

# Aspects of CP violation in electroweak baryogenesis

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## Outline :

1. Introduction
2. Relationship between the BAU and EDMs
3. Numerical results
4. Summary

# Introduction

# Baryon Asymmetry of the Universe (BAU)

Observational facts indicates that our Universe is baryon-asymmetric.

$$Y_B \equiv \frac{n_B}{s} = (8.59 \pm 0.11) \times 10^{-11}$$

P. A. R. Ade et al. [Planck Collaboration], arXiv:1303.5076

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One attractive scenario is electroweak baryogenesis (EWBG).

Kuzmin, Rubakov, Shaposhnikov, PLB155,36 ('85)



Relevant energy scale is  $\sim O(100)$  GeV.

$\sim$  Collider experiments can verify it.

# Electroweak Baryogenesis

Kuzmin, Rubakov, Shaposhnikov, PLB155,36 ('85)

Sakharov's criteria are satisfied as follows:

- Baryon number violation : Sphaleron process
- C and CP violation : Chiral gauge theory and CP phase
- Out of equilibrium : First order electroweak phase transition

# Electroweak Baryogenesis

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Sakharov's criteria are satisfied as follows:

- Baryon number violation : Sphaleron process
- C and CP violation : Chiral gauge theory and CP phase
  - (SM) CP violation is too small to generate the BAU.
- Out of equilibrium : First order electroweak phase transition
  - (SM) EWPT becomes crossover for  $m_h > 73$  GeV.

Gavela et al, NPB430, 382 (1994) ; Huet and Nelsen, PRD51, 379 (1995)

Kajantie et al, PRL77, 2887 (1996) ; Rummukainen et al, NPB 542, 283 (1998)  
Csikor et al, PRL 82, 21 (1999) ; Aoki et al, PRD 60, 013001 (1999)  
Laine et al, NPB 73, 180 (1999)

# New physics

For the successful EWBG, physics beyond the SM is needed.

New Physics for EWBG

First order EWPT

CP phase

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## New Physics for EWBG

First order EWPT

Extended Higgs sector

e.g. Real singlet,  
2 Higgs doublets

Collider

CP phase

\* Given the current situation in the LHC, the colored particles may not play a role in achieving EWBG.

# New physics

For the successful EWBG, physics beyond the SM is needed.

## New Physics for EWBG

First order EWPT

Extended Higgs sector

Real singlet

CP phase

EW-interacting fermions

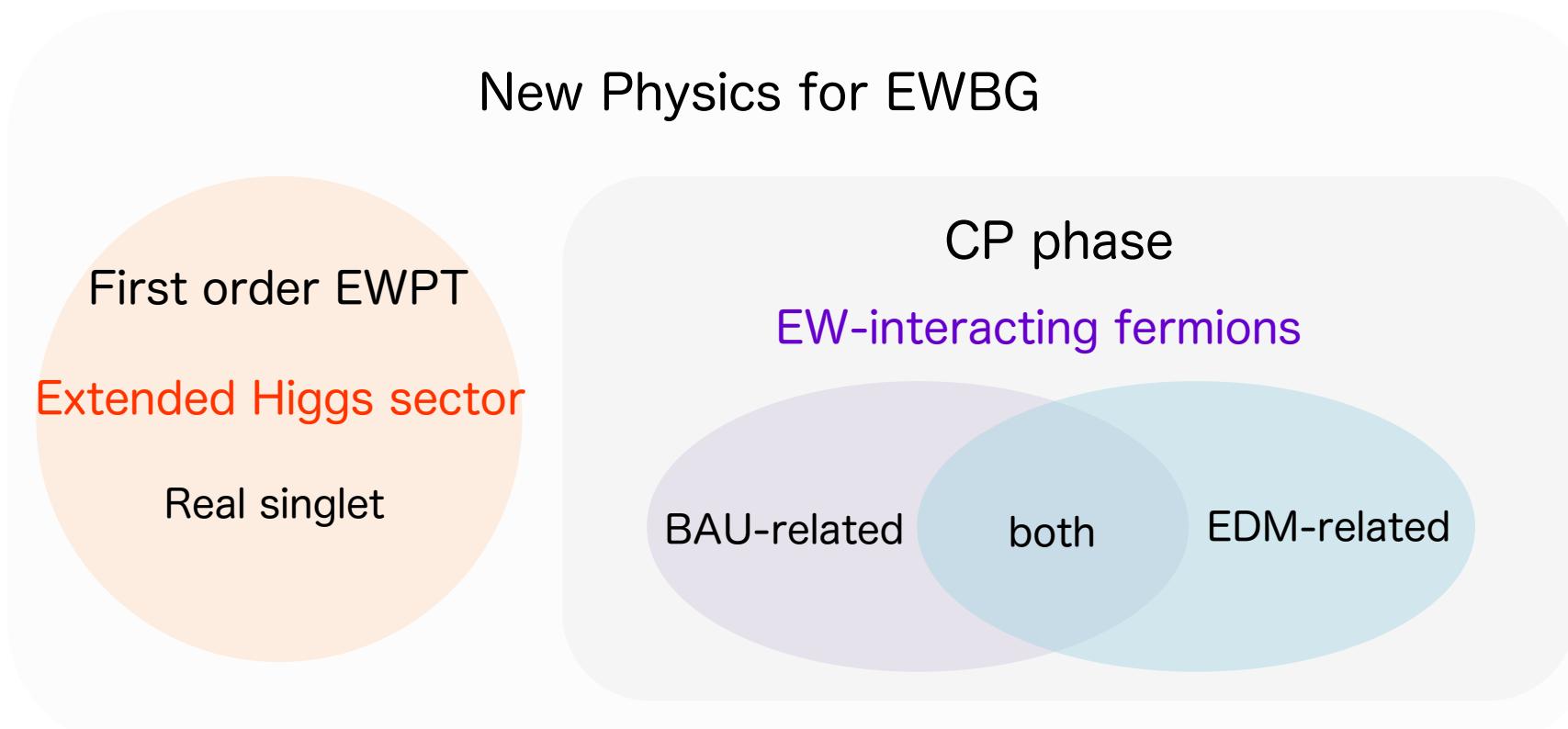
$$\begin{pmatrix} \psi^+ \\ \psi_i \end{pmatrix} \quad \psi_j$$

EDMs

\* It becomes possible to discuss both the first order EWPT and the CP violation, separately.

