



“Tell me that you have found no sign of  
New Physics again, I dare you.  
I double dare you. Tell me  
one more goddamn **time!**”

# SUSY Higgs Bosons in the Light of LHC Data

*Sven Heinemeyer, IFT/IFCA (CSIC, Madrid/Santander)*

Taipei, 09/2017

- Motivation
- SUSY Higgs mass and rate measurements
- MSSM Higgs results (CMSSM, NUHM1/2, pMSSM8)
- $\phi_{95} \rightarrow \gamma\gamma$
- Conclusions

# 1. Motivation

Fact:

The SM cannot be the ultimate theory!

1. gravity is not included
2. the hierarchy problem
3. Dark Matter is not included
4. neutrino masses are not included
5. anomalous magnetic moment of the muon shows a  $\sim 4\sigma$  discrepancy

⇒ Time to get ready for BSM physics

## Which model should we focus on?

Some “recent” measurements:

- top quark mass
- Higgs boson mass
- Higgs boson “couplings”
- Dark Matter (properties)

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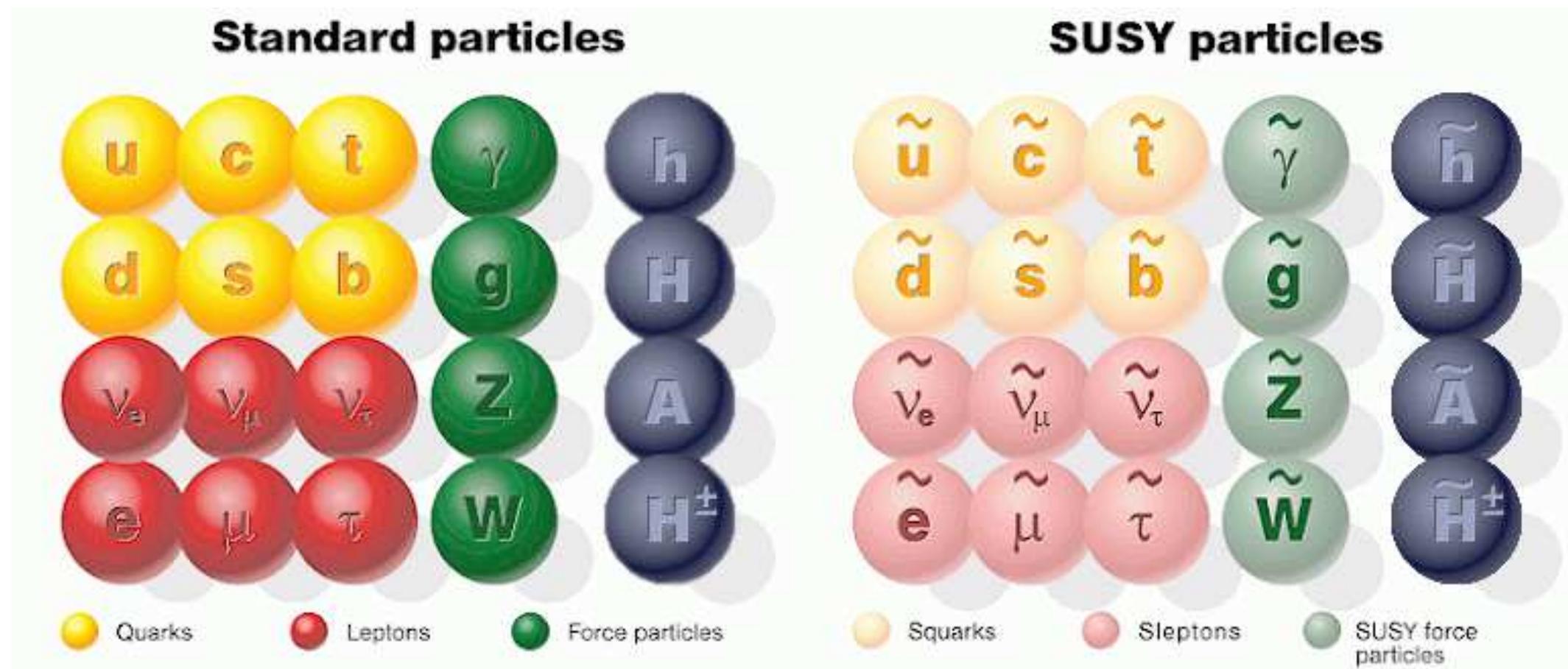
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⇒ **good motivation to look at SUSY! :-)**

# The Minimal Supersymmetric Standard Model (MSSM)

Superpartners for Standard Model particles



Problem in the MSSM: more than 100 free parameters

Nobody(?) believes that a model describing nature has so many free parameters!

## GUT based models: 1.) CMSSM (sometimes wrongly called mSUGRA):

⇒ Scenario characterized by

$$m_0, m_{1/2}, A_0, \tan\beta, \text{sign } \mu$$

$m_0$  : universal scalar mass parameter

$m_{1/2}$  : universal gaugino mass parameter

$A_0$  : universal trilinear coupling

$\tan\beta$  : ratio of Higgs vacuum expectation values

$\text{sign}(\mu)$  : sign of supersymmetric Higgs parameter

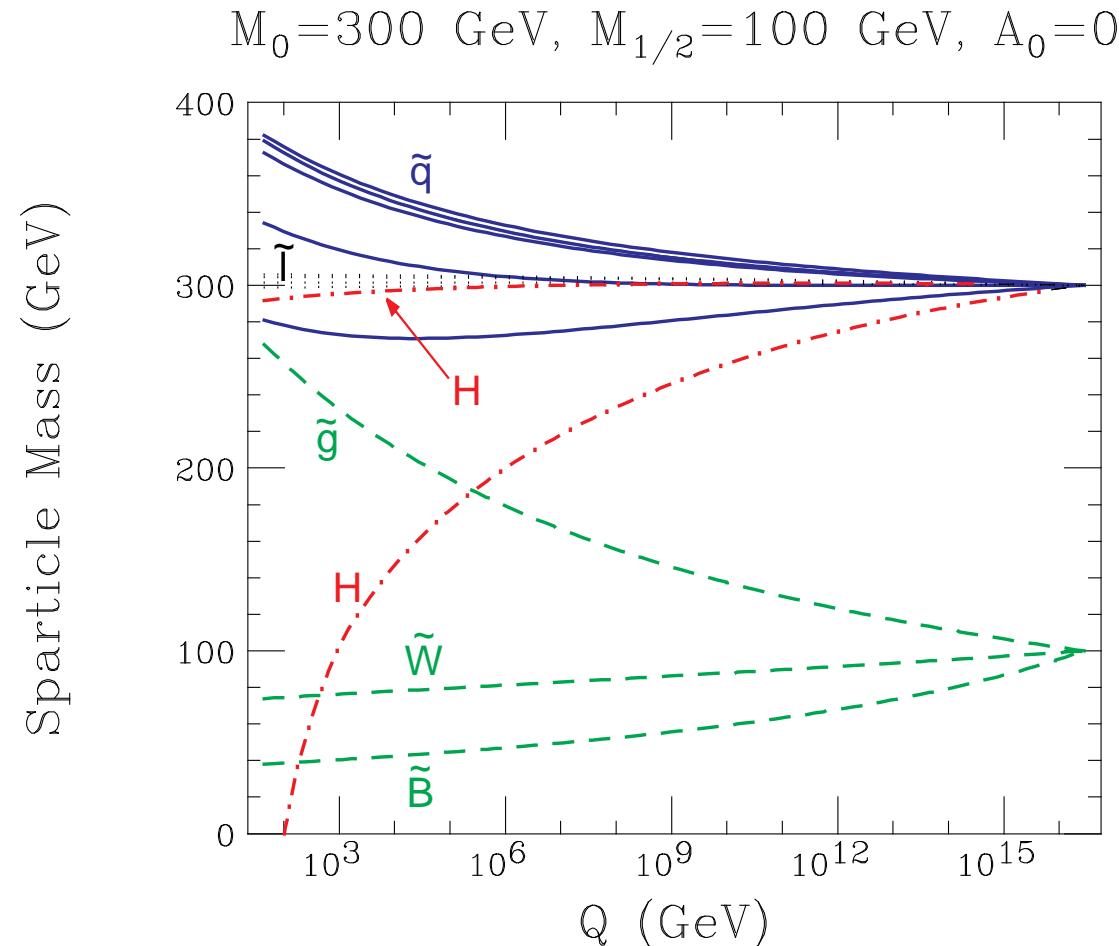
} at the GUT scale

⇒ particle spectra from renormalization group running to weak scale

⇒ Lightest SUSY particle (LSP) is the lightest neutralino ⇒ DM!

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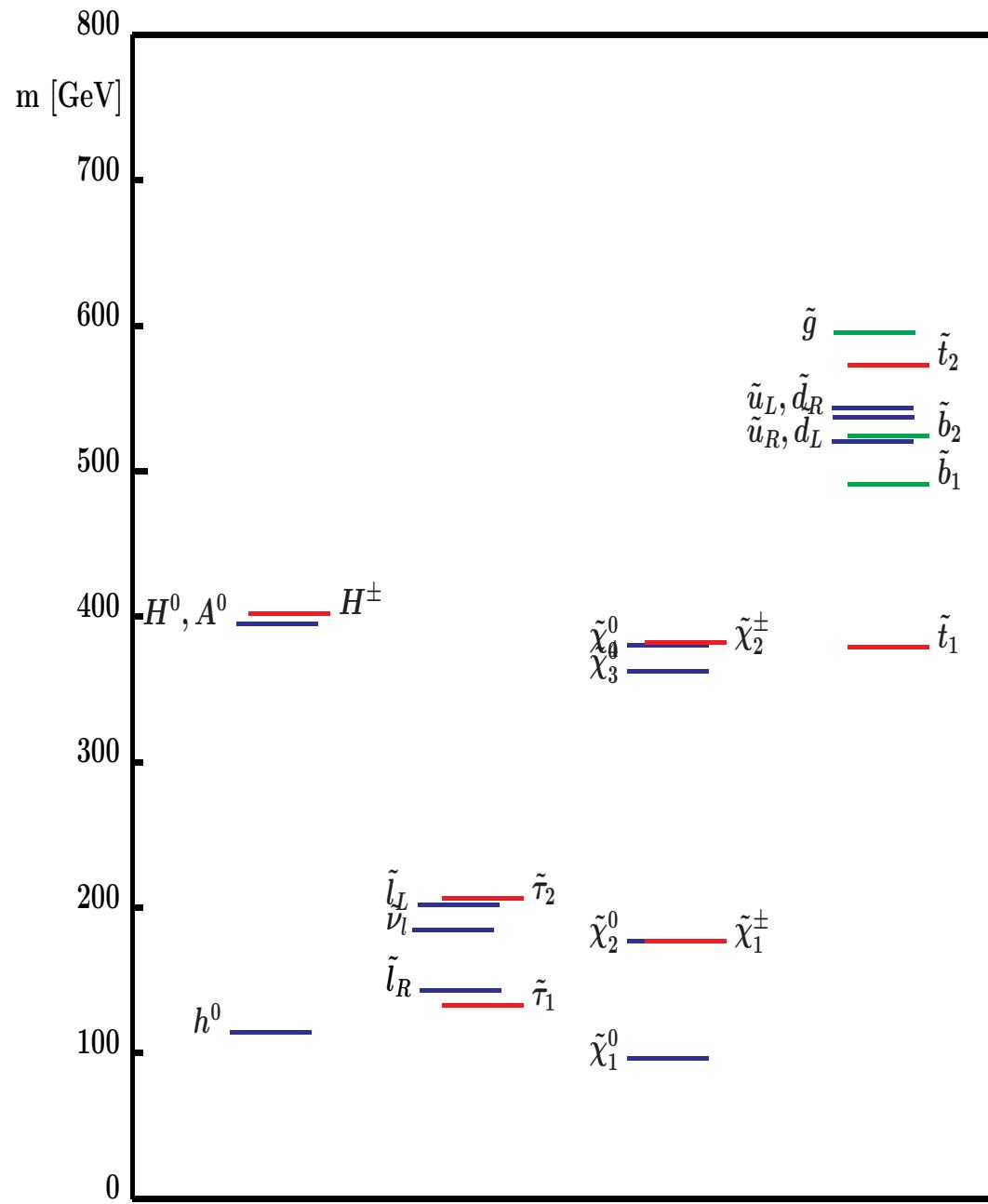
⇒ particle spectra from renormalization group running to weak scale



⇒ one parameter turns negative ⇒ Higgs mechanism for free

“Typical” CMSSM scenario  
(SPS 1a benchmark scenario):

Strong connection between  
all the sectors



## GUT based models: 2.) NUHM1: (Non-universal Higgs mass model)

Assumption: no unification of scalar fermion and scalar Higgs parameter at the GUT scale

⇒ effectively  $M_A$  as free parameters at the EW scale

⇒ Scenario characterized by

$m_0, m_{1/2}, A_0, \tan\beta, \text{sign } \mu$  and  $M_A$

## GUT based models: 3.) NUHM2: (Non-universal Higgs mass model 2)

Assumption: no unification of scalar Higgs parameter at the GUT scale

⇒ effectively  $M_A$  and  $\mu$  as free parameters at the EW scale

⇒ Scenario characterized by

$m_0, m_{1/2}, A_0, \tan\beta, \mu$  and  $M_A$

## Problem: We cannot be sure about the SUSY-breaking mechanism

- ⇒ it is possible that with the CMSSM, NUHM1, NUHM2, . . . we missed the “correct” mechanism
- ⇒ **hint: strong connection between colored and uncolored sector**  
tension between low-energy EW effects and (colored) LHC searches

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## Solution: investigate also the “general MSSM”

⇒ 8 parameters are manageable ⇒ pMSSM8

- 3rd gen. squark mass parameters:  $m_{\tilde{q}_3}$
- slepton mass parameter:  $m_{\tilde{l}_3}, m_{\tilde{l}_{1,2}}$
- gaugino masses:  $M_2$
- trilinear coupling:  $A_t$
- Higgs sector parameters:  $M_A, \tan \beta$
- Higgs mixing parameter:  $\mu$

⇒ Note: other “8 parameter selections” possible

## Data we have:

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- SUSY searches (LHC)  $\Rightarrow$  own re-cast
- electroweak precision data  $\Rightarrow$  FeynWZ, FeynHiggs
- flavor data  $\Rightarrow$  SuperIso, SuFla
- astrophysical data (DM properties)  $\Rightarrow$  MicrOMEGAs, SSARD

## 2. SUSY Higgs mass and rate measurements



## The Higgs mass accuracy: experiment vs. theory:

### Experiment:

ATLAS:  $M_h^{\text{exp}} = 125.36 \pm 0.37 \pm 0.18 \text{ GeV}$

CMS:  $M_h^{\text{exp}} = 125.03 \pm 0.27 \pm 0.15 \text{ GeV}$

combined:  $M_h^{\text{exp}} = 125.09 \pm 0.21 \pm 0.11 \text{ GeV}$

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### MSSM theory:

LHCHXSWG adopted **FeynHiggs** for the prediction of MSSM Higgs boson masses and mixings (considered to be the code containing the most complete implementation of higher-order corrections)

**FeynHiggs:**  $\delta M_h^{\text{theo}} \sim 3 \text{ GeV}$  (maybe 2 GeV?)

→ rough estimate, FeynHiggs contains algorithm to evaluate uncertainty, depending on parameter point

# Katharsis of Ultimate Theory Standards

7th meeting: 17.-19. July 2017, KIT (Karlsruhe, Germany)

Precise Calculation of  
 $(N)$

Higgs Boson masses

Local organizers: M. Muhlleitner, F. Staub, M. Steinhauser

Organized by:  
M. Carena, H. Haber  
R. Harlander, S. Heinemeyer  
W. Hollik, P. Slavich, G. Weiglein

⇒ next meeting: 01/2018 in Paris

## Enlarged Higgs sector: Two Higgs doublets

$$H_1 = \begin{pmatrix} H_1^1 \\ H_1^2 \end{pmatrix} = \begin{pmatrix} v_1 + (\phi_1 + i\chi_1)/\sqrt{2} \\ \phi_1^- \end{pmatrix}$$

$$H_2 = \begin{pmatrix} H_2^1 \\ H_2^2 \end{pmatrix} = \begin{pmatrix} \phi_2^+ \\ v_2 + (\phi_2 + i\chi_2)/\sqrt{2} \end{pmatrix}$$

$$V = m_1^2 H_1 \bar{H}_1 + m_2^2 H_2 \bar{H}_2 - m_{12}^2 (\epsilon_{ab} H_1^a H_2^b + \text{h.c.})$$

$$+ \underbrace{\frac{g'^2 + g^2}{8}}_{\text{gauge couplings, in contrast to SM}} (H_1 \bar{H}_1 - H_2 \bar{H}_2)^2 + \underbrace{\frac{g^2}{2}}_{\text{gauge couplings, in contrast to SM}} |H_1 \bar{H}_2|^2$$

gauge couplings, in contrast to SM

physical states:  $h^0, H^0, A^0, H^\pm$

Goldstone bosons:  $G^0, G^\pm$

Input parameters: (to be determined experimentally)

$$\tan \beta = \frac{v_2}{v_1}, \quad M_A^2 = -m_{12}^2(\tan \beta + \cot \beta)$$

## $\tilde{t}$ sector of the MSSM:

Stop mass matrices

$$M_{\tilde{t}}^2 = \begin{pmatrix} M_{\tilde{t}_L}^2 + m_t^2 + DT_{t_1} & m_t X_t \\ m_t X_t & M_{\tilde{t}_R}^2 + m_t^2 + DT_{t_2} \end{pmatrix} \xrightarrow{\theta_{\tilde{t}}} \begin{pmatrix} m_{\tilde{t}_1}^2 & 0 \\ 0 & m_{\tilde{t}_2}^2 \end{pmatrix}$$

with

$$X_t = A_t - \mu / \tan \beta$$

⇒ mixing important in stop sector!

Simplifying abbreviation:

$$M_{\text{SUSY}} := M_{\tilde{t}_L} = M_{\tilde{t}_R}$$

## Higgs boson mass scales from rate measurements?

We have a  $\sim 125$  GeV SM-like Higgs boson

⇒ What are the options?

1. Decoupling limit:

$M_A \gg M_Z$  ⇒ the light Higgs becomes SM-like

2. Alignment without decoupling:

⇒ a  $\mathcal{CP}$ -even Higgs becomes SM-like due to an “accidental” cancellation

3. Heavy Higgs SM-like: (see above!)

⇒ is the case with the heavy  $\mathcal{CP}$ -even Higgs being SM-like still a viable solution?

