New supersymmetric localizations from topological gravity

Dario Rosa

Seoul National University

Based on: 1510.00006 with J. Bae, S. J. Rey and C. Imbimbo (see also 1411.6635 with C. Imbimbo)

8th Taiwan String Workshop, November 16-20, 2015

Main results of the last few years

Exact results for SUSY QFTs in different dimensions and for different manifolds



Main results of the last few years

Exact results for SUSY QFTs in different dimensions and for different manifolds

Two main technical ingredients:

Localization — Exact results [Pestun,...]

Coupling to gravity — Manifolds with SUSY

[Festuccia-Seiberg,...]

An algorithmical recipe

6

The two methods combined give a clear recipe:
1. Flat space SUSY relevant SUGRA
2. SUSY backgrounds from SUGRA

3. Rigid SUSY SUGRA_{freezed} + matter

4. Localization — exact results

A couple of remarks

Localization is deeply related to topological field theories

Examples:

1) Nekrasov instanton partition function

2) Chern-Simons theory using SUSY3) 2d GLSM to compute GW invariants

Localization relies on a cohomological formulation of SUSY

--- This is a very useful property!!



A couple of remarks

Localization is deeply related to topological field theories

Examples:

- 1) Nekrasov instanton partition function
- 2) Chern-Simons theory using SUSY3) 2d GLSM to compute GW invariants
- Localization relies on a cohomological formulation of SUSY

--- This is a very useful property!!

Unfortunately SUGRA is NOT cohomological Much more complicated!

A different perspective

Another, cohomological, gravity theory can be found



Remark: TG ≠ twisted SUGRA





A different perspective

Another, cohomological, gravity theory can be found



Many advantages:

1) Easier equations

2) Cohomological

3) 1 to 1 mapping to SUGRA



A different perspective

Another, cohomological, gravity theory can be found



Many advantages:

1) Easier equations

2) Cohomological



3) 1 to 1 mapping to SUGRA

TG is the best choice for localization's users!



Plan of the talk

1) Our rigid topological starting point: 2d Yang-Mills

2d Yang-Mills coupled to TG

- 2) Classification of TG invariant backgrounds
- 3) From TG backgrounds to SUGRA backgrounds





Plan of the talk

1) Our rigid topological starting point: 2d Yang-Mills

2d Yang-Mills coupled to TG

- 2) Classification of TG invariant backgrounds
- 3) From TG backgrounds to SUGRA backgrounds

A full classification of 2d N=(2,2) SUGRA localizing backgrounds



Our starting point: 2d Yang-Mills

"Naïve" thought:

2d gauge inv. kills two photons \longrightarrow Topological theory!



Our starting point: 2d Yang-Mills

"Naïve" thought:

2d gauge inv. kills two photons \longrightarrow Topological theory!

• Why Naïve? [Witten] $I_{YM}(g,\epsilon) = \int_{\Sigma} \operatorname{Tr} \phi F + \epsilon \int_{\Sigma} d^2x \sqrt{g} \frac{1}{2} \operatorname{Tr} \phi^2$ Auxiliary scalar. Integrating it out we standard action $I_{YM}(g,\epsilon) = \int_{\Sigma} \operatorname{Tr} \phi F + \epsilon \int_{\Sigma} d^2x \sqrt{g} \frac{1}{2} \operatorname{Tr} \phi^2$ Auxiliary scalar. Integrating it out we standard action $I_{YM}(g,\epsilon) = \int_{\Sigma} \operatorname{Tr} \phi F + \epsilon \int_{\Sigma} d^2x \sqrt{g} \frac{1}{2} \operatorname{Tr} \phi^2$ Auxiliary scalar.

Our starting point: 2d Yang-Mills

"Naïve" thought:

2d gauge inv. kills two photons \longrightarrow Topological theory!

