

Nonperturbative Beta Function of the SU(3) Gauge Theory with $N_f=10$ Massless Fermions

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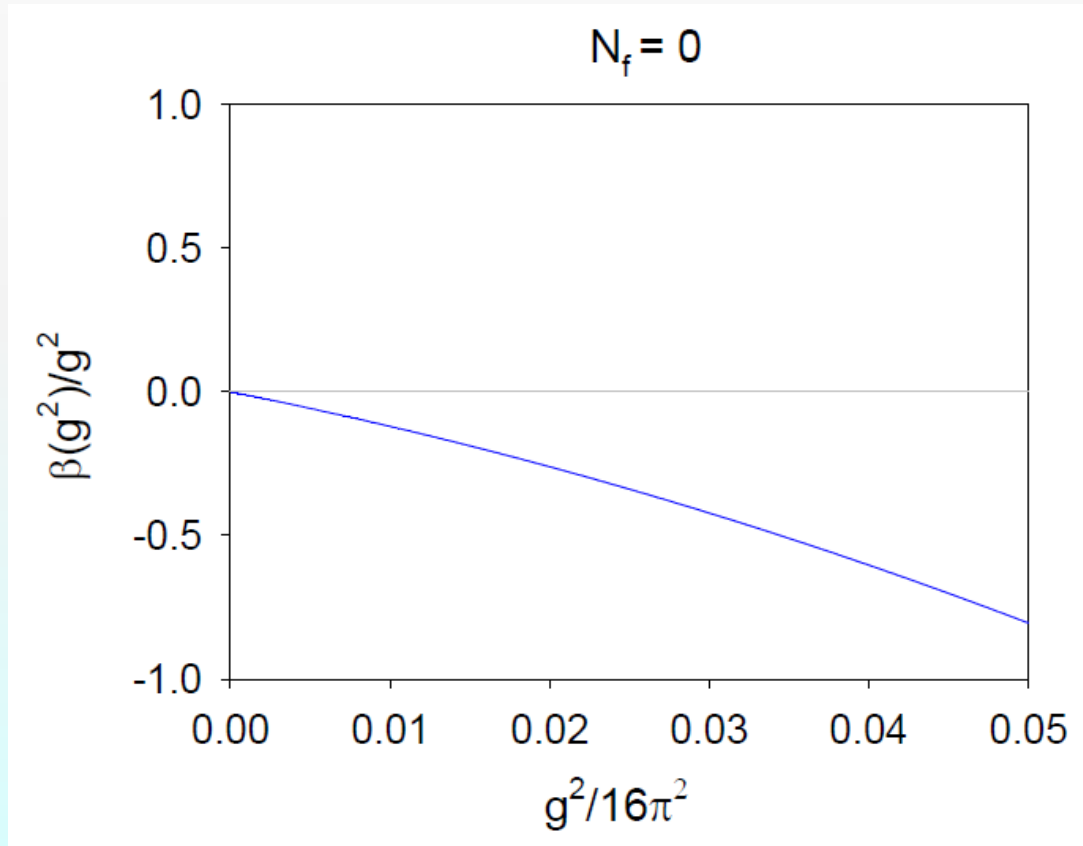
Outline

- **Introduction**
- **Finite Volume Gradient Flow Scheme**
- **Simulation of SU(3) gauge theory with $N_f=10$**
- **Renormalized Couplings**
- **Discrete β -function**
- **Concluding Remarks**