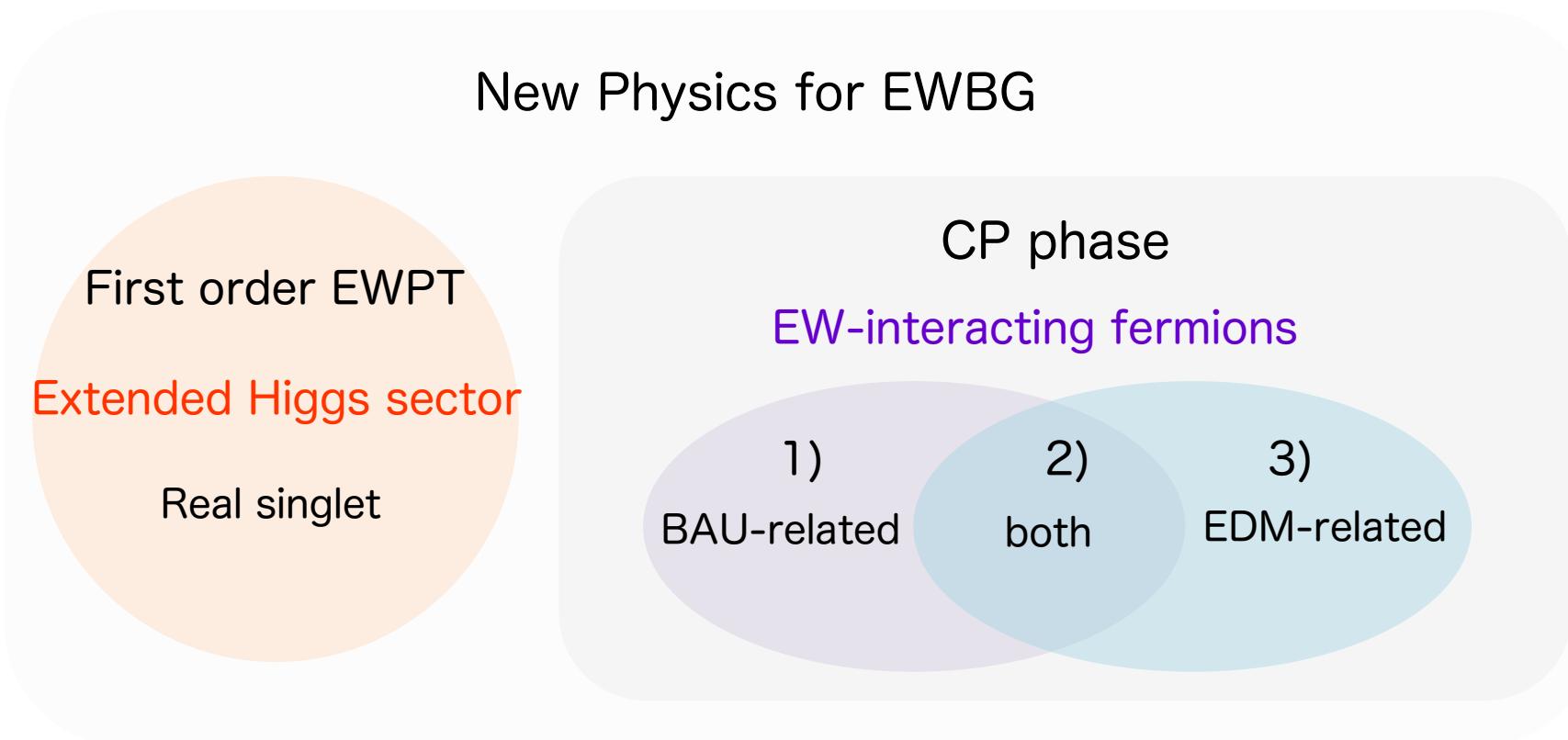
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- 1) What kind of CP phase is the BAU-related one ?
- 2) Clarify the relationship between the BAU and EDMs
- 3) Consider the situation where the other CP phase exists

# Relationship between the BAU and EDMs

# BAU

Two fermions  $\psi_{i,j}$  have the following interactions.

$$\mathcal{L} = \frac{1}{\sqrt{2}} \bar{\psi}_i [c_L v_b(x) P_L + c_R v_a(y) P_R] \psi_j + \text{h.c.}$$

$c_{L,R}$  : Complex numbers

$v_{a,b}(x)$  : Space-dependent VEVs

# BAU

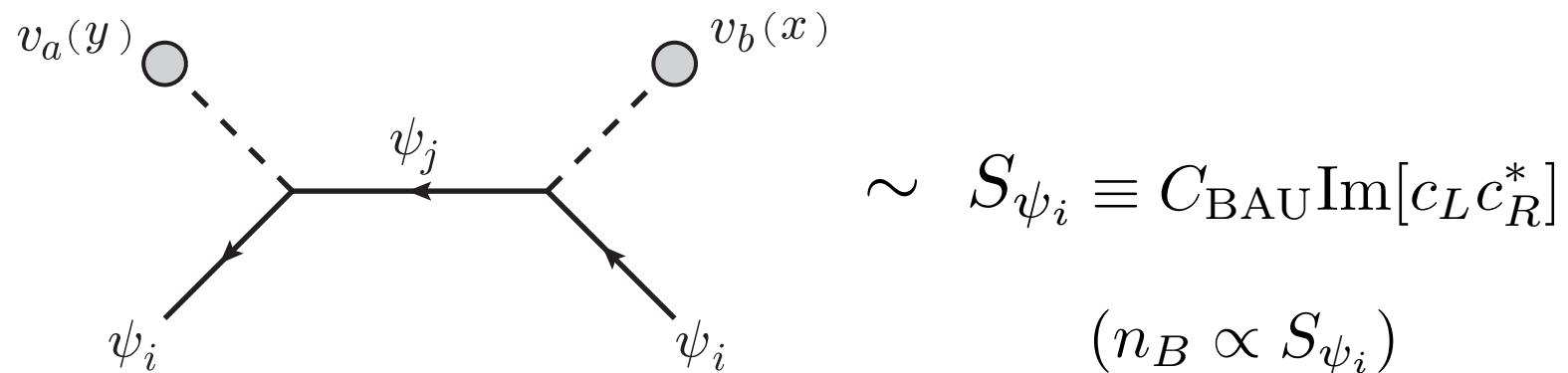
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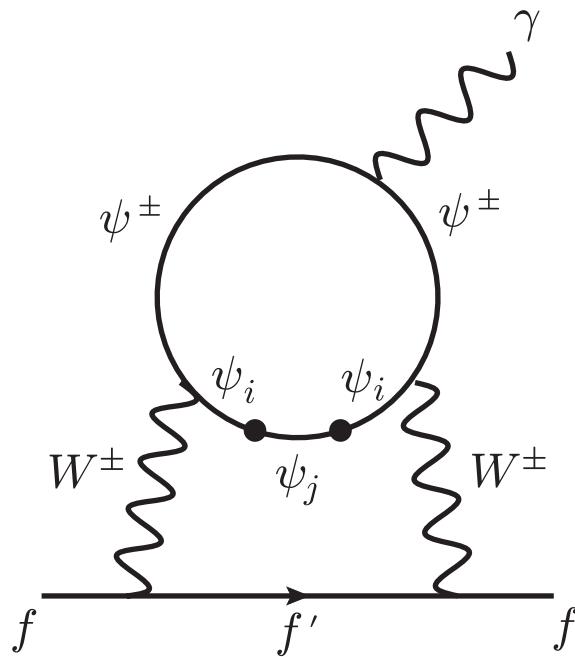
$v_{a,b}(x)$  : Space-dependent VEVs

CP-violating source is supplied to the BAU by the following diagram.



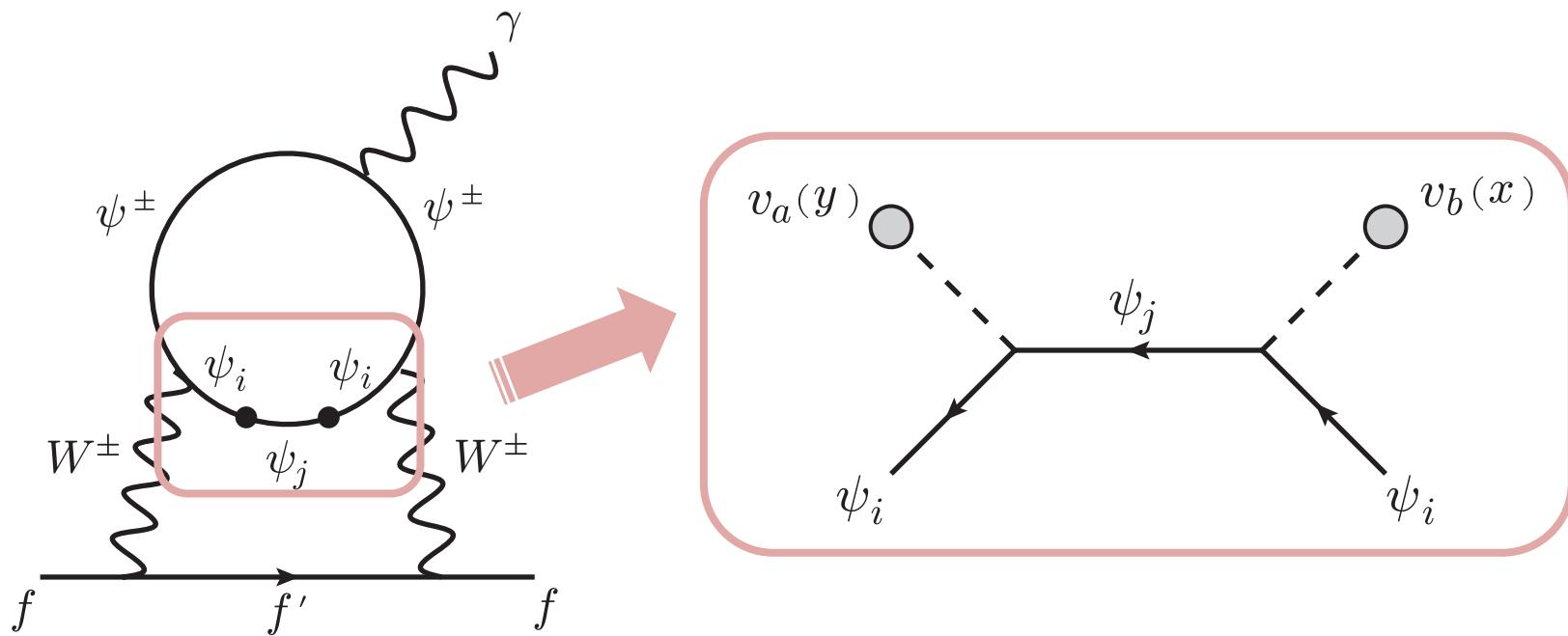
# EDM via Barr-Zee diagram

New fermions can induce EDMs via the Barr-Zee diagram.