• Why Naïve? [Witten] $I_{\rm YM}(g,\epsilon) = \int_{\Sigma} {\rm Tr} \, \phi \, F + \epsilon \int_{\Sigma} {\rm d}^2 x \, \sqrt{g} \, \frac{1}{2} {\rm Tr} \, \phi^2 \qquad {\rm standard\ action}$ $I_{\rm YM}(g,\epsilon) = \int_{\Sigma} {\rm Tr} \, \phi \, F + \epsilon \int_{\Sigma} {\rm d}^2 x \, \sqrt{g} \, \frac{1}{2} {\rm Tr} \, \phi^2 \qquad {\rm standard\ action}$ $I_{\rm The\ action\ is\ not\ fully\ topological.}$ It depends on the volume form

Making 2d Yang-Mills topological

We want to control the metric-dependence •

Replacement

Background 2-form $\epsilon \sqrt{g} \, \mathrm{d}^2 x \to -f^{(2)}$ Background 2-form It replaces the metric



Making 2d Yang-Mills topological

We want to control the metric-dependence .

Replacement

Background 2-form $\epsilon \sqrt{g} \, \mathrm{d}^2 x \to -f^{(2)}$ It replaces the metric

AND we require that the physics is unchanged

 $f^{(2)} = \Omega^{(2)} + \mathrm{d} \,\Omega^{(1)}$

Hödge decomposition

This part gives a -spurious dependence: We must remove it



Making 2d Yang-Mills topological

• We want to control the metric-dependence

Replacement

 $= \epsilon \sqrt{g} \, \mathrm{d}^2 x \to -f^{(2)} = f^{(2)} =$

AND we require that the physics is unchanged

 $f^{(2)} = \Omega^{(2)} + \mathrm{d}\,\Omega^{(1)}$

Hödge decomposition

This part gives a ---spurious dependence: We must remove it



The topological U(1) background

1. We make the spurious part BRST-trivial

0

Top. gaugino

 $s f^{(2)} = -d \psi^{(1)}$

The topological U(1) background

1. We make the spurious part BRST-trivial

Top. gaugino

$$s f^{(2)} = -d \psi^{(1)}$$

Topological U(1) background

2. This is not enough

$$s f^{(2)} = -d \psi^{(1)}$$

 $s \psi^{(1)} = -d \gamma^{(0)}$
 $s \gamma^{(0)} = 0$

0

The topological U(1) background

1. We make the spurious part BRST-trivial

Top. gaugino

$$s f^{(2)} = -d \psi^{(1)}$$

2. This is not enough

 \longrightarrow

This is the first actor of our new gravity background theory:

A topological abelian multiplet

$$s f^{(2)} = -d \psi^{(1)}$$

 $s \psi^{(1)} = -d \gamma^{(0)}$
 $s \gamma^{(0)} = 0$











By trying to couple 2d YM to TG we found a consistent coupling between TG and the topological background abelian multiplet. Let us call it "equivariant topological gravity"

Equivariant topological gravity

By coupling the two actors:

Equivariant topological gravity BRST operator

$$\begin{split} s \, g_{\mu\nu} &= -\mathcal{L}_{\xi} g_{\mu\nu} + \psi_{\mu\nu} & s \, f^{(2)} &= -\mathrm{d} \, \psi^{(1)} \\ s \, \psi_{\mu\nu} &= -\mathcal{L}_{\xi} \psi_{\mu\nu} + \mathcal{L}_{\gamma} g_{\mu\nu} & s \, \psi^{(1)} &= -\mathrm{d} \, \gamma^{(0)} + \iota_{\gamma} f^{(2)} & \text{Additional pieces} \\ s \, \xi^{\mu} &= -\frac{1}{2} \mathcal{L}_{\xi} \xi^{\mu} + \gamma^{\mu} & s \, \gamma^{(0)} &= \iota_{\gamma} \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{and} & s \, \psi^{(1)} &= -\mathrm{d} \, \psi^{(1)} & \text{$$



Equivariant topological gravity

By coupling the two actors:

Equivariant topological gravity BRST operator

$$\begin{split} s \, g_{\mu\nu} &= -\mathcal{L}_{\xi} g_{\mu\nu} + \psi_{\mu\nu} & s \, f^{(2)} = -\mathrm{d} \, \psi^{(1)} \\ s \, \psi_{\mu\nu} &= -\mathcal{L}_{\xi} \psi_{\mu\nu} + \mathcal{L}_{\gamma} g_{\mu\nu} & s \, \psi^{(1)} = -\mathrm{d} \, \gamma^{(0)} + \iota_{\gamma} f^{(2)} & \text{Additional pieces} \\ s \, \xi^{\mu} &= -\frac{1}{2} \mathcal{L}_{\xi} \xi^{\mu} + \gamma^{\mu} & s \, \gamma^{(0)} = \iota_{\gamma} \psi^{(1)} & \text{actors} \\ s \, \gamma^{\mu} &= -\mathcal{L}_{\xi} \gamma^{\mu} & \text{actors} \end{split}$$