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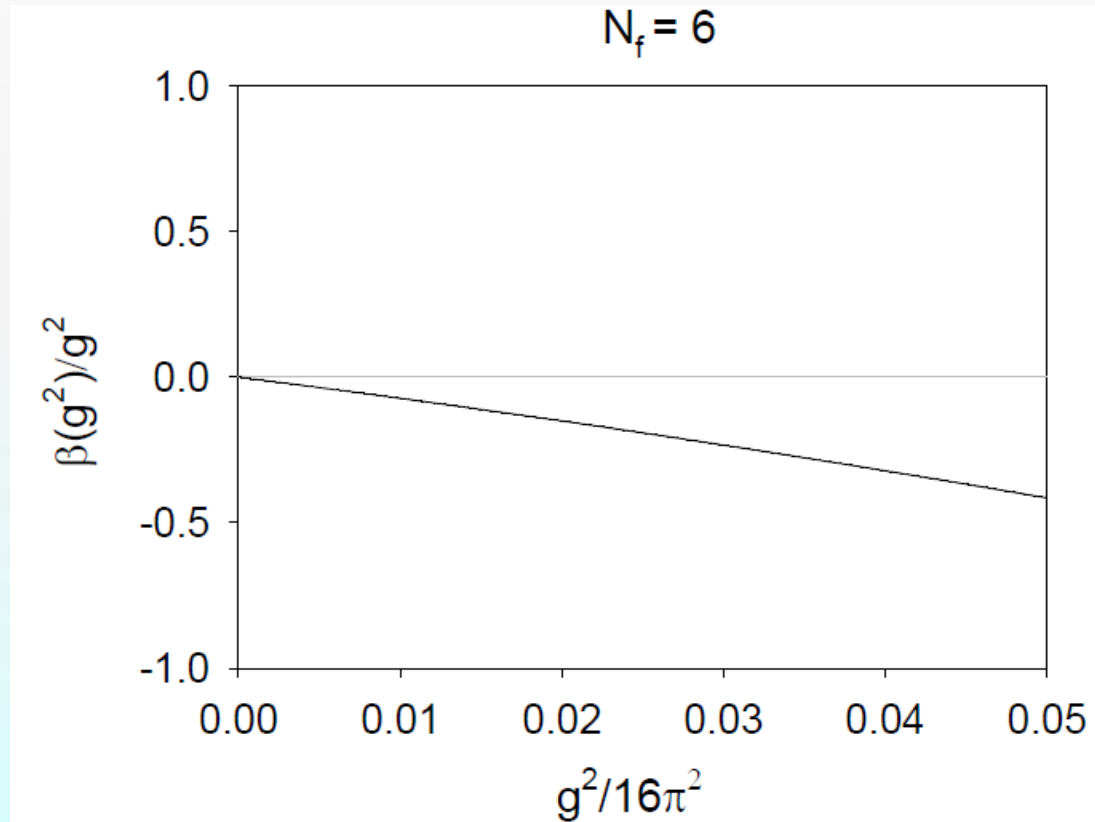
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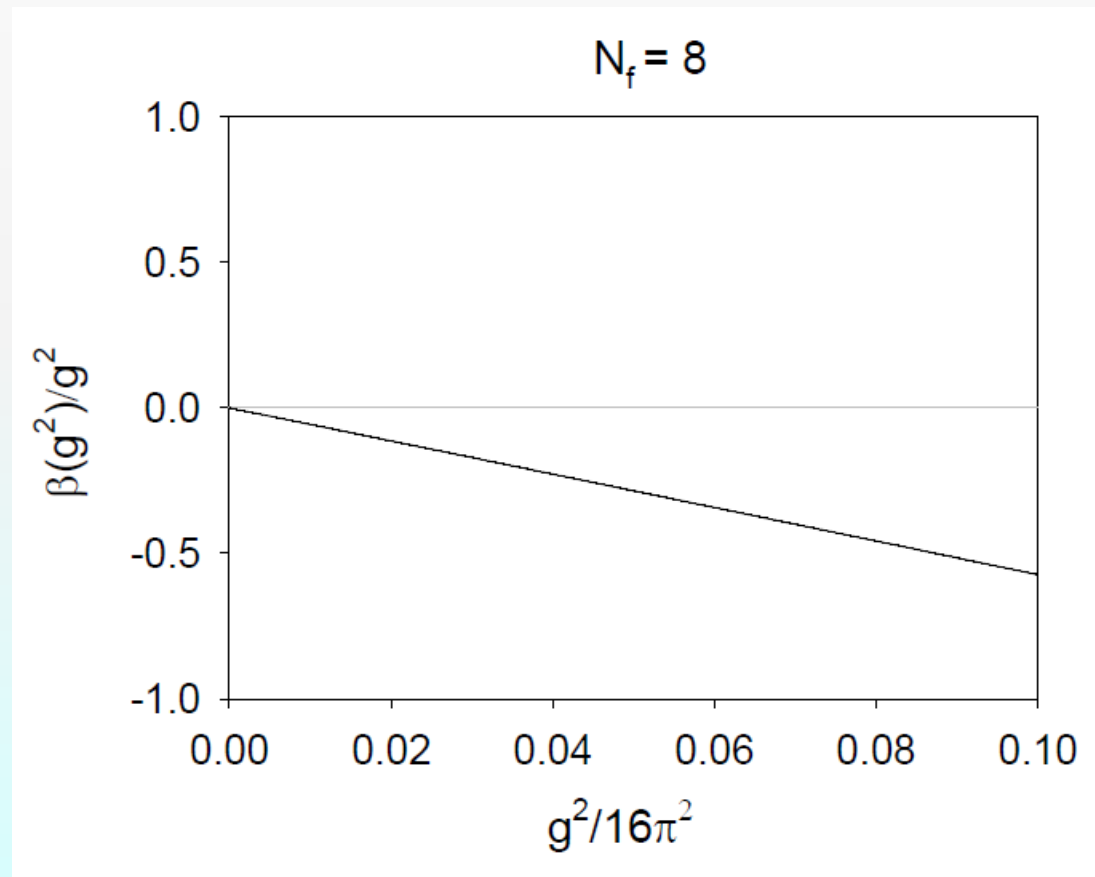
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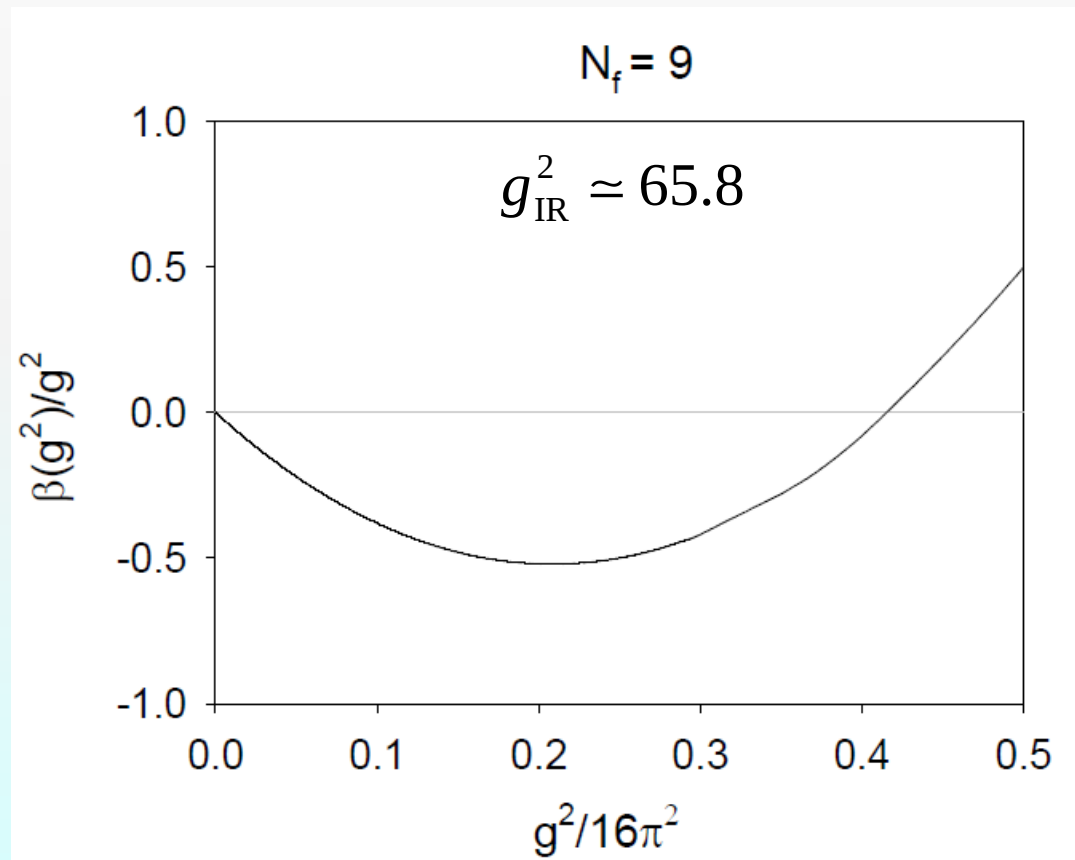
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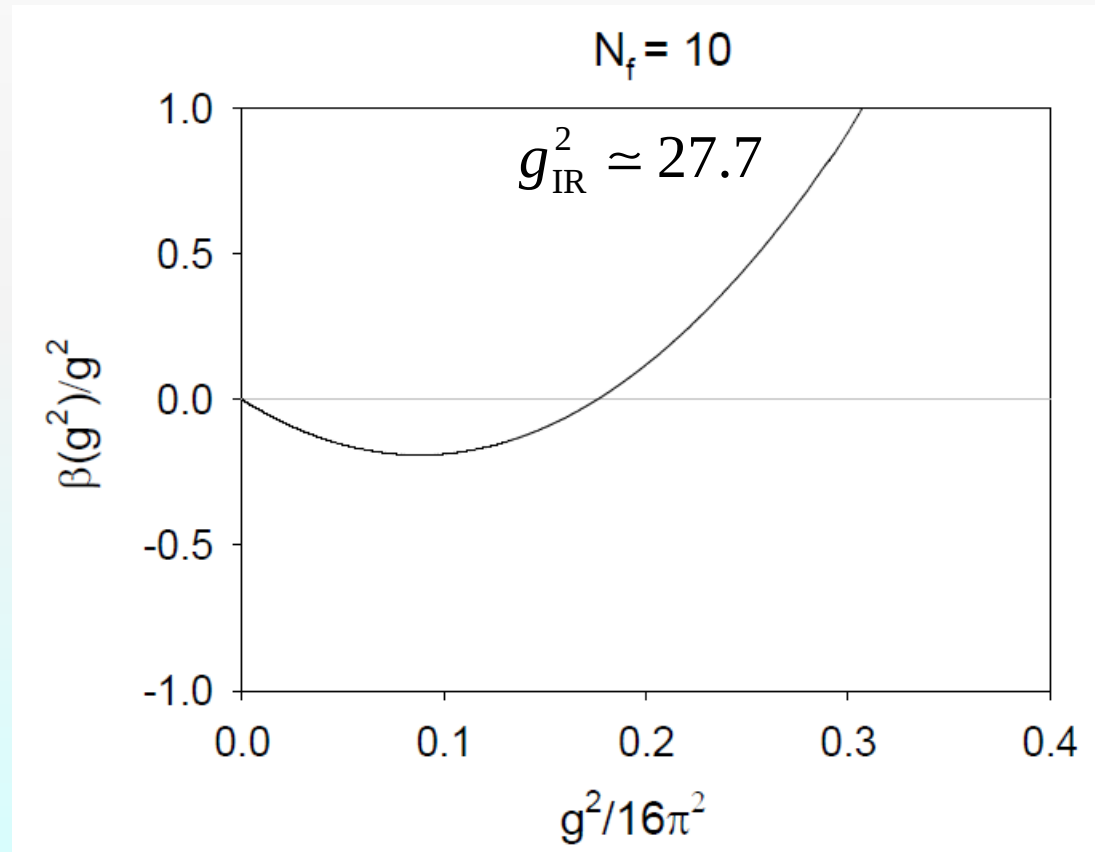