## Obtaining a light Higgs with SM-like couplings

[J. Gunion, H. Haber, [hep-ph/0207010](#)]

→  $\mathcal{CP}$  conserving 2HDM in the Higgs basis ( $\langle H_1 \rangle = v/\sqrt{2}$ ,  $\langle H_2 \rangle = 0$ )

$$\mathcal{V} = \dots + \frac{1}{2} Z_1 (H_1^\dagger H_1)^2 + \dots + \left[ \frac{1}{2} Z_5 (H_1^\dagger H_2)^2 + Z_6 (H_1^\dagger H_1)(H_1^\dagger H_2) + \text{h.c.} \right] + \dots$$

⇒  $\mathcal{CP}$ -even mass matrix:

$$\mathcal{M}^2 = \begin{pmatrix} Z_1 v^2 & Z_6 v^2 \\ Z_6 v^2 & M_A^2 + Z_5 v^2 \end{pmatrix}$$

with mixing angle  $\cos(\beta - \alpha) \equiv c_{\beta-\alpha}$

Decoupling limit:  $M_A^2 \gg Z_i v^2$

⇒  $m_h^2 \sim Z_1 v^2$ ,  $|c_{\beta-\alpha} \ll 1|$ ,  $h$  is SM-like

Alignment limit:  $Z_6 = 0$  and  $Z_1 < Z_5 + M_A^2/v^2$

⇒  $h$  is identical to the SM Higgs,  $c_{\beta-\alpha} = 0$

$Z_6 = 0$  and  $Z_1 > Z_5 + M_A^2/v^2$

⇒  $H$  is identical to the SM Higgs,  $c_{\beta-\alpha} = 1$

## Alignment limit: see e.g.

[M. Carena, I. Low, N. Shah, C. Wagner '13][M. Carena, H. Haber, I. Low, N. Shah, C. Wagner '14]

In the **MSSM**  $Z_6 = 0$  can be obtained through an “accidental” cancellation between tree-level and loop contribution, roughly at:

$$\tan \beta \sim \left[ M_h^2 + M_Z^2 + \frac{3m_t^2\mu^2}{4\pi^2 v^2 M_S^2} \left( \frac{A_t^2}{2M_S^2} - 1 \right) \right] / \left[ \frac{3m_t^2}{4\pi^2 v^2} \frac{\mu A_t}{M_S^2} \left( \frac{A_t^2}{6M_S^2} - 1 \right) \right]$$

Compare:  $m_h^{\text{mod+}}$  and  $m_h^{\text{alt}}$ :

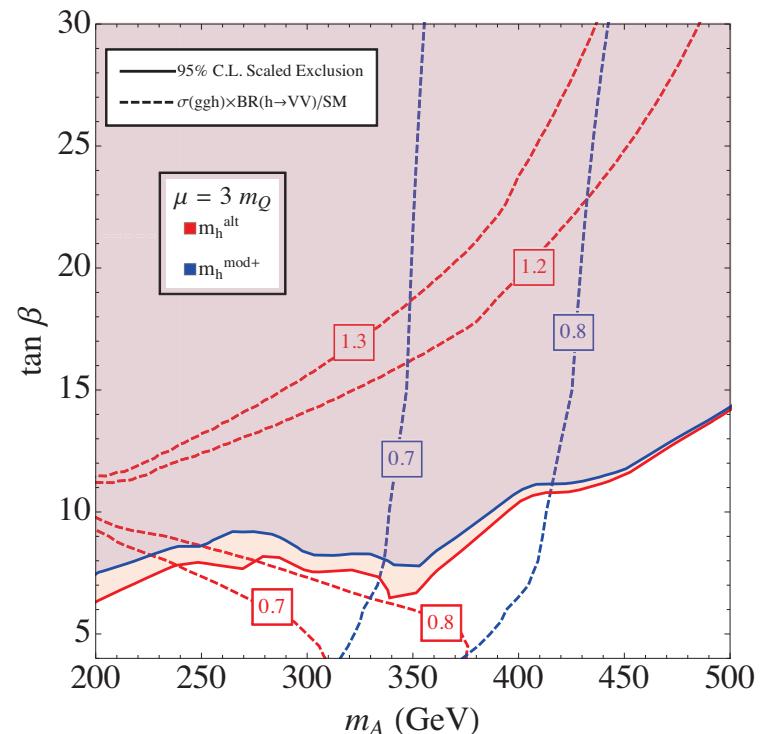
$$A_t/M_S = 2.45, \quad A_t = A_f,$$

$$M_S = m_{\tilde{f}} \geq 1 \text{ TeV}, \quad m_{\tilde{g}} = 1.5 \text{ TeV},$$

$$M_2 = 2M_1 = 200 \text{ GeV}, \quad \mu \text{ adjustable}$$

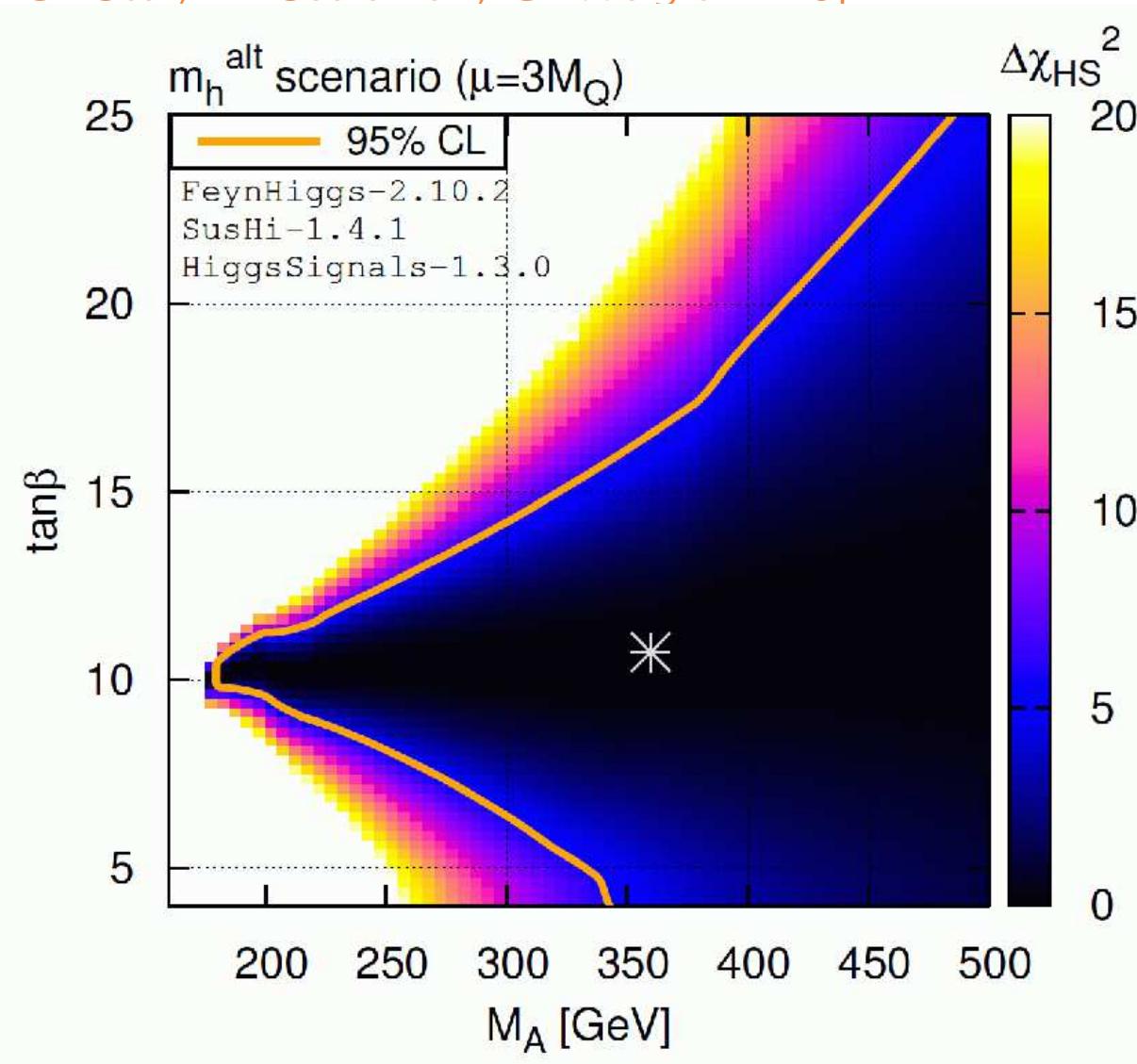
(low  $M_A$  and  $\tan \beta$ : tune  $M_S \geq 1 \text{ TeV}$  to obtain  $M_h \geq 122 \text{ GeV}$ )

⇒ SM-like Higgs for all  $M_A$



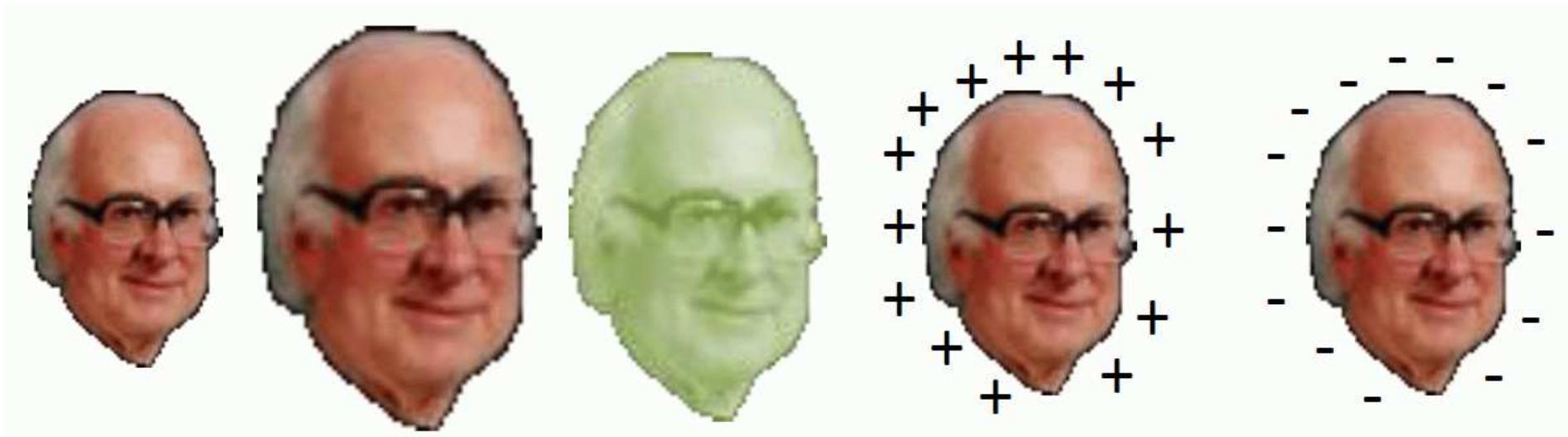
## Preferred $m_h^{\text{alt}}$ parameter space from HiggsSignals:

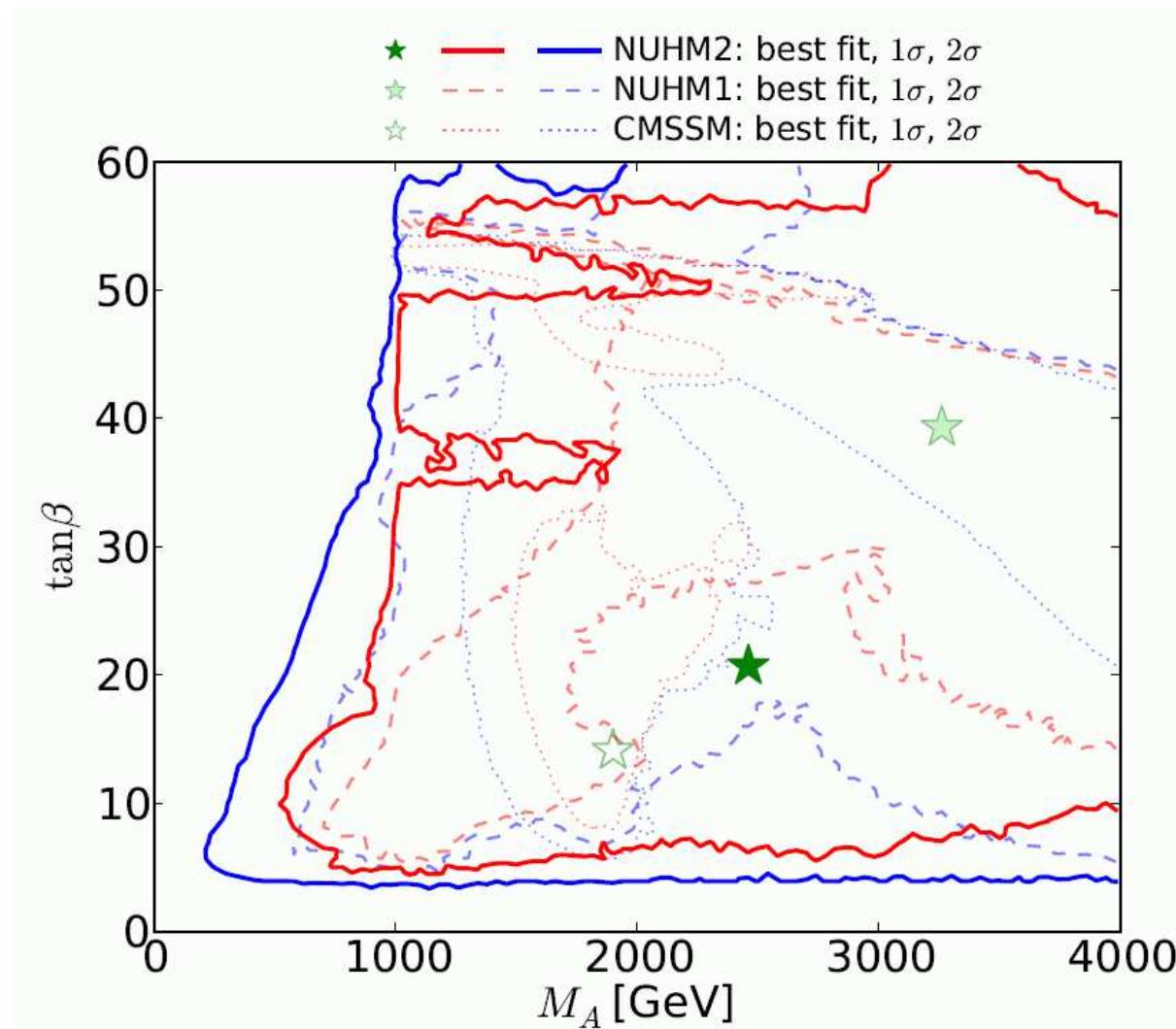
[P. Bechtle, S.H., O. Stål, T. Stefaniak, G. Weiglein '15]



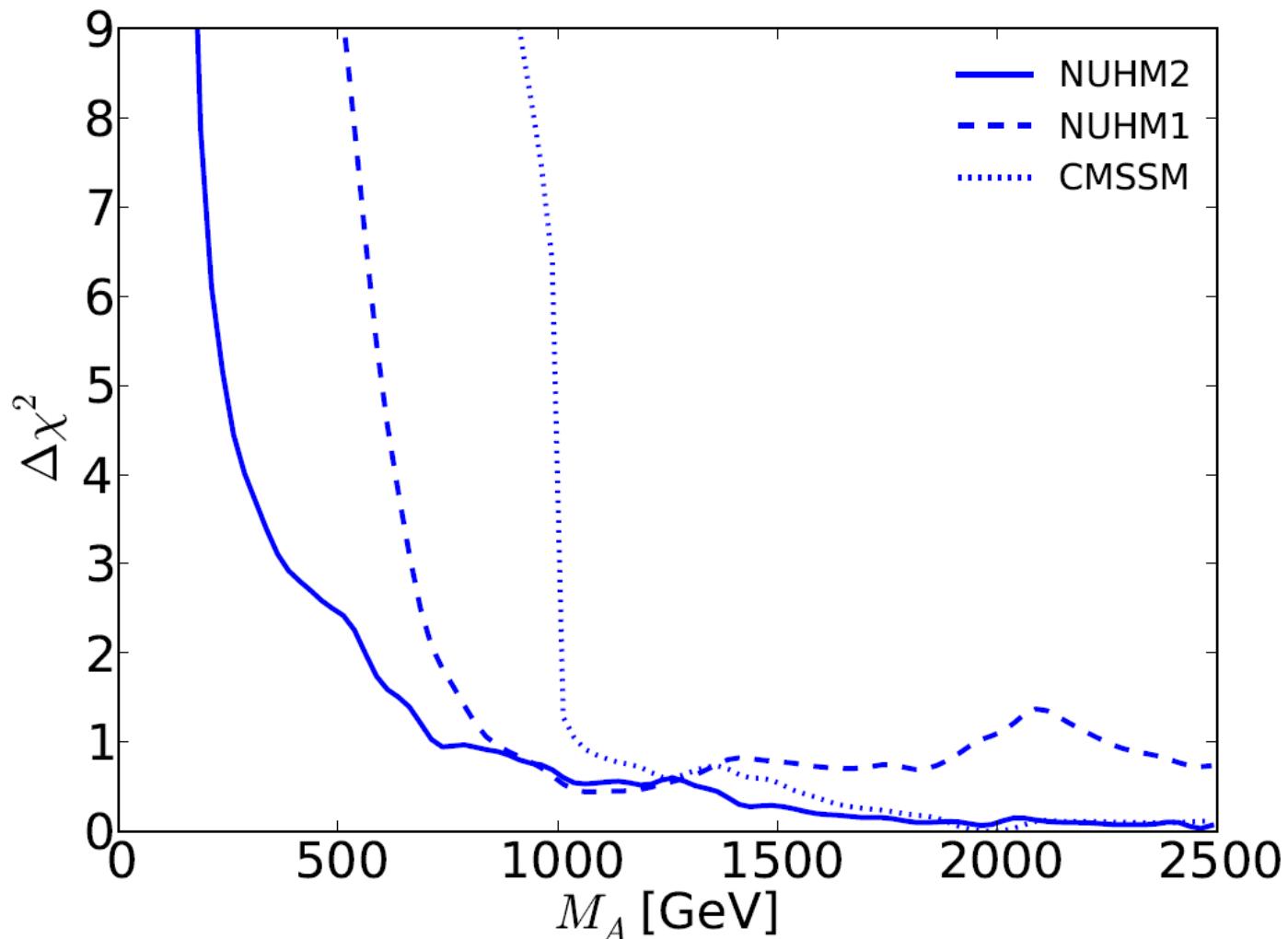
→ no Higgs mass scale restrictions from rates (in general)

### 3. MSSM Higgs results (CMSSM, NUHM1/2, pMSSM8)



$M_A$ - $\tan\beta$  plane in CMSSM, NUHM1, NUHM2:


→ high mass scales, only in NUHM2 lighter Higgs bosons . . .

$M_A$ - $\Delta\chi^2$  in CMSSM, NUHM1, NUHM2:


⇒ high mass scales, only in NUHM2 lighter Higgs bosons . . .