We will classify the bosonic BRST invariant backgrounds

1 to 1 correspondence with N = (2,2) SUGRA backgrounds



 \bigcirc

• As usual, we look for bosonic backgrounds

$$\xi^{\mu} = \psi_{\mu\nu} = \psi^{(1)} = 0$$

• As usual, we look for bosonic backgrounds

$$\xi^{\mu} = \psi_{\mu\nu} = \psi^{(1)} = 0$$

AND which solve the equations:

$$s \psi_{\mu\nu} = 0$$
 \Box $\mathcal{L}_{\gamma} g_{\mu\nu} = 0$

The superghost must be a - Killing vector. The interesting case is the 2-sphere



 \bigcirc

• As usual, we look for bosonic backgrounds

$$\xi^{\mu} = \psi_{\mu\nu} = \psi^{(1)} = 0$$

AND which solve the equations:

 $s\psi$

 $s \psi^{(}$

The superghost must be a - Killing vector. The interesting case is the 2-sphere



• As usual, we look for bosonic backgrounds

$$\xi^{\mu} = \psi_{\mu\nu} = \psi^{(1)} = 0$$

AND which solve the equations:

 $s\,\psi^{(1)}$

0

The superghost must be a - Killing vector. The interesting case is the 2-sphere

- Well-known equation: Cartan equivariant cohomology

In this form the problem has been already solved!!

Classifying the solutions

Take S² (round or squashed)







Classifying the solutions

V

Take S² (round or squashed) •



$$\mathcal{L}_{\gamma}g_{\mu\nu} = 0$$
U(1) isometry

$$V^{\mu} = \partial_{\phi}$$
implies

$$\gamma^{\mu} = \epsilon_{\Omega} V^{\mu}$$
equivariant
parameter - ---

$$d\,\gamma^{(0)} - \iota_{\gamma} f^{(2)} = 0$$

The solutions are classified Up to topological equivalences

integer for flux quantization

$$f^{(2)} = \operatorname{vol}_2|_g \cdot n$$
$$D^2 \gamma^{(0)} = \sqrt{g} \epsilon_{\mu\nu} D^{\mu} \gamma^{\nu} + A$$



Classifying the solutions

Take S² (round or squashed)

 \bigcirc



Killing spinor equations in SUGRA

We need to solve the Killing spinor equation (KSE)

gravitino

$$\delta \psi_{\mu} \equiv (D_{\mu} - i\mathcal{A}_{\mu})\zeta + \frac{1}{2}H\Gamma_{\mu}\zeta - \frac{i}{2}G\Gamma_{\mu}\Gamma_{3}\zeta = 0$$
U(1) R-symmetry
gauge field
graviphotons
field-strengths

Killing spinor equations in SUGRA

We need to solve the Killing spinor equation (KSE)

gravitino

$$\delta \psi_{\mu} \equiv (D_{\mu} - i\mathcal{A}_{\mu})\zeta + \frac{1}{2} H \Gamma_{\mu}\zeta - \frac{i}{2} G \Gamma_{\mu}\Gamma_{3}\zeta = 0$$

$$U(1) \text{ R-symmetry}_{qauge field}$$

$$graviphotons_{field-strengths}$$

Solutions known in the literature [Witten, Gomis et al., Benini and Cremonesi, Closset and Cremonesi]

- A-Twist $A_{\mu} = \pm \frac{1}{2}\omega_{\mu}$, H = G = 0
- **Q-background** $A_{\mu} = \pm \frac{1}{2} \omega_{\mu}$, $H = G = \frac{\epsilon_{\Omega}}{2} \sqrt{g} \epsilon_{\mu\nu} D^{\mu} V^{\nu}$
- "No twist"

$$A_{\mu} = 0 , \qquad H = -i , \qquad G = 0$$

From spinors to forms: Fierzing

• A simple Fierz computation tells us:

 $\zeta_a(x)\,\zeta_b^{\dagger}(x) = c_0(x)\,\frac{1}{2}\mathrm{Id}_{ab} + c_{\mu}(x)\frac{1}{2}\Gamma_{ab}^{\mu} + \widetilde{c}_0(x)\,\frac{1}{2}\Gamma_{ab}^3$

 $c_0(x) \equiv \zeta^{\dagger}(x) \zeta(x)$ $c_{\mu}(x) \equiv \zeta^{\dagger}(x) \Gamma_{\mu} \zeta(x)$ $\widetilde{c}_0(x) \equiv \zeta^{\dagger}(x) \Gamma^3 \zeta(x)$

Not all independent

$$c^{\mu}c_{\mu} = c_0^2(x) - \widetilde{c}_0^2(x)$$

The KSE get rewritten:

$$D_{\mu} c_{\nu} + D_{\nu} c_{\mu} = 0$$

$$D_{\mu} \widetilde{c}_{0} = -i H \sqrt{g} \epsilon_{\mu\nu} c^{\nu}$$

$$D_{\mu} c_{\nu} = \sqrt{g} \epsilon_{\mu\nu} (G c_{0} + i H \widetilde{c}_{0})$$

$$D_{\mu} c_{0} = G \sqrt{g} \epsilon_{\mu\nu} c^{\nu}$$



 \bigcirc

From spinors to forms: Fierzing

• A simple Fierz computation tells us:

 $\zeta_a(x)\,\zeta_b^{\dagger}(x) = c_0(x)\,\frac{1}{2}\mathrm{Id}_{ab} + c_{\mu}(x)\frac{1}{2}\Gamma_{ab}^{\mu} + \widetilde{c}_0(x)\,\frac{1}{2}\Gamma_{ab}^3$

 $c_0(x) \equiv \zeta^{\dagger}(x) \zeta(x)$ $c_{\mu}(x) \equiv \zeta^{\dagger}(x) \Gamma_{\mu} \zeta(x)$ $\widetilde{c}_0(x) \equiv \zeta^{\dagger}(x) \Gamma^3 \zeta(x)$

The KSE get rewritten:

$$D_{\mu} c_{\nu} + D_{\nu} c_{\mu} = 0$$

$$D_{\mu} \widetilde{c}_{0} = -i H \sqrt{g} \epsilon_{\mu\nu} c^{\nu}$$

$$D_{\mu} c_{\nu} = \sqrt{g} \epsilon_{\mu\nu} (G c_{0} + i H \widetilde{c}_{0})$$

$$D_{\mu} c_{0} = G \sqrt{g} \epsilon_{\mu\nu} c^{\nu}$$

$$c^{\mu}c_{\mu} = c_0^2(x) - \widetilde{c}_0^2(x)$$

These equations
completely solve
the KSE



The topological system and the KSE are equivalent with the identifications $f \equiv *f^{(2)} = -i H, \quad \gamma^{(0)} = \tilde{c}_0, \quad \gamma^{\mu} = c^{\mu}$



C

C

The topological system and the KSE are equivalent with the identifications $f \equiv *f^{(2)} = -iH, \quad \gamma^{(0)} = \widetilde{c}_0, \quad \gamma^{\mu} = c^{\mu}$

Remark The topological approach has the notion of topological equivalence We have exported this notion also to the SUGRA side. Complete classification also in SUGRA

All the inequivalent backgrounds are parametrized by the fluxes of the graviphotons field-strengths



All the inequivalent backgrounds are parametrized by the fluxes of the graviphotons field-strengths



All the inequivalent backgrounds are parametrized by the fluxes of the graviphotons field-strengths



All the inequivalent backgrounds are parametrized by the fluxes of the graviphotons field-strengths



All the inequivalent backgrounds are parametrized by the fluxes of the graviphotons field-strengths

Infinite new solutions! Inequivalent localizations!







A new gravity theory for finding localizing backgrounds

cohomological classification

Infinite new localizing backgrounds in two dimensions

great extension of the known results

An intrinsic topological reformulation of localization