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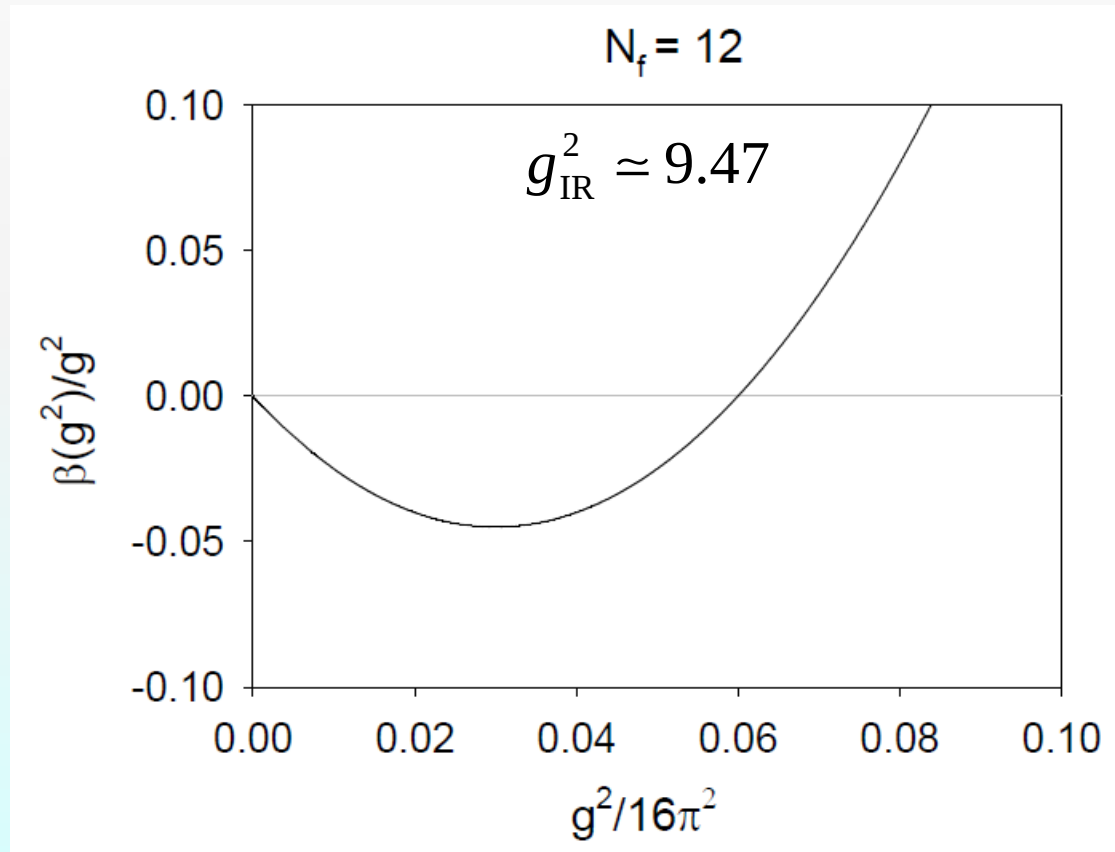
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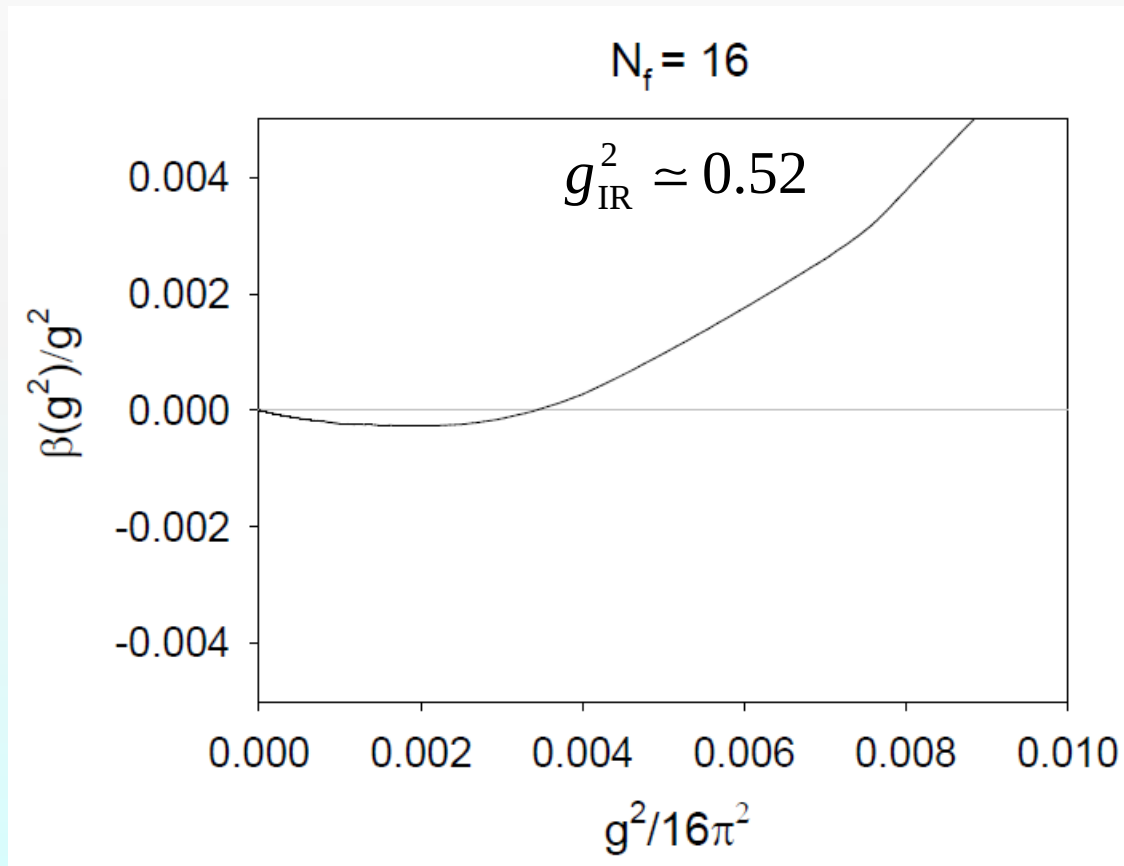
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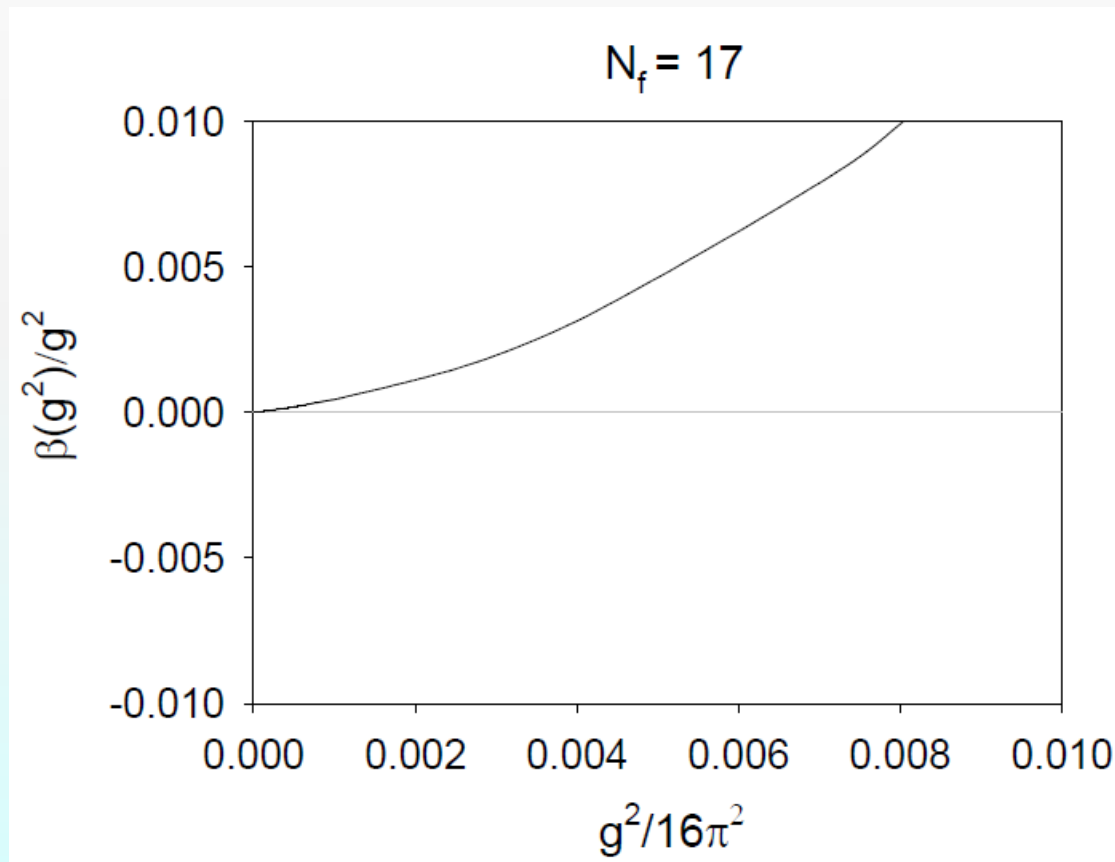
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- **The strong interaction physics around IRFP may have significant impacts to the phenomenological models beyond the Standard Model, e.g., models of composite Higgs boson.**
- **Lattice field theory provides a viable framework for nonperturbative study of vector gauge theories.**
- **It is vital to use lattice fermions with exact chiral symmetry, having exactly the same flavor symmetry as their counterparts in the continuum.**