# EDM via Barr-Zee diagram

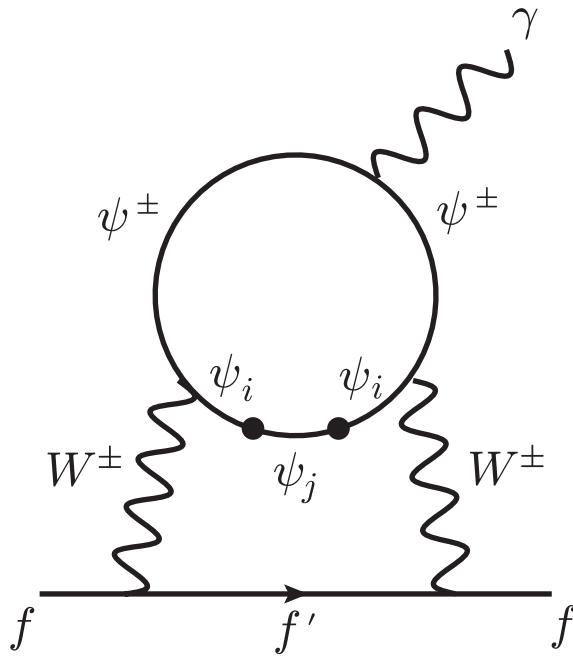
New fermions can induce EDMs via the Barr-Zee diagram.



EDMs are directly connected to the CP-violating source.

# EDM via Barr-Zee diagram

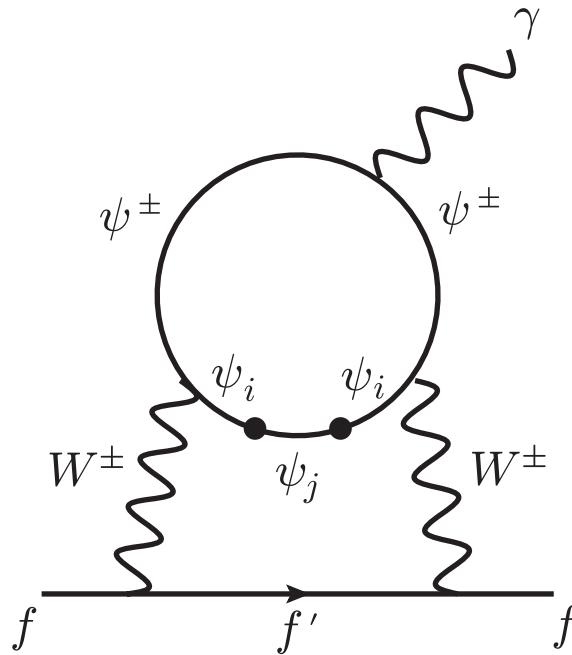
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$$\frac{d_f^{WW}}{e} \equiv C_{\text{EDM}}^{WW} \text{Im}[c_L c_R^*]$$

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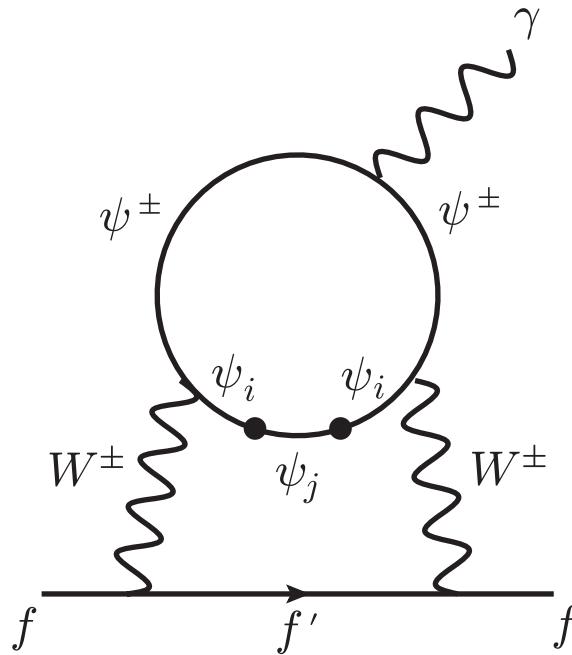
$$\frac{d_f^{WW}}{e} \equiv C_{\text{EDM}}^{WW} \text{Im}[c_L c_R^*]$$

We finally obtain

$$S_{\psi_i} = \frac{C_{\text{BAU}}}{C_{\text{EDM}}^{WW}} \left( \frac{d_f^{WW}}{e} \right)$$

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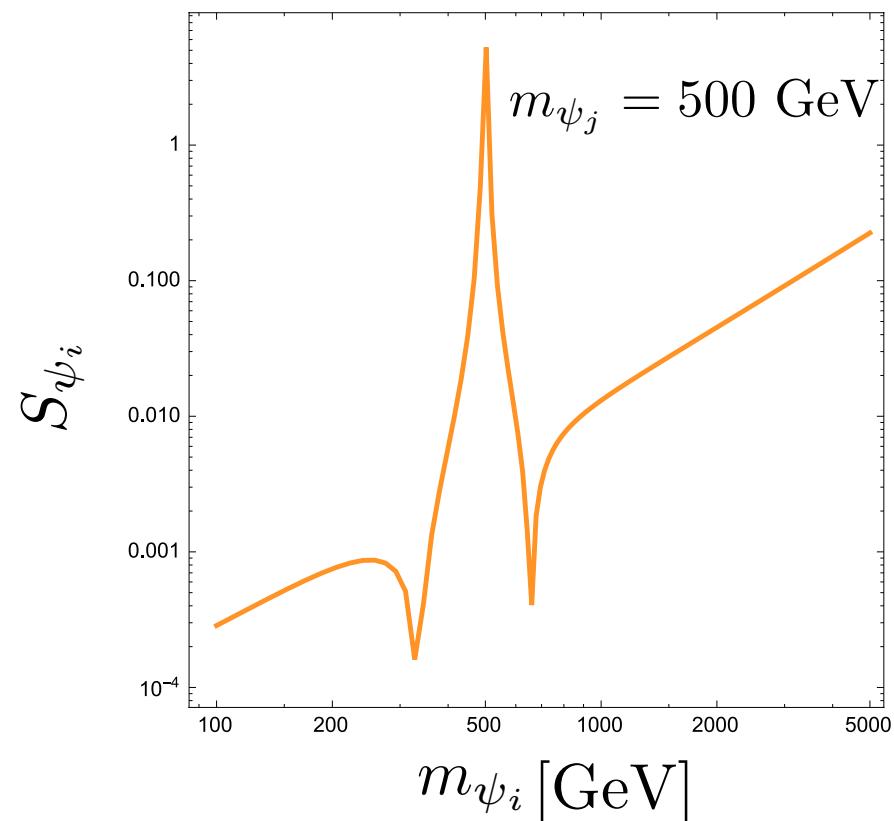
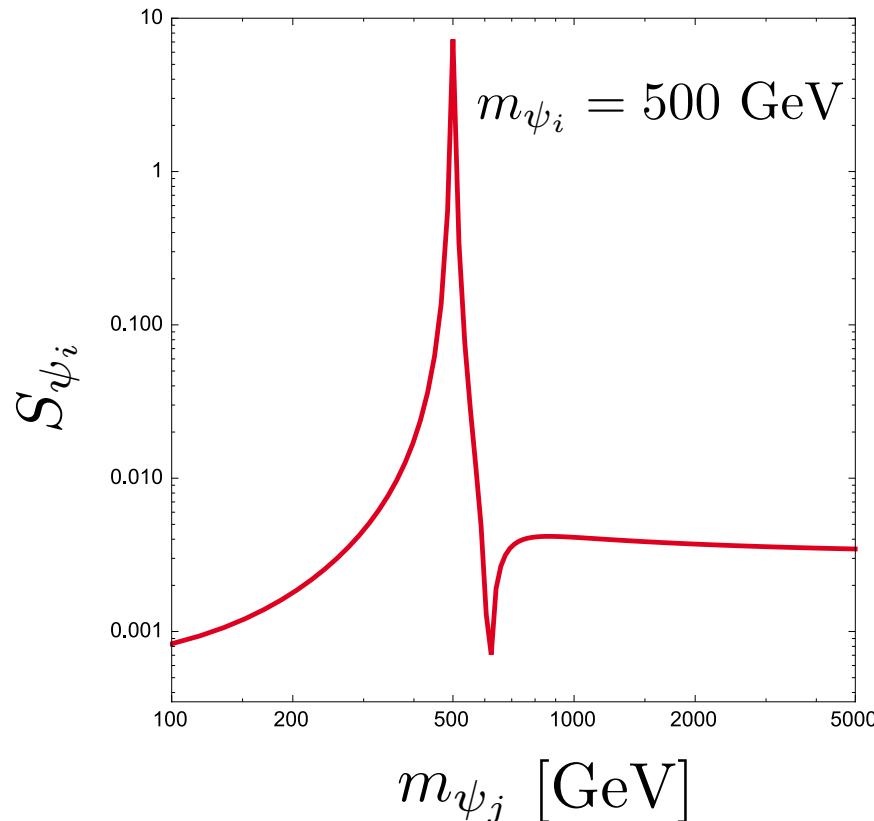
$$S_{\psi_i} = \frac{C_{\text{BAU}}}{C_{\text{EDM}}^{WW}} \left( \frac{d_e^{WW}}{e} \right)_{\text{exp}}$$