## Results in the pMSSM8

[*P. Bechtle, H. Haber, S.H., O. Stål, T. Stefaniak, G. Weiglein, L. Zeune '16*]

- decoupling,  $M_h = 125 \text{ GeV}$
- alignment without decoupling,  $M_h = 125 \text{ GeV}$
- “heavy Higgs” case,  $M_H = 125 \text{ GeV}$ ,  $h$  lighter

	Min	Max
$M_A$	90 GeV	1000 GeV
$\tan \beta$	1	60
$M_{Q_3}$	200 GeV	5000 GeV
$A_t$	$-3M_{Q_3}$	$+3M_{Q_3}$
$\mu$	$-3M_{Q_3}$	$+3M_{Q_3}$
$M_{L_3}$	200 GeV	1000 GeV
$M_{L_{1,2}}$	200 GeV	1000 GeV
$M_2$	200 GeV	500 GeV

$$M_{Q_{1,2}} = M_{U_{1,2}} = M_{D_{1,2}} = 1.5 \text{ TeV}$$

$$M_{D_3} = M_{U_3} = M_{Q_3}$$

$$M_{L_{1,2}} = M_{E_{1,2}}$$

$$A_b = A_\tau = A_t$$

$$M_3 = 1.5 \text{ TeV}$$

$M_1$  fixed by GUT relation

$10^7$  random points

$$R_{XX}^\phi := \frac{\sum_i [\sigma_i(\phi) \times \text{BR}(\phi \rightarrow XX)]_{\text{MSSM}}}{\sum_i [\sigma_i(\phi) \times \text{BR}(\phi \rightarrow XX)]_{\text{SM}}}$$

use [FeynHiggs-2.10.2](#) and [SuperIso-3.3](#) for MSSM predictions.

Construct global  $\chi^2$  from observables:

- Higgs mass and signal rates ([HiggsSignals-1.4.0](#))
- Low energy observables (LEO):  $b \rightarrow s\gamma$ ,  $B_s \rightarrow \mu\mu$ ,  $B_u \rightarrow \tau\nu$ ,  $(g - 2)_\mu$ ,  $M_W$
- exclusion likelihood from CMS  $\phi \rightarrow \tau\tau$  search ([HiggsBounds-4.2.0](#))
- LEP Higgs exclusion likelihood,  $\chi^2_{\text{LEP}}$ , if relevant. ([HiggsBounds-4.2.0](#))

Further constraints:

- 95% CL Higgs exclusion limits (w/o MSSM  $\phi \rightarrow \tau\tau$  limits) ([HiggsBounds-4.2.0](#))
- Sparticle mass limits from LEP, (fixed  $m_{\tilde{q}_{1,2}} = m_{\tilde{g}} = 1.5$  TeV to evade LHC limits)
- Require neutral lightest supersymmetric particle (LSP).

Newly included: [CheckMate](#) to check SUSY exclusion limits

⇒ “naive”  $\chi^2$  calculation (heavily relying on [HiggsSignals](#))

## The best-fit points:

Case	full fit			fit without $a_\mu$			fit without all LEOs		
	$\chi^2/\nu$	$\chi^2_\nu$	$p$	$\chi^2/\nu$	$\chi^2_\nu$	$p$	$\chi^2/\nu$	$\chi^2_\nu$	$p$
SM	83.7/91	0.92	0.69	72.4/90	0.80	0.91	70.2/86	0.82	0.89
$h$	68.5/84	0.82	0.89	68.2/83	0.82	0.88	67.9/79	0.86	0.81
$H$	73.7/85	0.87	0.80	71.9/84	0.86	0.82	70.0/80	0.88	0.78

## Best-fit points parameters:

Case	$M_A$ (GeV)	$\tan \beta$	$\mu$ (GeV)	$A_t$ (GeV)	$M_{\tilde{q}_3}$ (GeV)	$M_{\tilde{\ell}_3}$ (GeV)	$M_{\tilde{\ell}_{1,2}}$ (GeV)	$M_2$ (GeV)
$h$	929	21.0	7155	4138	2957	698	436	358
$H$	172	6.6	4503	-71	564	953	262	293

⇒ SM and both MSSM cases provide similar fit to the Higgs data

⇒ Including LEOs, SM fit becomes worse

## The “exotic” solution:

the discovery is interpreted as the heavy  $\mathcal{CP}$ -even Higgs

In principle also possible:

$$M_h < 125 \text{ GeV}$$

$$M_H \approx 125 \text{ GeV}$$

Consequences:

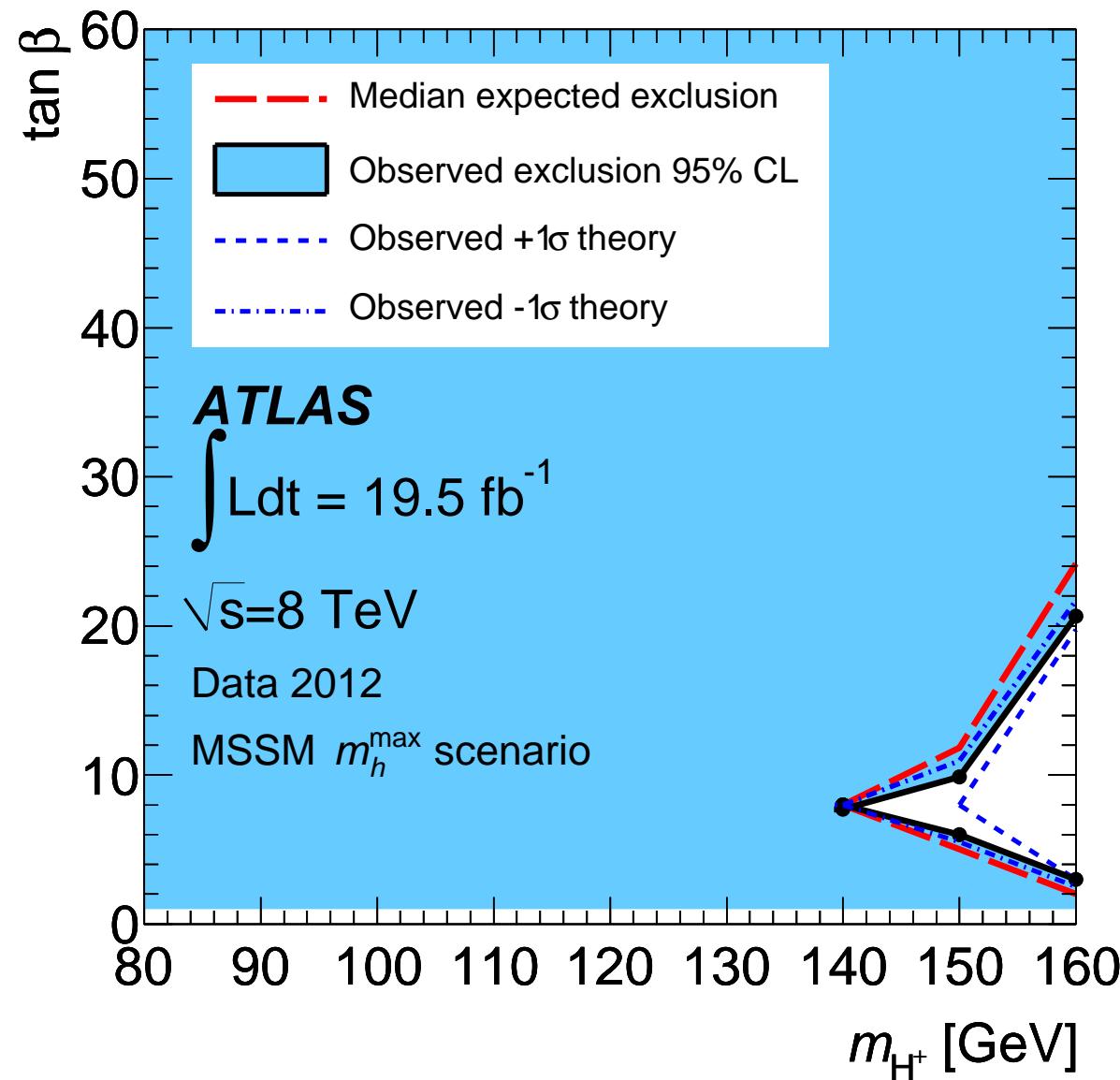
- all Higgs bosons very light
- easy(?) discovery of additional Higgs bosons at the LHC

Constraints:

- direct searches for the lightest  $\mathcal{CP}$ -even Higgs
- direct searches for the heavy neutral Higgses
- direct searches for the charged Higgses
- flavor constraints ( $\text{BR}(B_s \rightarrow \mu^+ \mu^-)$  etc.)

⇒ original scenario: low- $M_H$

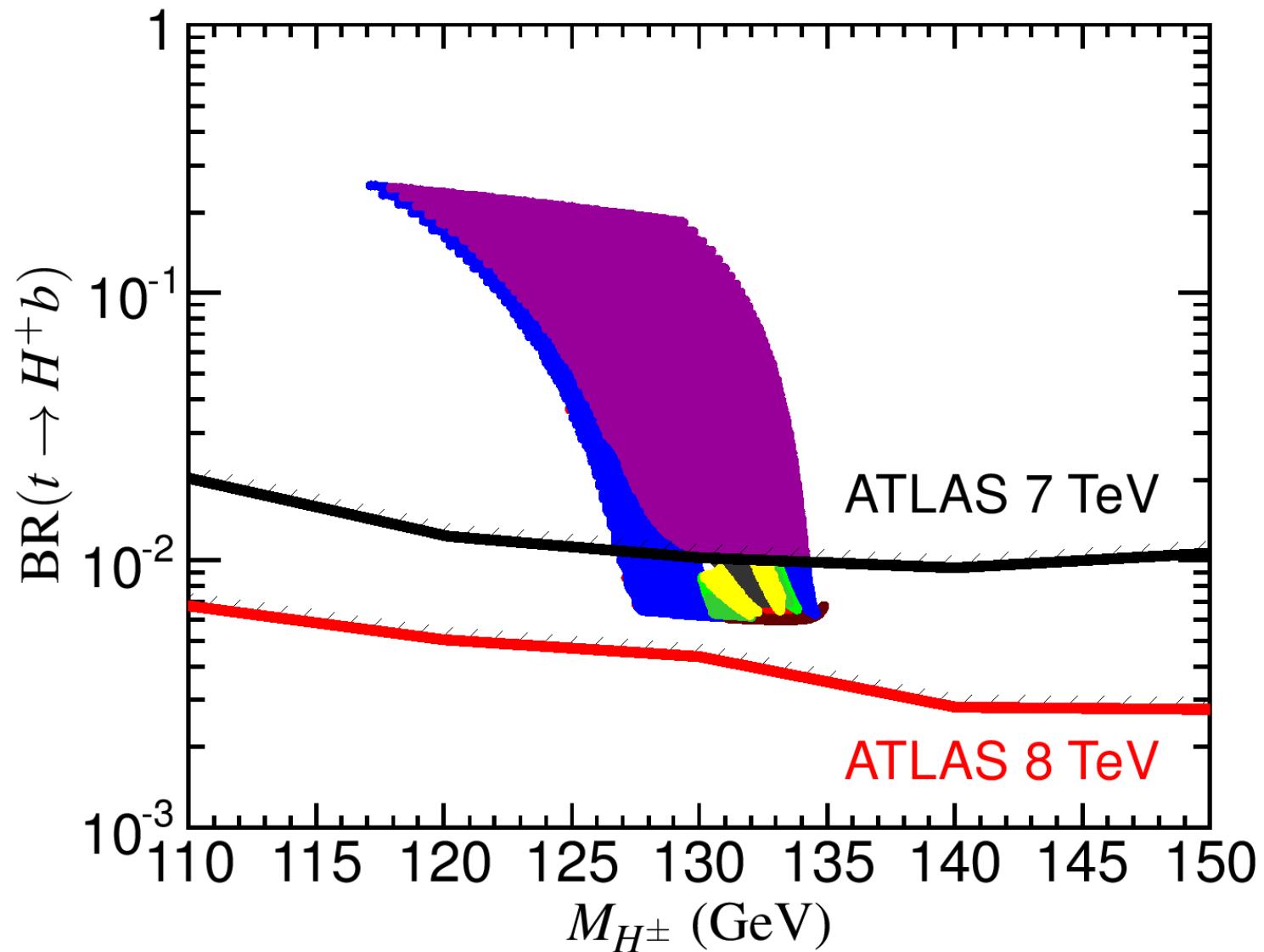
[M. Carena, S.H., O. Stål, C. Wagner, G. Weiglein '13]



→ exclusion of light  $M_{H^\pm}$  in the  $m_h^{\max}$  scenario! . . . low- $M_H$ ?

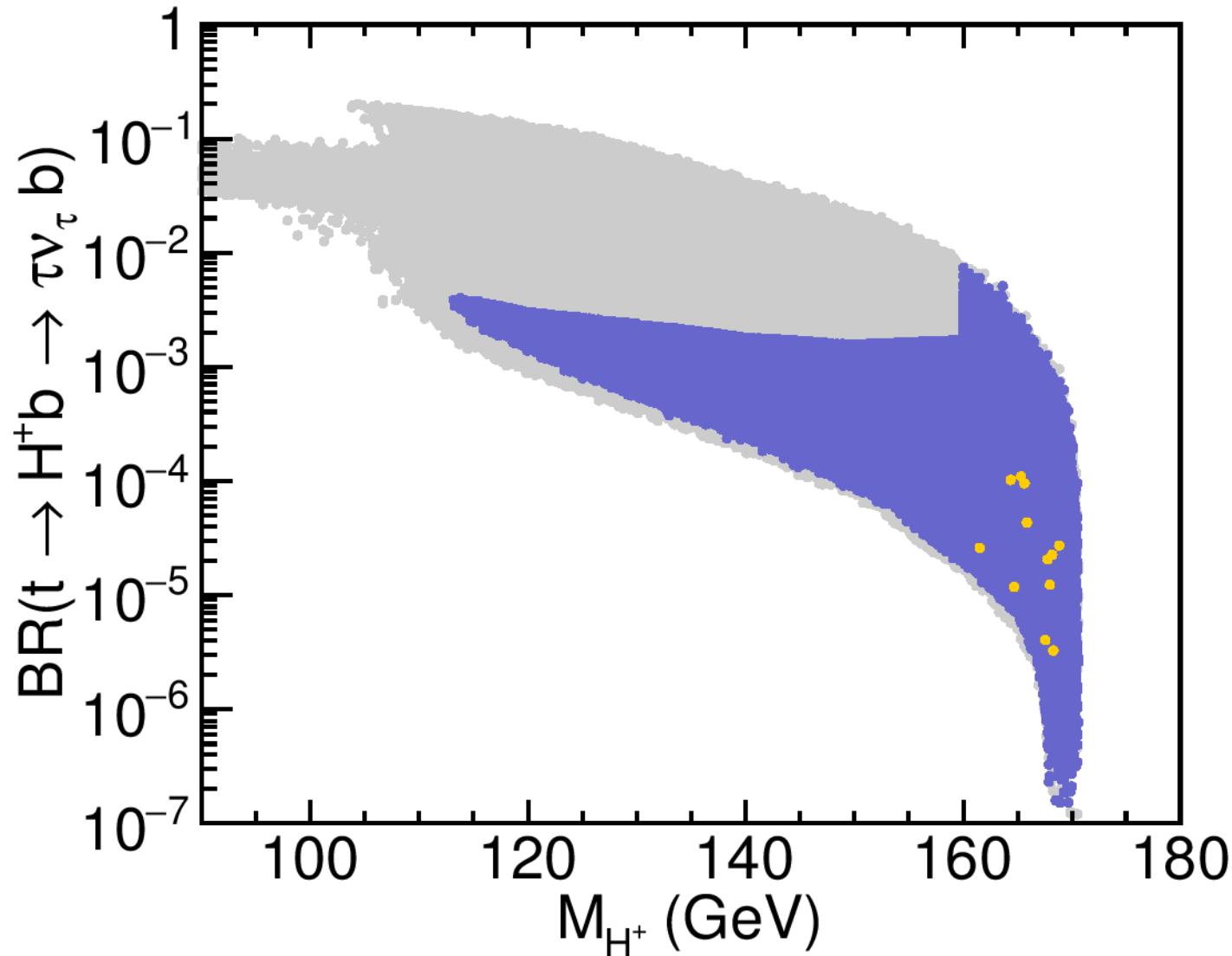
## Application of charged Higgs limits on low- $M_H$ scenario:

[HiggsBounds 4.1]



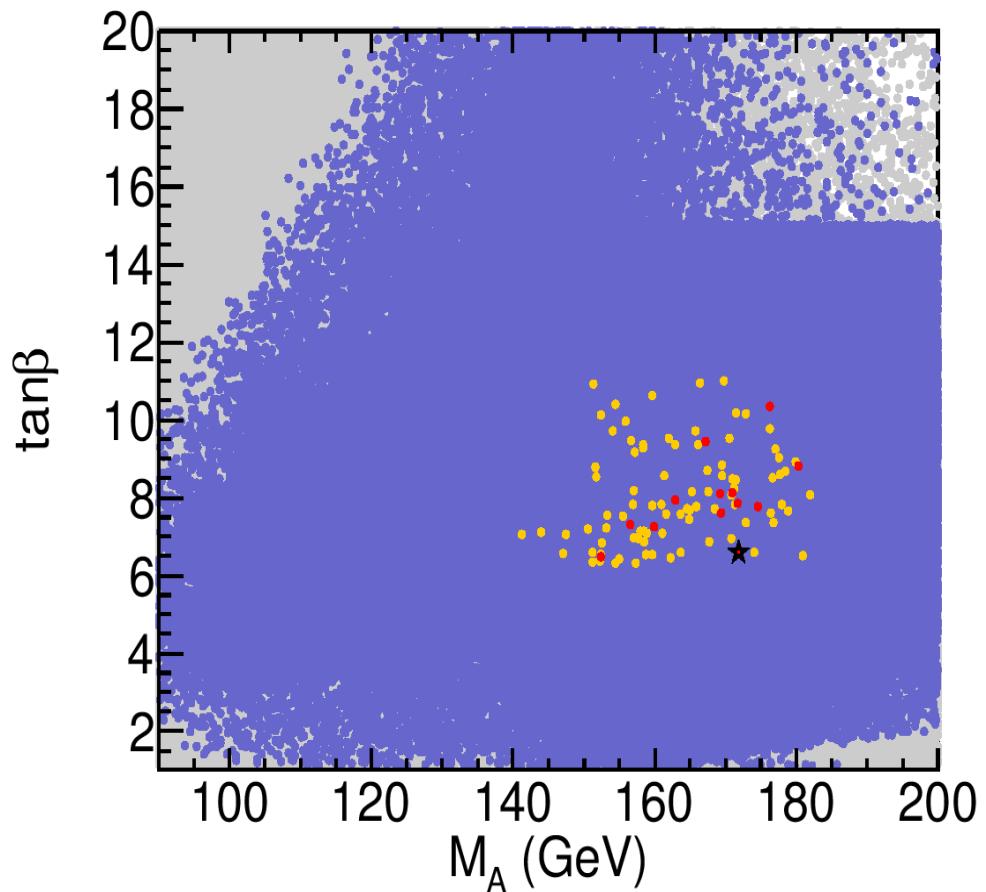
⇒ that (particular incarnation of the) low- $M_H$  scenario is excluded!