Finite Volume Gradient Flow Scheme

The gradient flow (Wilson flow) amounts to solving

$$\frac{dB_\mu}{dt} = D_\nu G_{\nu\mu}, \quad G_{\mu\nu} = \partial_\mu B_\nu - \partial_\nu B_\mu + [B_\mu, B_\nu], \quad D_\nu = \partial_\mu + [B_\mu, \],$$

in the **flow time t** with the initial condition $B_\mu|_{t=0} = A_\mu$

The gradient flow is a process to average the gauge field over a 4D spherical range with $R_{\text{rms}} = \sqrt{8t}$

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The renormalized coupling on the lattice L^4 is obtained from

$$g_c^2(L, a) \propto \langle t^2 E(t) \rangle, \quad c = \frac{\sqrt{8t}}{L} = \text{constant},$$

$$E(t) = \frac{1}{4} G_{\mu\nu}^a G_{\mu\nu}^a(t), \text{ energy density}$$

The proportional constant is fixed by requiring $g_c^2(L, a) = g_{\overline{MS}}^2$ to the 1-loop order.

Simulation of $SU(3)$ gauge theory with $N_f=10$

- Gauge Action: Wilson plaquette action

$\beta = 6/g_0^2 = 6.45, 6.5, 6.6, 6.7, 6.8, 7.0, 7.5, 8.0, 9.0, 10.0, 12.0, 15.0$

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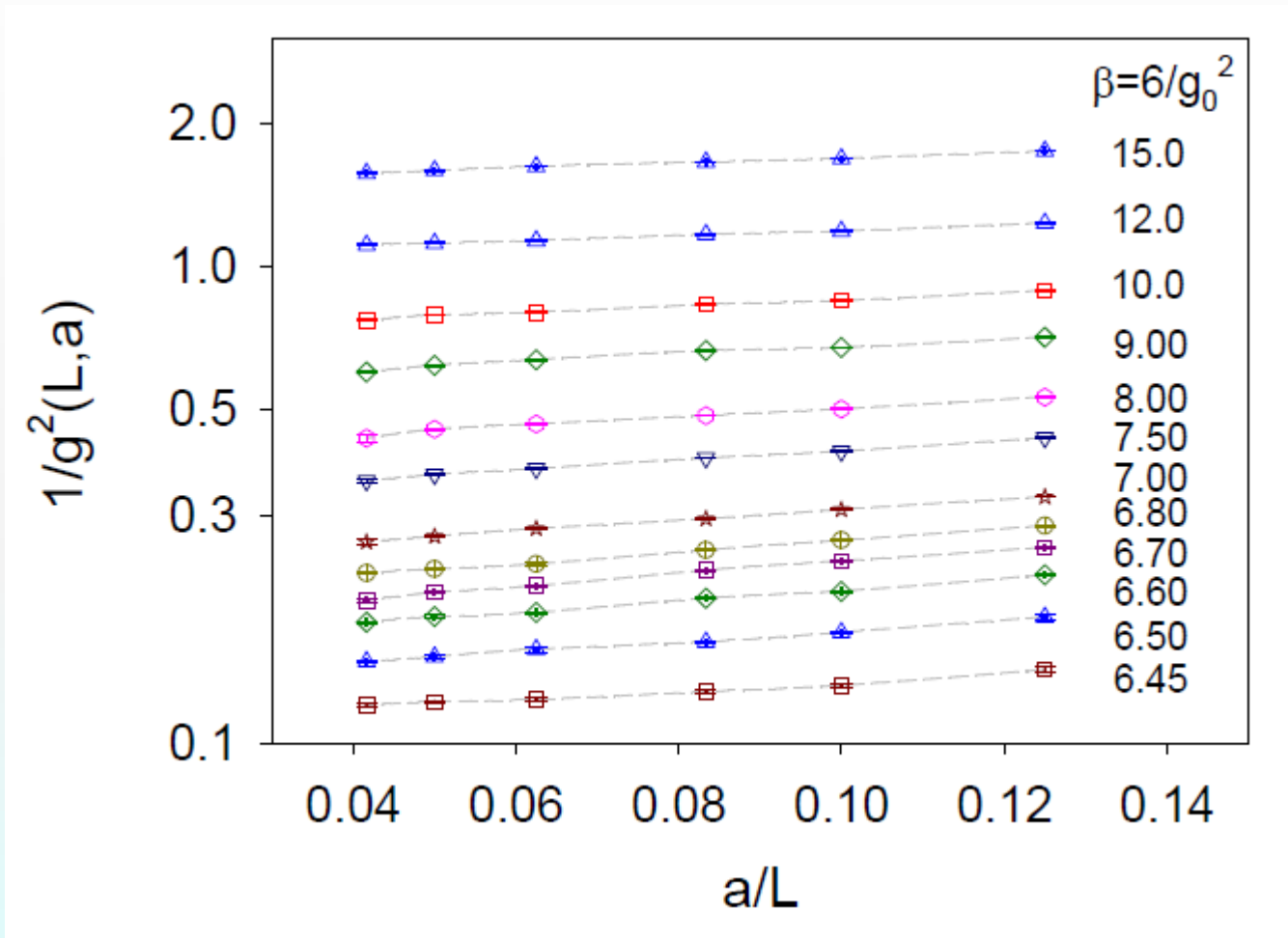
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- The residual mass of any gauge ensemble is less than 5×10^{-5}

Renormalized Couplings with $c = \frac{\sqrt{8t}}{L} = 0.3$



Other values of $g^2(L, a)$ are obtained by cubic spline interpolation.