# Numerical results

# CP-violating source as a function of $m_{\psi_{i,j}}$

For numerical calculations,  $|d_e^{\text{exp}}| = 8.7 \times 10^{-29} e \cdot \text{cm}$

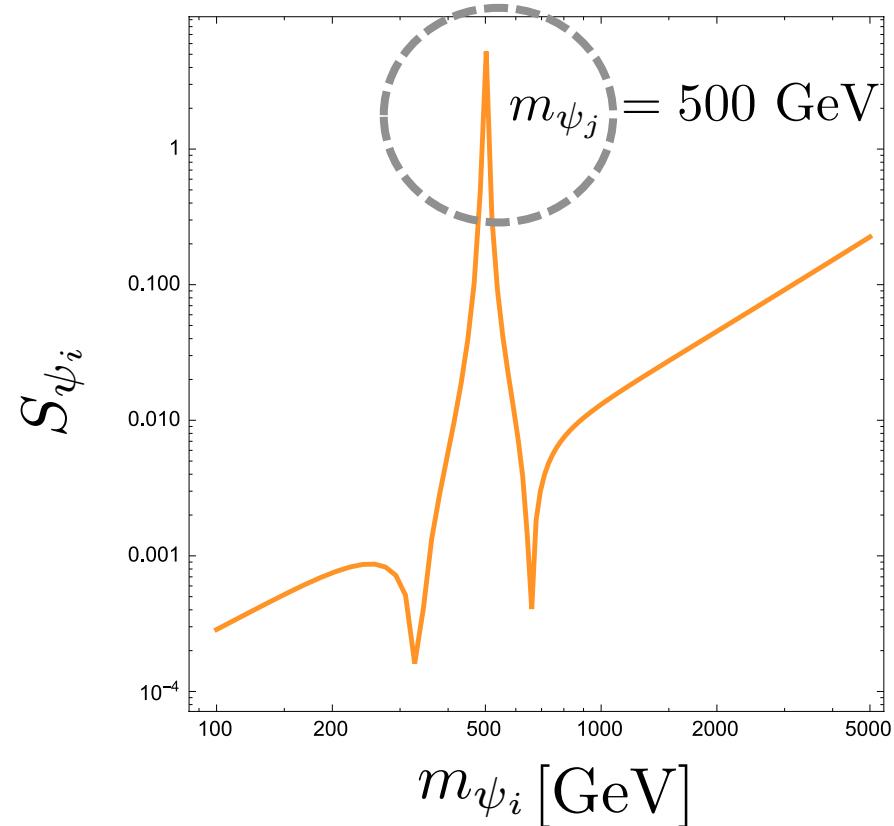
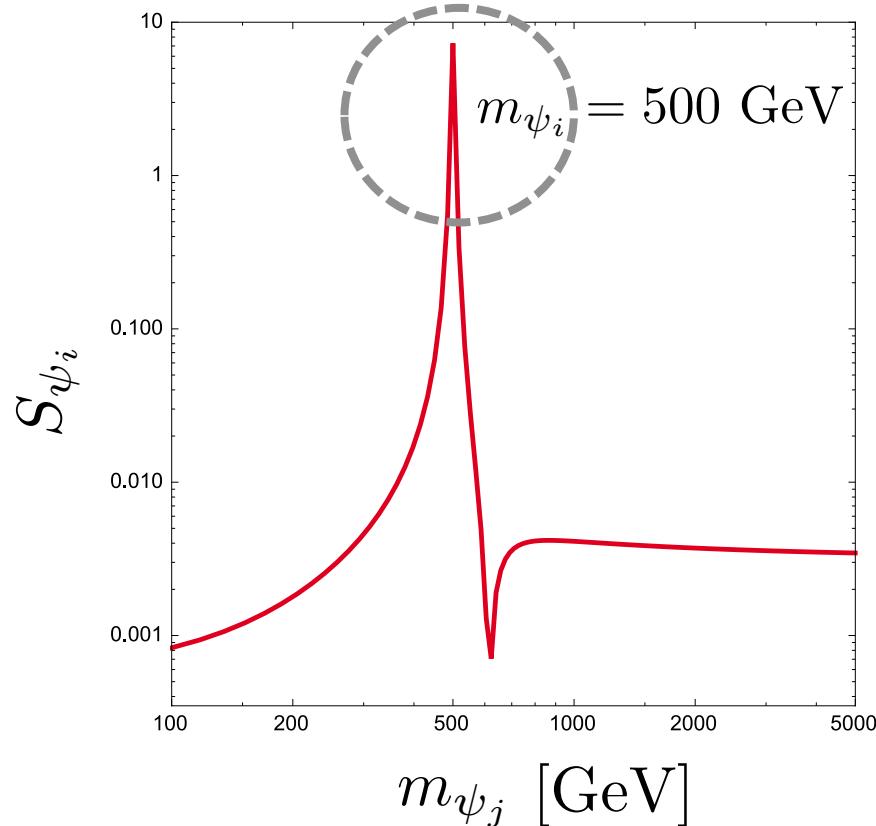
\* We get rid of model-dependent parameters, VEVs and  $\beta$ .



# CP-violating source as a function of $m_{\psi_{i,j}}$

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Enhancement shows up due to thermal effect in  $C_{\text{BAU}}$ .