How to avoid  $\text{BR}(t \rightarrow H^\pm b)$  bounds:  $\Rightarrow$  higher  $M_{H^\pm}$ !



$\Rightarrow$  “tricky” region below and beyond the top threshold!

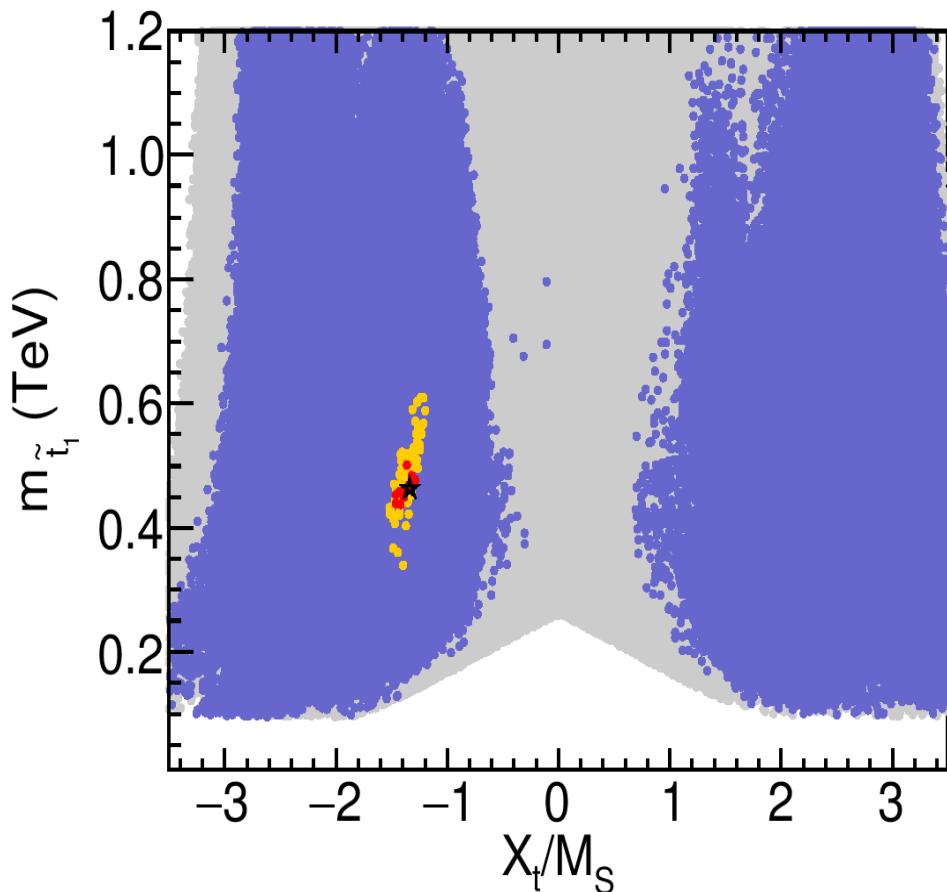
## Heavy-Higgs case: preferred parameters



$\Rightarrow M_A \sim 140 \dots 180 \text{ GeV}$

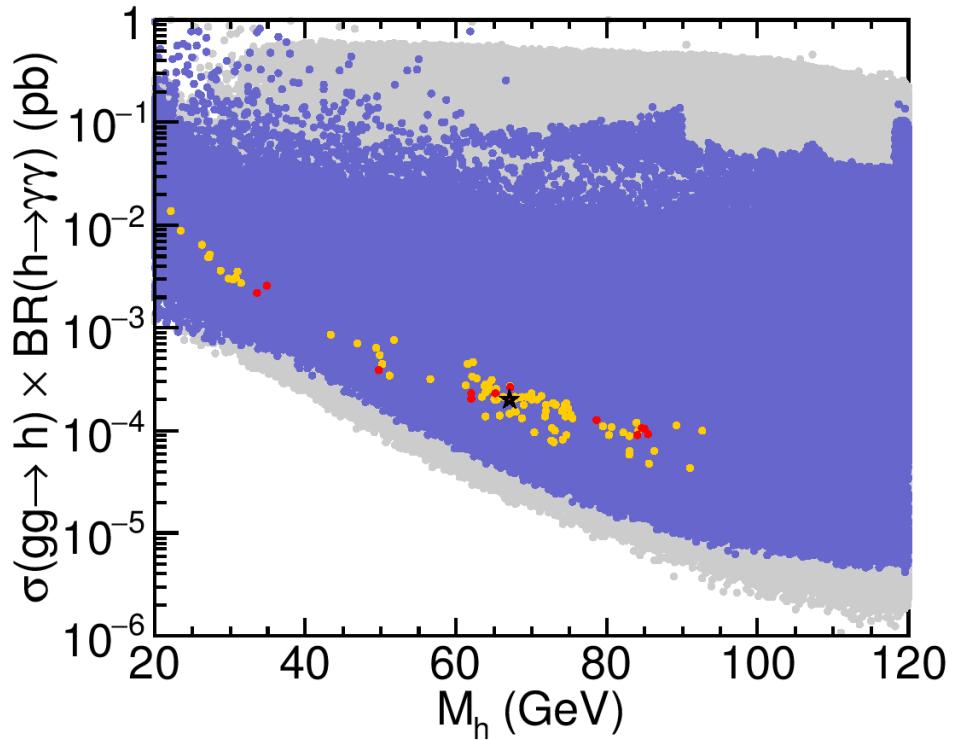
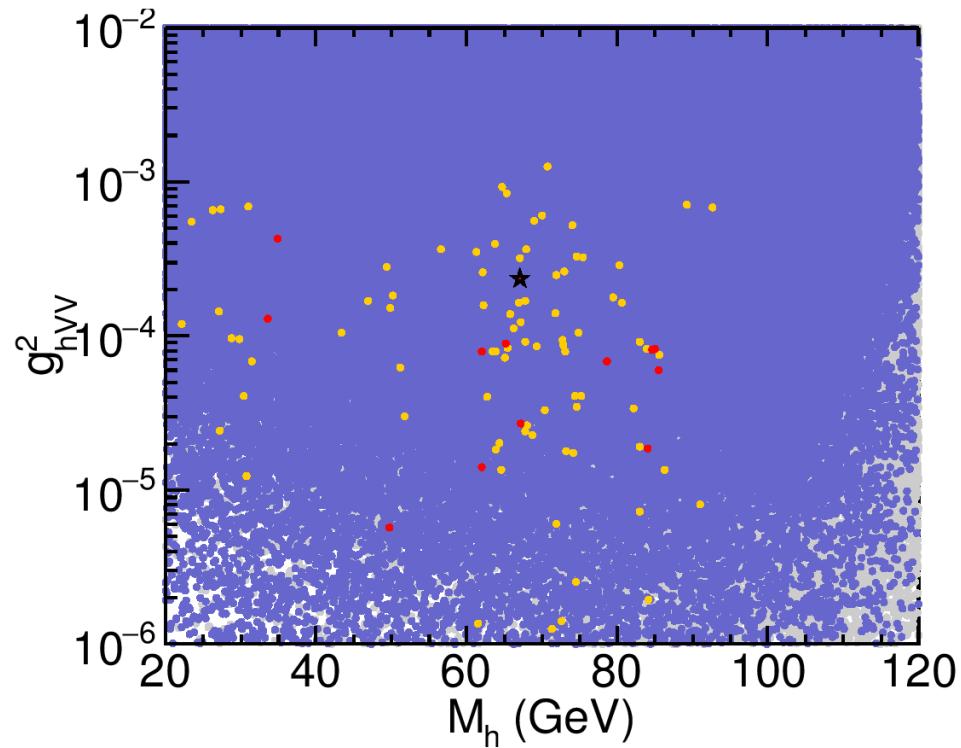
$R_{VV}^H = [0.95, 1.13]$ ,  $R_{\gamma\gamma}^H = [0.81, 0.94]$ ,  $R_{bb}^{VH} = [0.94, 1.03]$ ,  $R_{\tau\tau}^H = [0.78, 0.90]$

$\Rightarrow$  not fully SM-like ...



$\Rightarrow m_{\tilde{t}_1} \sim 350 \dots 650 \text{ GeV}$

## Where is the light Higgs?



⇒ strongly reduced couplings to gauge bosons ⇒ beyond LEP reach!

⇒  $M_h > M_H/2$  (mostly) to avoid  $H \rightarrow hh$  (or  $BR(H \rightarrow hh) \lesssim 10\%$ )

⇒ visible in  $gg \rightarrow h \rightarrow \gamma\gamma$ ?

⇒  $M_h \sim 95$  GeV possible?!

## New low- $M_H$ benchmark scenarios

⇒ currently investigated by the LHCHXSWG!

Based on our best-fit region:

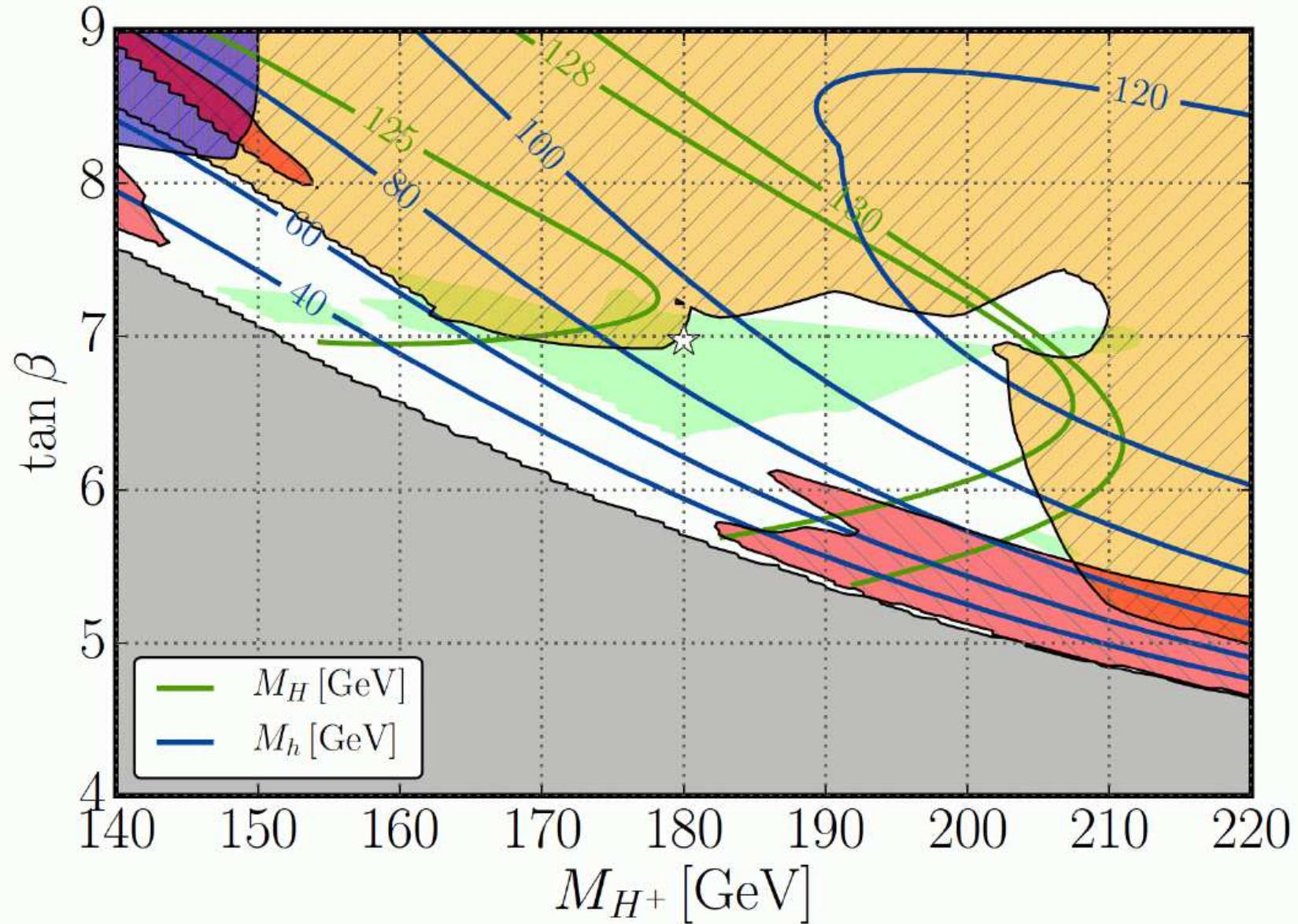
Benchmark scenario	$M_{H^\pm}$ [GeV]	$\mu$ [GeV]	$\tan \beta$
low- $M_H^{\text{alt},-}$	155	3800 – 6500	4 – 9
low- $M_H^{\text{alt},+}$	185	4800 – 7000	4 – 9
low- $M_H^{\text{alt},v}$	140 – 220	6000	4 – 9
fixed parameters:	$m_t = 173.2$ GeV, $A_t = A_\tau = A_b = -70$ GeV, $M_2 = 300$ GeV, $M_{\tilde{q}_L} = M_{\tilde{q}_R} = 1500$ GeV ( $q = c, s, u, d$ ), $m_{\tilde{g}} = 1500$ GeV, $M_{\tilde{q}_3} = 750$ GeV, $M_{\tilde{\ell}_{1,2}} = 250$ GeV, $M_{\tilde{\ell}_3} = 500$ GeV		

low- $M_H^{\text{alt}-}$ : fixed  $M_{H^\pm} < m_t$

low- $M_H^{\text{alt}+}$ : fixed  $M_{H^\pm} > m_t$

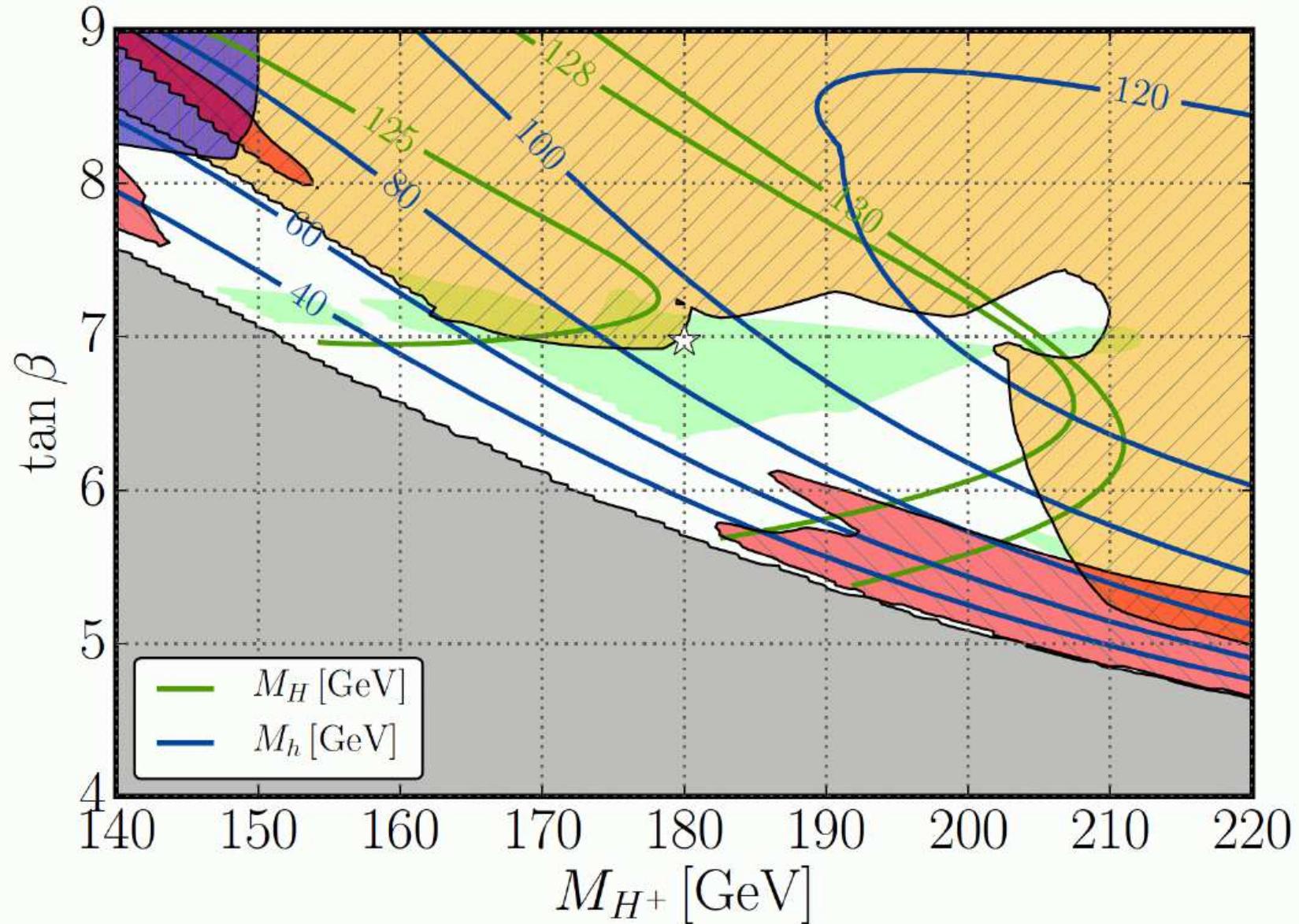
low- $M_H^{\text{alt}v}$ : varied  $M_{H^\pm}$  ( $\mu$  fixed)

low- $M_H^{\text{alt}\nu}$  ( $140 \text{ GeV} \leq M_{H^\pm} \leq 220 \text{ GeV}$ ):



⇒ green area in agreement with all data!

low- $M_H^{\text{alt}\nu}$  ( $140 \text{ GeV} \leq M_{H^\pm} \leq 220 \text{ GeV}$ ):



⇒ green area in agreement with all data!