Discrete β -function (DBF)

$$\beta(s, a/L) = \frac{g^2(sL, a) - g^2(L, a)}{\ln(s^2)}, \text{ for each pair of lattices } (L, sL).$$

Given a set of lattice pairs with a fixed ratio s , $\lim_{a \rightarrow 0} \beta(s, a/L)$

can be obtained by extrapolation. Moreover, if $\lim_{a \rightarrow 0} \beta(s, a/L)$

is available for several s , then the limit $s \rightarrow 1$ can be taken,

$$\lim_{s \rightarrow 1} \lim_{a \rightarrow 0} \beta(s, a/L) = -\beta(g^2) = -\frac{dg^2}{d \ln \mu^2}$$

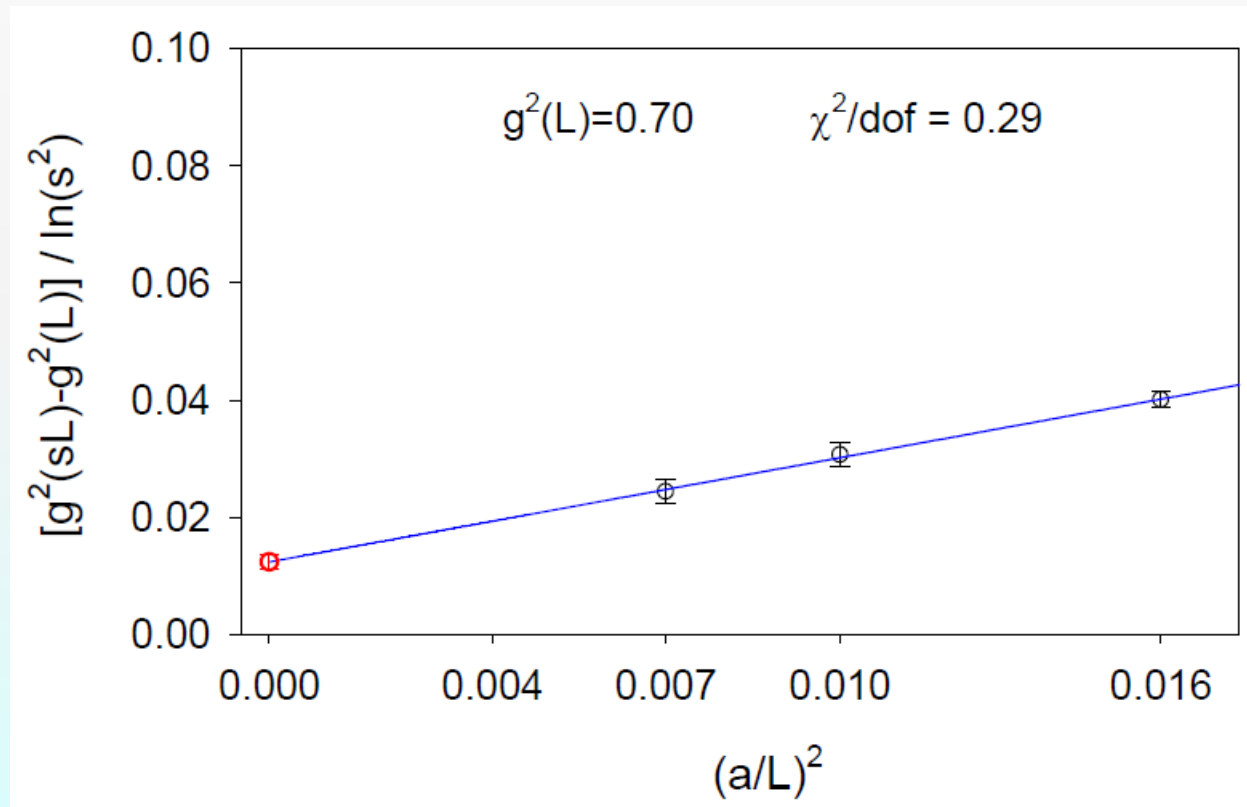
If $\lim_{a \rightarrow 0} \beta(s, a/L)$ has IRFP, then $\beta(g^2)$ also has IRFP, and vice versa.