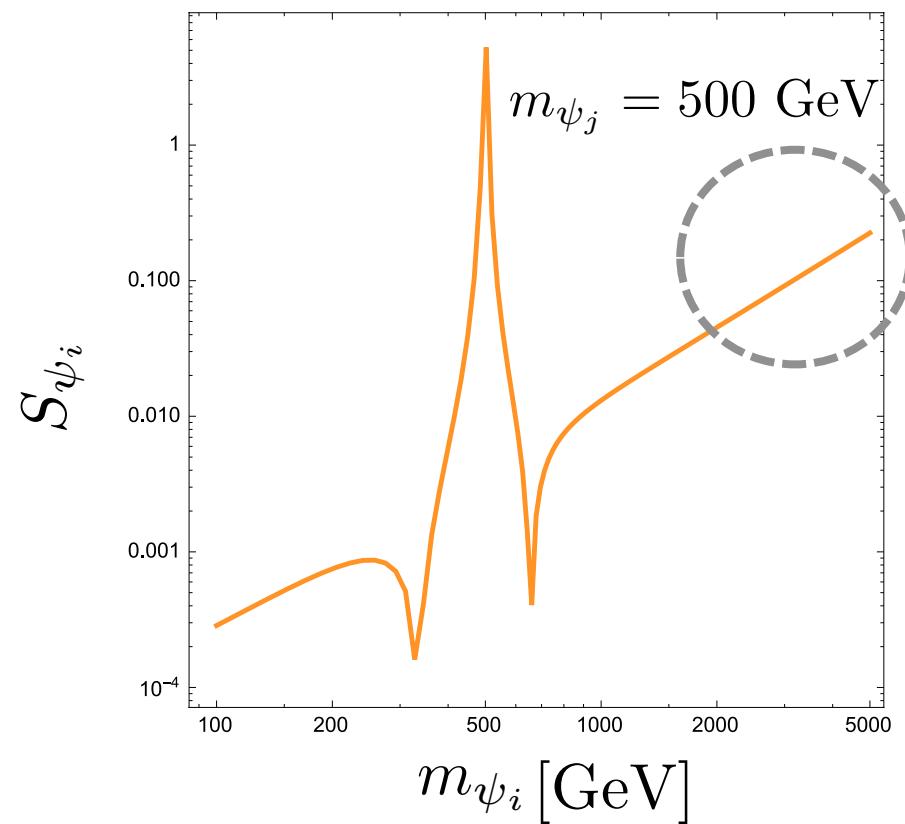
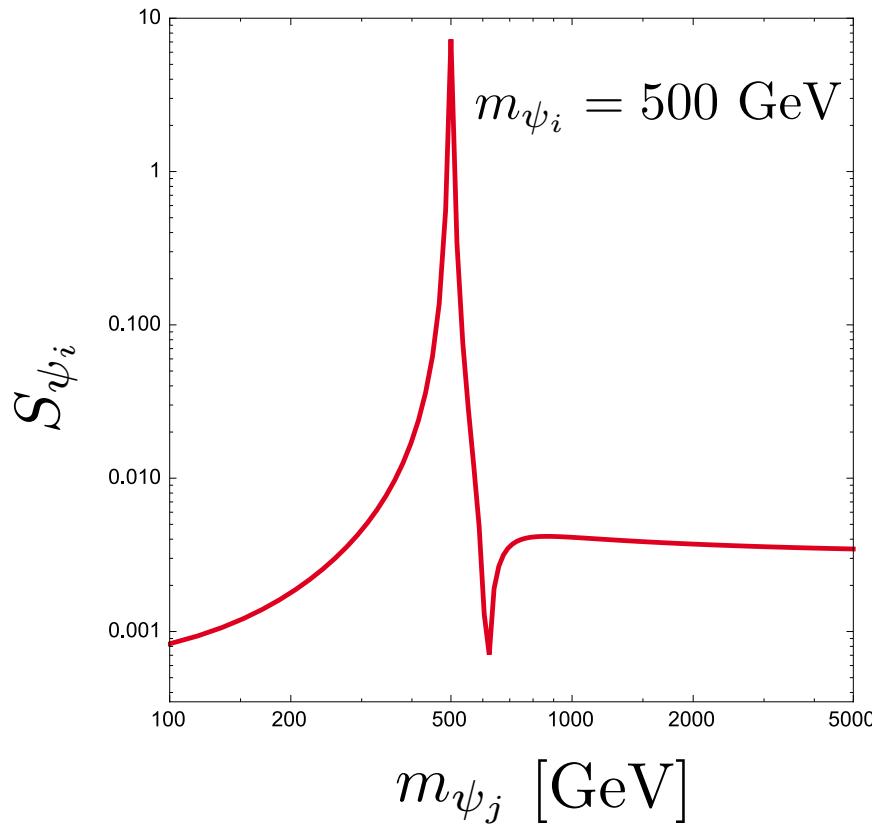


# CP-violating source as a function of $m_{\psi_{i,j}}$

For numerical calculations,  $|d_e^{\text{exp}}| = 8.7 \times 10^{-29} e \cdot \text{cm}$

Due to the rapid suppression of

$$C_{\text{EDM}}^{WW} \sim \frac{m_e m_j}{m_i^3}$$



# Particle contents

Particle contents are NMSSM-like.

Scalars	$SU(3)_C \times SU(2)_L \times U(1)_Y$	$Z_2$
$\Phi_1$	(1, 2, 1/2)	-
$\Phi_2$	(1, 2, 1/2)	+
$S$	(1, 1, 0)	-

Fermions	$SU(3)_C \times SU(2)_L \times U(1)_Y$	$Z_2$
$\tilde{\Phi}_1$	(1, 2, 1/2)	-
$\tilde{\Phi}_2$	(1, 2, 1/2)	+
$\tilde{S}^0$	(1, 1, 0)	-

Parameters are set in such a way as to take the limit of the real-singlet model.

# BAU vs EDM

$$|c_L| = |c_R| = 0.42, \phi = \phi_L - \phi_R = 225^\circ$$

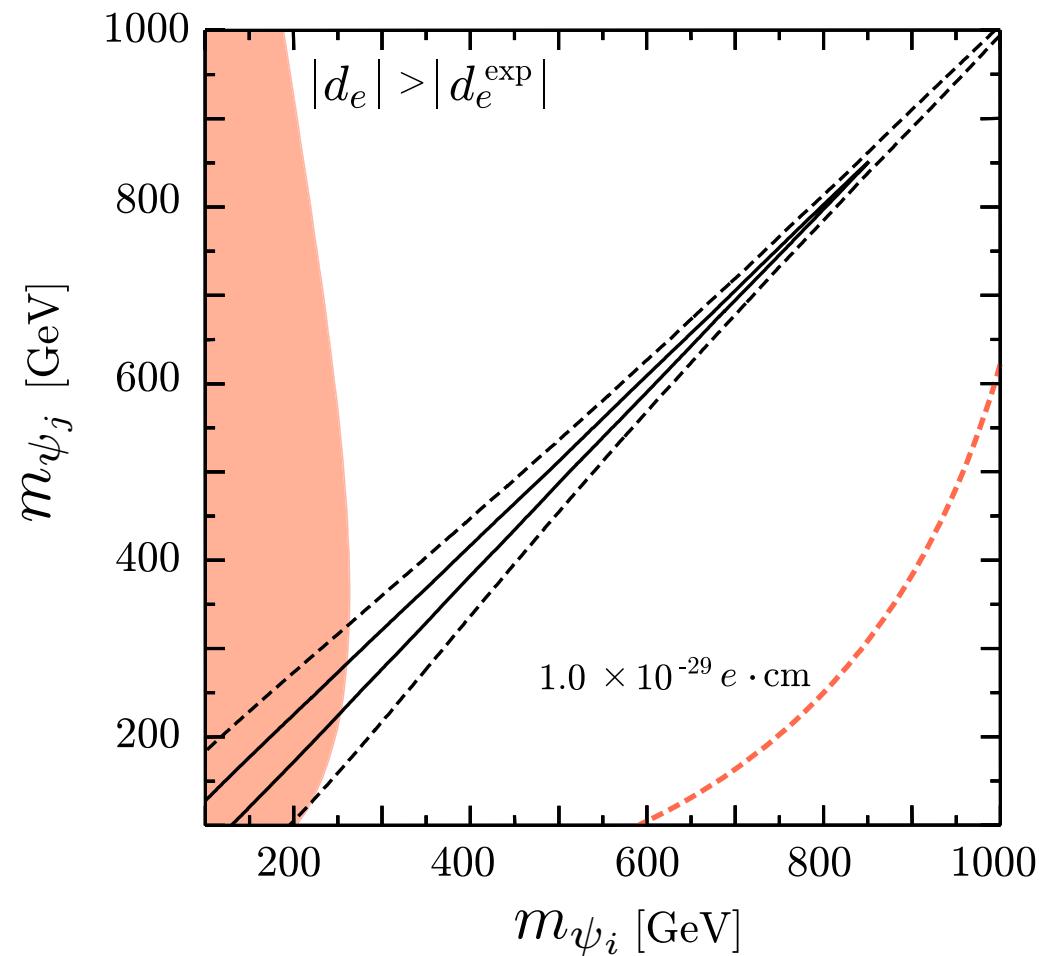
Solid line :