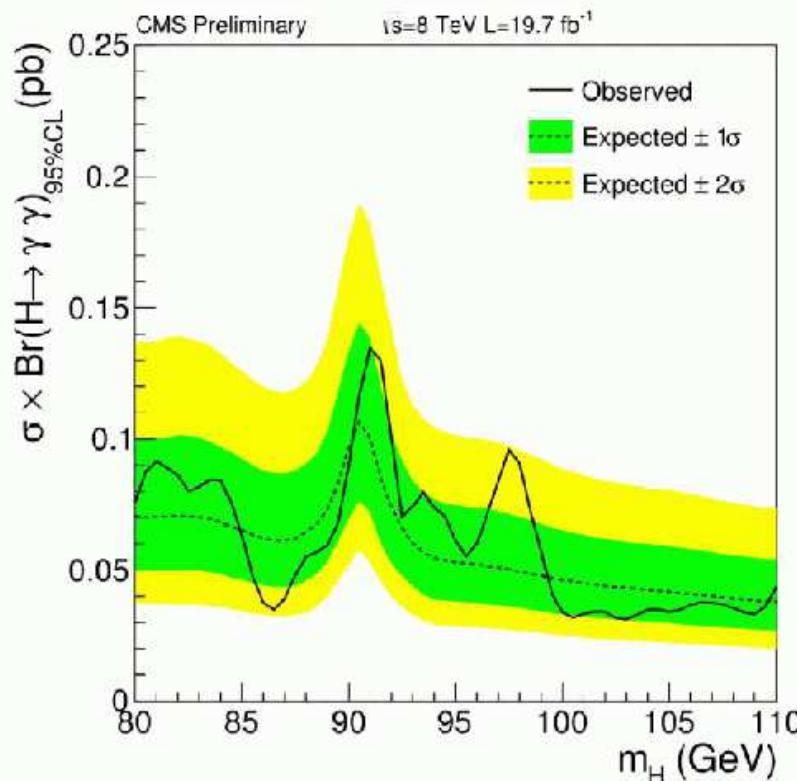
Go and exclude it!

## 4. $\phi_{95} \rightarrow \gamma\gamma$

- What was seen in Run I?
- What was seen in Run II?
- What was seen at LEP?
- Should we get excited?
- Which model fits?

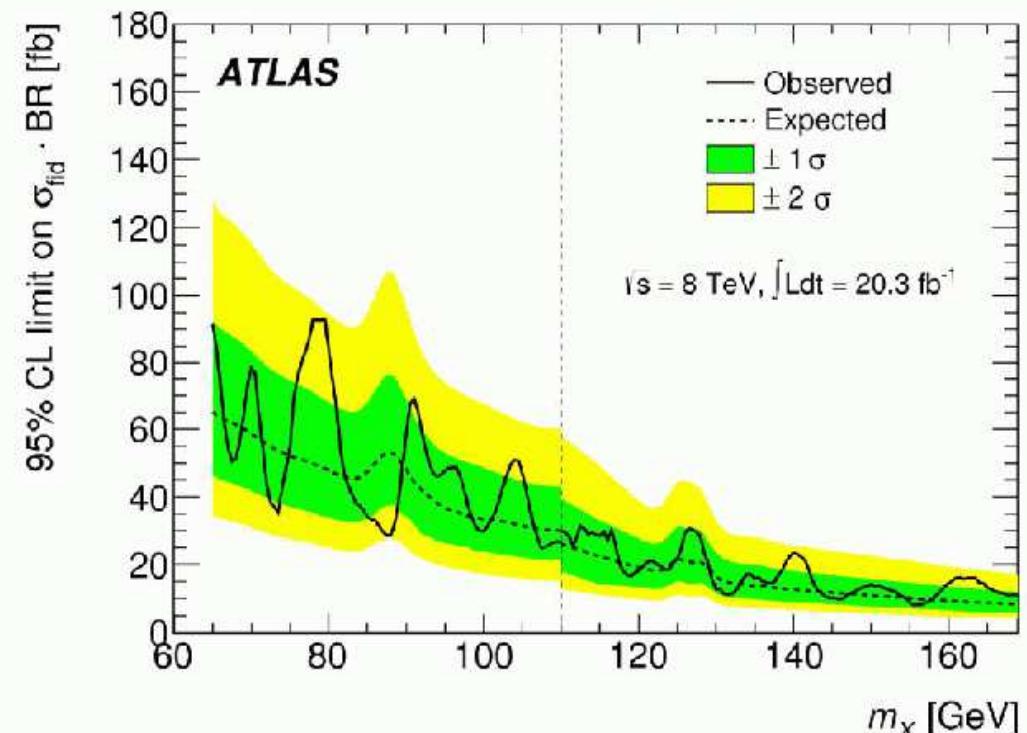


CMS PAS HIG-14-037



# $h \rightarrow \gamma\gamma$ (65-110GeV) Run 1

PRL 113 171801 (2014)



- $\sim 2\sigma$  excursion @  $\sim 97.5 \text{ GeV}$

- $\sim 2\sigma$  excursion @  $\sim 80 \text{ GeV}$

18

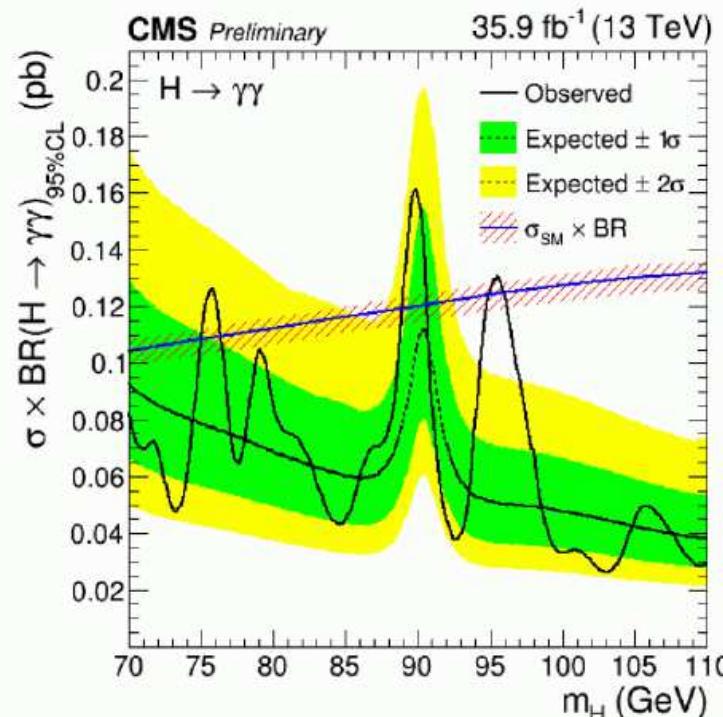
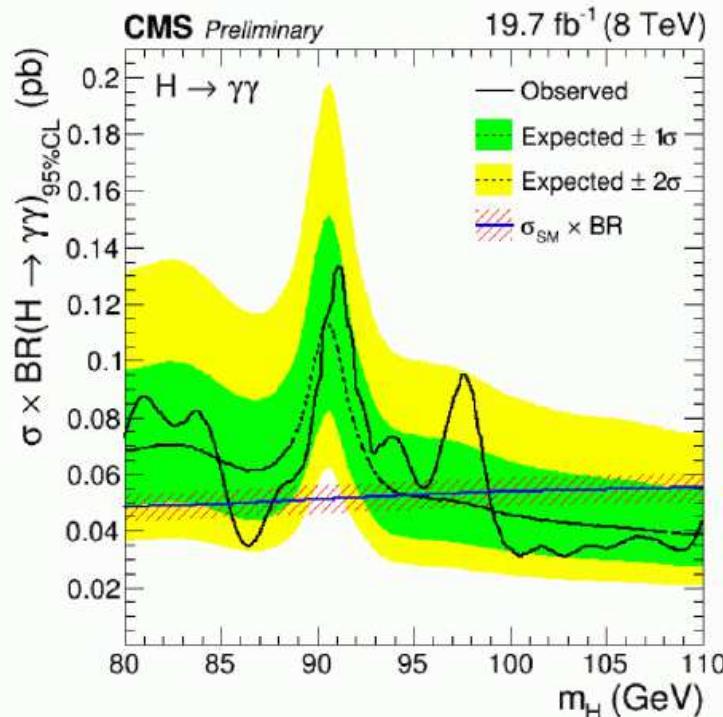
S. Gascon-Shotkin HDays17, Santander, ES Sept. 22 2017



# $h \rightarrow \gamma\gamma$ (70-110 GeV) Runs 1+2

New!

CMS PAS HIG-17-013



8 TeV:  
minimum(maximum)  
limit on  $\sigma \times \text{Br}$  :  
 $31(133) \text{ fb}$  at  
 $m=102.8(91.1) \text{ GeV}$

13 TeV:  
minimum(maximum)  
limit on  $\sigma \times \text{Br}$  :  
 $26(161) \text{ fb}$  at  
 $m=103.0(89.9) \text{ GeV}$

- 8 TeV limits on  $\sigma \times \text{Br}$  redone with 0.1 GeV step. Production processes assumed in SM proportions. No significant excess with respect to expected limits observed.

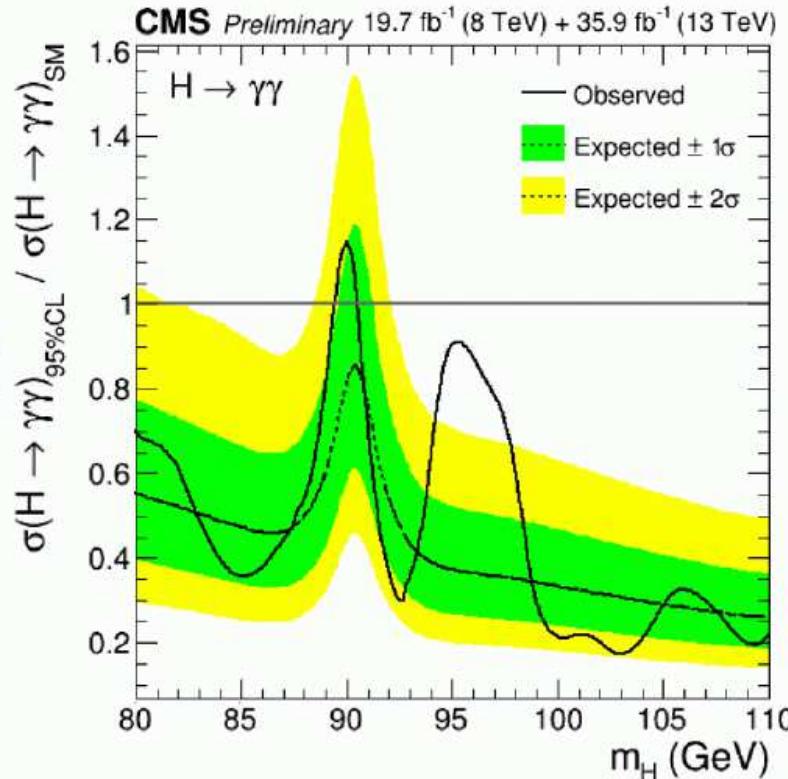
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# $h \rightarrow \gamma\gamma$ (70-110 GeV) Runs 1+2

All experimental + theoretical systematic uncertainties assumed uncorrelated except for those on signal acceptance due to scale variations + those on production cross sections (assumed 100% correlated).



- Combined 8 TeV+13 TeV  $\sigma \times BR$  limit normalized to SM expectation (production processes assumed in SM proportions). No significant excess with respect to expected limits observed.



CMS PAS HIG-17-013

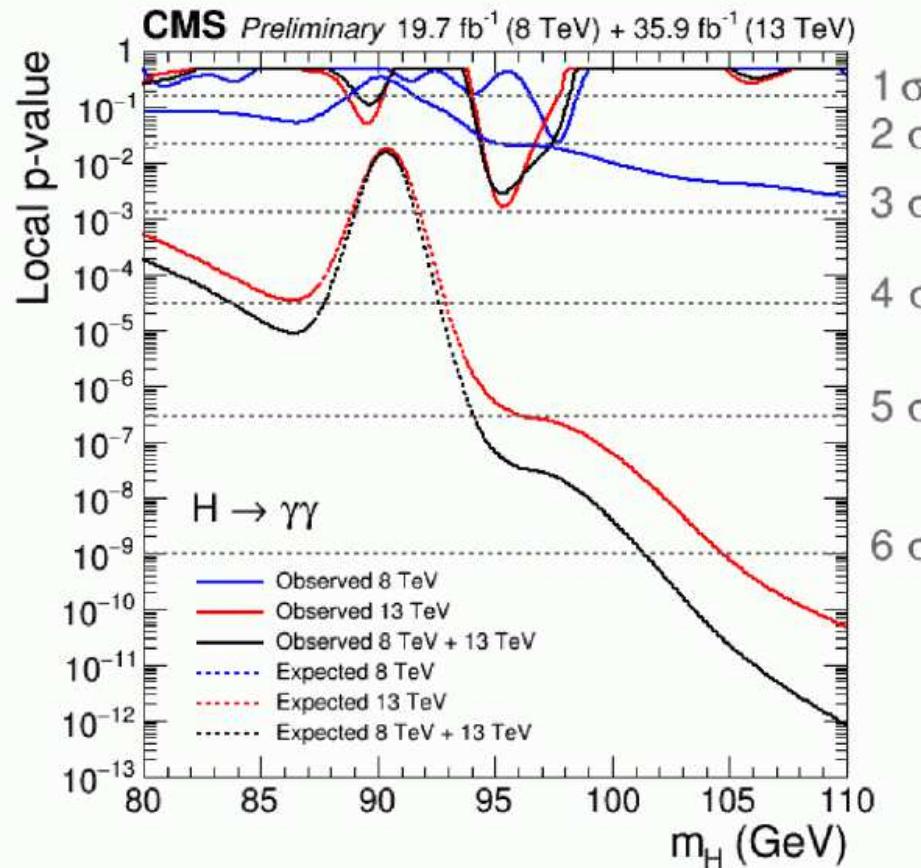
8 TeV+13 TeV:  
minimum(maximum) limit  
on  $(\sigma \times Br) / (\sigma \times Br)_{SM}$  :  
0.17(1.15) at  
 $m=103.0(90.0)$ GeV

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# $h \rightarrow \gamma\gamma$ (70-110 GeV) Runs 1+ 2



- Expected and observed local p-values for **8 TeV**, **13 TeV** and their combination

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8 TeV: Excess with  $\sim 2.0 \sigma$  local significance at  $m=97.6$  GeV

13 TeV: Excess with  $\sim 2.9 \sigma$  local ( $1.47 \sigma$  global) significance at  $m=95.3$  GeV

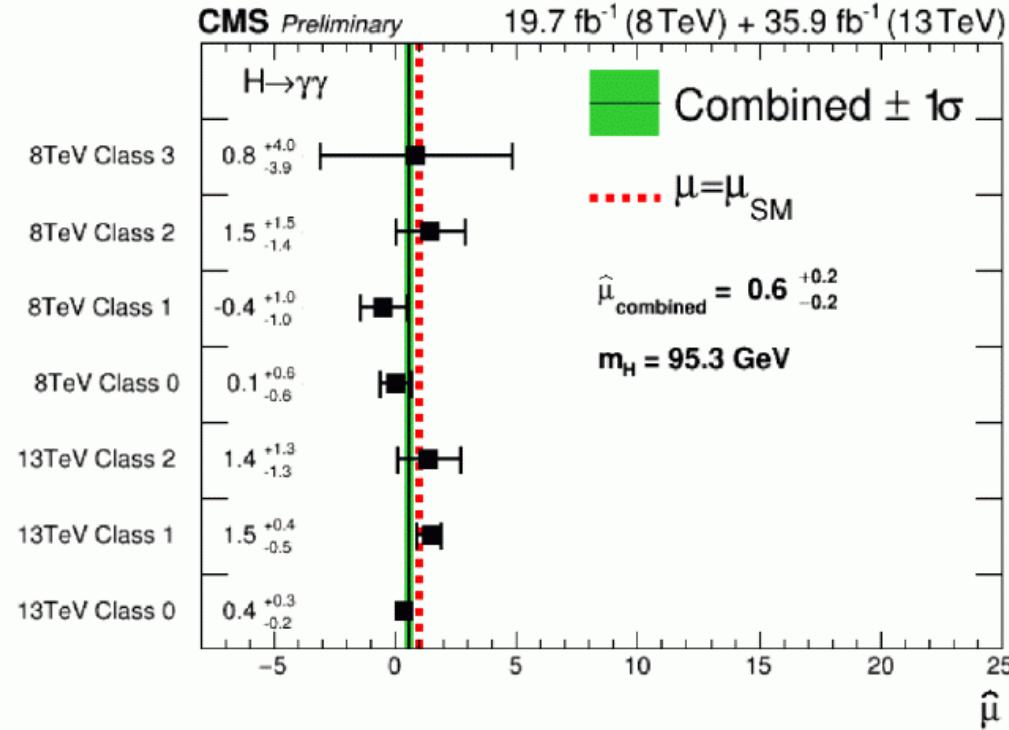
8TeV+13 TeV: Excess with  $\sim 2.8 \sigma$  local ( $1.3 \sigma$  global) significance at  $m=95.3$  GeV

More data are required to ascertain the origin of this excess

30



# $h \rightarrow \gamma\gamma$ (70-110 GeV) Runs 1+2

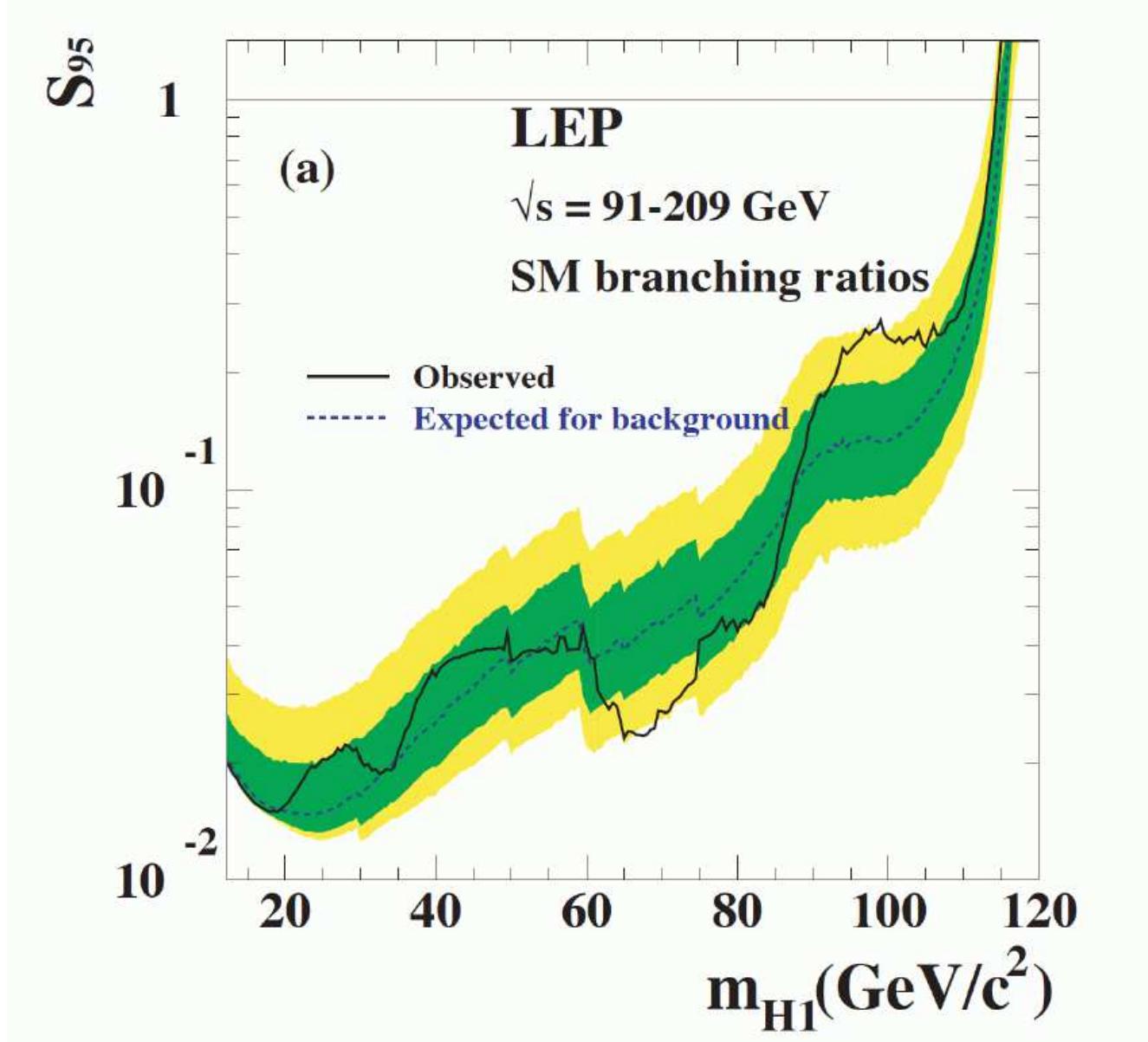


Excess here mostly driven by class 1 (&2) at 13 TeV

$\chi^2$  probability for the seven individual values to be compatible with a single signal hypothesis: 41%