In the following, use lattice pairs (8,16), (10,20), (12,24) with $s = 2$,

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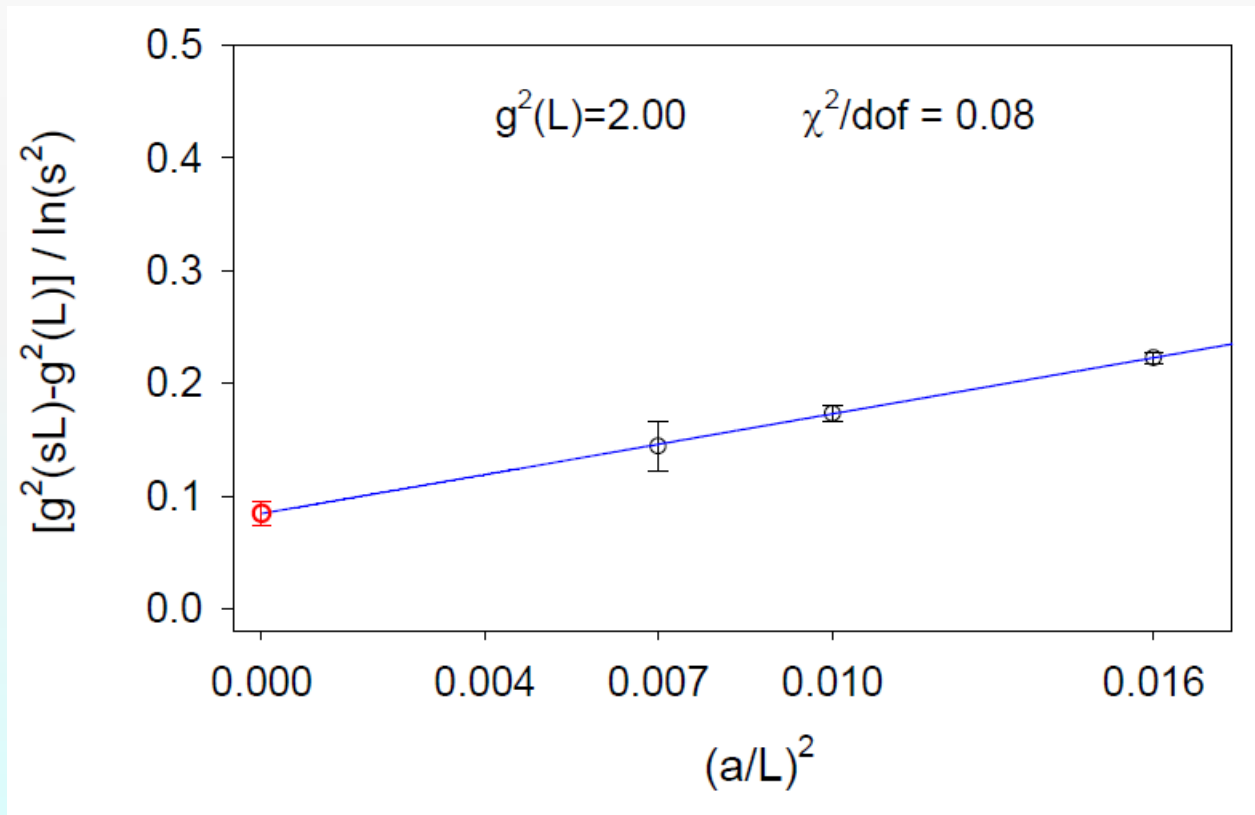
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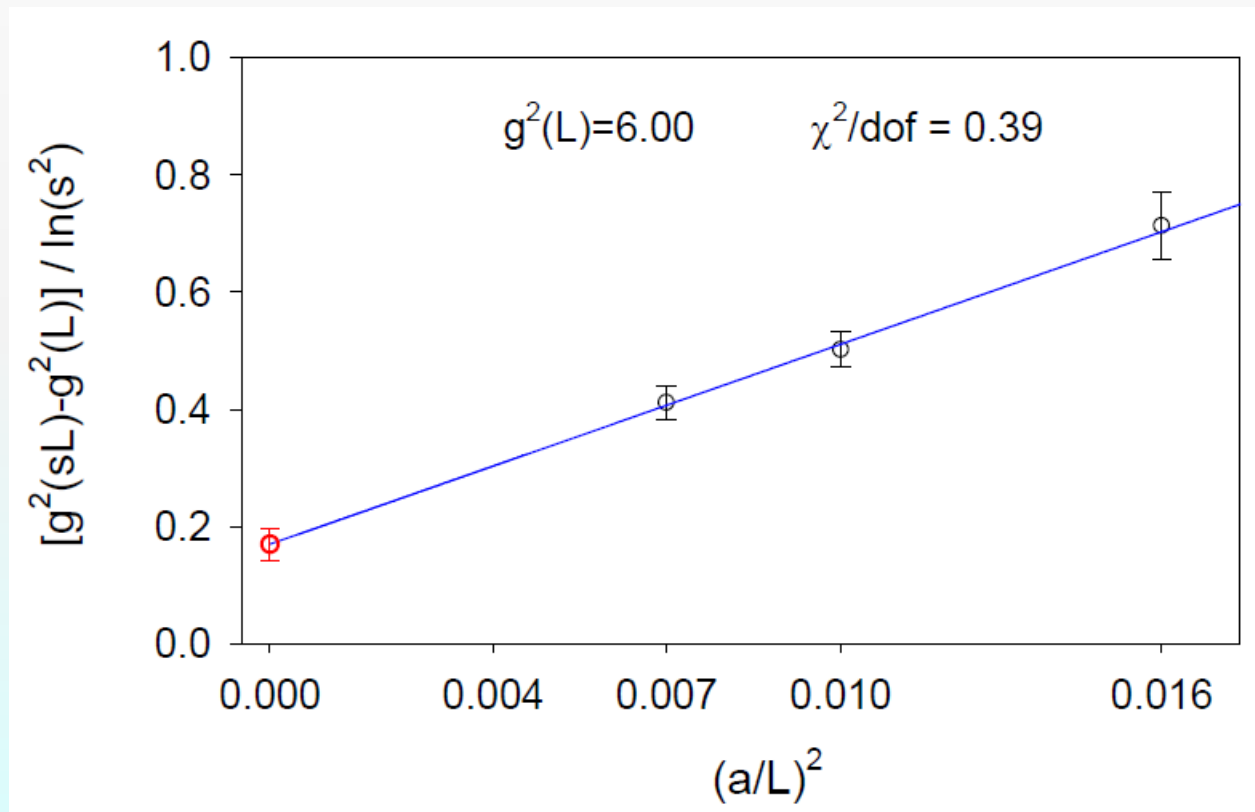
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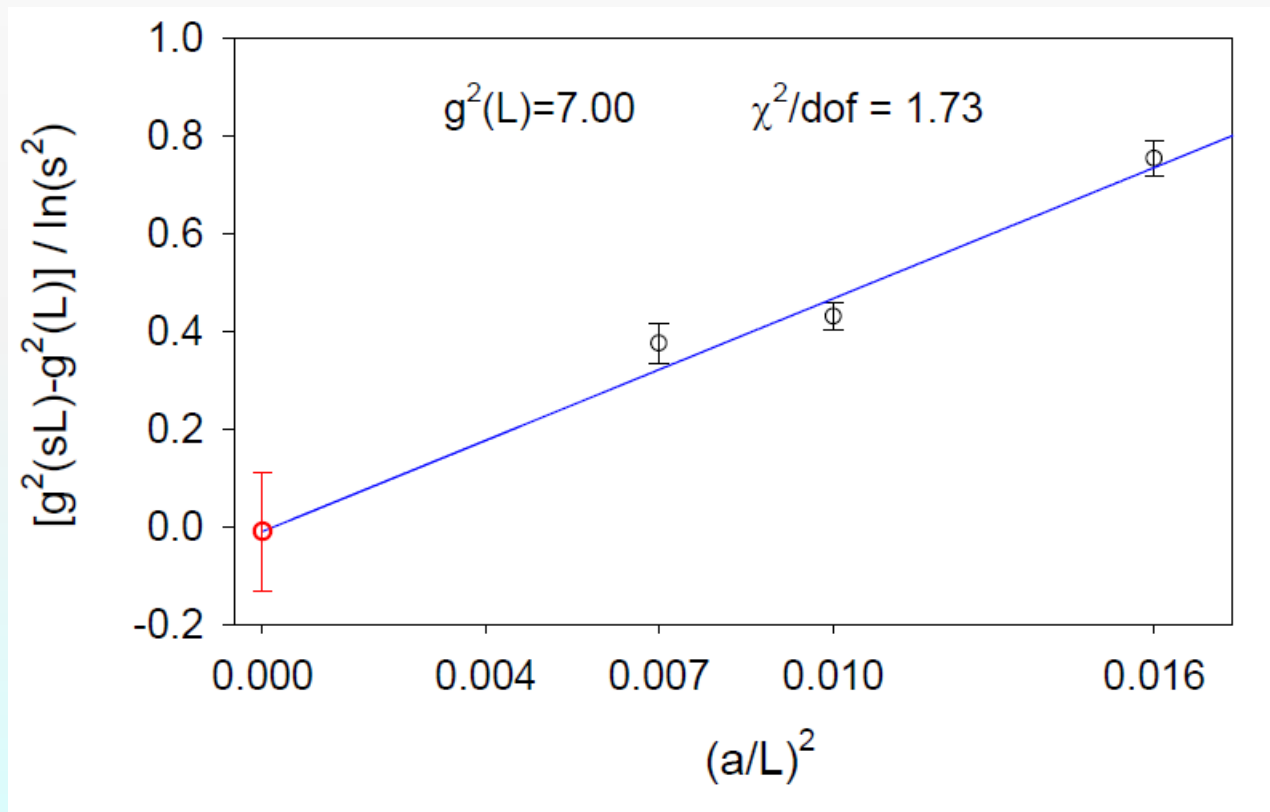
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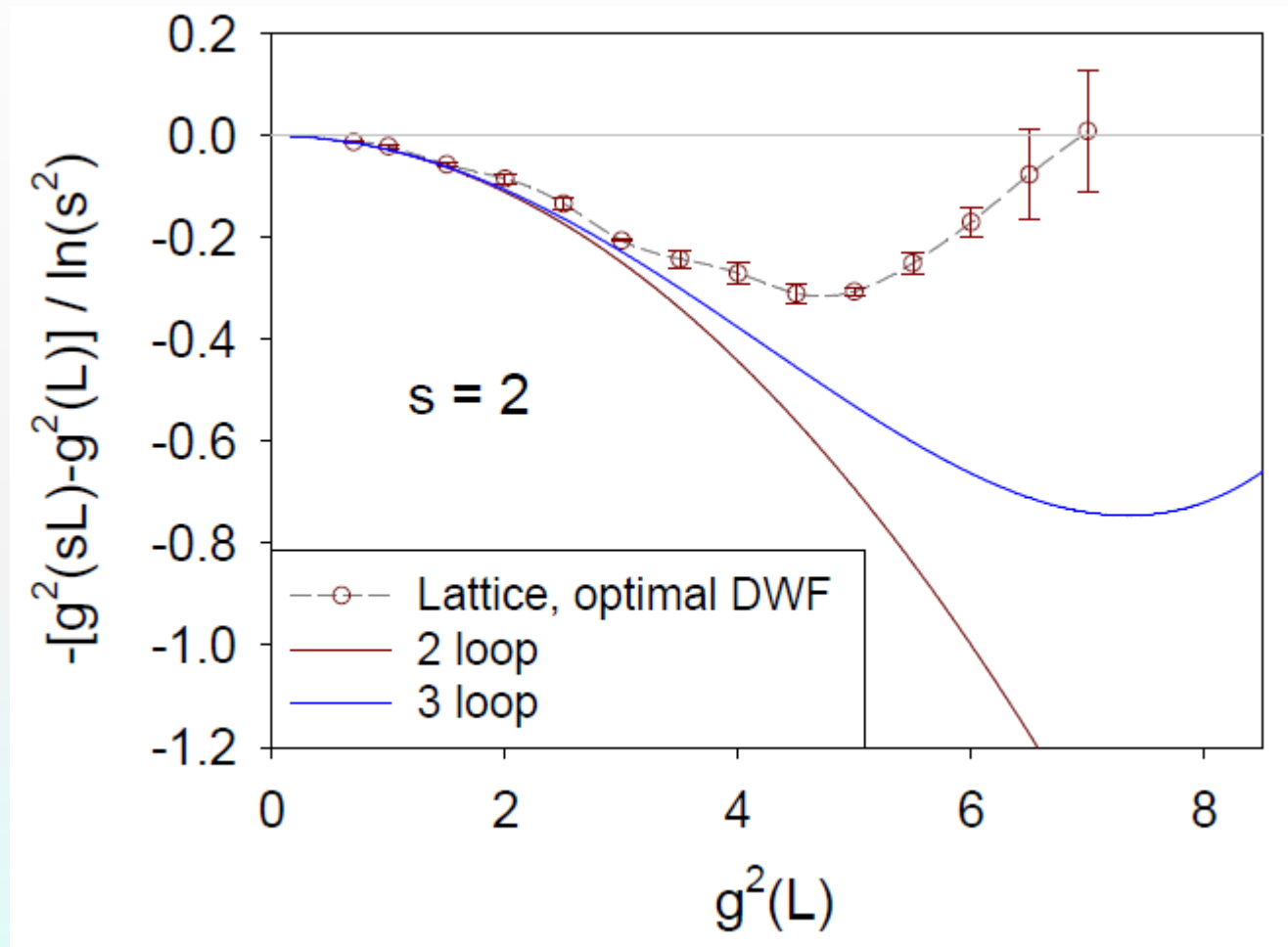
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Discrete β -function of $SU(3)$ Gauge Theory with $N_f=10$



Concluding Remarks

- The present data of the discrete β -function (with $s=2$) of $SU(3)$ gauge theory with 10 massless optimal DWF suggests that the theory may possess IRFP at $g_c^2 \sim 7.0$, and its long-distance behavior is conformal.

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- The statistical+systematic errors of the discrete β -function around the IRFP can be further reduced by increasing the statistics as well as more data points around the regime of strong couplings. Moreover, it is instructive to simulate $L=32$ (a difficult task !) and add the fourth lattice pair (16,32) to the existing data of DBF.