$$Y_B/Y_B^{\text{obs}} = 1$$

Dashed line :

$$Y_B/Y_B^{\text{obs}} = 0.1$$

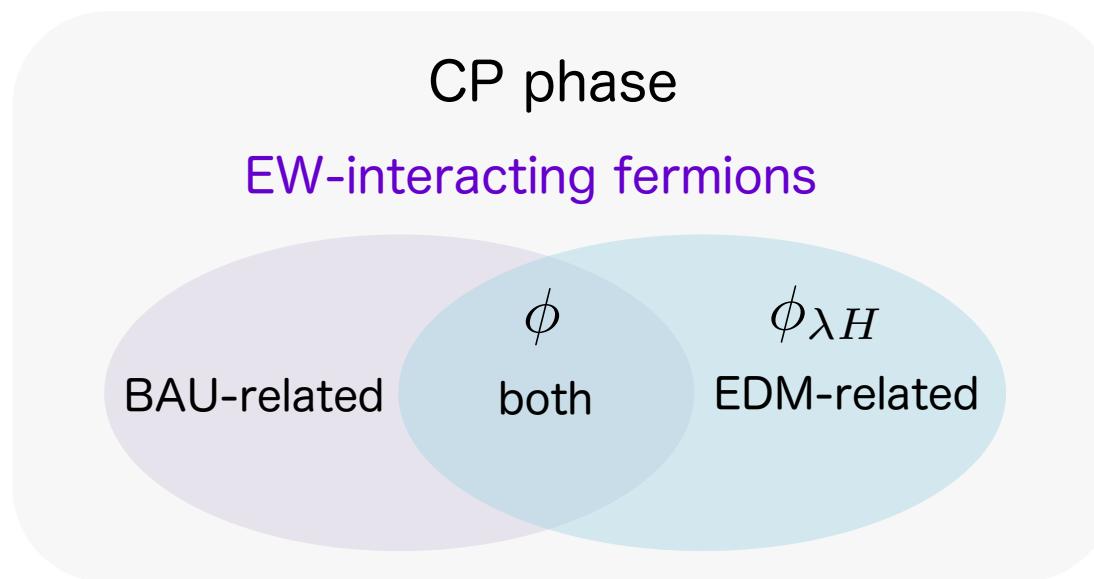
If  $d_e = 10^{-29} e \cdot \text{cm}$  can be achieved, the successful region would be tested !



Other CP phase

# Other CP phase

If the other CP phase exists, the situation is different.



The other Barr-Zee diagram can be induced.

# The other Barr-Zee diagram

Real singlet  $h_S$  has the following interactions.

$$\mathcal{L} \ni h_S \bar{\psi}^+ (g^S + i\gamma_5 g^P) \psi^+$$

with  $g^S = |\lambda| \cos \phi_{\lambda H}$

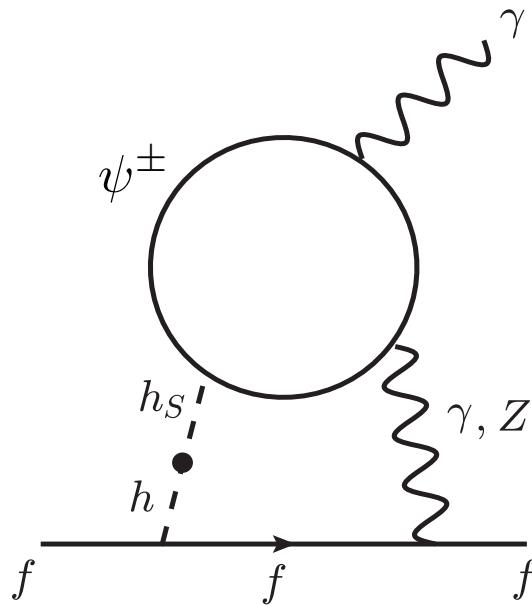
$$g^P = -|\lambda| \sin \phi_{\lambda H}$$

# The other Barr-Zee diagram

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$$\mathcal{L} \ni h_S \bar{\psi}^+ (g^S + i\gamma_5 g^P) \psi^+$$

This interaction induces the other Barr-Zee diagram,  $d_e^{H\gamma}$ .



The electron EDM becomes

$$d_e^{\text{sum}} = d_e^{WW} + d_e^{H\gamma}$$

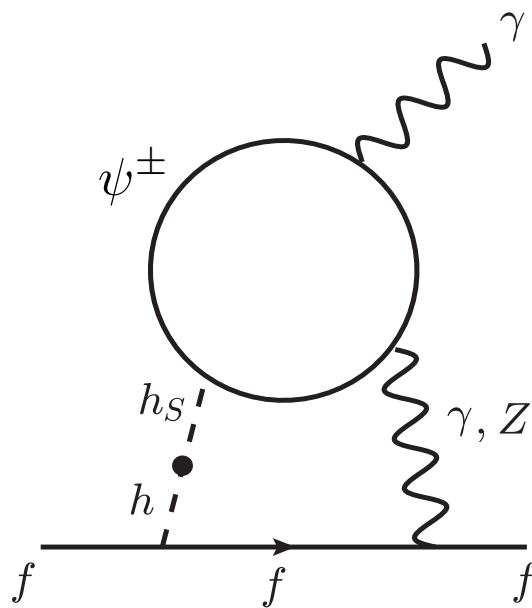
It is possible that  $d_e^{\text{sum}} = 0$ .

# The other Barr-Zee diagram

Real singlet  $h_S$  has the following interactions.

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This interaction induces the other Barr-Zee diagram,  $d_e^{H\gamma}$ .



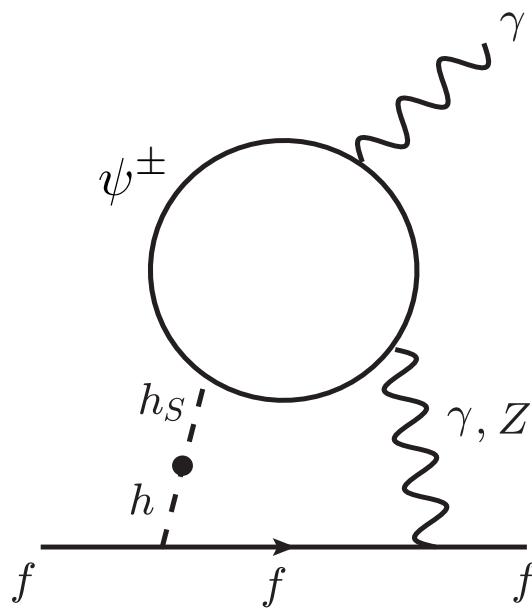
Even if  $d_e^{\text{sum}} = 0$ ,  
the Higgs physics helps  
the verifiability in this model.

# The other Barr-Zee diagram

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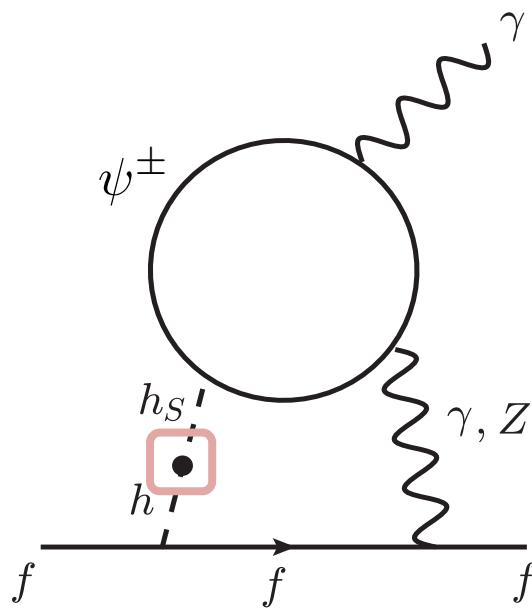
The first loop contributes to the signal strength of the Higgs to two gammas  $\mu_{\gamma\gamma}$ .

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The mixing angle  $\gamma$  is required to be a certain size for EWPT.