- ‘Signal’ strengths for the 7 event classes and overall, in the 8 TeV+13TeV combination, fixing  $m_H=95.3$  GeV
- More data are required to ascertain the origin of this excess

## What was seen at LEP?

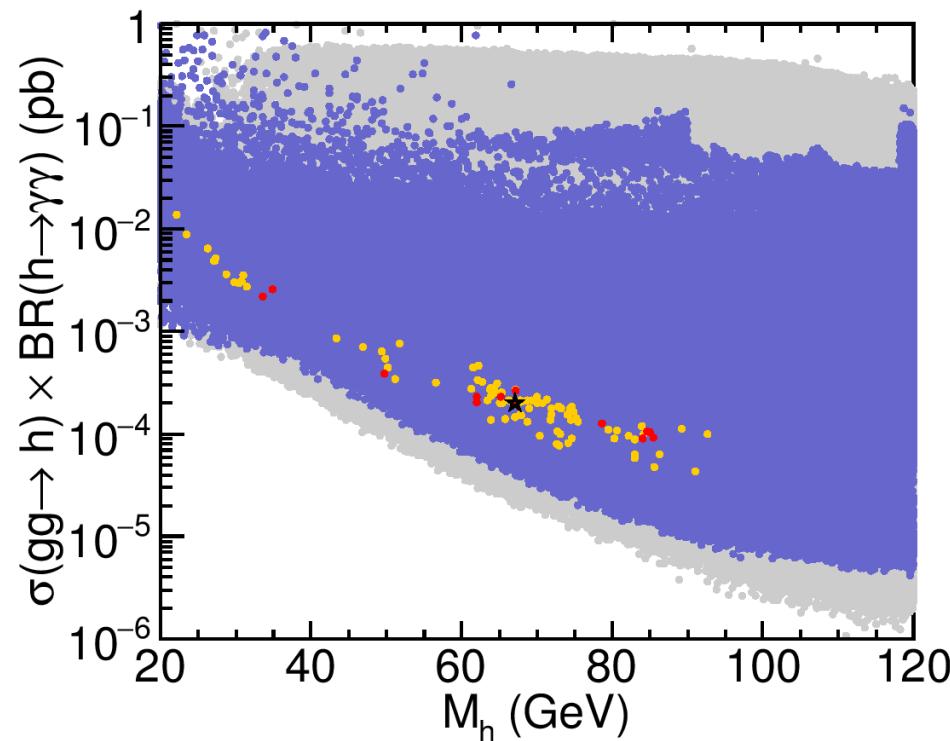


## Should we get excited?

⇒ according to CMS no!

⇒ let's wait for ATLAS, ETA summer '18

## Which model fits?



⇒ not the MSSM

⇒ 2HDM? NMSSM? . . . ? HAVE FUN!

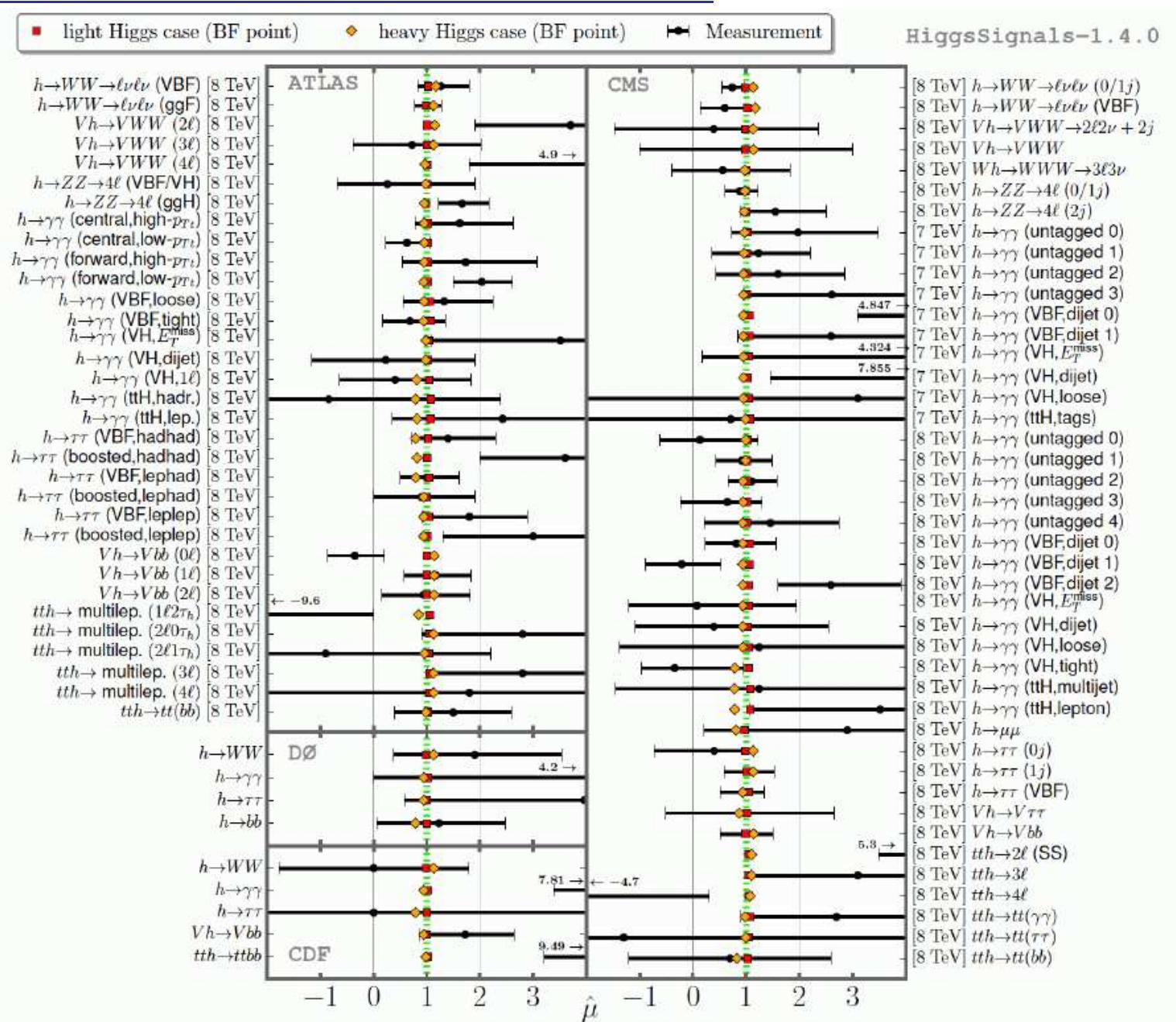
## 5. Conclusions

- SUSY is (still) the best motivated BSM theory
- Higgs rate measurements can be fulfilled by
  - the light  $\mathcal{CP}$ -even Higgs in the decoupling regime
  - the light  $\mathcal{CP}$ -even Higgs in the alignment w/o decoupling regime
  - the heavy  $\mathcal{CP}$ -even Higgs with  $M_h < 125$  GeV
- MSSM results in GUT based models:
  - CMSSM, NUHM1: relatively high Higgs mass scales favored
  - NUHM2: somewhat lower values possible, but still high . . .
- pMSSM8:
  - light  $\mathcal{CP}$ -even Higgs for “all”  $M_A$   
⇒ alignment without decoupling possible
  - heavy  $\mathcal{CP}$ -even Higgs at 125 GeV  
⇒ new benchmark scenarios
- $\phi_{95} \rightarrow \gamma\gamma$ : new CMS result possibly interesting!

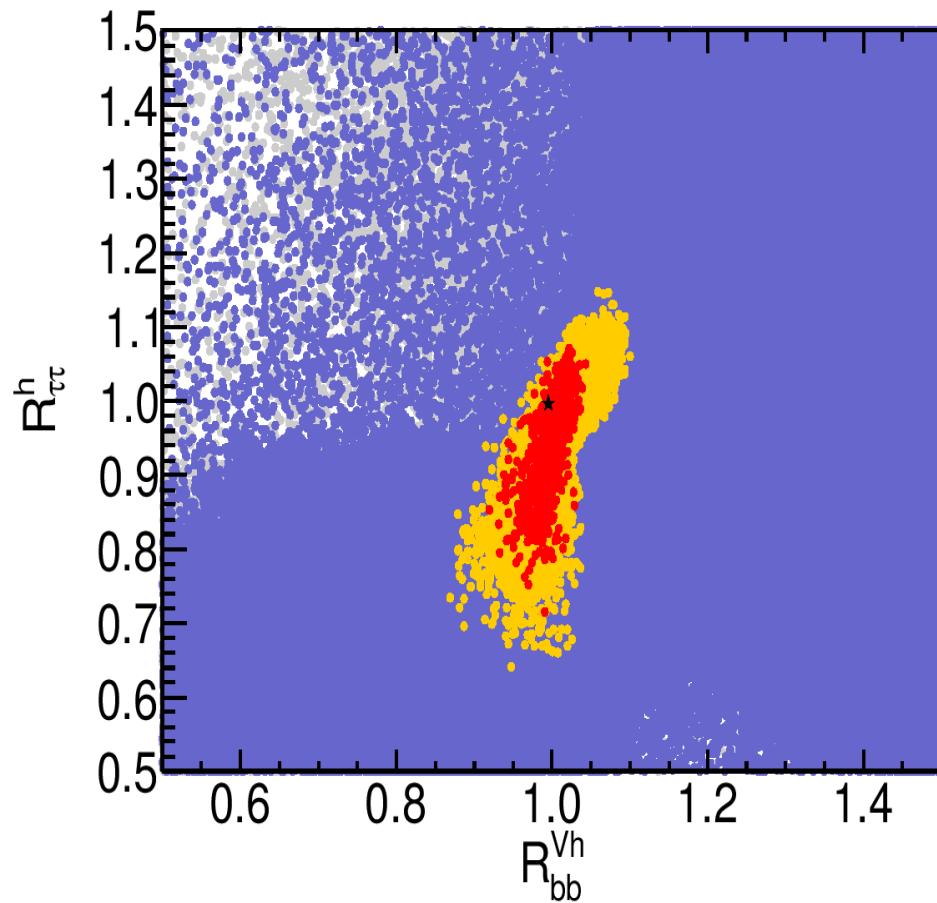
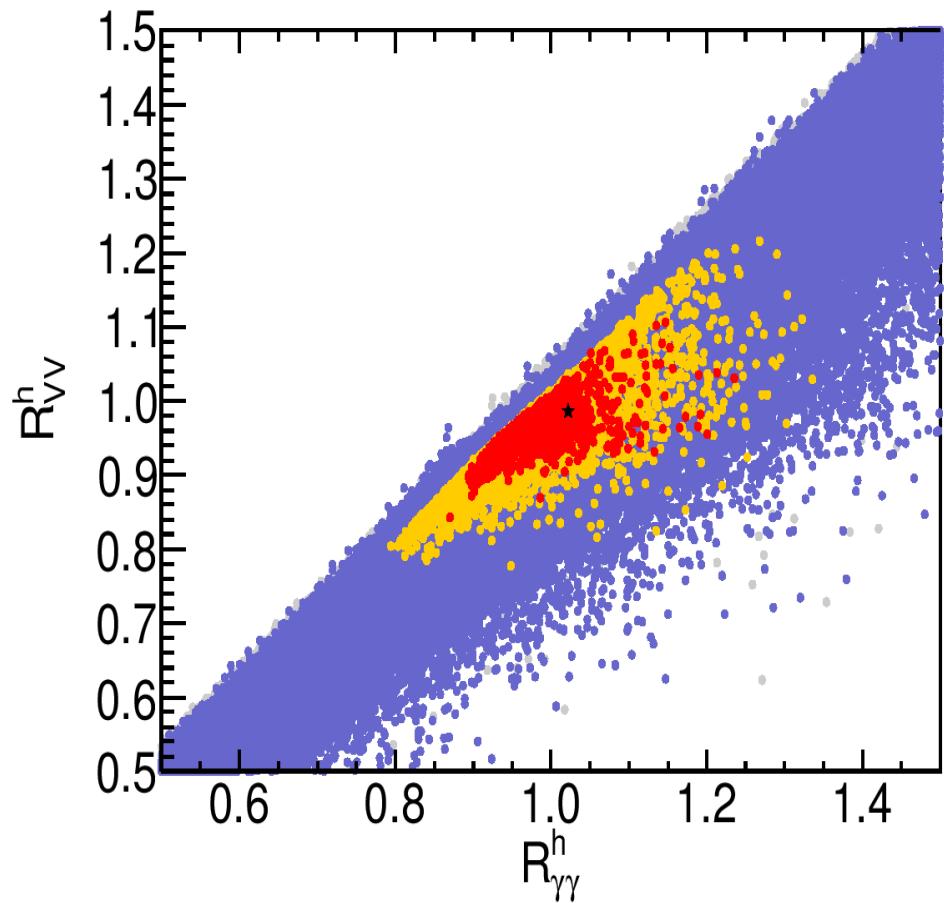


Further Questions?

## Best-fit point rates in the two Higgs cases:



## Light-Higgs case: preferred rates

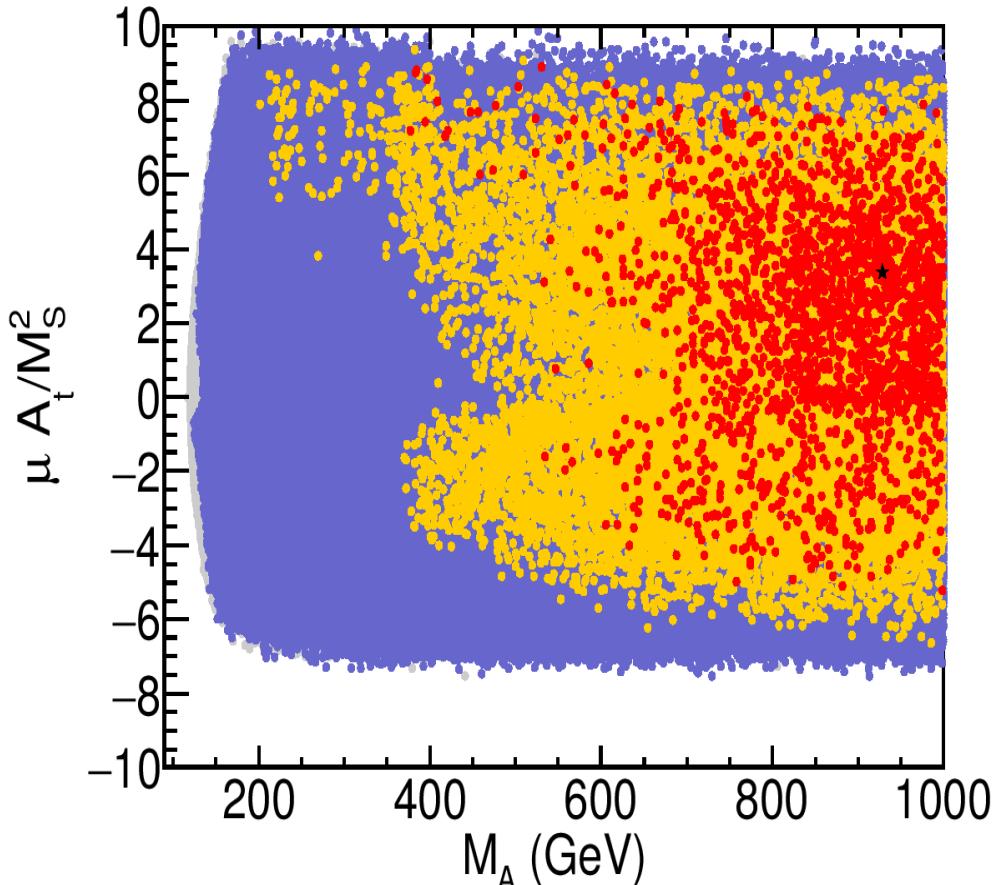
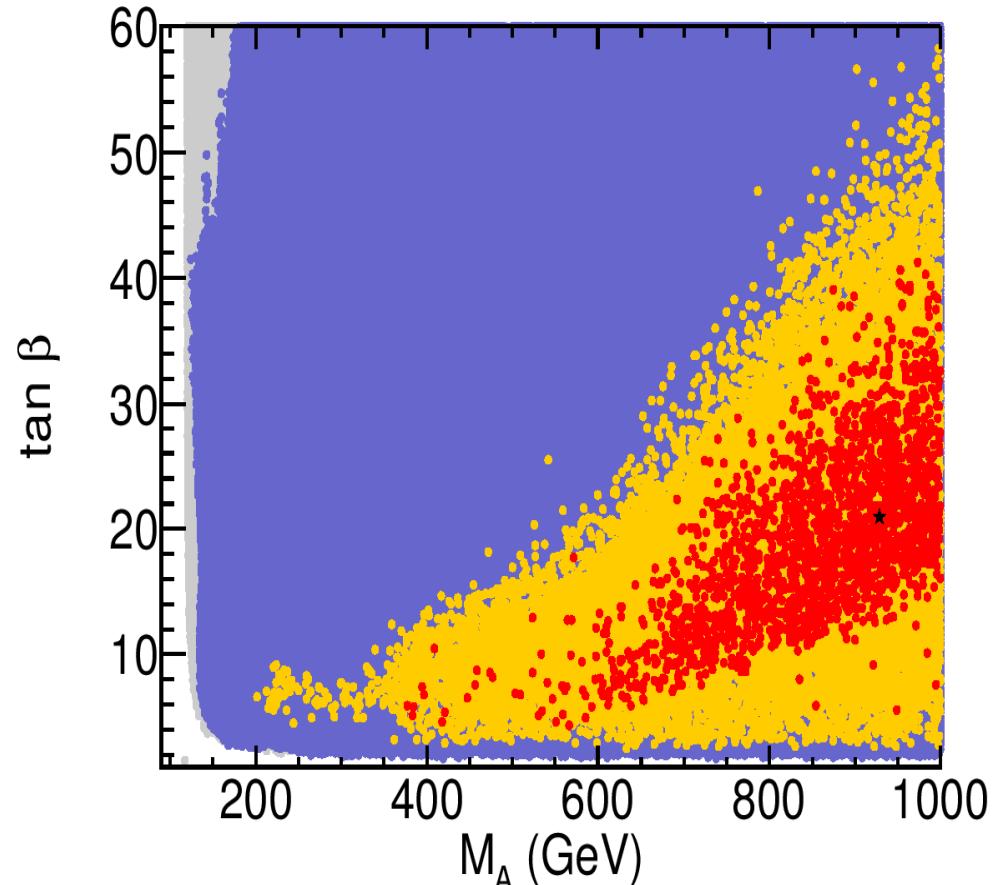


$$R_{VV}^h = 0.99^{+0.09}_{-0.08}, \quad R_{\gamma\gamma}^h = 1.02^{+0.16}_{-0.10},$$

$$R_{bb}^{Vh} = 1.00^{+0.02}_{-0.05}, \quad R_{tt}^h = 1.00^{+0.06}_{-0.20}$$

⇒ all very SM-like (no surprise . . .)  
 ⇒ but some (BSM) spread is allowed!

