rel}^2 \rangle \equiv \frac{\alpha_{3\to 2}^3}{M_{\rm DM}^5}$$

DM

single scale theory

 $\sim rac{\kappa\lambda}{M^2}$

 $\kappa \sim g M, \ \lambda \sim g^2 \quad \Longrightarrow \quad \mathcal{M} \sim \frac{g^3}{M}$

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DM

Hot DM ?



Need connection with SM to keep it in Thermal Equilibrium

$$\begin{array}{ccc} \mathrm{DM} & & & & \\ & & & \\ \mathrm{SM} & & & \\ & & & \\ & & & \\ \end{array} & & & \\ & & & \\ \mathrm{SM} & & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ \end{array} & \begin{array}{c} \mathrm{New \ portal \ int.} & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ &$$

SIMP Condition





Neutrino

<u>Neutrino</u>

 $M_{\nu} < < < < < < v_{\rm EW}$ $O(10^{12})$ $= 2.46 \times 10^{11} \, eV$ $\sim 0.1\,\mathrm{eV}$

 $M_{\nu} << < M_e \quad \mathcal{O}(10^7)$ = 5.11 × 10⁵ eV



 $\sim 0.1\,\mathrm{t}$

 $\sim 2.9 imes 10^{12} \, {
m t}$ <u>Mt. Fuji</u>

(a big mountain)



Seesaw Mechanism

Neutrinos can be Majorana Particle



Neutrinos Meet DM

Scoto-genic Model (Scotos = Darkness)

Ma (06) [Ma (98)]

Z₂ odd Imposed Discrete Sym. → WIMP DM



Z₃ charged

Ma (07)



Can we construct a model with SIMP DM?

NR

η

η

S-Y Ho, T. Toma, KT, JHEP1708, 101 (2017) [arXiv:1705.00592] "A Radiative Neutrino Mass Model with SIMP Dark Matter"

	vSIMP ^g ₃→2											
						New Particle						
	E	Φ	N	η	χ	γ η χ ^ο η Υποτεί						
SU(2)	2	2	1	2	1	• portal N N						
$\mathrm{U}(1)_Y$	-1/2	1/2	0	1/2	0							
\mathbb{Z}_3	1	1	ω	ω	ω	Mediator for Cooling ω=exp(2πi/3)						

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 $M_{\rm DM} \simeq \alpha_{3 \to 2} \times 100 {\rm MeV}$

<u>Relic Abundance</u> $\kappa_{\chi} \leftrightarrow g_{3 \rightarrow 2} M_{\text{DM}}$





Requiring the correct DM relic abundance







Resonant SIMP .Choi, Lee '16

Large $\sigma_{3\rightarrow 2}$ while keeping $\sigma_{2\rightarrow 2}$ small by Resonance



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Radiative M_v

From SIMP Condition

≈ 100 MeV ≈ 0.1 ξ≈ 0.05

$$(M_{\nu})_{rs} = \frac{\mu_2 Y_{rj} Y_{sk} s_{2\xi}^2}{4(4\pi)^2} \begin{pmatrix} \bullet \approx 10^{-3} \\ \mathcal{Y}_{ji}^L \mathcal{C}_{ji}^L + \mathcal{Y}_{ji}^R \mathcal{C}_{ji}^R \end{pmatrix}$$

$$\approx 0.1 \text{eV}$$
Loop integrals $\approx 0.1-10$



✓ Direct Search of η

Relic Abundance



Stability : $\mu_2 \lesssim 100 \text{MeV} \rightarrow \text{Y s}_{2\xi} \lesssim 0.01$ Koji Tsumura (Kyoto U.)

Lepton Flavor Violation

$$\mathcal{B}_{\ell_r \to \ell_s \gamma} \approx \frac{\alpha \, \mathcal{B}_{\ell_r \to \ell_s \overline{\nu_s} \nu_r}}{768 \pi G_F^2 m_\eta^4} \Big| \sum_k \mathcal{Y}_{rk}^{\star} \mathcal{Y}_{sk} \Big|^2$$

Without considering the specific structure of Yukawa





Conflict with SIMP condition

An easy solution :

$$\mathcal{Y} = \begin{pmatrix} Y_{e1} & Y_{e2} & Y_{e3} \\ Y_{\mu 1} & Y_{\mu 2} & Y_{\mu 3} \\ Y_{\tau 1} & Y_{\tau 2} & Y_{\tau 3} \end{pmatrix} \longrightarrow \begin{pmatrix} Y_{e1} & & \\ & Y_{\mu 2} & \\ & & Y_{\tau 3} \end{pmatrix}$$



 η^-

 ℓ_r^-

 η^-

 N_k

 ℓ_s^-

Structure of M_v can be adjusted by Y^L , Y^R

Energy/Precision Frontier

Higgs Invisible Decay

 $[h_{125} \rightarrow \chi^0 \chi^0]$ is induced by mixing

$$\mathcal{L}_{\text{eff}} = -\frac{1}{2} \kappa \, s_{2\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0 \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0} \, z_{\xi} \, z_{\xi} \, h \, (\chi^0)^2 \implies \Gamma_{h \to \chi^0} \, z_{\xi} \, z_{\xi}$$

$$\frac{(\kappa s_{2\xi})^2}{64\pi m_h}$$

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The latest data : $B(h \rightarrow \chi^0 \chi^0) < 0.16$

 $|s_{2\xi}| < 0.11 \times \left(\frac{\mathcal{B}_{inv}^{up}}{0.16}\right)^{1/4} \left(\frac{300 \,\mathrm{GeV}}{m_n}\right)$

 $n^{0}-\chi^{0}$ mixing

$$\kappa_{2\xi} = \frac{\kappa v}{m_{\eta}^2 - m_{\chi}^2}$$

SIMP condition \rightarrow Y s₂ \leq 0.01

Slepton Direct Search



<u>Summary</u>

- v & DM are evidences of BSM
 vSIMP
 - M_v may be generated by SIMP
 - SIMP may be thermalized by v
 - Collider search is powerful probe of vSIMP

		Φ	N	η	χ	$egin{array}{c} S \end{array}$
SU(2)	2	2	1	2	1	1
$\mathrm{U}(1)_{Y}$	-1/2	1/2	0	1/2	0	0
\mathbb{Z}_5	1	1	ω	ω	ω	ω^3