Mixing < LHC constraint

# Electron EDM

$$d_e^{\text{sum}} = d_e^{WW} + d_e^{H\gamma} \quad \text{in } (\lambda, \phi_{\lambda H}) \text{ plane.}$$

## □ Parameters

$$m_{\psi_i} = 300 \text{ GeV}$$

$$m_{\psi_j} = 277 \text{ GeV}$$

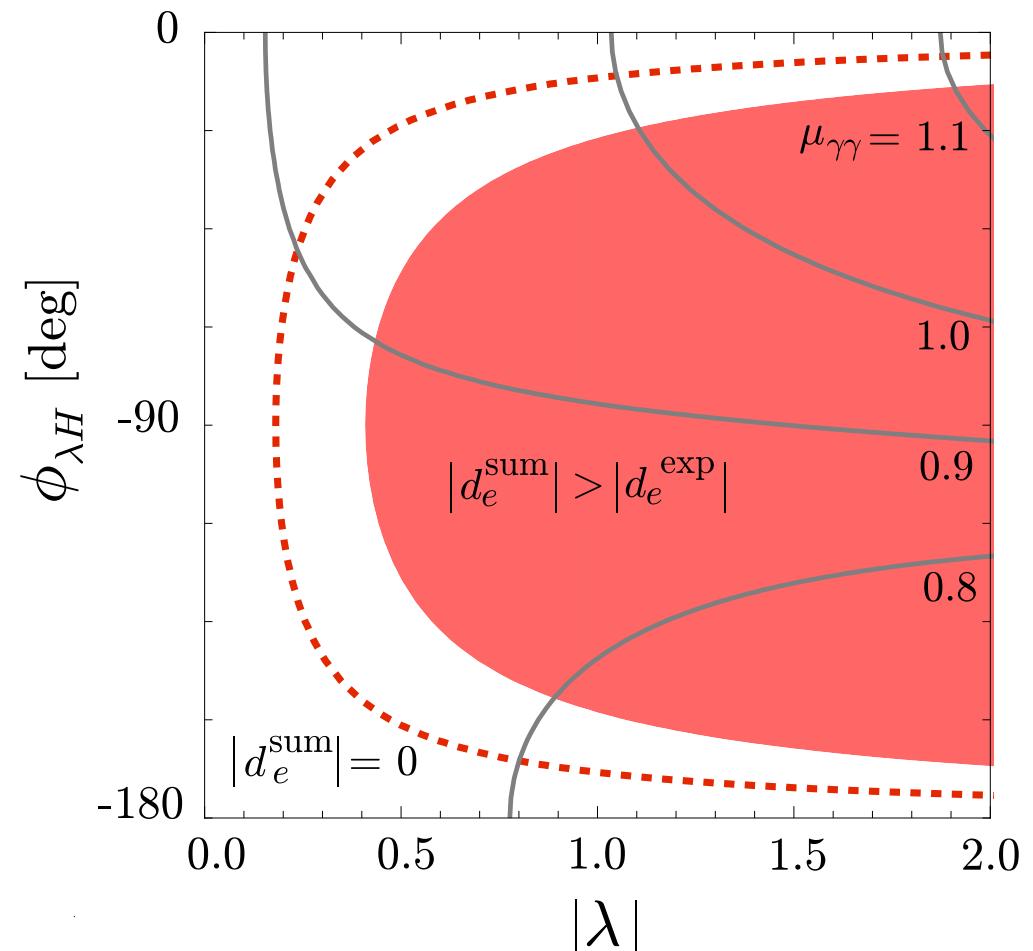
$$\gamma = -20^\circ$$

which leads to

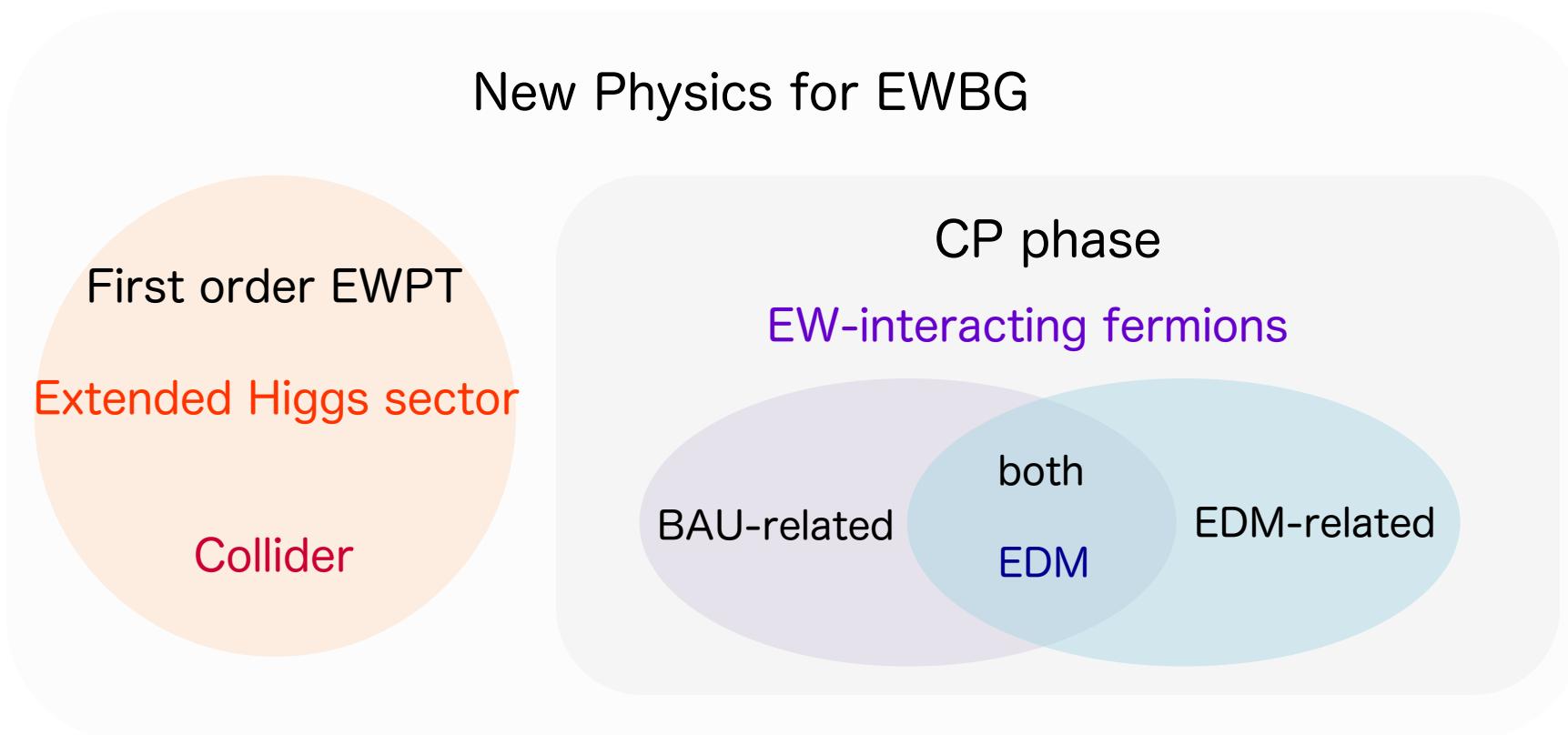
$$Y_B/Y_B^{\text{obs}} = 1$$

■ Gray line :  $\mu_{\gamma\gamma}$

Signal strength of the Higgs decay to two gammas



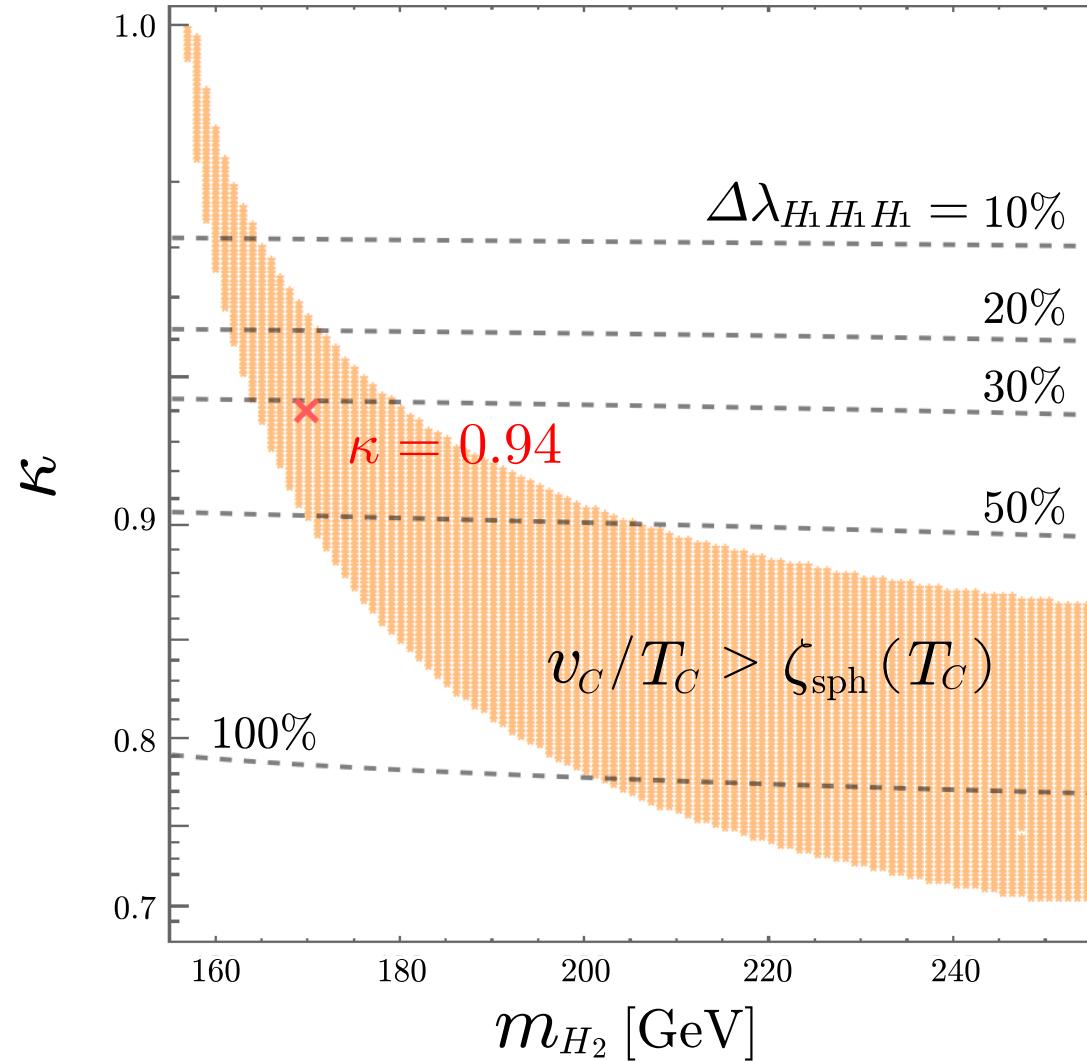
# Summary



As long as only the BAU-related CP phase exists in the model,  
it can be verified by the electron EDM.

Buckup

# EWPT in the real-singlet model



# Coupling measurements

Higgs couplings :

$$\kappa_V = 1.09 \pm 0.07 \text{ (ATLAS)}, \quad \kappa_V = 1.11 \pm 0.16 \text{ (ATLAS)},$$

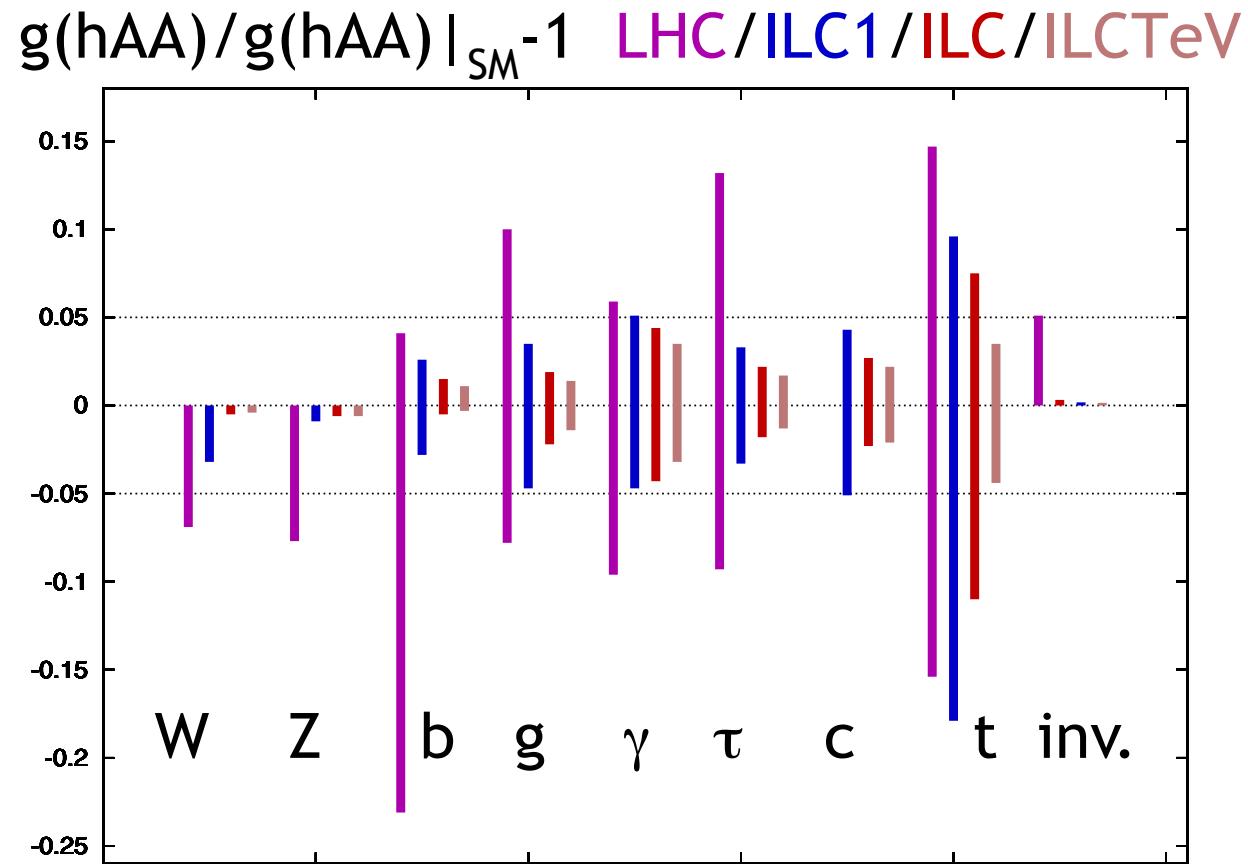