## Light-Higgs case: preferred parameters



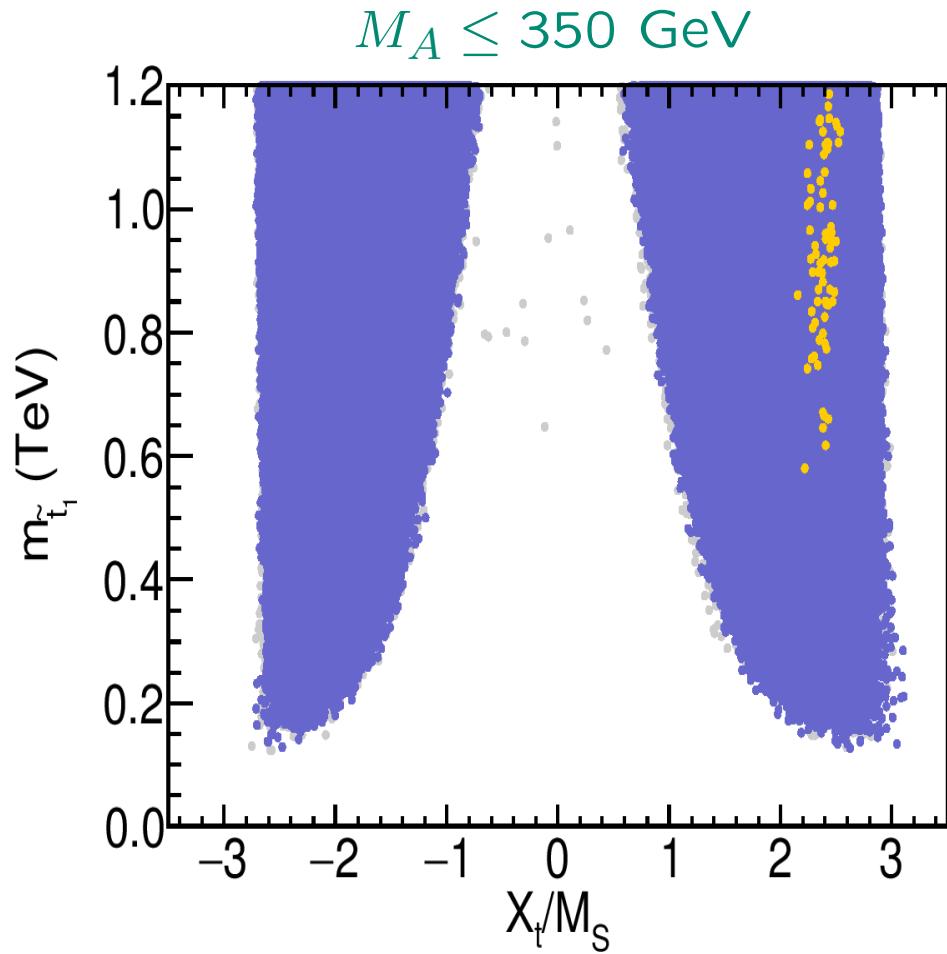
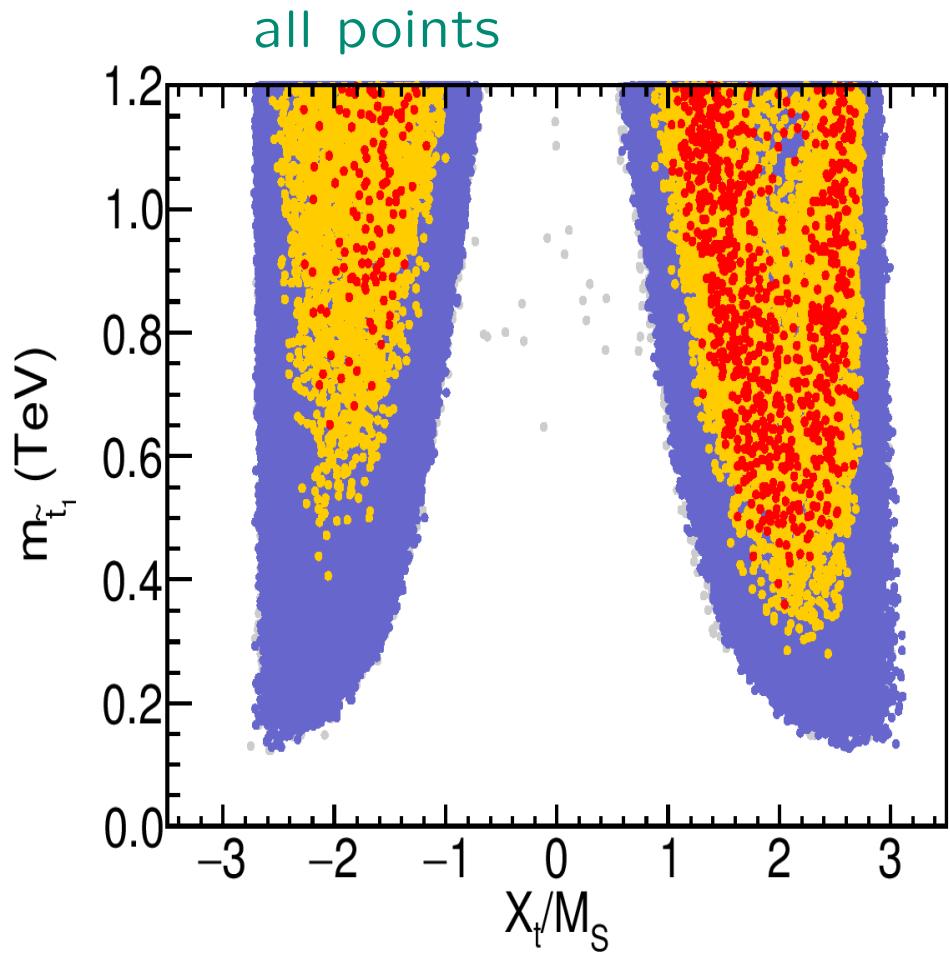
Favored points with  $M_A \gtrsim 500$  GeV  $\Rightarrow$  decoupling limit  
 $M_A \gtrsim 200$  GeV  $\Rightarrow$  alignment limit

$$\text{Alignment: } \tan \beta \sim 1 / \left[ \frac{\mu A_t}{M_S^2} \left( \frac{A_t^2}{6M_S^2} - 1 \right) \right]$$

$\Rightarrow$  small(er)  $\tan \beta$  needed to avoid  $\tau\tau$  limits  $\Rightarrow \mu A_t / M_S^2$  larger

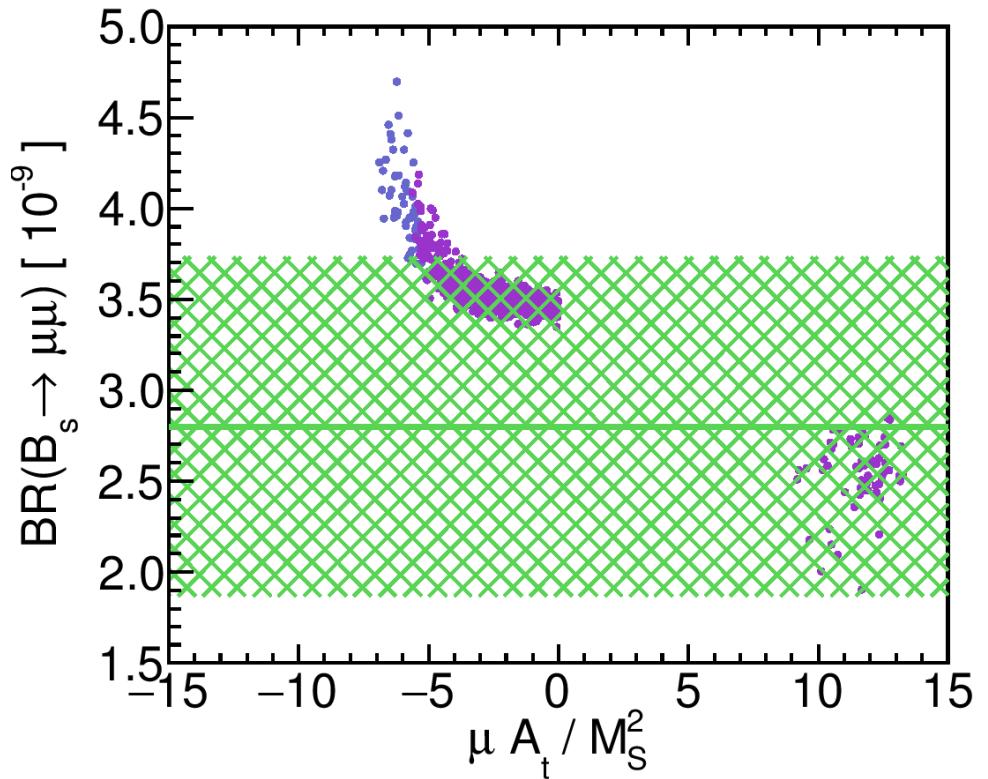
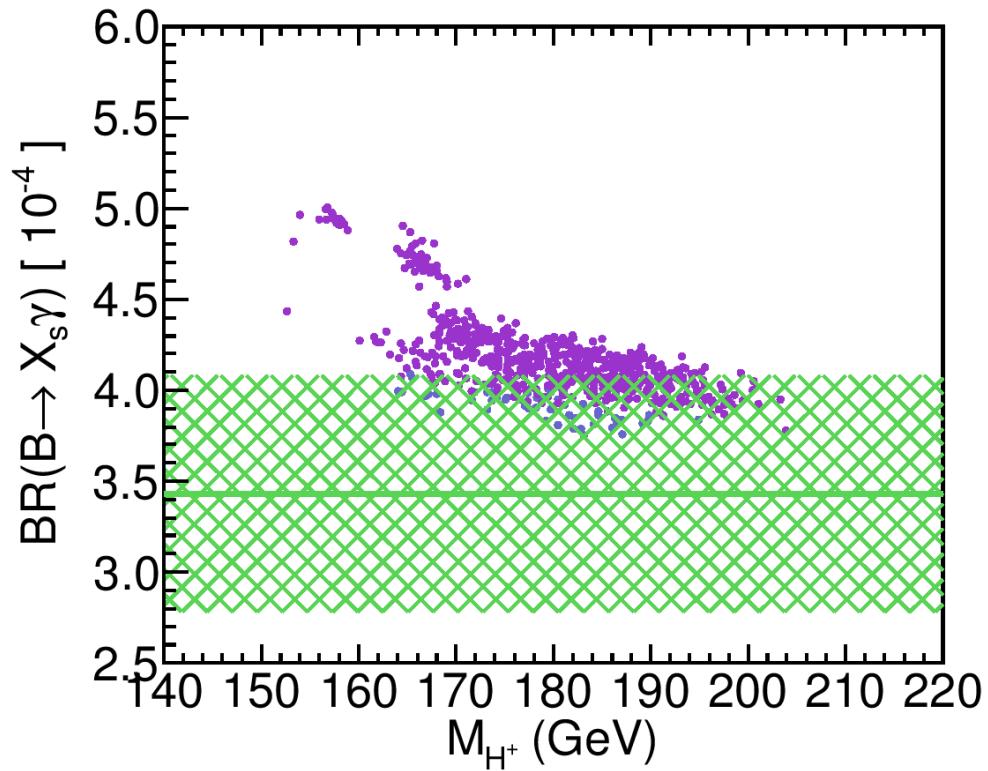
$\Rightarrow$  positive  $A_t$  preferred (for  $\mu > 0$ )

## Light-Higgs case: preferred parameters in the $\tilde{t}$ sector



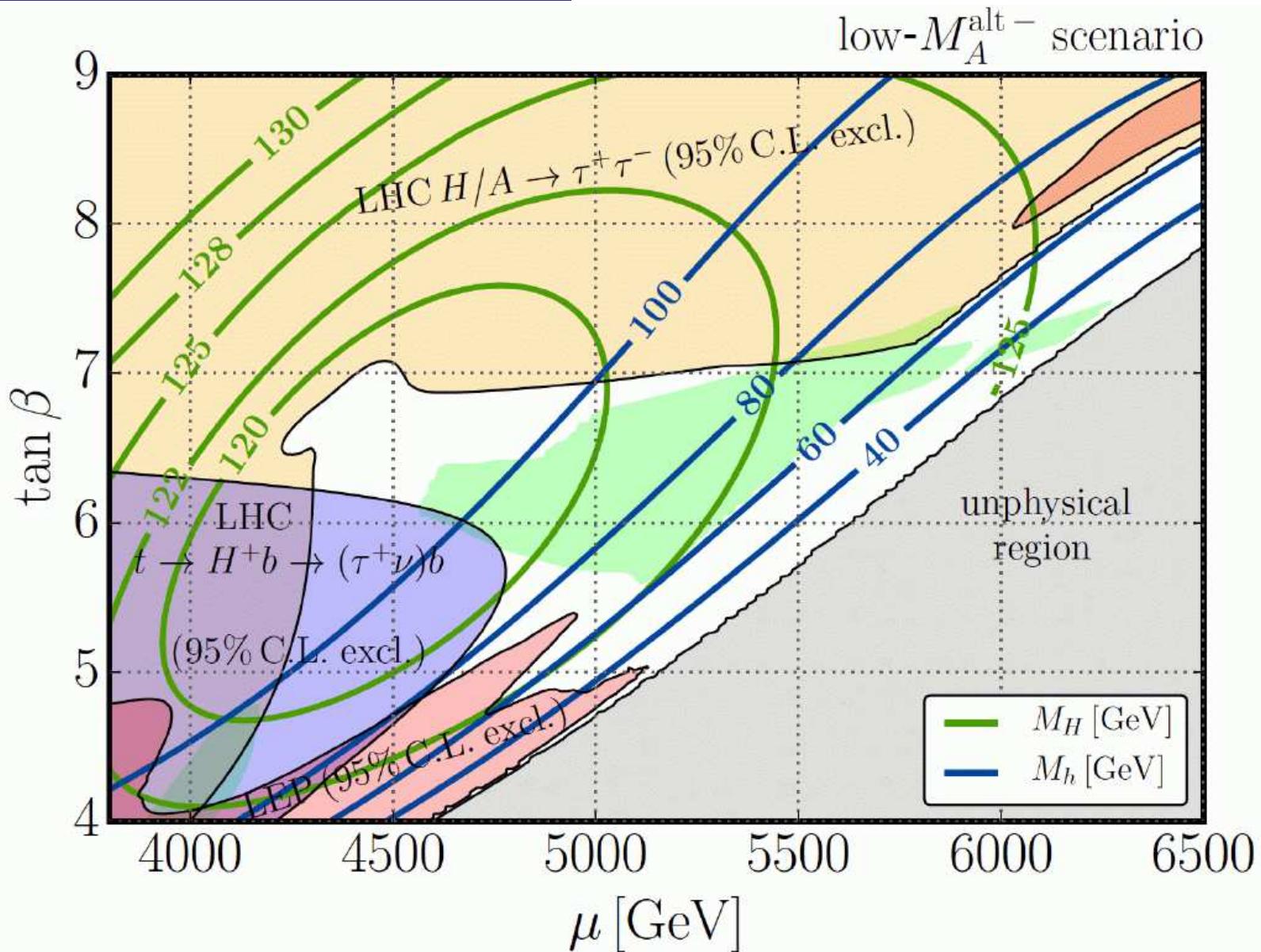
→ light stops down to  $m_{\tilde{t}_1} \sim 300 \text{ GeV}$  possible  
(even lighter stops possible with  $M_{\tilde{t}_L} \neq M_{\tilde{t}_R}$ )

## B-physics constraints?



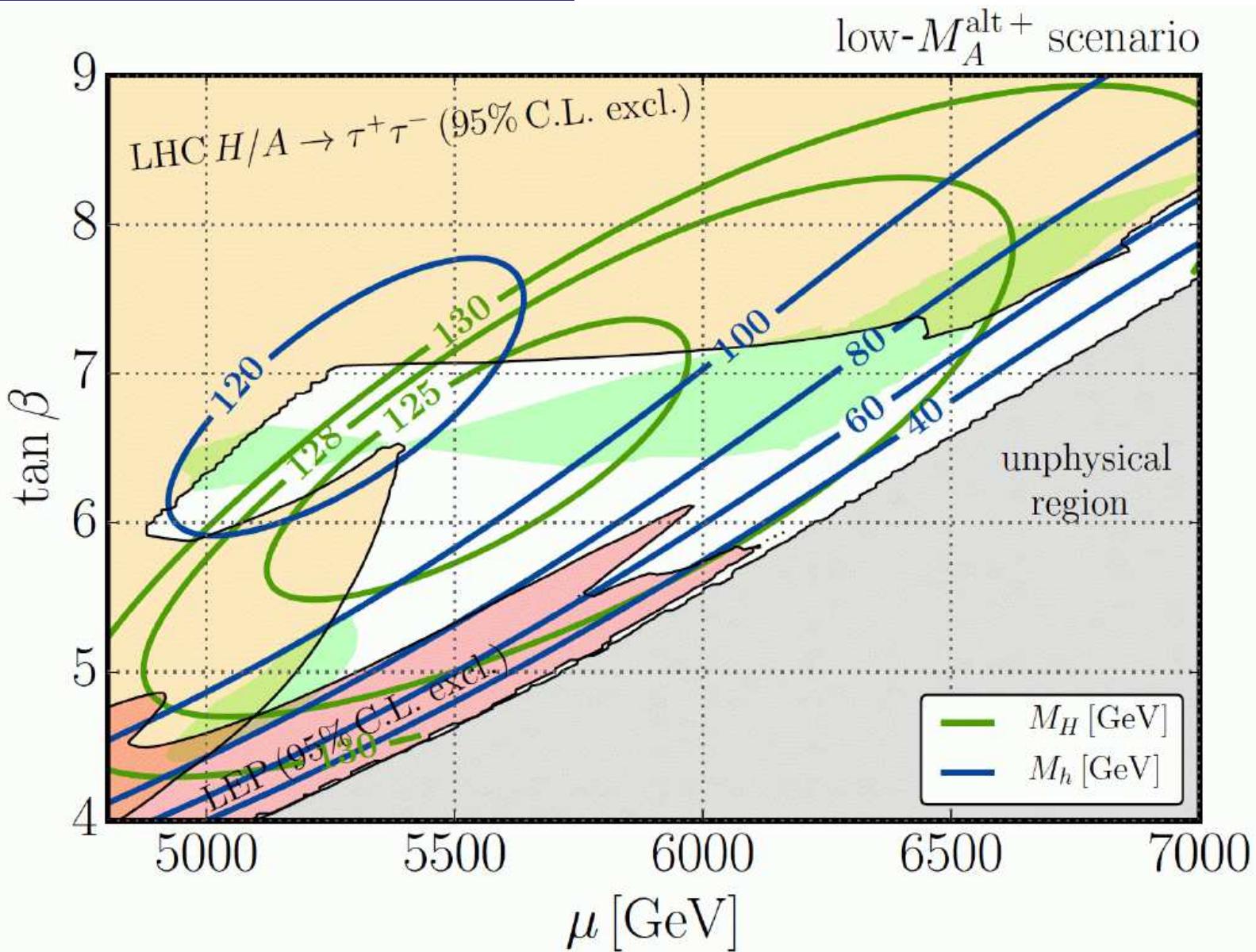
⇒ flavor constraints fulfilled!