$$\kappa_V = 1.01^{+0.07}_{-0.07} \text{ (CMS)}, \quad \kappa_F = 0.89^{+0.14}_{-0.13} \text{ (CMS)},$$

Signal strength :  $\mu_{\gamma\gamma} = 1.17 \pm 0.27 \text{ (ATLAS)},$

$$\mu_{\gamma\gamma} = 1.14^{+0.26}_{-0.23} \text{ (CMS)},$$

# Coupling measurements

M. Peskin, 1207.2516



# Coupling measurements

	ILC(250)	ILC(500)	ILC(1000)	ILC(LumUp)
$\sqrt{s}$ (GeV)	250	250+500	250+500+1000	250+500+1000
$L$ ( $\text{fb}^{-1}$ )	250	250+500	250+500+1000	1150+1600+2500
$\gamma\gamma$	18 %	8.4 %	4.0 %	2.4 %
$gg$	6.4 %	2.3 %	1.6 %	0.9 %
$WW$	4.8 %	1.1 %	1.1 %	0.6 %
$ZZ$	1.3 %	1.0 %	1.0 %	0.5 %
$t\bar{t}$	—	14 %	3.1 %	1.9 %
$b\bar{b}$	5.3 %	1.6 %	1.3 %	0.7 %
$\tau^+\tau^-$	5.7 %	2.3 %	1.6 %	0.9 %
$c\bar{c}$	6.8 %	2.8 %	1.8 %	1.0 %
$\mu^+\mu^-$	91%	91%	16 %	10 %
$\Gamma_T(h)$	12 %	4.9 %	4.5 %	2.3 %
$hhh$	—	83 %	21 %	13 %
BR(invis.)	< 0.9 %	< 0.9 %	< 0.9 %	< 0.4 %