low- $M_H^{\text{alt-}}$  (155 GeV =  $M_{H^\pm} < m_t$ ):



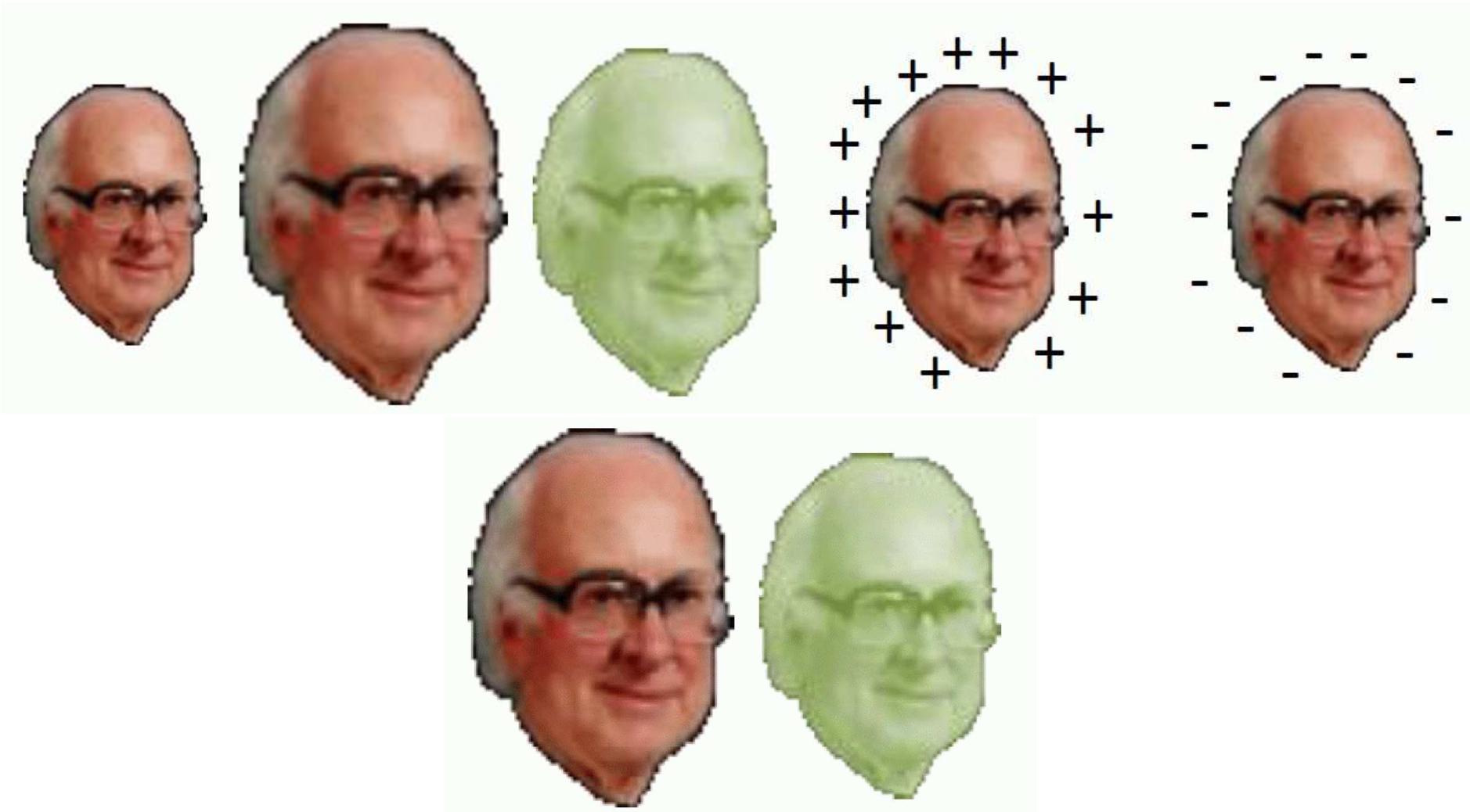
⇒ green area in agreement with all data!

low- $M_H^{\text{alt}+}$  (180 GeV =  $M_{H^\pm} > m_t$ ):



⇒ green area in agreement with all data!  $M_H \sim M_h \sim 125$  GeV possible!

## Results in the NMSSM



## Some NMSSM Higgs theory ( $Z_3$ invariant NMSSM)

MSSM Higgs sector: Two Higgs doublets

$$H_1 = \begin{pmatrix} H_1^1 \\ H_1^2 \end{pmatrix} = \begin{pmatrix} v_1 + (\phi_1 + i\chi_1)/\sqrt{2} \\ \phi_1^- \end{pmatrix}$$
$$H_2 = \begin{pmatrix} H_2^1 \\ H_2^2 \end{pmatrix} = \begin{pmatrix} \phi_2^+ \\ v_2 + (\phi_2 + i\chi_2)/\sqrt{2} \end{pmatrix}$$

$$\begin{aligned} V = & (\tilde{m}_1^2 + |\mu_1|^2) H_1 \bar{H}_1 + (\tilde{m}_2^2 + |\mu_2|^2) H_2 \bar{H}_2 - m_{12}^2 (\epsilon_{ab} H_1^a H_2^b + \text{h.c.}) \\ & + \frac{g'^2 + g^2}{8} (H_1 \bar{H}_1 - H_2 \bar{H}_2)^2 + \frac{g^2}{2} |H_1 \bar{H}_2|^2 \end{aligned}$$

## Some NMSSM Higgs theory ( $Z_3$ invariant NMSSM)

NMSSM Higgs sector: Two Higgs doublets + one Higgs singlet

$$H_1 = \begin{pmatrix} H_1^1 \\ H_1^2 \end{pmatrix} = \begin{pmatrix} v_1 + (\phi_1 + i\chi_1)/\sqrt{2} \\ \phi_1^- \end{pmatrix}$$

$$H_2 = \begin{pmatrix} H_2^1 \\ H_2^2 \end{pmatrix} = \begin{pmatrix} \phi_2^+ \\ v_2 + (\phi_2 + i\chi_2)/\sqrt{2} \end{pmatrix}$$

$$S = v_s + S_R + IS_I$$

$$V = (\tilde{m}_1^2 + |\mu\lambda S|^2)H_1\bar{H}_1 + (\tilde{m}_2^2 + |\mu\lambda S|^2)H_2\bar{H}_2 - m_{12}^2(\epsilon_{ab}H_1^aH_2^b + \text{h.c.})$$

$$+ \frac{g'^2 + g^2}{8}(H_1\bar{H}_1 - H_2\bar{H}_2)^2 + \frac{g^2}{2}|H_1\bar{H}_2|^2$$

$$+ |\lambda(\epsilon_{ab}H_1^aH_2^b) + \kappa S^2|^2 + m_S^2|S|^2 + (\lambda A_\lambda(\epsilon_{ab}H_1^aH_2^b)S + \frac{\kappa}{3}A_\kappa S^3 + \text{h.c.})$$

Free parameters:

$$\lambda, \kappa, A_\kappa, M_{H^\pm}, \tan\beta, \mu_{\text{eff}} = \lambda v_s$$

## Higgs spectrum:

$\mathcal{CP}$ -even :  $h_1, h_2, h_3$

$\mathcal{CP}$ -odd :  $a_1, a_2$

charged :  $H^+, H^-$

Goldstones :  $G^0, G^+, G^-$

## Neutralinos:

$$\mu \rightarrow \mu_{\text{eff}}$$

compared to the MSSM: one singlino more

$$\rightarrow \tilde{\chi}_1^0, \tilde{\chi}_2^0, \tilde{\chi}_3^0, \tilde{\chi}_4^0, \tilde{\chi}_5^0$$

## Mass of the lightest $\mathcal{CP}$ -even Higgs: (no singlet mixing)

$$m_{h,\text{tree},\text{NMSSM}}^2 = m_{h,\text{tree},\text{MSSM}}^2 + M_Z^2 \frac{\lambda^2}{g^2} \sin^2 2\beta$$

## Mass of the $\mathcal{CP}$ -odd Higgs:

$$\text{MSSM} : M_A^2 = -m_{12}^2(\tan \beta + \cot \beta) = \mu B(\tan \beta + \cot \beta)$$

$$\text{NMSSM} : "M_A^2" = \mu_{\text{eff}} B_{\text{eff}} (\tan \beta + \cot \beta)$$

with  $B_{\text{eff}} = A_\lambda + \kappa v_s$ ,  $\mu_{\text{eff}} = \lambda v_s$   $\Rightarrow$  one very light  $a_1$

## Mass of the charged Higgs:

$$\text{MSSM} : M_{H^\pm}^2 = M_A^2 + M_W^2 = M_A^2 + \frac{1}{2} v^2 g^2$$

$$\text{NMSSM} : M_{H^\pm}^2 = M_A^2 + v^2 \left( \frac{g^2}{2} - \lambda^2 \right)$$

Mass of the lightest  $\mathcal{CP}$ -even Higgs: (no singlet mixing)

$$m_{h,\text{tree,NMSSM}}^2 = m_{h,\text{tree,MSSM}}^2 + M_Z^2 \frac{\lambda^2}{g^2} \sin^2 2\beta$$

Mass of the  $\mathcal{CP}$ -odd Higgs:

$$\text{MSSM} : M_A^2 = -m_{12}^2(\tan \beta + \cot \beta) = \mu B(\tan \beta + \cot \beta)$$

$$\text{NMSSM} : "M_A^2" = \mu_{\text{eff}} B_{\text{eff}} (\tan \beta + \cot \beta)$$

with  $B_{\text{eff}} = A_\lambda + \kappa v_s$ ,  $\mu_{\text{eff}} = \lambda v_s$   $\Rightarrow$  one very light  $a_1$

Mass of the charged Higgs:

$$\text{MSSM} : M_{H^\pm}^2 = M_A^2 + M_W^2 = M_A^2 + \frac{1}{2}v^2 g^2$$

$$\text{NMSSM} : M_{H^\pm}^2 = M_A^2 + v^2 \left( \frac{g^2}{2} - \lambda^2 \right)$$

$$\Rightarrow M_{h_1}^{\text{MSSM,tree}} \leq M_{h_1}^{\text{NMSSM,tree}}, \text{ one light } a_1, M_{H^\pm}^{\text{MSSM,tree}} \geq M_{H^\pm}^{\text{NMSSM,tree}}$$

## Interesting case: light singlet

Singlet does not couple to SM particles!

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“Non-interacting particles are hard to detect.”



[F. Klinkhamer]

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Singlet does not couple to SM particles!



[F. Klinkhamer]

“Non-interacting particles are hard to detect.”

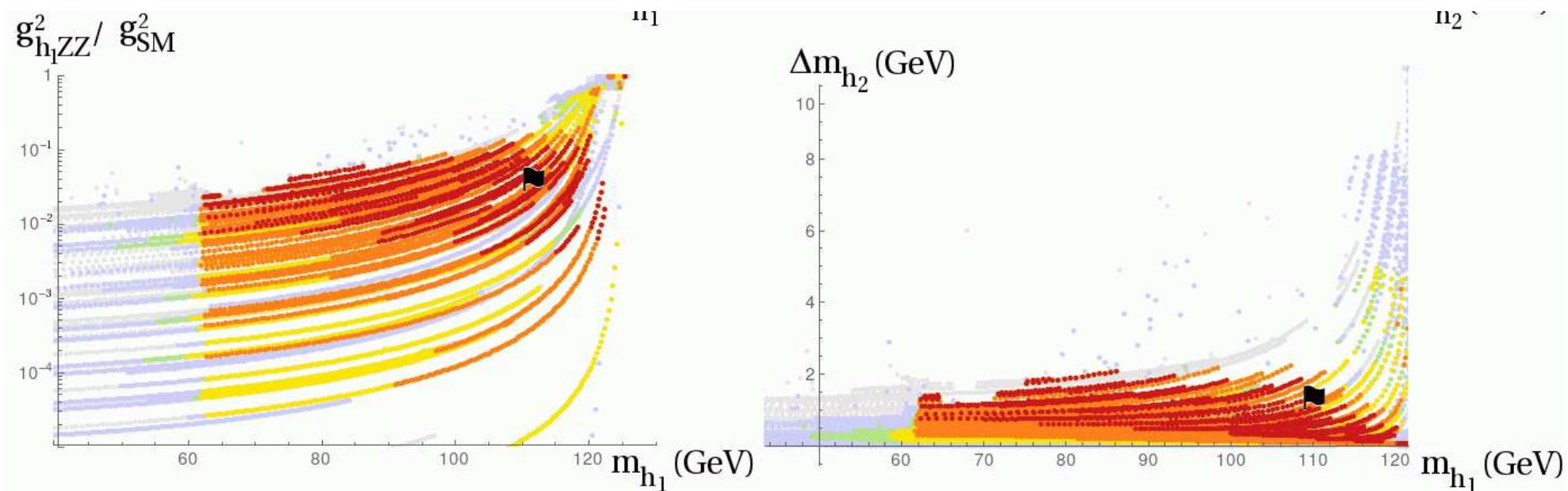
“Easily” possible in the NMSSM:

Light, singlet-like Higgs below 125 GeV

Can the LHC find them?

Parameters:

$\tan \beta = 8$ ,  $M_A = 1 \text{ TeV}$ ,  $A_\kappa = -2...0 \text{ TeV}$ ,  $\mu = 120...2000 \text{ GeV}$ ,  
 $2M_1 = M_2 = 500 \text{ GeV}$ ,  $M_3 = 1.5 \text{ TeV}$ ,  $m_{\tilde{Q}_3} = 1 \text{ TeV}$ ,  $m_{\tilde{Q}_{1,2}} = 1.5 \text{ TeV}$ ,  
 $A_t = -2 \text{ TeV}$ ,  $A_{b,\tau} = -1.5 \text{ TeV}$

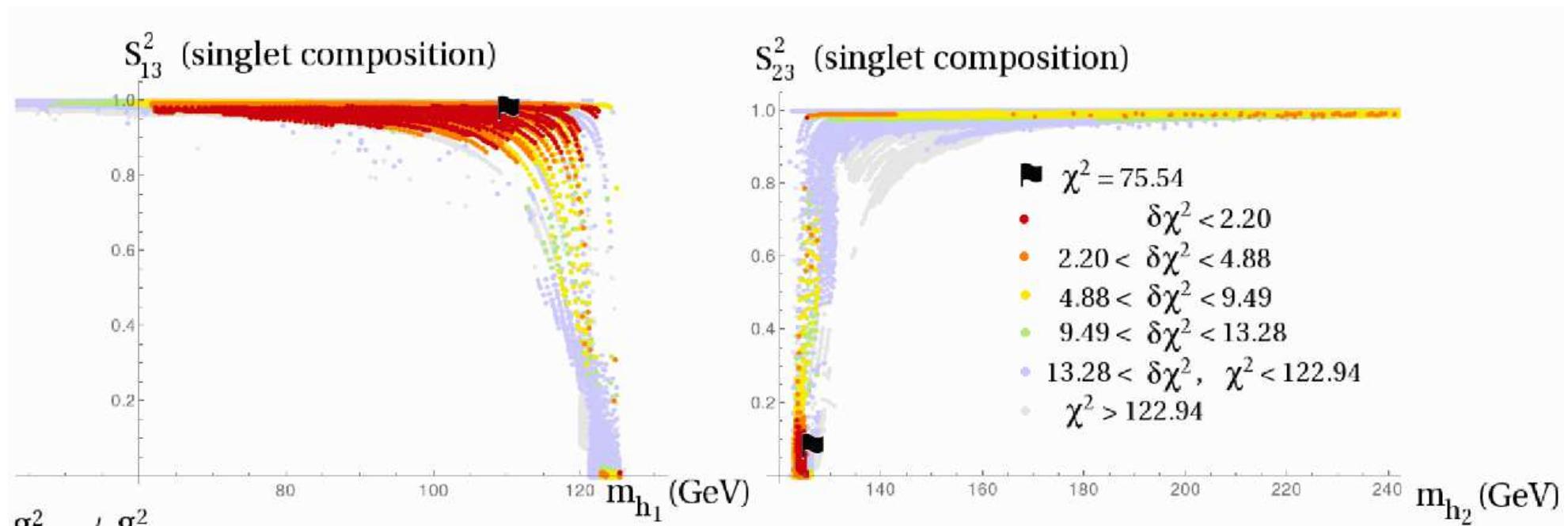


⇒ light Higgs below 125 GeV

⇒ strongly reduced couplings to gauge bosons!

## Parameters:

$\tan \beta = 8$ ,  $M_A = 1 \text{ TeV}$ ,  $A_\kappa = -2...0 \text{ TeV}$ ,  $\mu = 120...2000 \text{ GeV}$ ,  
 $2M_1 = M_2 = 500 \text{ GeV}$ ,  $M_3 = 1.5 \text{ TeV}$ ,  $m_{\tilde{Q}_3} = 1 \text{ TeV}$ ,  $m_{\tilde{Q}_{1,2}} = 1.5 \text{ TeV}$ ,  
 $A_t = -2 \text{ TeV}$ ,  $A_{b,\tau} = -1.5 \text{ TeV}$



→ light Higgs below 125 GeV has large singlet component  
 → second Higgs is SM-like